

New physics effects on the W -boson mass from a doublet extension of the SM Higgs sector

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based on 2204.05269 and 2103.07484

In collaboration with

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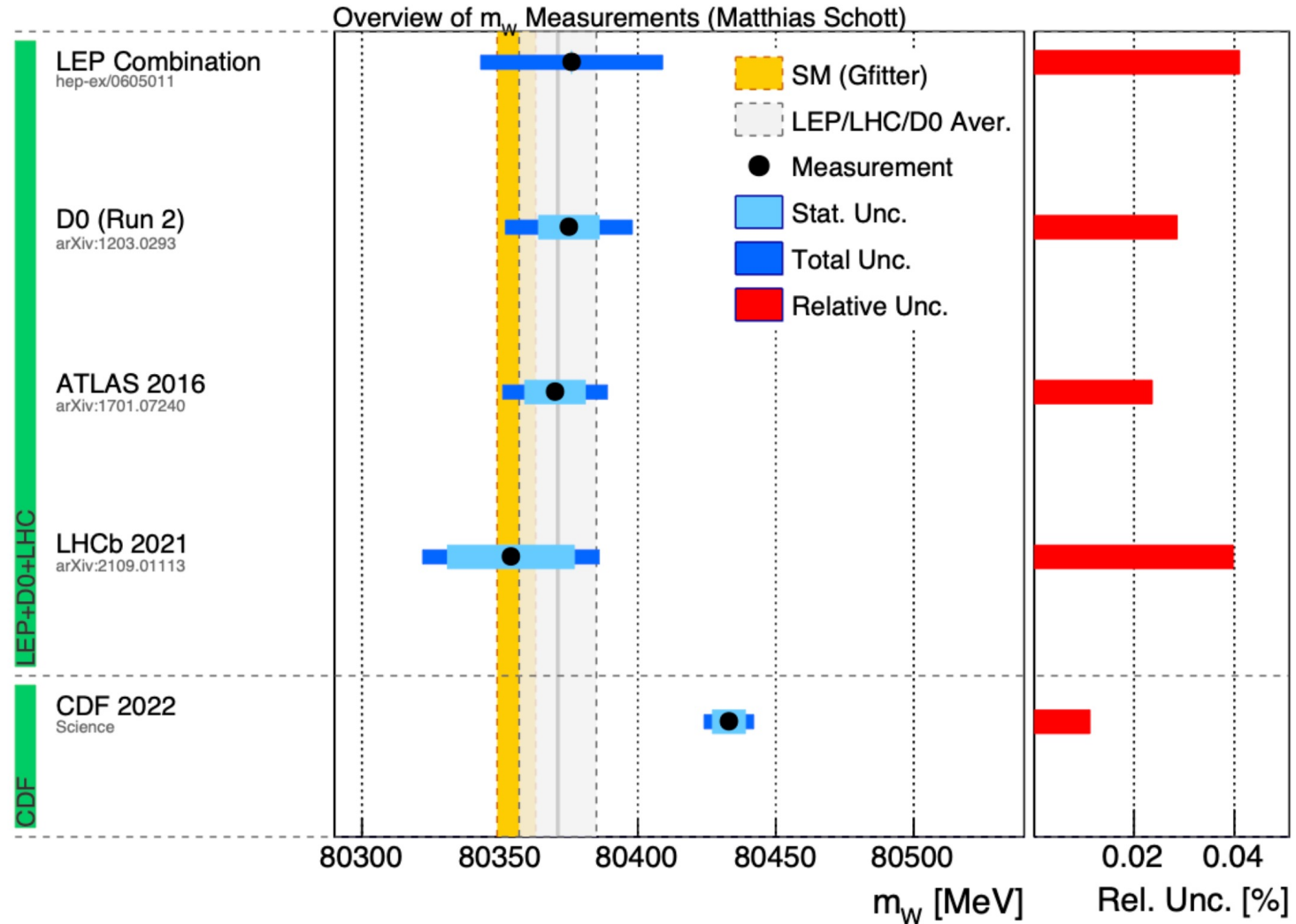
EFI lunch seminar, 5/02/2022

Talk outline

- Introduction
- M_W in the 2HDM
- Implications for LHC searches
- Conclusions

Introduction

M_W measurements



Predicting M_W in the SM

1. Measure Fermi constant G_F via muon decay.

$$\frac{1}{\tau_\mu} = \frac{G_F^2 m_\mu^5}{192\pi^3} F\left(\frac{m_e^2}{m_\mu^2}\right) \left(1 + \frac{3}{5} \frac{m_\mu^2}{M_W^2}\right) (1 + \Delta q),$$

2. Relate M_W to G_F via

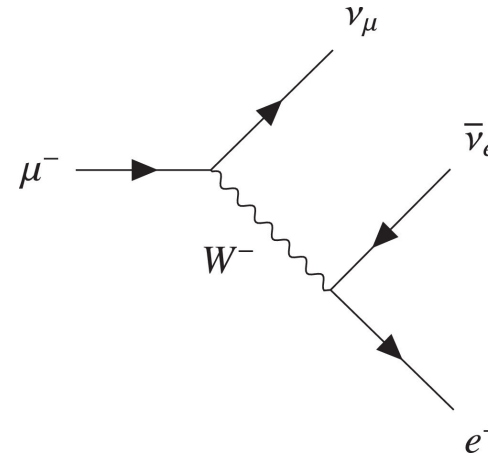
$$M_W^2 \left(1 - \frac{M_W^2}{M_Z^2}\right) = \frac{\pi\alpha_{em}}{\sqrt{2}G_F} (1 + \Delta r).$$

QED corrections

Loop corrections

3. Extract M_W^{SM} via iteration.

→ $M_W^{SM} = 80.357 \pm 6 \text{ GeV}$



Theoretical calculation/prediction is more precise than experimental determination; no large shifts expected.

New physics in M_W ?

How can we shift M_W w.r.t. the SM value? CDF paper has already 82 citations...

- Change relation between M_W and G_F at the **tree level**:
 - Extensions of the SM gauge group: Z' , W' models, ...;
 - Extensions of the SM Higgs sector: $Y = 0$ triplet, $Y = 3/2$ quadruplet, ...;
 - ...
- Change relation between M_W and G_F at the **loop level**:
 - Extensions of the SM Higgs sector: **2nd doublet**, $Y = 1$ triplet, ...;
 - Additional fermions: 4th generation, vector-like fermions, ...;
 - Leptoquarks;
 - SUSY;
 - ...
- Change relation between **muon lifetime** and G_F :
 - Right-handed neutrinos;
 - ...

Focus of this talk: M_W in the 2HDM

M_W in the 2HDM

The 2-Higgs-doublet-model (2HDM)

- Focus first on **2HDM type I** in the alignment limit (similar results expected for other types/models).

$$V_{2\text{HDM}}(\Phi_1, \Phi_2) = m_{11}^2 \Phi_1^\dagger \Phi_1 + m_{22}^2 \Phi_2^\dagger \Phi_2 - m_{12}^2 (\Phi_1^\dagger \Phi_2 + \Phi_2^\dagger \Phi_1) \\ + \frac{1}{2} \lambda_1 (\Phi_1^\dagger \Phi_1)^2 + \frac{1}{2} \lambda_2 (\Phi_2^\dagger \Phi_2)^2 + \lambda_3 (\Phi_1^\dagger \Phi_1) (\Phi_2^\dagger \Phi_2) + \lambda_4 (\Phi_1^\dagger \Phi_2) (\Phi_2^\dagger \Phi_1) + \frac{1}{2} \lambda_5 \left((\Phi_1^\dagger \Phi_2)^2 + (\Phi_2^\dagger \Phi_1)^2 \right).$$

- 2 Higgs doublets \rightarrow 5 physical Higgs bosons: CP-even h, H ; CP-odd A ; charged H^\pm .

SM-like Higgs boson 

- Convenient choice of input parameters: $m_h, m_H, m_A, m_{H^\pm}, t_\beta = \tan\beta \equiv v_2/v_1, M = m_{12}/(s_\beta c_\beta)$.
- Higgs masses are related to λ s via: $m_i^2 = M^2 + \sum_j c_{ij} \lambda_j v^2$

M_W in the 2HDM

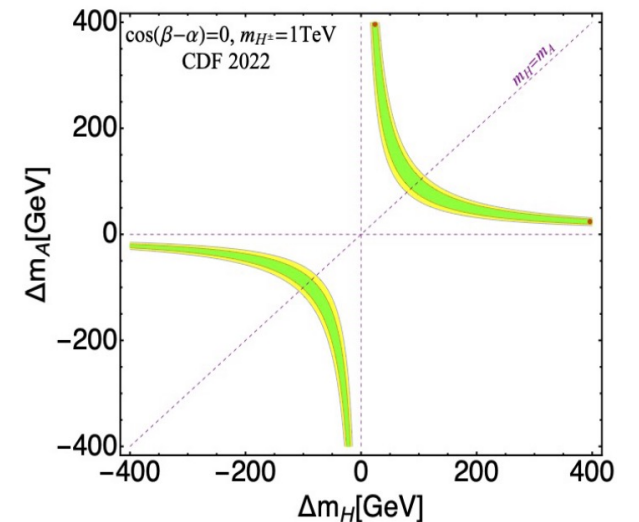
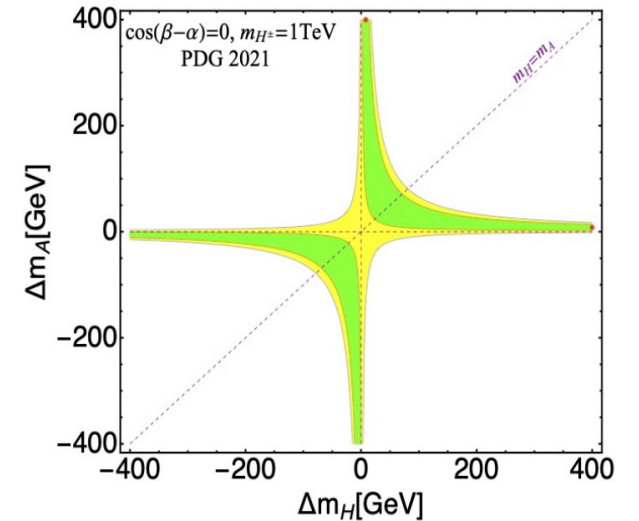
- Non-SM one-loop corrections to muon decay result in

$$\Delta\rho_{\text{non-SM}}^{(1)} = \frac{\alpha}{16\pi^2 s_W^2 M_W^2} \left\{ \frac{m_A^2 m_H^2}{m_A^2 - m_H^2} \ln \frac{m_A^2}{m_H^2} - \frac{m_A^2 m_{H^\pm}^2}{m_A^2 - m_{H^\pm}^2} \ln \frac{m_A^2}{m_{H^\pm}^2} - \frac{m_H^2 m_{H^\pm}^2}{m_H^2 - m_{H^\pm}^2} \ln \frac{m_H^2}{m_{H^\pm}^2} + m_{H^\pm}^2 \right\},$$

$$\Delta M_W \simeq \frac{1}{2} M_W \frac{c_W^2}{c_W^2 - s_W^2} \Delta\rho.$$

- Restoration of custodial symmetry (i.e. $\Delta M_W \simeq 0$) if $m_A = m_{H^\pm}$ or $m_H = m_{H^\pm}$.

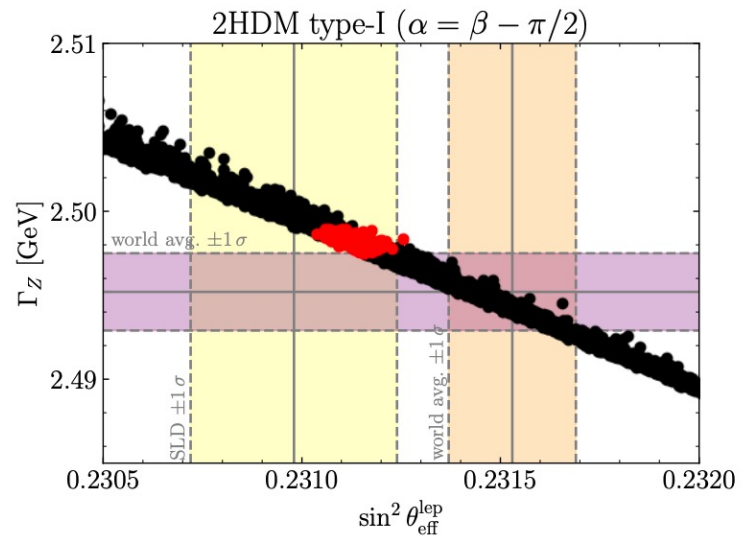
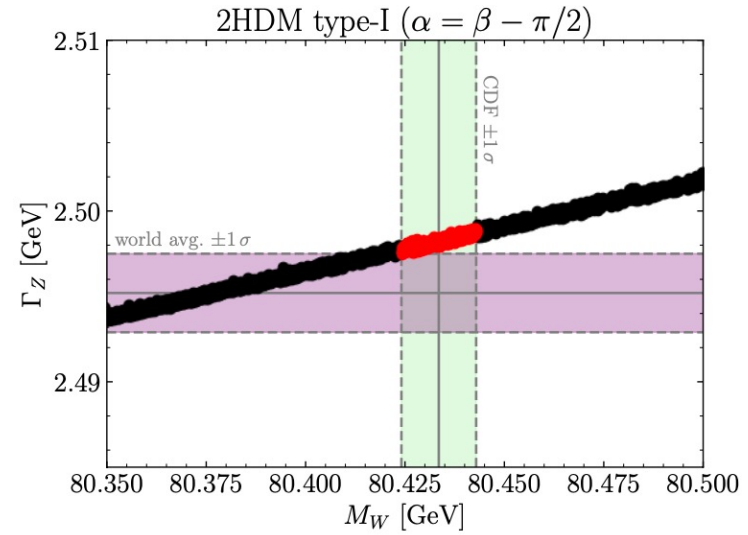
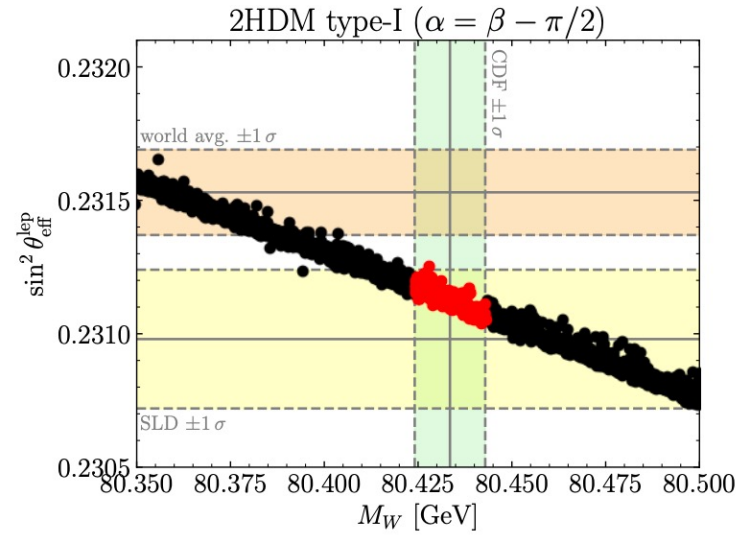
Is this compatible with other constraints?



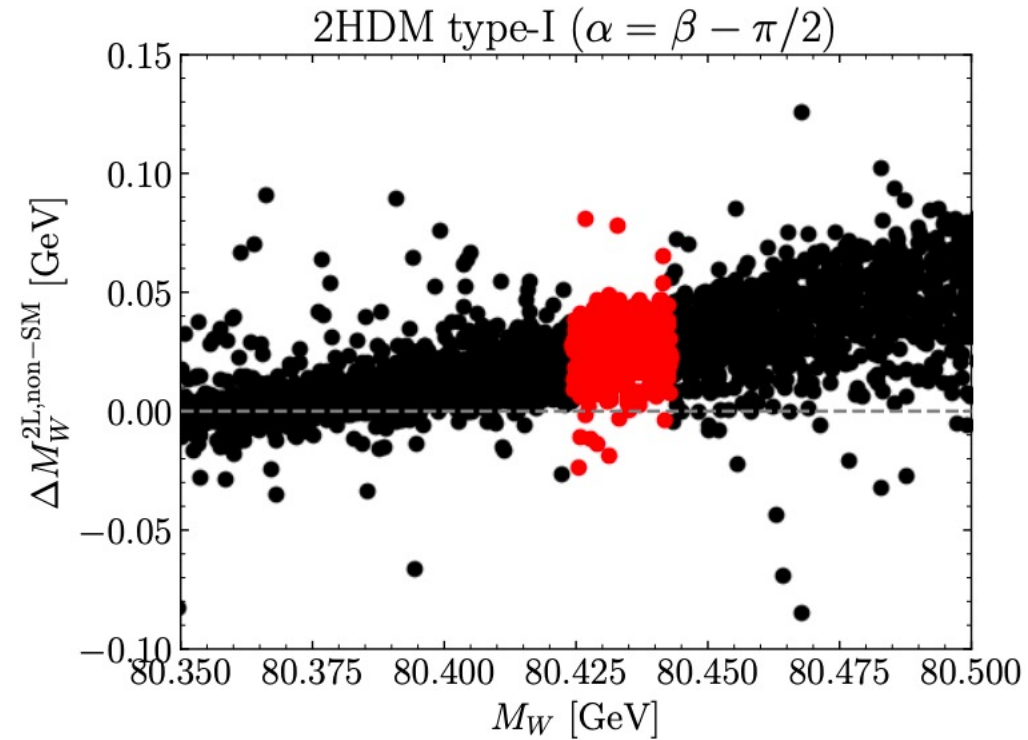
2HDM parameter scan

- We checked for:
 - Vacuum stability and boundedness-from-below.
 - NLO perturbative unitarity.
 - SM-like Higgs measurements via HiggsSignals.
 - Direct searches for BSM scalars via HiggsBounds.
 - b-physics constraints.
- Most constraints checked using ScannerS.
- Loop corrections affecting M_W will also affect other electroweak precision observables.
- Calculate electroweak precision observables — M_W , effective weak mixing angle $\sin^2 \theta_{\text{eff}}^{\text{lep}}$, Z decay width — at the 2L level using THDM_EWPOS [Hessenberger,Hollik,1607.04610].

Scan results — EWPOs

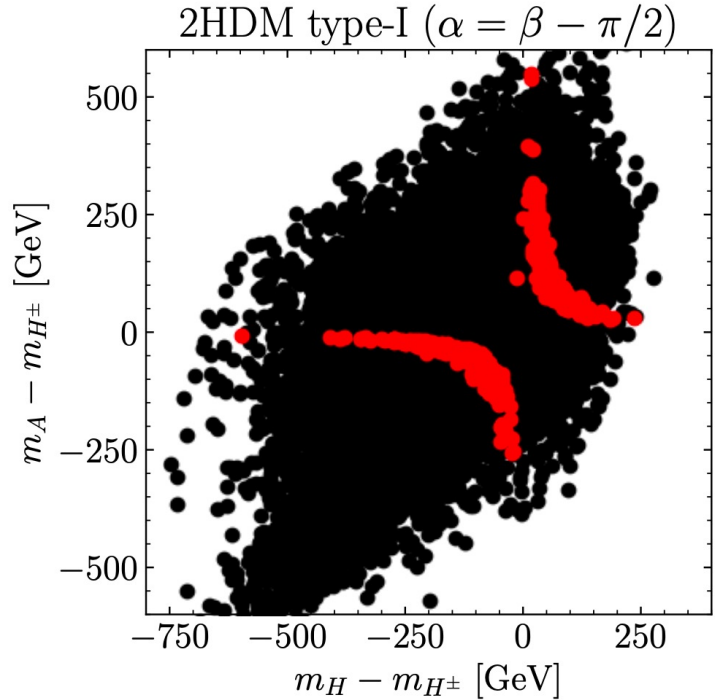


Scan results — 1L vs. 2L



Two-loop non-SM corrections important for precise results.

Scan results — mass hierarchies



- CDF M_W measurement favors

$$m_H, m_A < m_{H^\pm}, \quad m_{H^\pm} < m_H, m_A$$

- CDF M_W measurement disfavors

$$m_{H,A} < m_{H^\pm} < m_{A,H}, \quad m_H \sim m_A \sim m_{H^\pm}$$

How can we test this scenario at the LHC?

m_H [GeV]	m_A [GeV]	m_{H^\pm} [GeV]	$\tan \beta$ —	M^2 [GeV ²]	M_W [GeV] (non-SM@1L)	M_W [GeV] (non-SM@2L)	$\sin^2 \theta_{\text{eff}}^{\text{lep}}$ —	Γ_Z [GeV]
853.813	928.352	809.047	1.206	444.166×10^3	80.4001	80.4337	0.23113	2.4981
351.962	751.498	762.911	1.255	55.451×10^3	80.3990	80.4339	0.23109	2.4979

Implications for LHC searches

Bosonic Higgs couplings

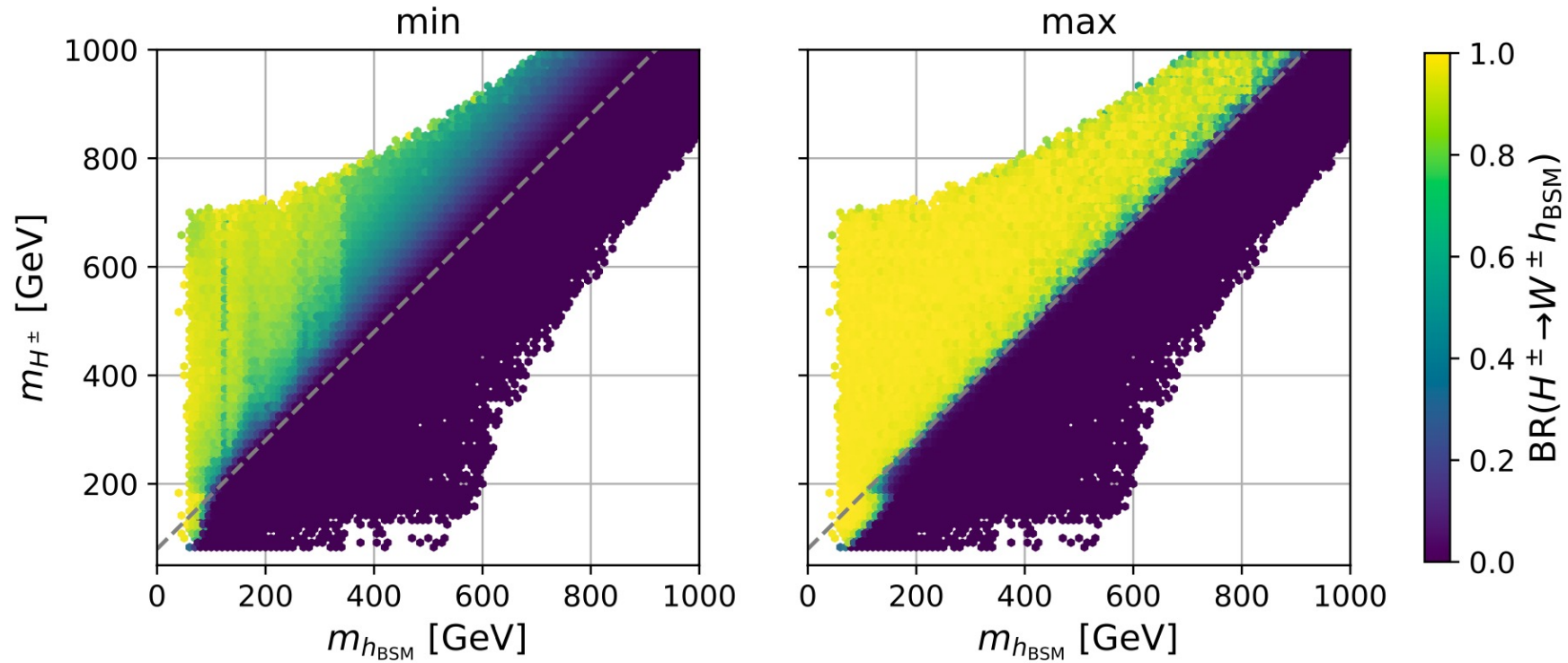
- Fermionic decay modes of BSM Higgs bosons comparably well explored.
- But also bosonic decay modes are important (α : mixing angle between CP-even Higgs bosons):

$$g(ZAh) \propto c_{\beta-\alpha}, \quad g(ZAH) \propto s_{\beta-\alpha}$$

$$g(H^\pm W^\mp h) \propto c_{\beta-\alpha}, \quad g(H^\pm W^\mp H) \propto s_{\beta-\alpha}, \quad g(H^\pm W^\mp A) = -\frac{g}{2}$$

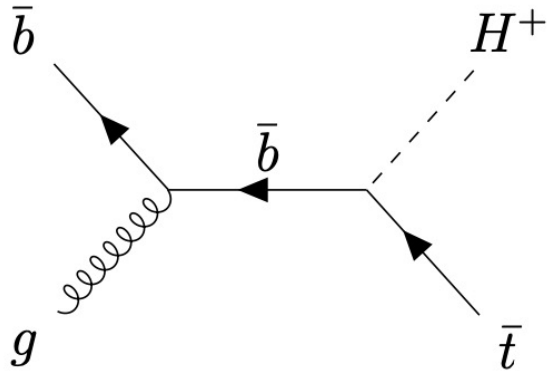
- In the alignment limit both the ZAH , $H^\pm W^\mp H$, and $H^\pm W^\mp A$ couplings are maximized.
- Mass hierarchies favored by CDF M_W measurement allow for $H^\pm \rightarrow W^\pm A/H$ or $A/H \rightarrow W^\pm H^\mp$ decays.
- Focus here on $H^\pm \rightarrow W^\pm A/H$ decays.

Bosonic Higgs couplings

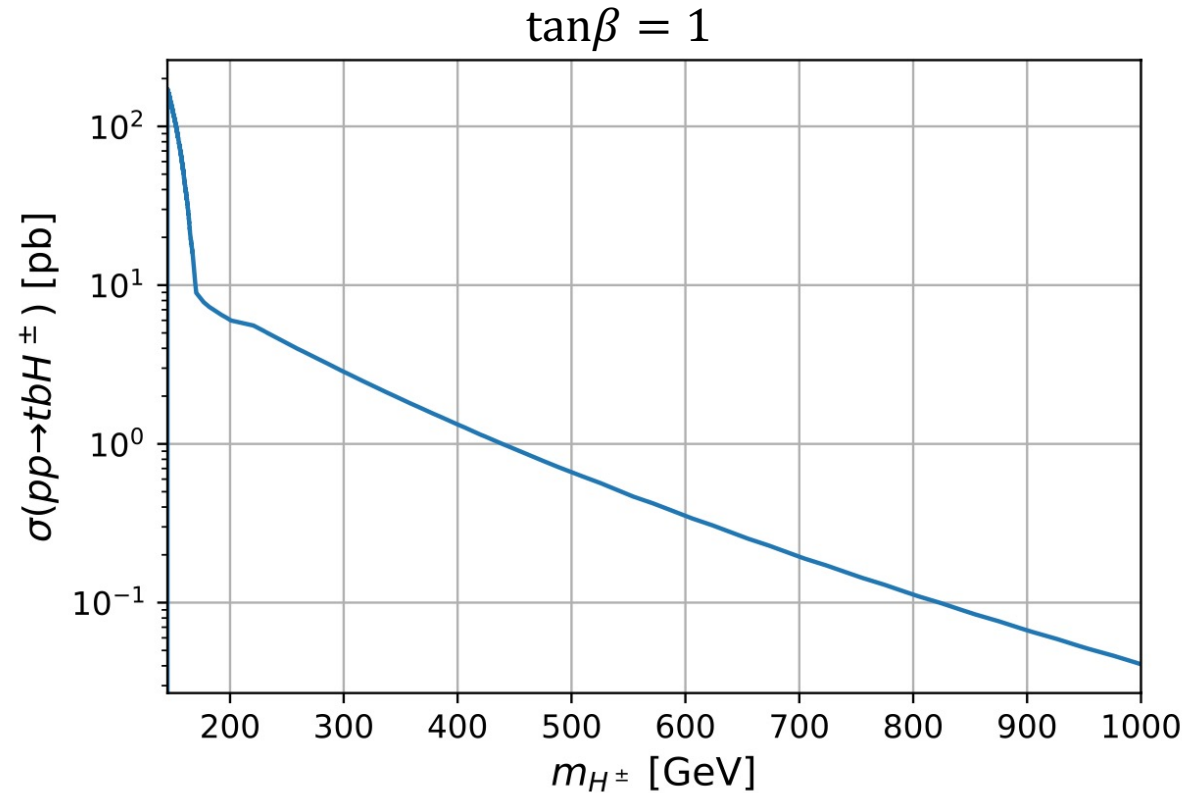


Very large $\text{BR}(H^\pm \rightarrow W^\pm A/H)$ expected if decay kinematically allowed.

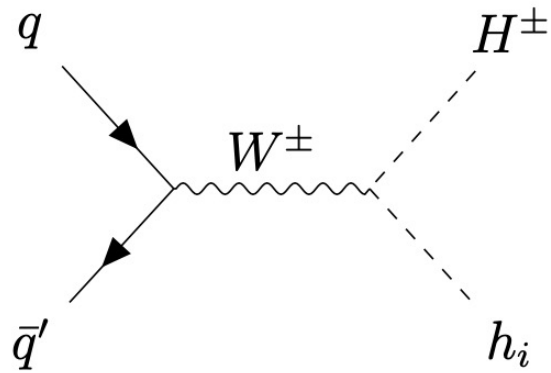
H^\pm production at the LHC — $pp \rightarrow H^\pm tb$



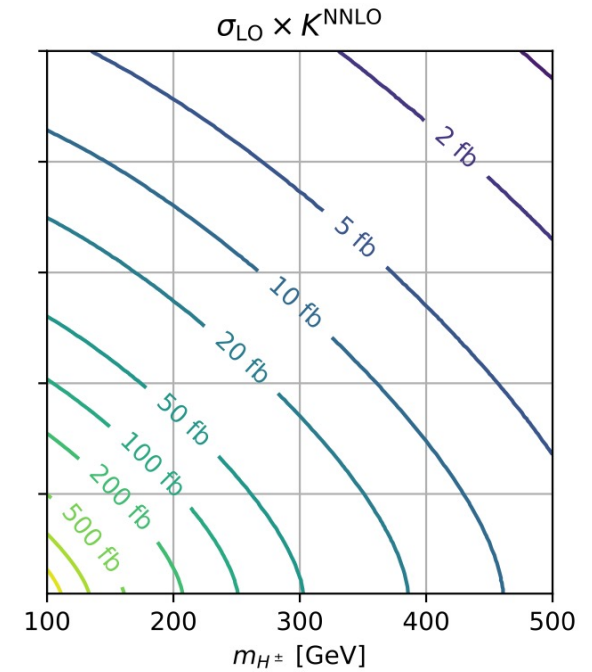
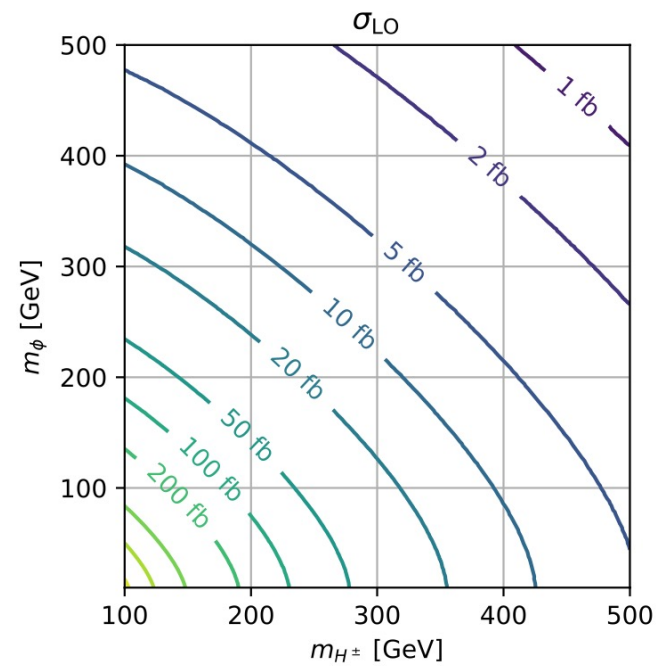
Suppressed by $\propto 1/\tan^2 \beta$.



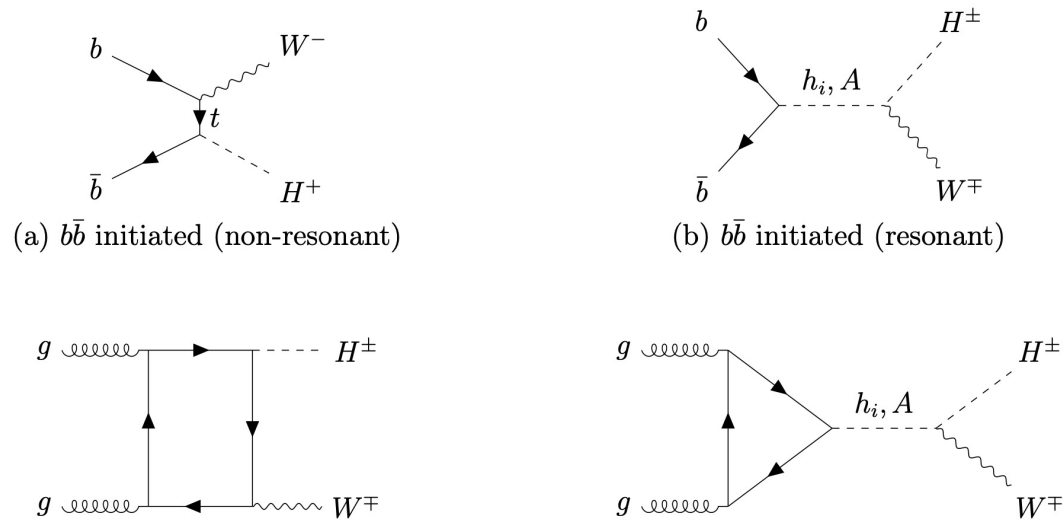
H^\pm production at the LHC — $pp \rightarrow H^\pm H/A$



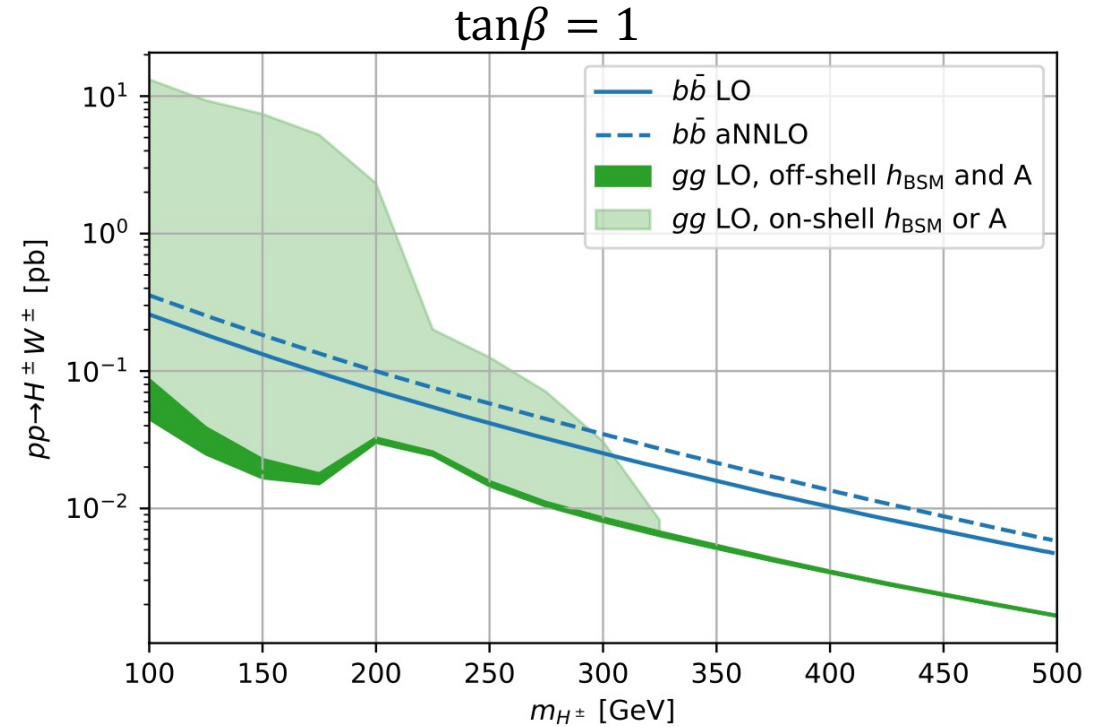
Not suppressed by $\tan\beta$.



H^\pm production at the LHC — $pp \rightarrow H^\pm W^\mp$



Dependence $\tan\beta$ varies between subchannels.



Existing charged Higgs boson searches

Production process	Higgs decay	Final state	# of exp. searches
$pp \rightarrow H^\pm tb$	$H^\pm \rightarrow \tau \nu_\tau$	$tb(\tau \nu_\tau)$	5
$pp \rightarrow H^\pm tb$	$H^\pm \rightarrow tb$	$tbtb$	4
$pp \rightarrow tt, t \rightarrow H^\pm b$	$H^\pm \rightarrow cb$	$tbc b$	1
$pp \rightarrow tt, t \rightarrow H^\pm b$	$H^\pm \rightarrow cs$	$tbc s$	2
$pp \rightarrow H^\pm qq'$ (VBF)	$H^\pm \rightarrow W^\pm Z$	$W^\pm Z qq'$	3
$pp \rightarrow tt, t \rightarrow H^\pm b$	$H^\pm \rightarrow W^\pm A$	$tbW^\pm \mu^+ \mu^-$	2
$pp \rightarrow H \rightarrow H^\pm W^\mp$	$H^\pm \rightarrow W^\pm h$	$W^\pm W^\mp bb$	1

Do these existing searches (as well as other constraint) still allow for large $H^\pm \rightarrow W^\pm H/A$ signals?

Benchmark scenarios for bosonic charged Higgs searches

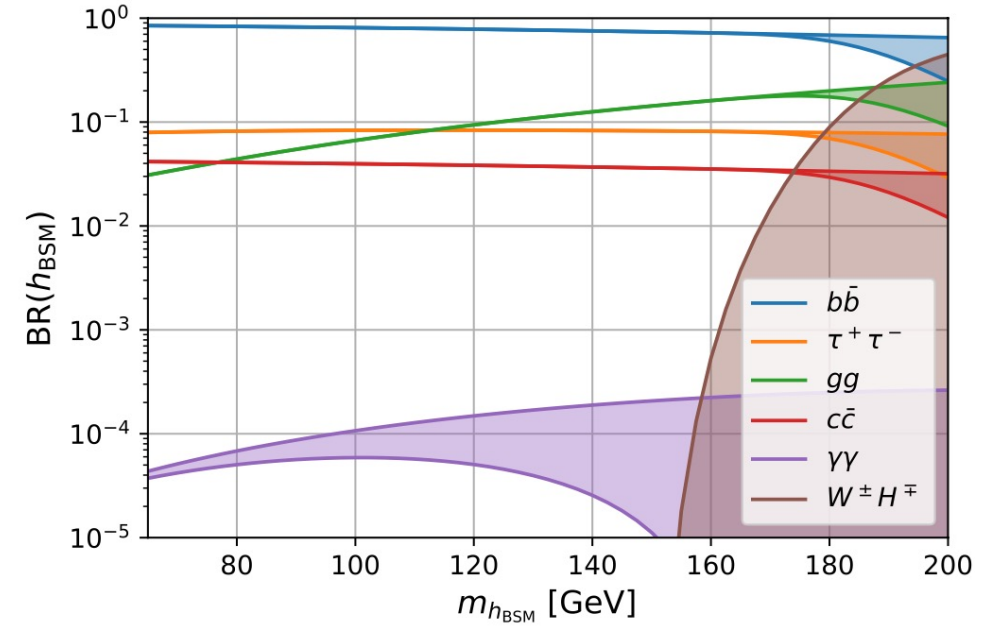
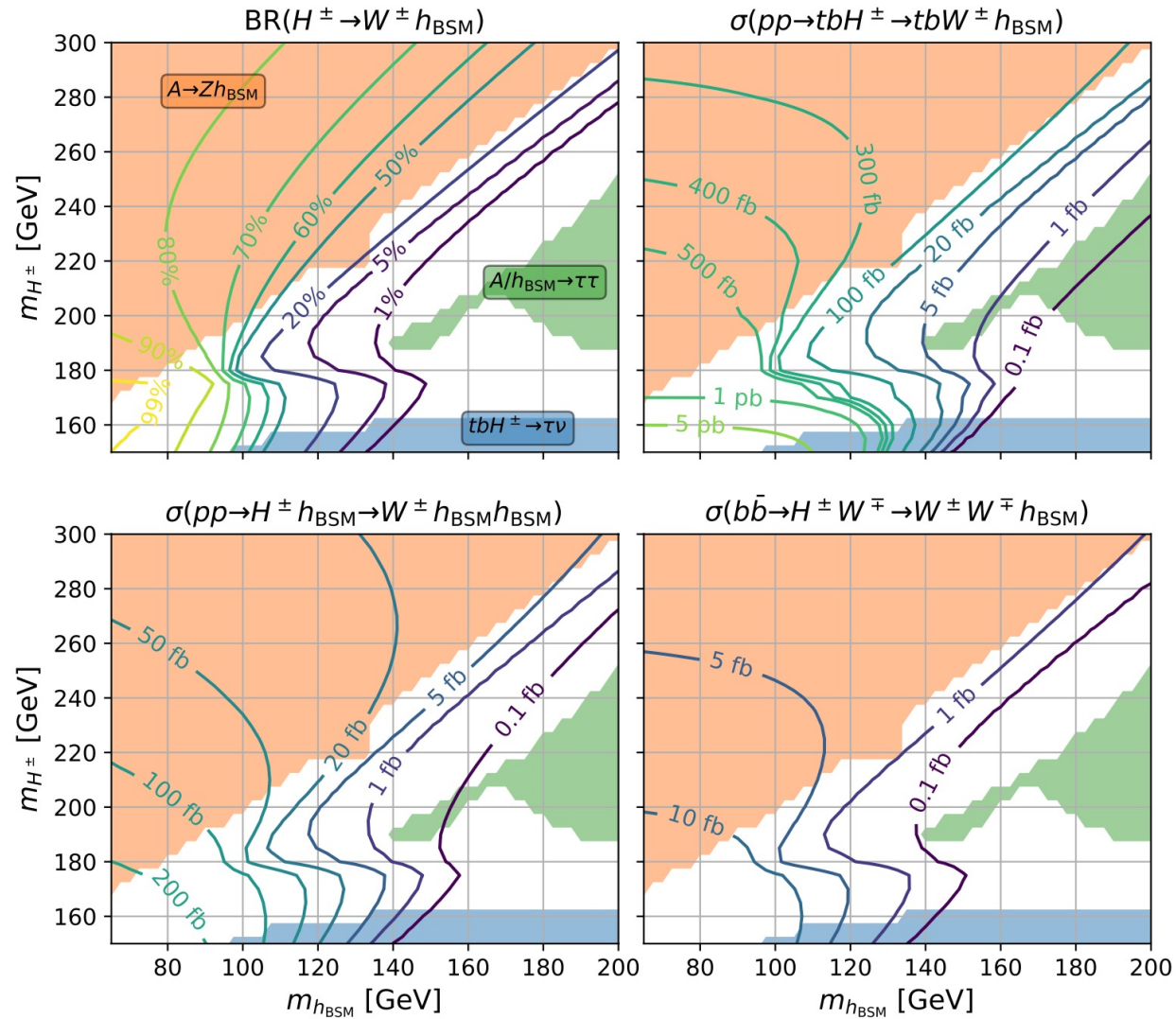
Concept: define $(m_{H,A}, m_{H^\pm})$ planes with sizeable unconstrained $H^\pm \rightarrow W^\pm H/A$ signal.

- Take into account all experimental and theoretical constraints.
- Benchmark scenarios cover different mass ranges and different decay modes of H/A .

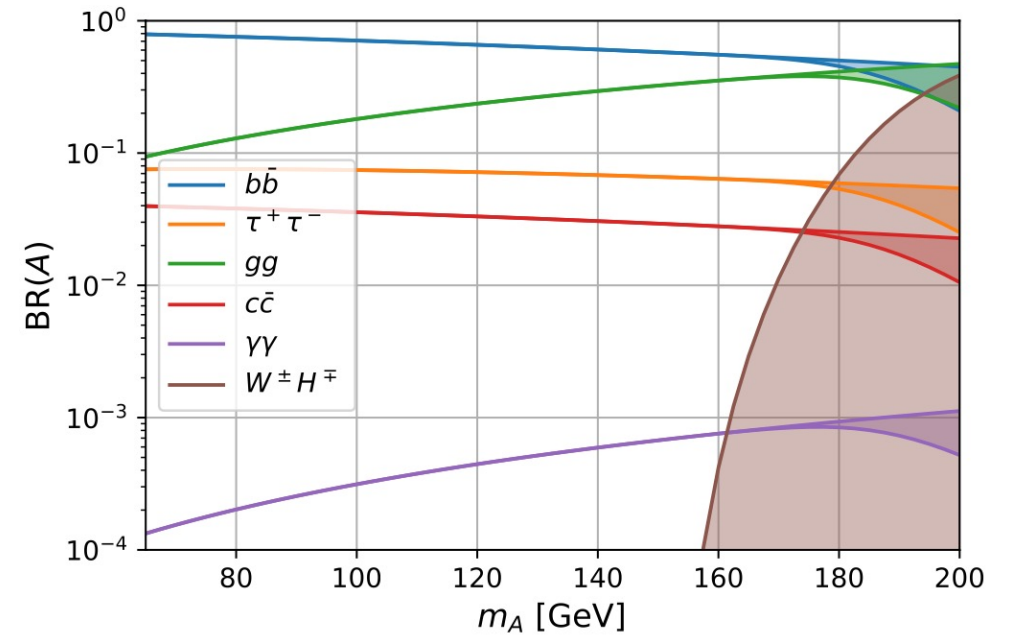
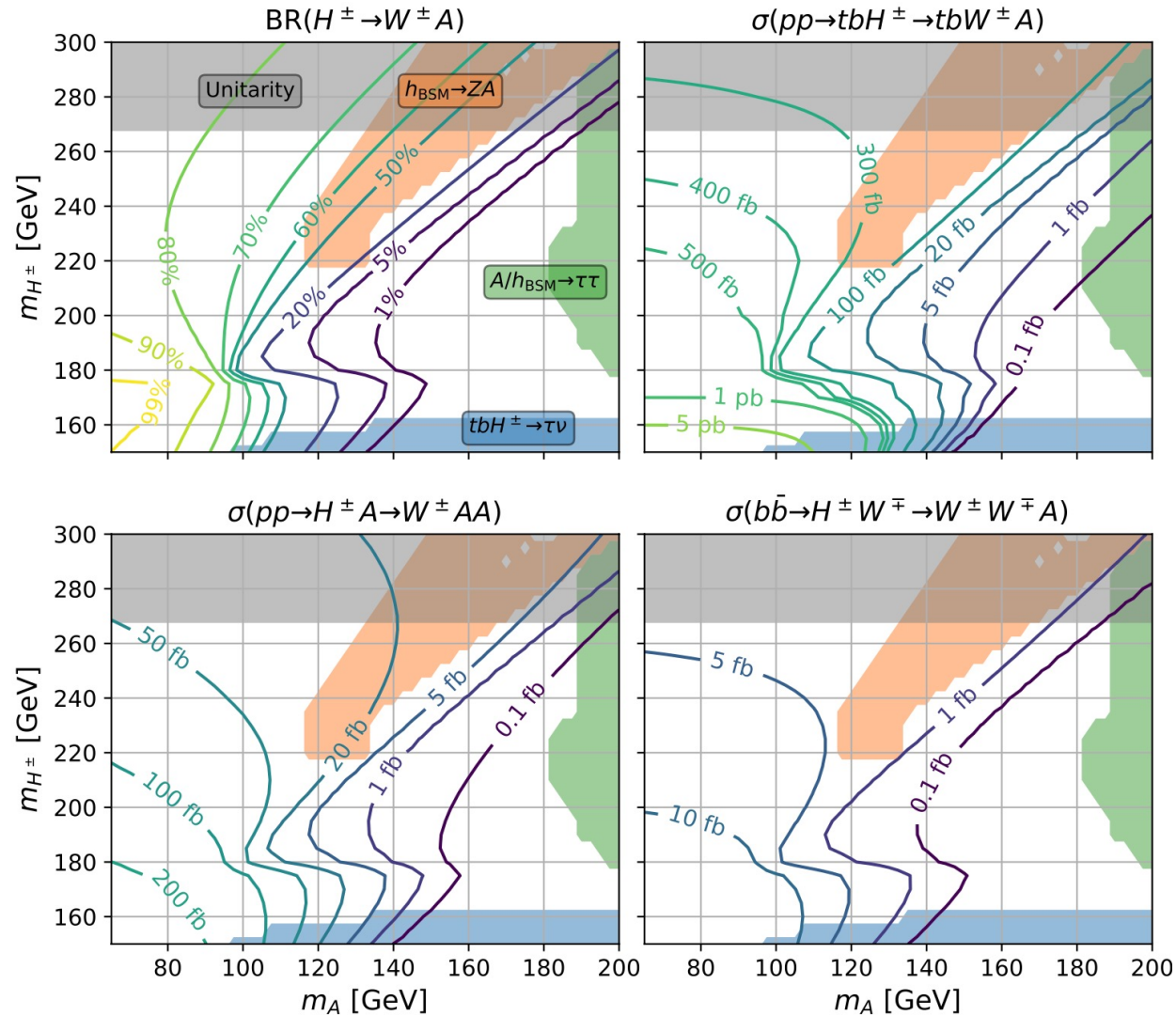
Benchmark scenarios:

- $\text{cH}(Wh_{\text{BSM}})$ scenario \rightarrow exact alignment, $m_A = m_{H^\pm}$, type I.
- $\text{cH}(WA)$ scenario \rightarrow same as $\text{cH}(Wh_{\text{BSM}})$ scenario, but $h_{\text{BSM}} \leftrightarrow A$.
- $\text{cH}(Wh_{\text{BSM}}^{\text{fpob}})$ scenario \rightarrow approximate alignment, fermiophobic h_{BSM} , $m_A = m_{H^\pm}$, type I.
- $\text{cH}(Wh_{\text{BSM}}^{\text{light}})$ scenario \rightarrow approximate alignment, light BSM Higgs, $m_A = m_{H^\pm}$, type I.
- $\text{cH}(Wh_{\text{BSM}}^{\text{lphil}})$ scenario \rightarrow same as $\text{cH}(Wh_{\text{BSM}}^{\text{light}})$ scenario but in lepton-specific 2HDM.

cH(Wh_{BSM}) scenario



cH(WA) scenario



Conclusions

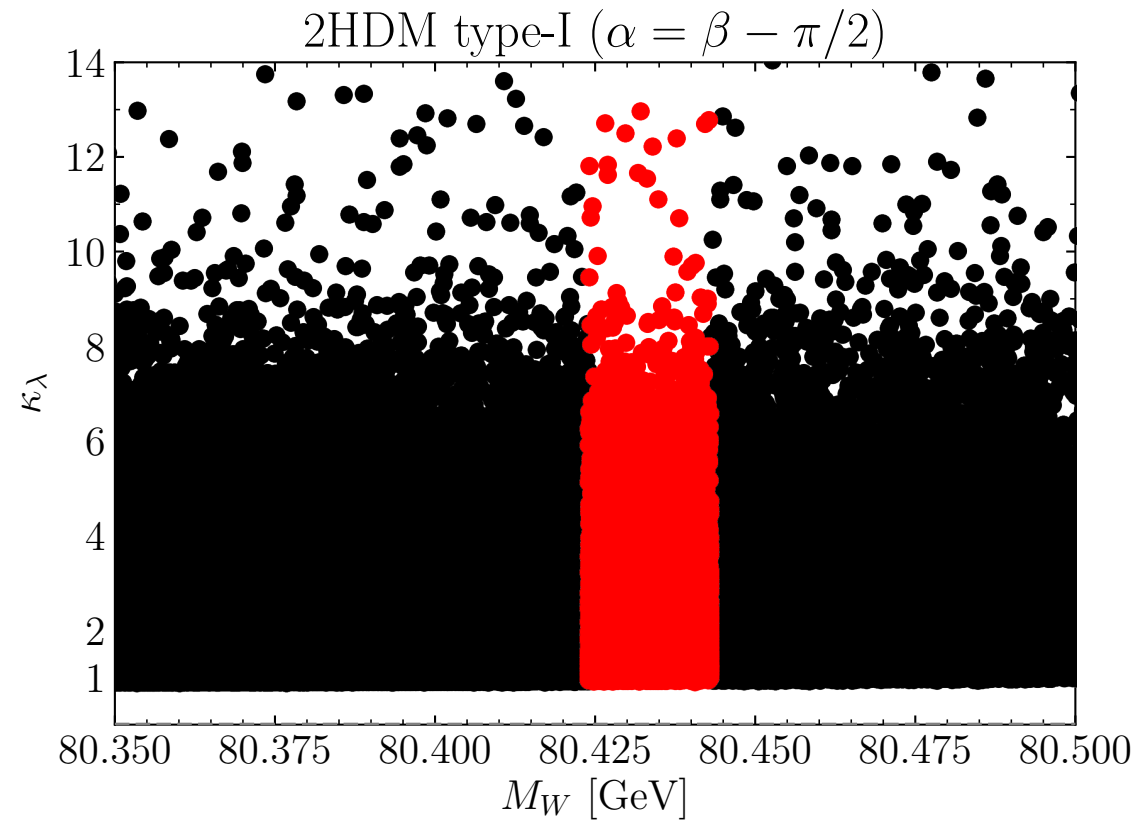
Conclusions

- Many BSM models can accommodate M_W measurement by CDF.
- Should check for every model that this is still consistent with other constraints.
- We showed that the 2HDM can explain the CDF M_W measurement while still satisfying other constraints.
- CDF M_W measurements favours large mass splitting between neutral and charged Higgs bosons.
- The 2HDM explanation of the CDF M_W measurement could be probed by LHC searches for $H^\pm \rightarrow W^\pm H/A$ or $H/A \rightarrow H^\pm W^\mp$.

Thanks for your attention!

Appendix

M_W in the 2HDM – correlation with κ_λ ?



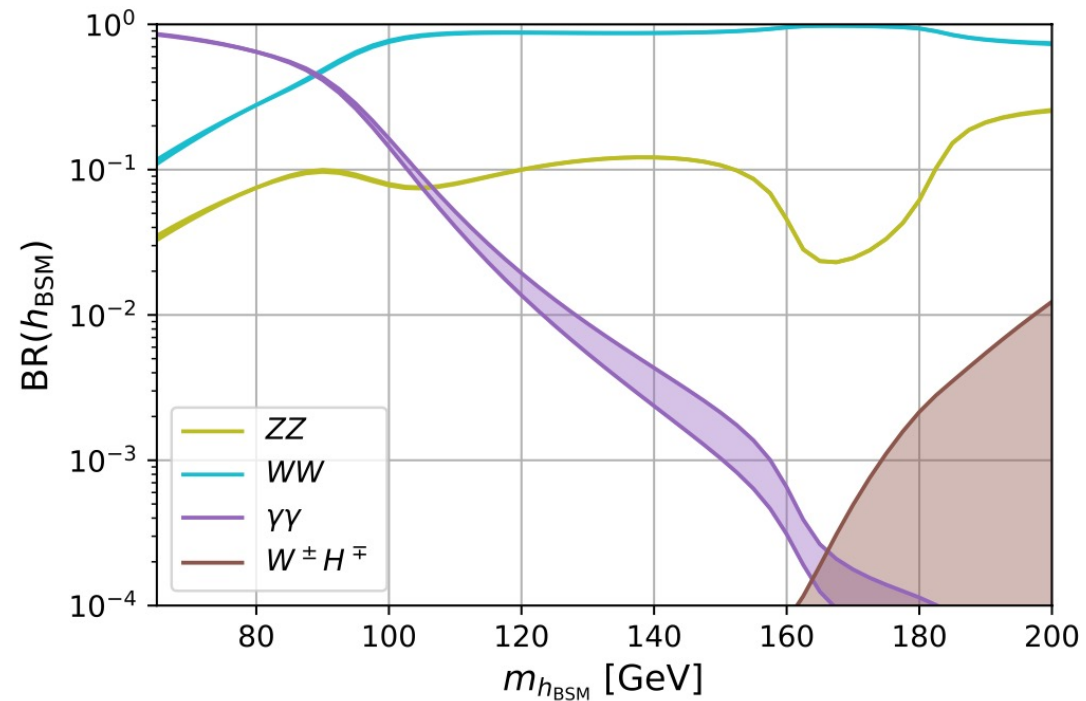
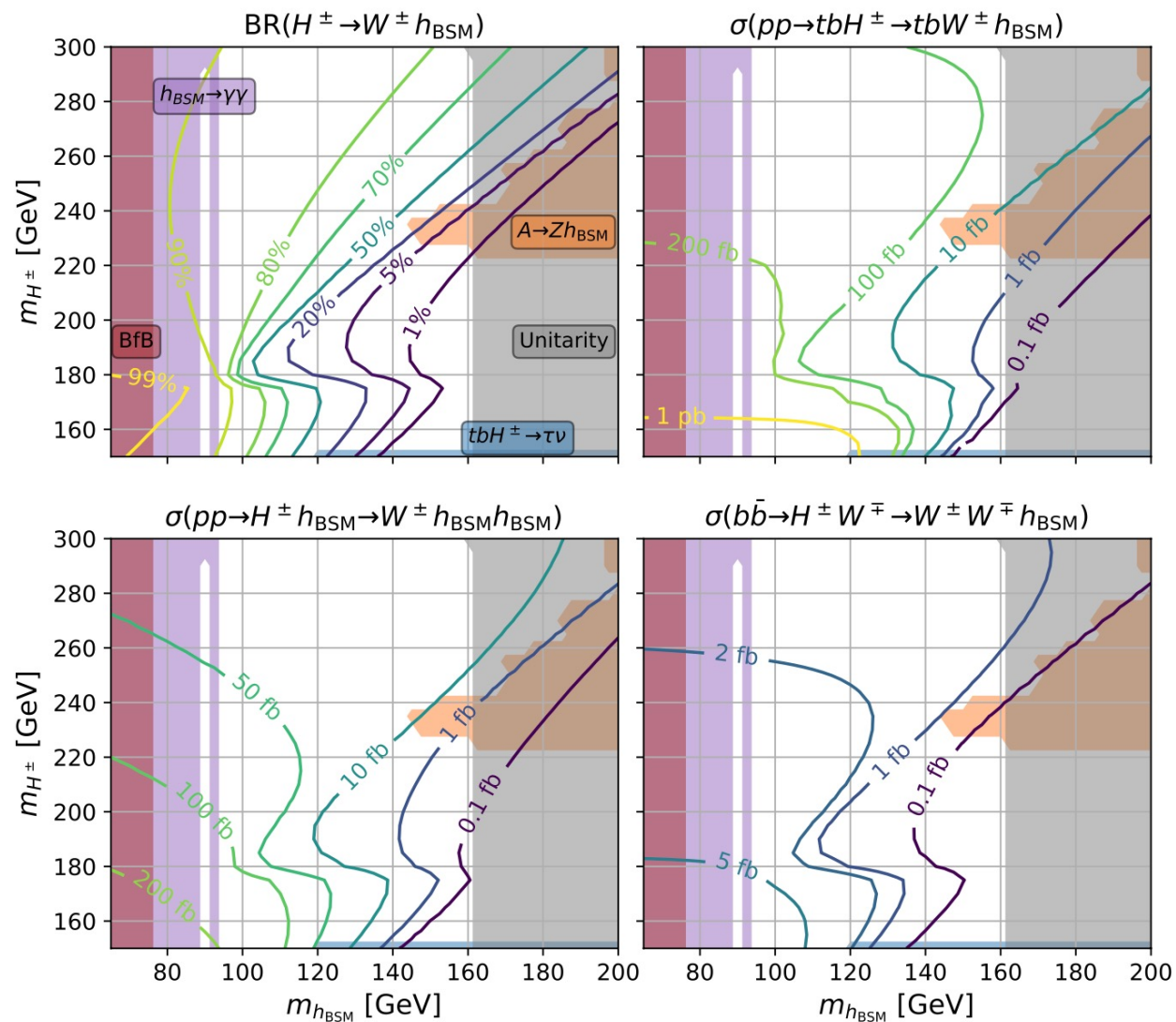
2HDM types

Type	u_R	d_R	l_R	λ_{uu}	λ_{dd}	λ_{ll}
I	+	+	+	$\cot \beta$	$\cot \beta$	$\cot \beta$
II	+	-	-	$\cot \beta$	$-\tan \beta$	$-\tan \beta$
Flipped	+	-	+	$\cot \beta$	$-\tan \beta$	$\cot \beta$
Lepton-specific	+	+	-	$\cot \beta$	$\cot \beta$	$-\tan \beta$

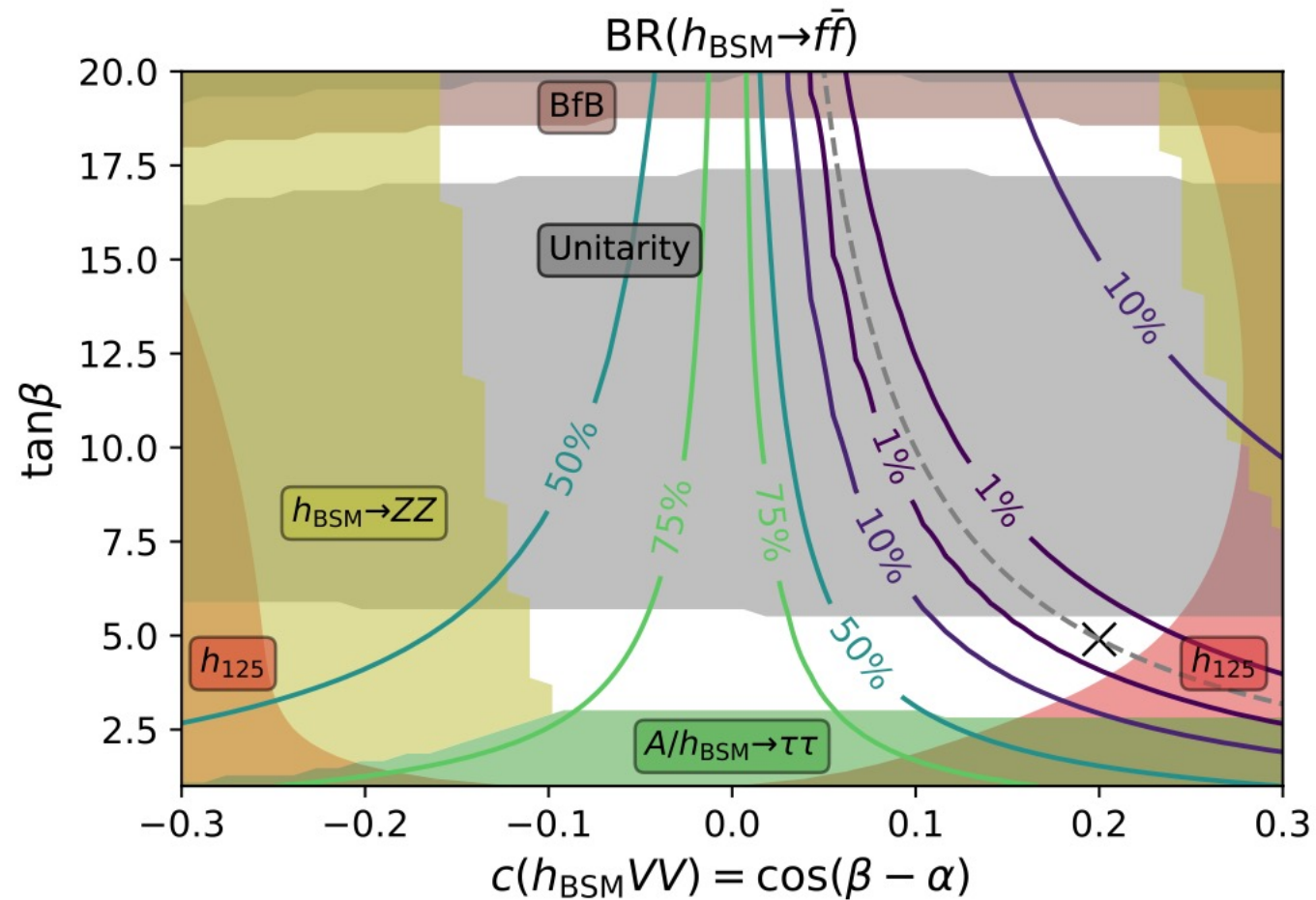
Benchmark scenarios

	$m_{h_{125}}$ [GeV]	m_{H^\pm} [GeV]	$m_{h_{\text{BSM}}}$ [GeV]	m_A [GeV]	$c(h_{\text{BSM}}VV)$	$\tan \beta$	m_{12}^2 [GeV ²]
cH(Wh_{BSM})	125.09	150–300	65–200	m_{H^\pm}	0	3	500
cH(WA)			m_{H^\pm}	65–200			5000
cH($Wh_{\text{BSM}}^{\text{fphob}}$)	125.09	150–300	65–200	m_{H^\pm}	0.2	Eq. (26)	1200
cH($Wh_{\text{BSM}}^{\text{light}}$)		100–300	10–62.5	m_{H^\pm}	-0.062	16.6	Eq. (33)
cH($Wh_{\text{BSM}}^{\ell\text{phil}}$)		same as cH($Wh_{\text{BSM}}^{\text{light}}$) but in the lepton-specific 2HDM					

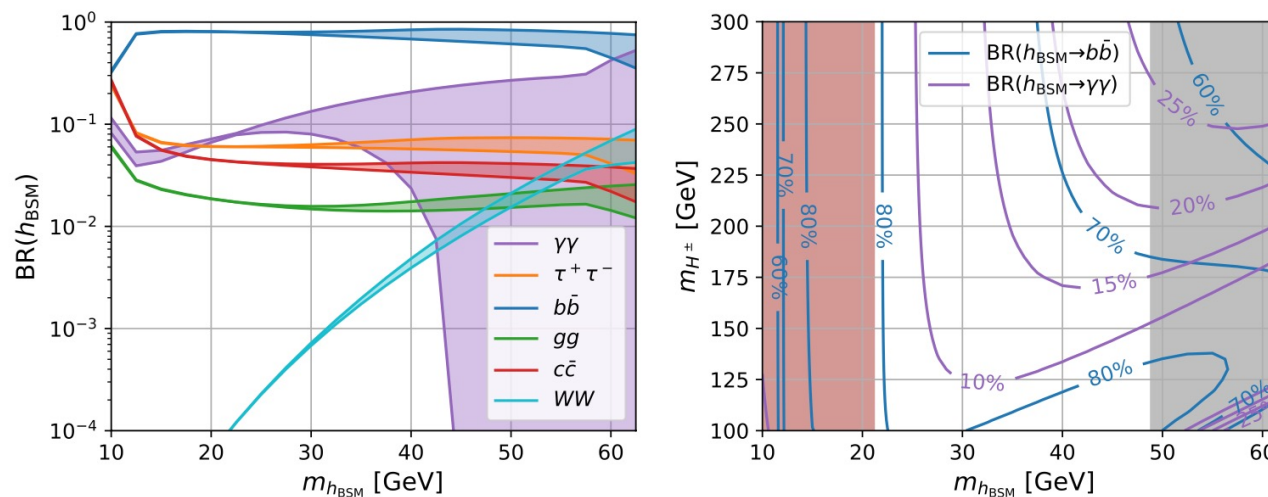
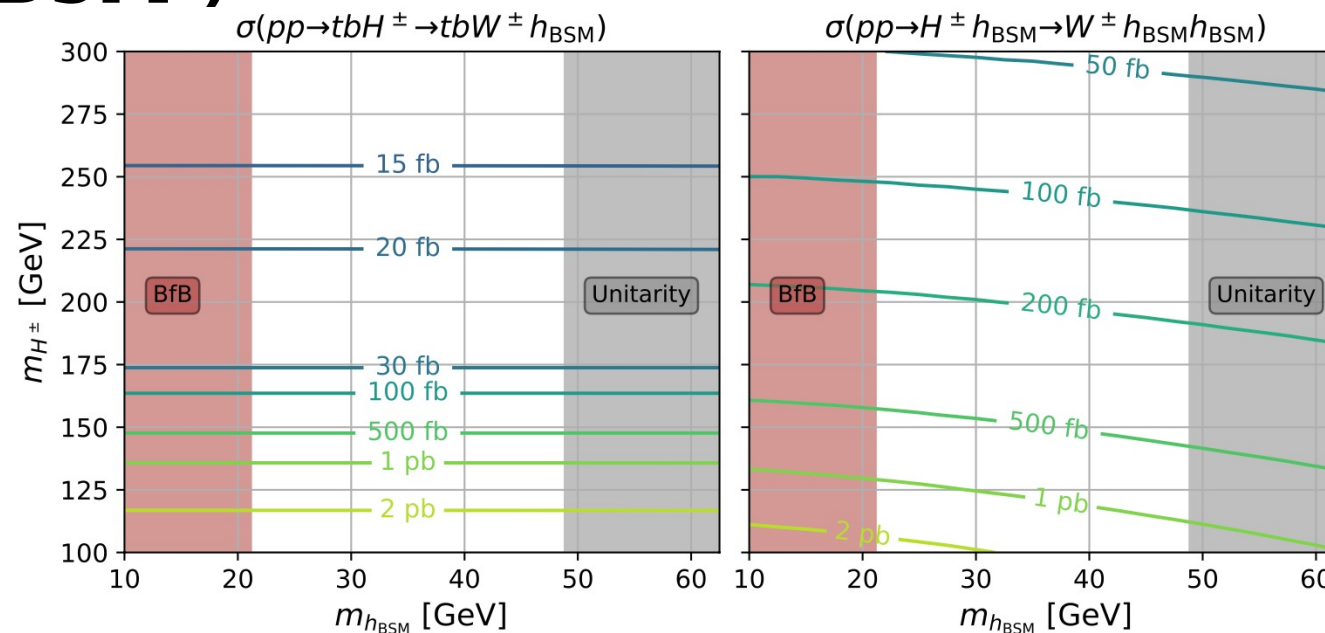
cH ($Wh_{\text{BSM}}^{\text{fphob}}$) scenario



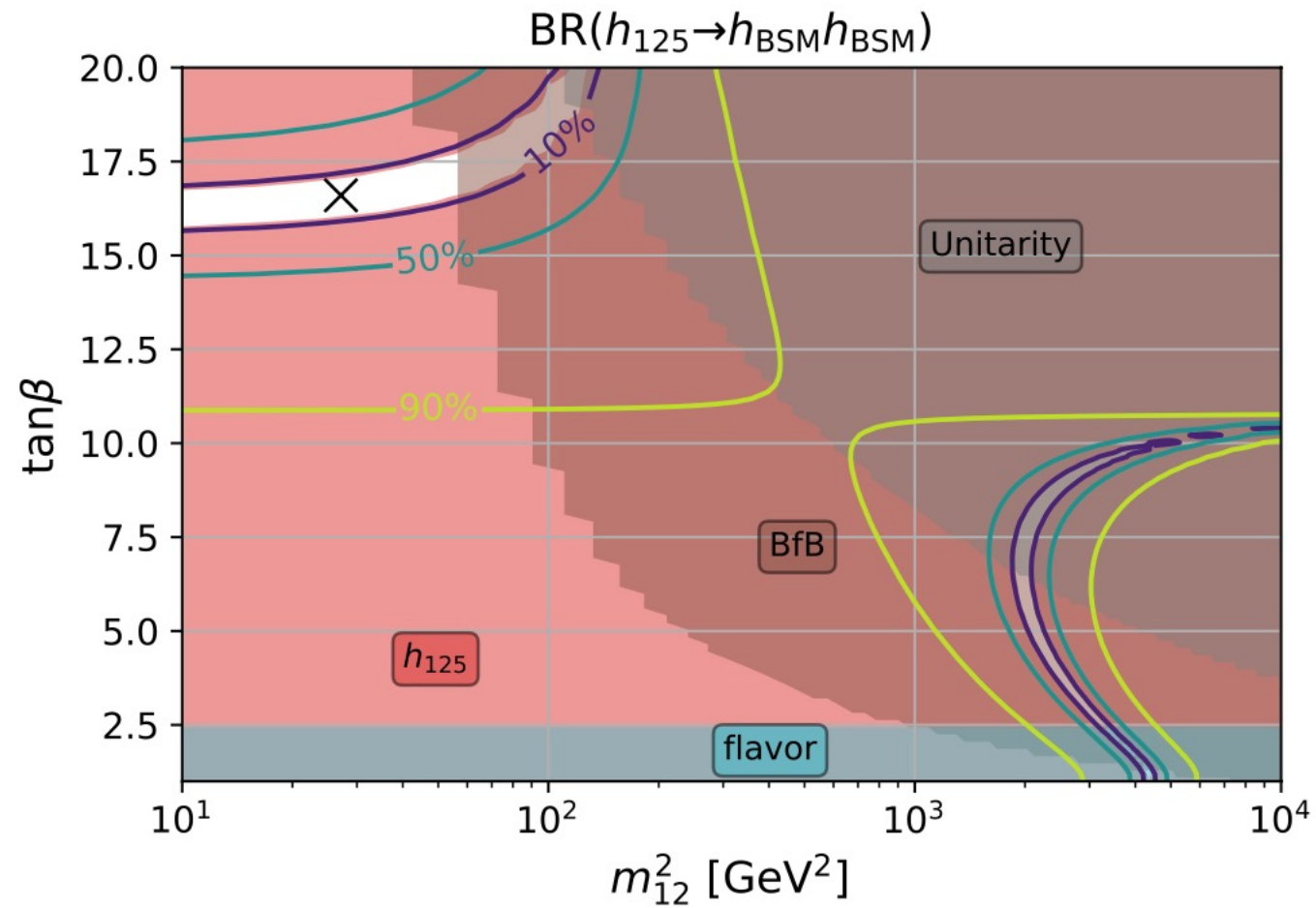
cH ($W h_{\text{BSM}}^{\text{phob}}$) scenario — tuning



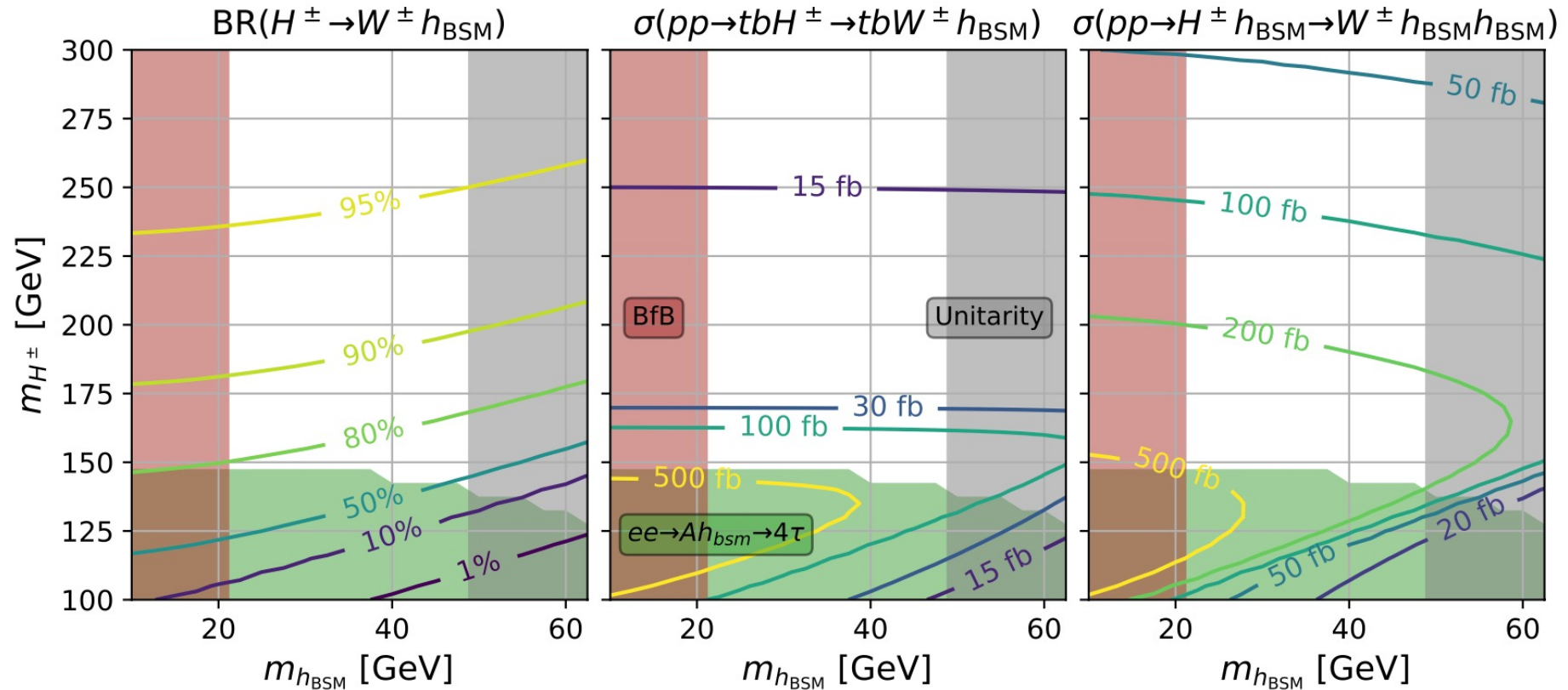
cH ($W h_{\text{BSM}}^{\text{light}}$) scenario



cH ($Wh_{\text{BSM}}^{\text{light}}$) scenario— tuning



cH ($W h_{\text{BSM}}^{\text{phil}}$) scenario



$$\text{BR}(h_{\text{BSM}} \rightarrow \tau^+ \tau^-) \simeq 1$$