Constraining CP violation in the Higgs–fermion interactions

Henning Bahl



Fermilab theory seminar, October 26th 2022

Talk based on

• 2007.08542:

Indirect CP probes of the Higgs-top-quark interaction: current LHC constraints and future opportunities; in collaboration with P. Bechtle, S. Heinemeyer, J. Katzy, T. Klingl, K. Peters, M. Saimpert, T. Stefaniak, G. Weiglein.

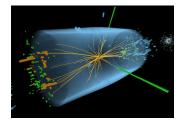
• **2110.10177**:

Constraining CP-violation in the Higgs-top-quark interaction using machine-learning-based inference; in collaboration with S. Brass.

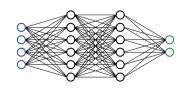
• **2202.11753**:

Constraining the CP structure of Higgs-fermion couplings with a global LHC fit, the electron EDM and baryogenesis; in collaboration with E. Fuchs, S. Heinemeyer, J. Katzy, M. Menen, K. Peters, M. Saimpert, G. Weiglein.

Outline of the talk



Global LHC fit



Constraining CP violation using ML-based inference





Complementarity with EDM and baryogenesis constraints

Introduction

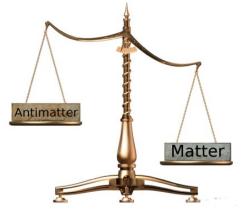
Why should we care about CP violation in the Higgs-fermion couplings?

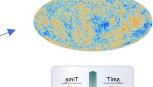
CP violation in the Higgs sector

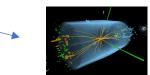
- New sources of CP violation are necessary to explain the baryon asymmetry of the Universe.
- One possibility: CP violation in the Higgs sector.

Is the SM-like Higgs boson a CP-admixed state?

- CP violation in the Higgs sector can be constrained by
 - demanding significant contribution to the baryon asymmetry (BAU)
 - electric dipole measurements,
 - collider measurements.







The CP nature of the Higgs boson

- CP violation in *HVV* couplings already tightly constrained via VBF and $pp \rightarrow VH$ production as well as $H \rightarrow 4l$ decay. [ATLAS,CMS:..,2002.05315, 2104.12152,2109.13808,2202.06923,2205.05120]
- CP-violating HVV coupling can only be induced at the loop level → expected to be small in most BSM theories.
- CP violation in Higgs—fermion couplings can be induced at the tree level.

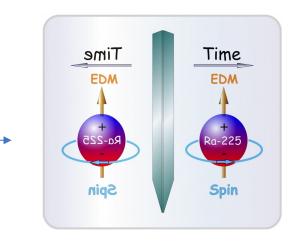
Focus of this talk: Constraining CP violation in the Higgs–fermion interactions.

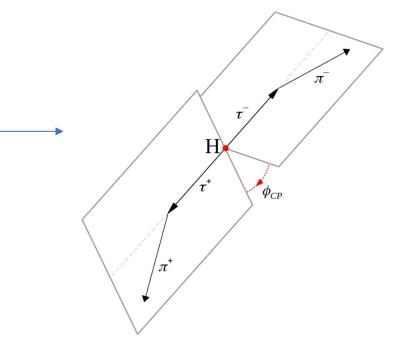
 $H^{\mathcal{CP}\text{-even}} - H^{\mathcal{CP}\text{-odd}}$

Constraining CP violation

CP violation in the Higgs sector can be constrained using:

- Pure CP-odd observables:
 - Unambiguous markers for CP violation: e.g.
 - EDM measurements,
 - decay angle in $H \rightarrow \tau^+ \tau^-$.
 - Experimentally difficult for some processes (i.e., top-associated Higgs production).

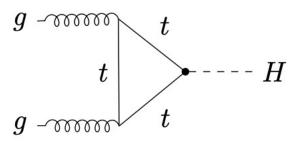


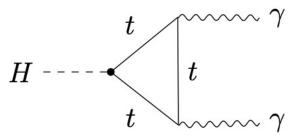


Constraining CP violation

CP violation in the Higgs sector can be constrained using:

- Pure CP-even observables:
 - Many rate measurements are indirectly sensitive: e.g.
 - Higgs production via gluon fusion,
 - $H \rightarrow \gamma \gamma$.
 - Deviations from SM need not be due to CP violation.

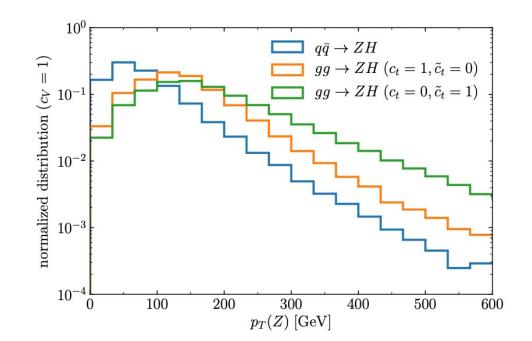




Constraining CP violation

CP violation in the Higgs sector can be constrained using:

- Kinematic information:
 - Effectively mixes CP-even and CP-odd observables.
 - High sensitivity expected since all available information is used.
 - Can be difficult to reinterpret if multivariate analysis is used.

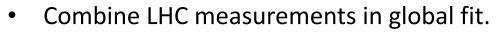




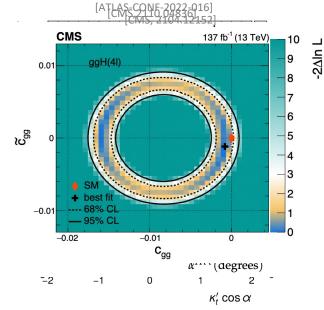
Exploit all three complementary approaches to learn as much as possible!

Why should we look for CP-violating $Hf\bar{f}$ couplings now?

- More and more experimental CP measurements:
 - *Htt* coupling: $t\bar{t}H$, tH, tWH production (with $H \rightarrow \gamma\gamma$, $\bar{b}b$) using kinematic analysis,
 - $H\tau\tau$ coupling: $H \rightarrow \tau\tau$ using CP-odd observable,
 - ggH coupling: H + 2j production using CP-odd observable.
- Increased precision on indirectly sensitive channels.
 - E.g., ggH, $H \rightarrow \gamma\gamma$, etc.
- Much more luminosity to be collected at HL-LHC.
- Tighter upper bounds on CP violation from EDM measurements.



- Propose ways to improve measurements in the future.
- Compare LHC measurements with EDM bounds and baryogenesis constraints.



Effective model

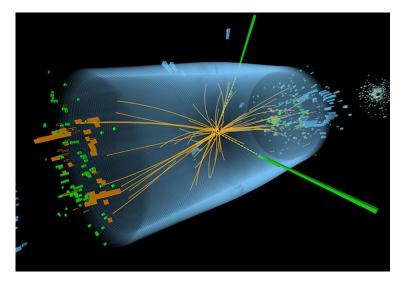
• Modify Yukawa interactions by (e.g. generated by dim-6 $(\phi^{\dagger}\phi)Q_L\tilde{\phi}t_R$ operator)

$$\mathcal{L}_{\text{yuk}} = -\sum_{f=u,d,c,s,t,b,e,\mu,\tau} \frac{y_f^{\text{SM}}}{\sqrt{2}} \bar{f} \left(c_f + i\gamma_5 \tilde{c}_f \right) fH,$$

• Allow moreover for CP-conserving modification of HVV couplings

$$\mathcal{L}_V = c_V H\left(\frac{M_Z^2}{v} Z_\mu Z^\mu + 2\frac{M_W^2}{v} W_\mu^+ W^{-\mu}\right)$$

- SM: $c_f = 1$, $\tilde{c}_f = 0$, $c_V = 1$.
- Parametrize effect of undiscovered colored and neutral BSM particles via effective Higgs–gluon and Higgs–photon interactions.

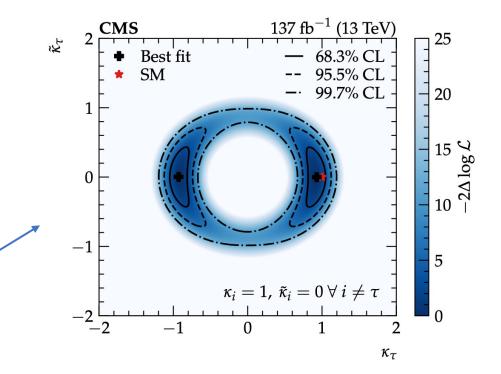


Global LHC fit

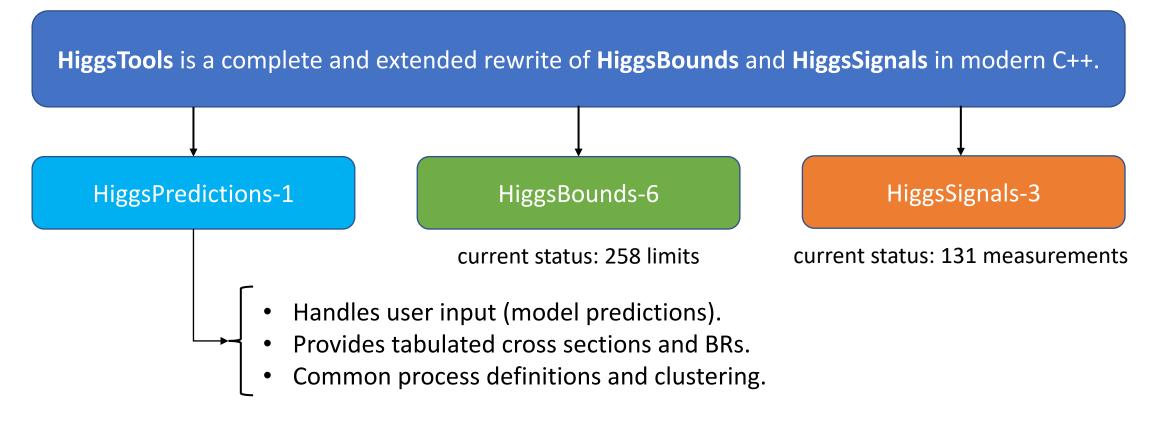
What can we learn from current LHC data?

LHC constraints — setup

- Experimental input:
 - All relevant Higgs measurements:
 - Rate measurements (production + decay),
 - ZH STXS measurements (p_T shape),
 - CMS $H \rightarrow \tau \tau$ CP analysis, [2110.04836]
 - did not include dedicated experimental top-Yukawa CP analyses (difficult to reinterpret in another model).
 - If available, included all uncertainty correlations.
- Scanning using either random scan or Markov-chain algorithm,
- χ^2 fit performed using HiggsSignals including ~ 100 different measurements.



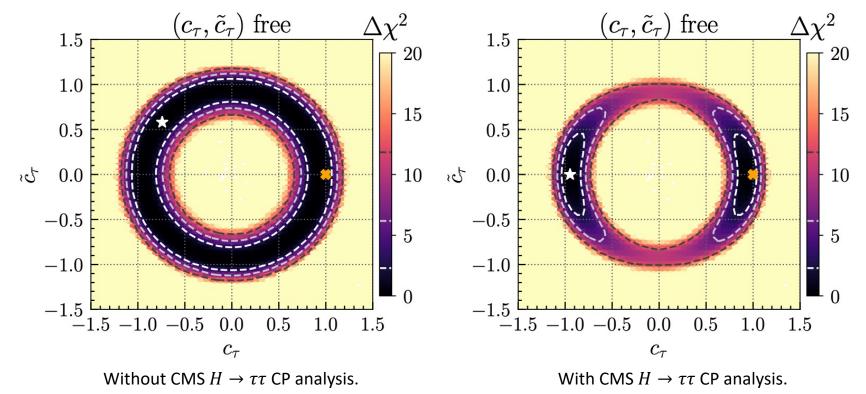
Interlude: HiggsTools



C++ interface for high performance; Python and Mathematica interfaces for ease of use.

Starting point — 1 flavor fits: au

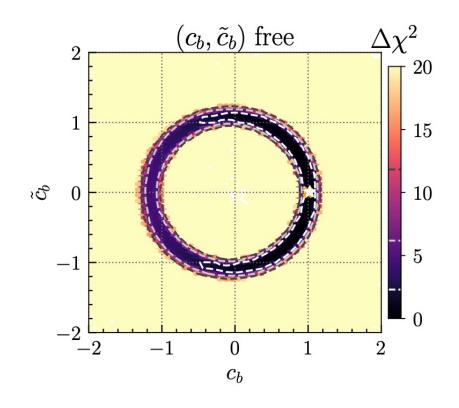
[HB et al.,2202.11753]



- Without CMS $H \to \tau \tau$ CP analysis ring-like structure since $\Gamma_{H \to \tau \tau} \propto c_{\tau}^2 + \tilde{c}_{\tau}^2$ (similar for muon-Yukawa coupling).
- With CMS $H \rightarrow \tau \tau$ CP analysis, we can differentiate between CP-even and CP-odd tau-Yukawa coupling.

1 flavor fits: *b*

[HB et al.,2202.11753]



- Ring-like structure since $\Gamma_{H \to bb} \propto c_b^2 + \tilde{c}_b^2$.
- Bottom-Yukawa coupling, however, also affects *ggH* rate:

•
$$\frac{\sigma_{gg \rightarrow H}}{\sigma_{gg \rightarrow H}^{SM}} \simeq 1.1c_t^2 + 2.6\tilde{c}_t^2 - 0.1c_tc_b + \cdots$$

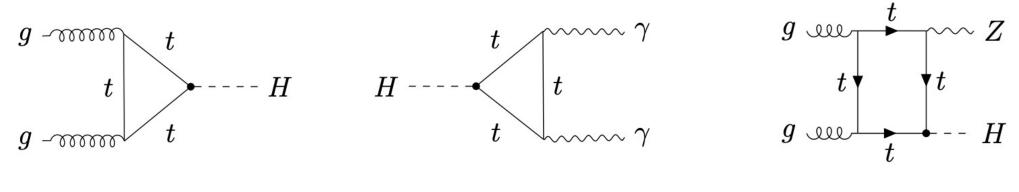
- Negative c_b values disfavored since ggH rate is enhanced by $\sim 20\%$.
- Direct bottom CP measurements very difficult.



Indirect CP constraints will remain important for the bottom-Yukawa coupling.

Top-Yukawa coupling

• Probe top-Yukawa coupling at the loop-level via $gg \rightarrow H, H \rightarrow \gamma\gamma, gg \rightarrow ZH$:



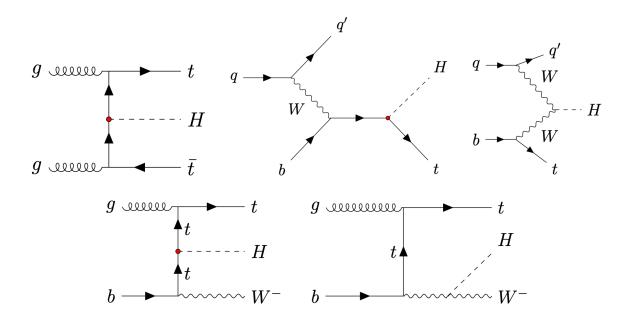
•
$$\kappa_g^2 \equiv \frac{\sigma_{gg \to H}}{\sigma_{gg \to H}^{SM}} \simeq 1.1c_t^2 + 2.6\tilde{c}_t^2 - 0.1c_tc_b - 0.2\tilde{c}_t\tilde{c}_b + \cdots$$
, disfavors large \tilde{c}_t .

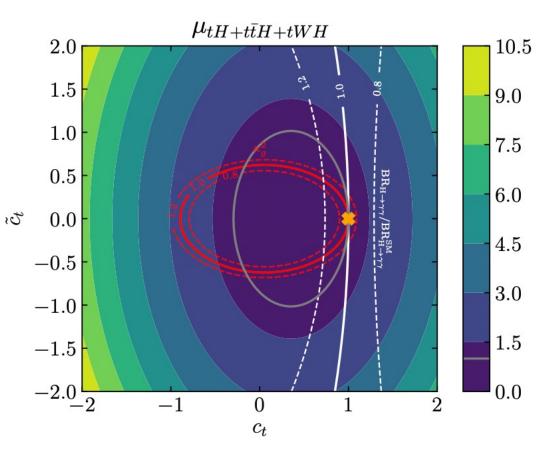
•
$$\kappa_{\gamma}^2 \equiv \frac{\Gamma_{H \to \gamma \gamma}}{\Gamma_{H \to \gamma \gamma}^{SM}} \simeq 1.6c_V^2 - 0.7c_Vc_t + 0.1c_t^2 + 0.2\tilde{c}_t^2 + \cdots$$
, disfavors negative/small c_t .

• $\frac{\sigma_{gg \to ZH}}{\sigma_{gg \to ZH}^{SM}} \simeq 0.5c_t^2 + 0.5\tilde{c}_t^2 + 2.4c_V^2 - 1.9c_Vc_t$..., disfavors negative c_t .

Top-Yukawa coupling

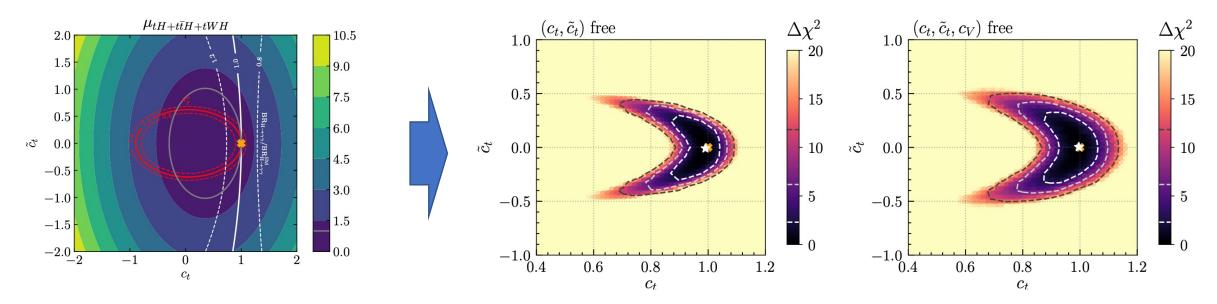
- Probe top-Yukawa coupling at the tree-level via top-associated Higgs production:
 - Three subchannels: $t\bar{t}H$, tH, tWH.
 - Difficult to disentangle experimentally.
 - Consider combined signal strength.





1 flavor fits: t

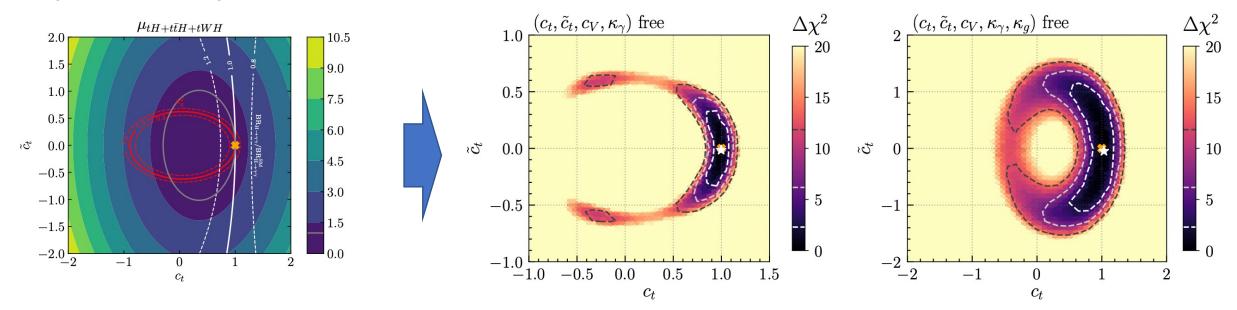
[HB et al.,2007.08542]



- ggH and $H \rightarrow \gamma\gamma$ total rates strongly constraint CP violation in top-Yukawa coupling.
- Relies on assumption that no other BSM physics affect ggH and $H \rightarrow \gamma\gamma$.
- What happens if we allow κ_{γ} and κ_{g} to float freely?

1 flavor fits: $t - \text{free } \kappa_{\gamma}, \kappa_{g}$

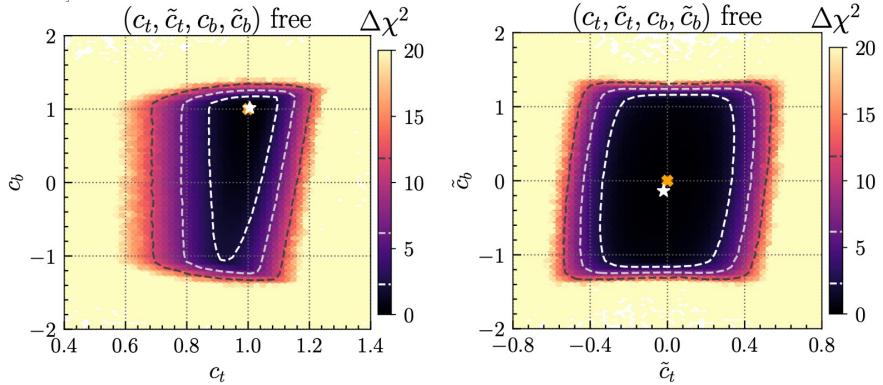
[HB et al.,2007.08542]



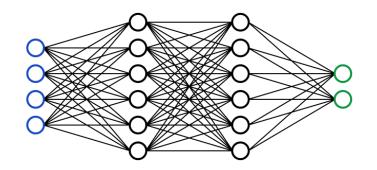
- Colored and charged BSM particles can cancel the effect of a modified top-Yukawa coupling.
- Top-associated Higgs production is a more model-independent but weaker probe.

2 flavor fits: *t* and *b*

[HB et al.,2202.11753]



- ggH rate correlates top and bottom Yukawa couplings: $\kappa_g^2 \simeq 1.1c_t^2 + 2.6\tilde{c}_t^2 0.1c_tc_b 0.2\tilde{c}_t\tilde{c}_b$.
- Correlation of CP-odd coupling modifiers weaker since bounds on \tilde{c}_t are stronger.



Constraining CP-violation

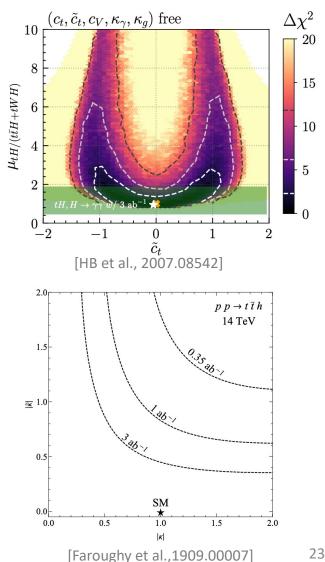
using ML-based inference

How can we improve the LHC bounds on the top-Yukawa coupling in the future?

Future probes of the top-Yukawa interaction

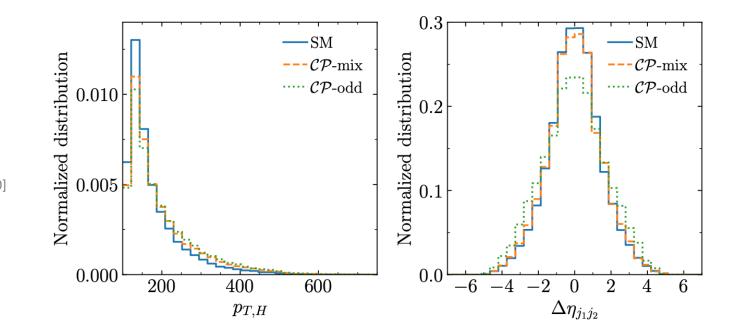
- Future rates measurements:
 - Need to disentangle $t\bar{t}H$ and tH to improve sensitivity on \tilde{c}_t . [HB et al., 2007.08542]
 - Possible alternative channels: $t\bar{t}$, $t\bar{t}t\bar{t}$. [Cao et al.,1901.04567;Martini et al.,2104.04277]
 - Even at HL-LHC comparably weak bounds expected.

- Future measurements of CP-odd observables: [e.g. Faroughy et al.,1909.00007; Bortolato et al.,2006.13110;Barman et al.,2110.07635]
 - Difficult since top quarks need to be reconstructed.
 - Resulting projected limits are relatively weak.



Alternative approach: kinematic analysis

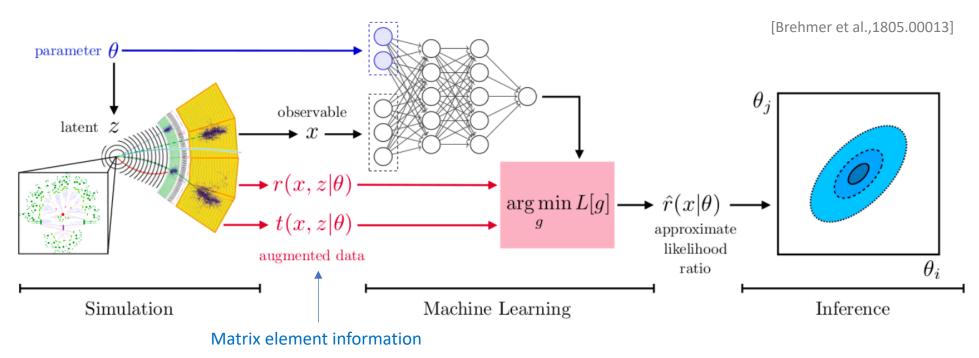
- Multivariate analyses exploiting kinematic information:
 - BDT analysis, [CMS,2003.10866;ATLAS,2004.04545]
 - matrix-element approach, [e.g. Goncalves et al,1804.05874;Kraus et al.,1908.09100]
 - High sensitivity expected.



How to best exploit the full available information to constraint top-Yukawa interaction?

Machine-learning based inference

[Brehmer et al., 1906.01578, 1805.12244, 1805.00013, 1805.00020, 1808.00973]



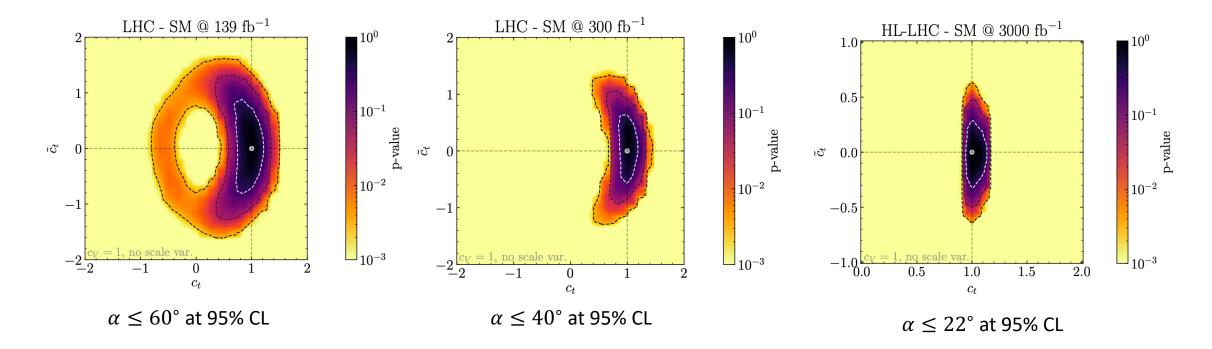
- Allows to extract the full available information (maximal sensitivity).
- No information loss due to binning (as for BDT analysis).
- No approximation of shower and detector effects (as for matrix-element approach).
- Use implementation in public code MadMiner designed to work with MadGraph + Pythia + Delphes. [Brehmer,Kling,Espejo,Cranmer,1907.10621]

ML-based inference: setup

- Focus on top-associated Higgs production $(t\bar{t}H,tH,tWH)$ with $H \rightarrow \gamma\gamma$.
- We require at least one lepton \rightarrow consider ZH, WH as backgrounds.
- Non-Higgs backgrounds are assumed to be subtracted by fit to smoothly falling $m_{\gamma\gamma}$ distribution.
- Free parameters: c_t , \tilde{c}_t , and c_V (+ renormalization scale μ_R).
- Defined 47 observables used by neural network (photon, jet, lepton momenta, Higgs p_T , etc.).
- Averaged over ensemble of six neural networks to minimize ML uncertainty.
- \Rightarrow Evaluate likelihoods for different luminosities at the LHC + HL-LHC.

Expected limits at the (HL-)LHC

[HB&Brass,2110.10177]



- Can also interpret result in terms of mixing angle $\tan \alpha = \tilde{c}_t/c_t$.
- Additional variation of c_V (and of the renormalization scale) only slightly weakens bounds (~ 5° for 300 fb⁻¹).

Which observables drive these constraints?

• Use Fisher matrix to evaluate information for different observables

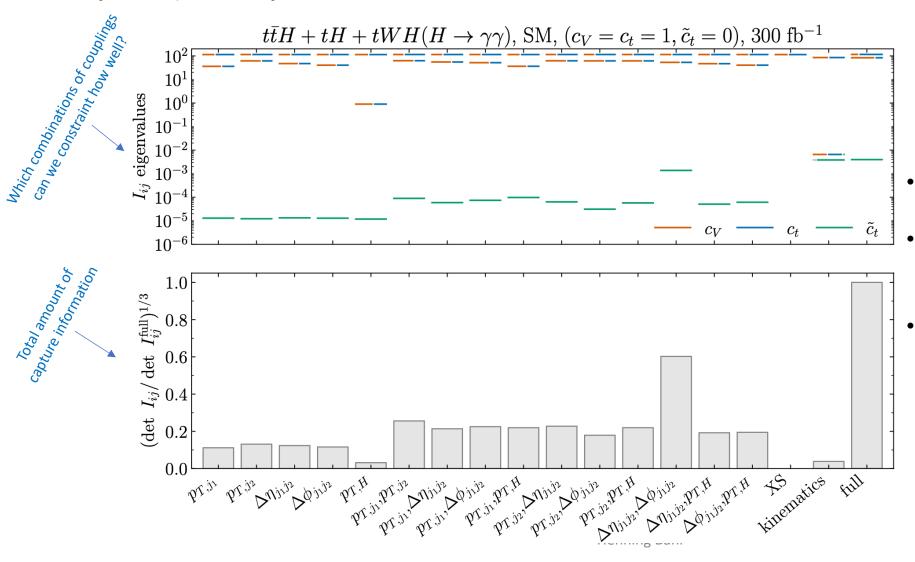
$$I_{ij}(\theta) = \mathbb{E}\left[\frac{\partial \log p_{\text{full}}(\{x\}|\theta)}{\partial \theta_i} \frac{\partial \log p_{\text{full}}(\{x\}|\theta)}{\partial \theta_j}\Big|_{\theta}\right], \quad \text{with} \quad \operatorname{cov}(\hat{\theta}|\theta)_{ij} \ge I_{ij}^{-1}(\theta),$$

 \rightarrow The higher the information, the more precise we can measure a parameter.

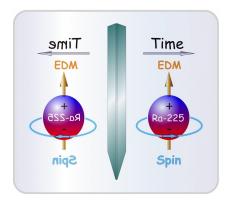
- E.g., for SM point we have $Information about c_V \qquad Information about c_t$ $I_{ij}^{\text{full}}(\text{SM}) \simeq \begin{pmatrix} 91.4 & 13.7 & 0.1 \\ 13.7 & 108.2 & -0.1 \\ 0.1 & -0.1 & 0.004 \end{pmatrix},$ Correlation of c_t and c_V Information about \tilde{c}_t
- Evaluate Fisher matrix for various 1D and 2D histograms, full likelihood, XS only, kinematics only.

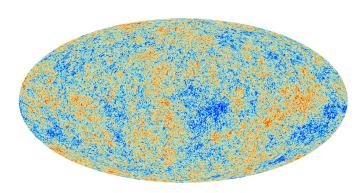
Fisher information for SM scenario

[HB&Brass,2110.10177]



- \tilde{c}_t not constrained by rate.
- Use of kinematic information mandatory.
- No single observable able to capture information about \tilde{c}_t .

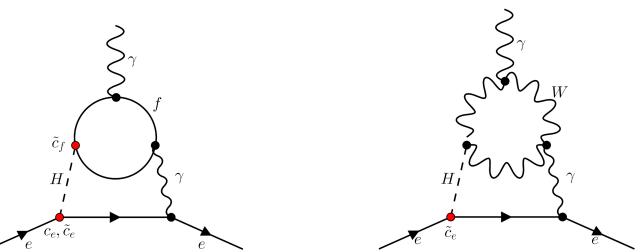




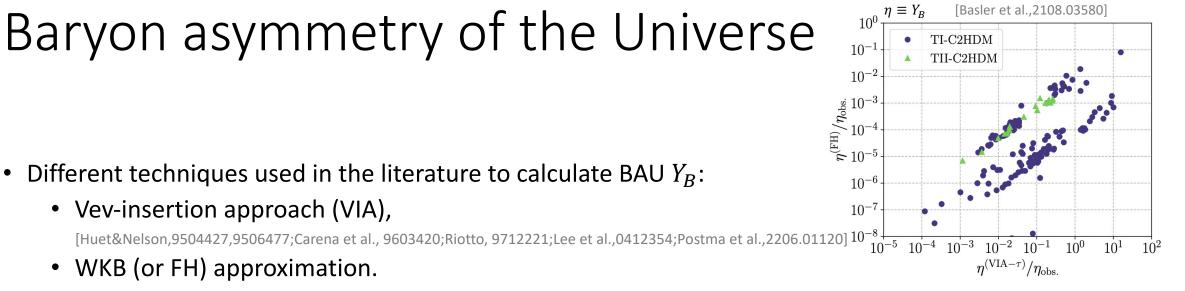
Complementarity with EDM and baryogenesis constraints

Can CP violation in the Higgs sector explain the BAU?

EDM constraints



- Several EDMs are sensitive to CP violation in the Higgs sector.
- We consider only constraints from theoretically cleanest EDM: the electron EDM. [Brod et al.,1310.1385,1503.04830, 1810.12303, 2203.03736;Panico et al.,1810.09413;Altmannshofer et al.,2009.01258]
- Strongest limit by ACME collaboration: $d_e^{\text{ACME}} = 1.1 \cdot 10^{-29} e \text{ cm}$ at 90% CL. [ACME, Nature 562 (2018) 7727, 355-360]
- $\frac{d_e}{d_e^{\text{ACME}}} \simeq \frac{c_e}{c_e} (870.0\tilde{c}_t + 3.9\tilde{c}_b + 3.4\tilde{c}_\tau + \dots) + \tilde{c}_e (610.1c_t + 3.1c_b + 2.8c_\tau 1082.6c_V + \dots)$
- Bounds strongly depend on assumptions about electron-Yukawa coupling.



[Joecy et al.,9410282;Kainulainen et al.,0105295, 0202177;Prokopec et al., 0312110, 0406140;Konstandin et al.,1302.6713, 1407.3132]

- VIA approach yields consistently higher results by orders of magnitude.
- We use VIA approach with bubble wall parameters close to optimal values for Y_B : [de Vries,1811.11104;Fuchs et al.,2003.00099,2007.06940;Shapira,2106.05338]

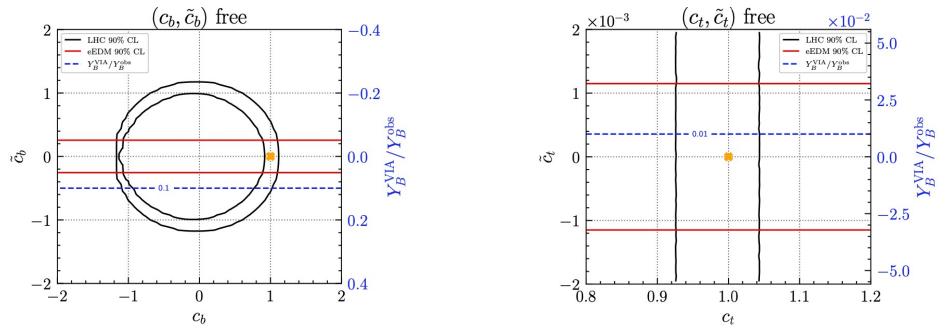
$$\frac{Y_B}{Y_B^{\rm obs}} \simeq 28\tilde{c}_t - 0.2\tilde{c}_b - 11\tilde{c}_\tau + \cdots$$



 Y_B values should be regarded as **upper bound** on what is theoretically achievable.

1 flavor results: t and b

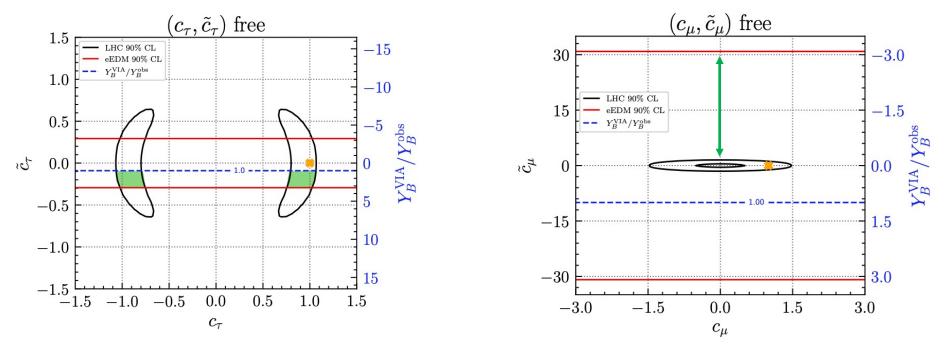
[HB et al.,2202.11753]



- CP-violating bottom-Yukawa coupling contributes too less to BAU.
- CP-violating top-Yukawa coupling strongly constrained by eEDM → not able produce sufficient BAU.

1 flavor results: au and μ

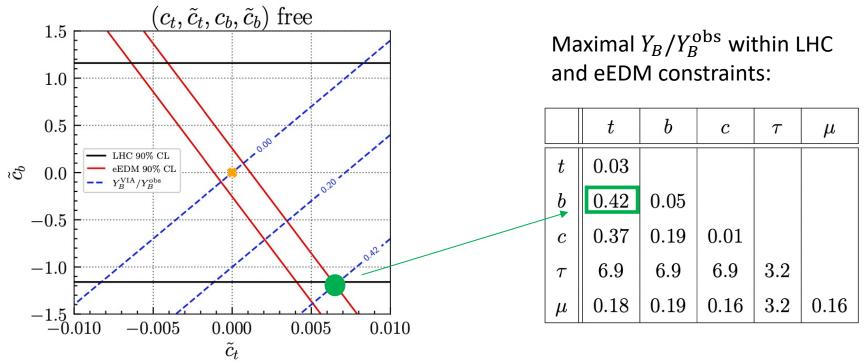
[HB et al.,2202.11753]



- CP-violating tau-Yukawa coupling can potentially explain BAU within LHC and eEDM constraints.
- LHC constraints on CP-violating muon-Yukawa coupling are stronger than eEDM bounds.

2 flavor results: *t* and *b*

[HB et al.,2202.11753]

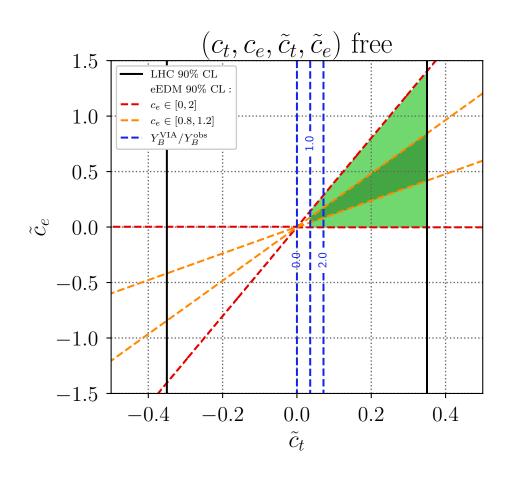


• Presence of more than one CP-violating coupling allows for cancellation in eEDM.

 \rightarrow Larger values for Y_B/Y_B^{obs} can be reached.

Dependence on electron-Yukawa coupling

[HB et al.,2202.11753]

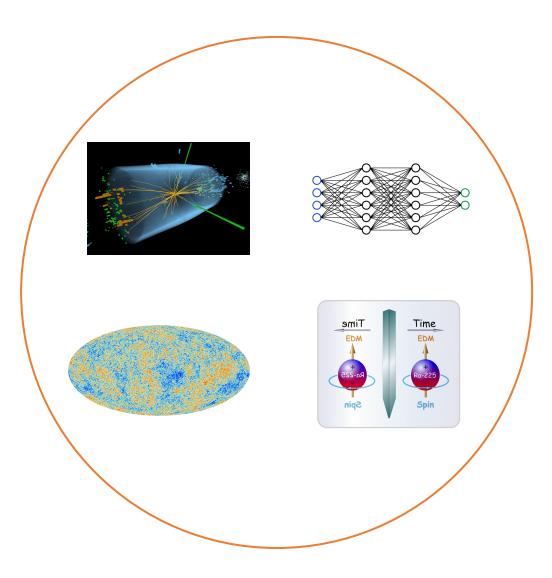


- Electron Yukawa-coupling only very weakly constrained ($g_e \leq 268$ at 95% CL).
- If *c_e* smaller, eEDM significantly weakened.
- Moreover, we can fine-tune CP-odd electron-Yukawa coupling such that $d_e < d_e^{ACME}$.
- Neutron EDM has similar dependence on firstgeneration quark-Yukawa couplings.



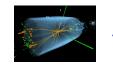
LHC bounds important since they do not depend on 1st gen. Yukawa couplings.

Conclusions

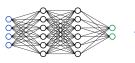


Conclusions

Initial question: how well can we constrain **CP violation in the Higgs–fermion interactions** and what are the implications for the BAU?



LHC already tightly constrains CP violation in the top- and tau-Yukawa couplings.



To improve bounds in the **future**, we need to exploit all available information using e.g. **machine-learning based inference**.



Complementarity of LHC, EDM, and baryogenesis constraints:

- CP violation in tau-Yukawa coupling remains viable source for electroweak baryogenesis.
- LHC allows to distinguish between CP violation in various Yukawa couplings (beginning to probe 2nd generation).
- EDM interpretation strongly depends on first generation Yukawa couplings.

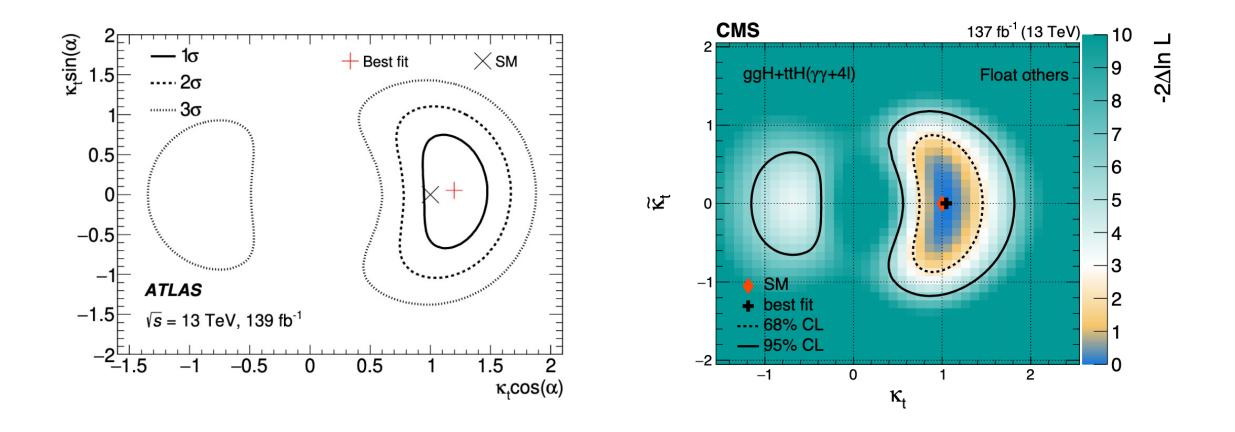
Thanks for your attention!

Appendix

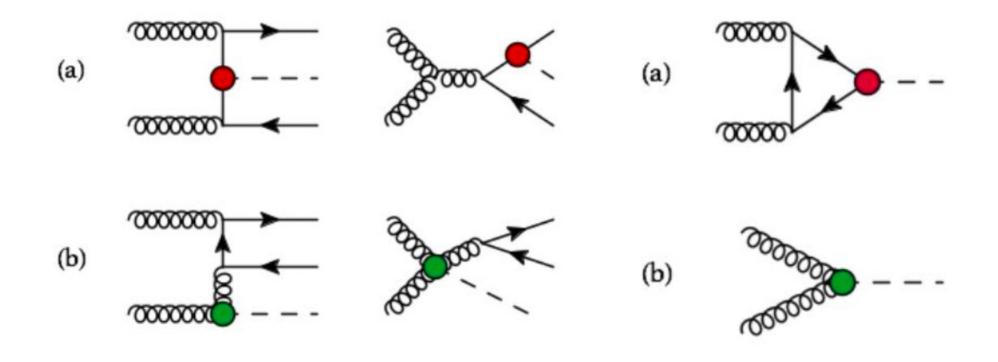
Reasons for not including ATLAS and CMS top CP studies

- CMS study: [2003.10866]
 - All Higgs production modes (apart form top-associated Higgs production) are constrained to their SM predictions.
 - No two-dimensional likelihood given when our study was published (now available in [CMS-PAS-HIG-19-009])
- ATLAS study: [2004.04545]
 - Two setups:
 - 1. κ_g (and κ_{γ}) constrained by other measurements (ggH) excluding $t\bar{t}H$ and tH but events generated at NLO \rightarrow top-associated Higgs production and gluon fusion cannot be regarded as independent.
 - *2.* κ_g and κ_γ calculated as function of c_t and \tilde{c}_t .
 - Assumed *HVV* couplings equal to SM value.

Experimental top CP studies [ATLAS, 2004.04545;CMS, 2104.12152]

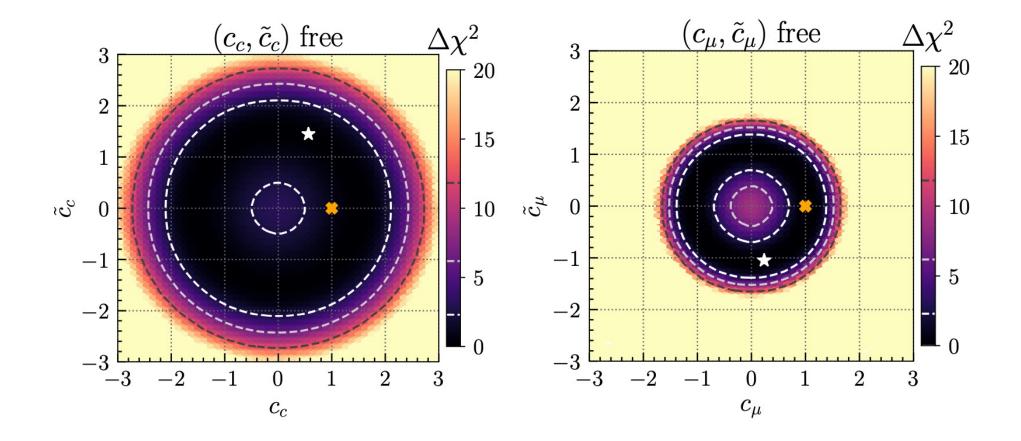


Correlation between ggH and $t\bar{t}H$ at NLO

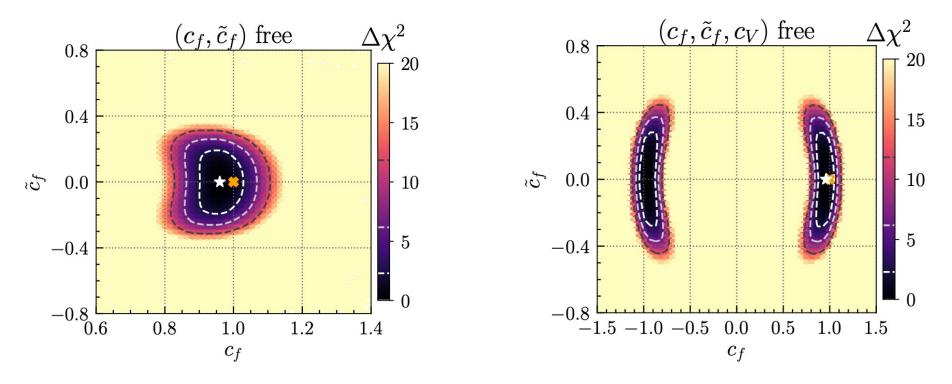


[Maltoni,Vryonidou,Zhang,1607.05330]

Charm- and muon-Yukawa couplings



Global modification fits

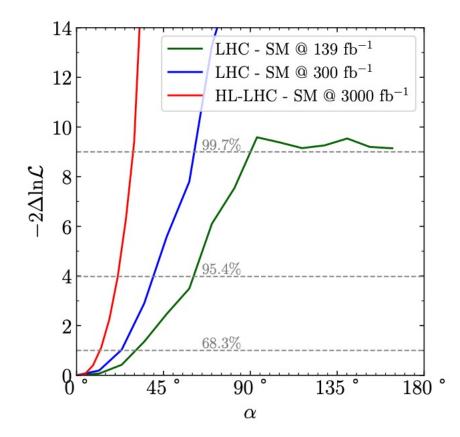


- Universal fermion coupling modifiers: $c_f = c_t = c_b = \cdots = c_{\tau}$, $\tilde{c}_f = \tilde{c}_t = \tilde{c}_b = \cdots = \tilde{c}_{\tau}$.
- Dominated by constraints on top-Yukawa coupling.
- Additional varying c_V reopens negative c_f range.

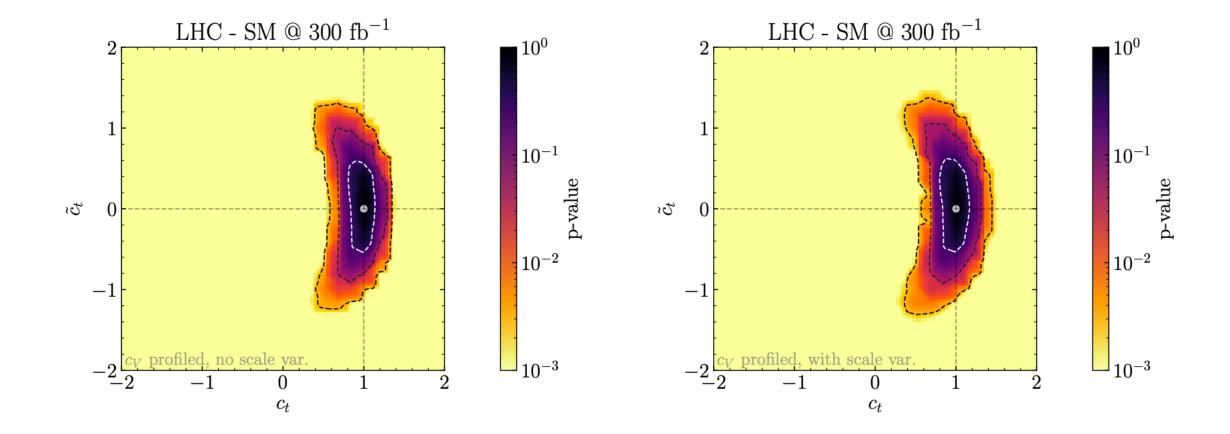
observable	condition
N_{γ}	$\geq 2 \text{ (with } \eta < 2.5 \text{ and } p_T > 25 \text{ GeV})$
$(p_{T,1}^\gamma, p_{T,2}^\gamma)$	$\geq (35,25)~{ m GeV}$
$m_{\gamma\gamma}$	$[105-160]~{\rm GeV}$
$(p_{T,1}^\gamma/m_{\gamma\gamma},p_{T,2}^\gamma/m_{\gamma\gamma})$	$\geq (0.35, 0.25)$
N_ℓ	$\geq 1 \text{ (with } \eta < 2.5 \text{ and } p_T > 15 \text{ GeV})$
$m_{\ell\ell}$	[80, 100] GeV vetoed if same flavour
N_{jet}	$\geq 1 \text{ (with } \eta < 2.5 \text{ and } p_T > 25 \text{ GeV})$

 Table 1: Summary of preselection cuts.

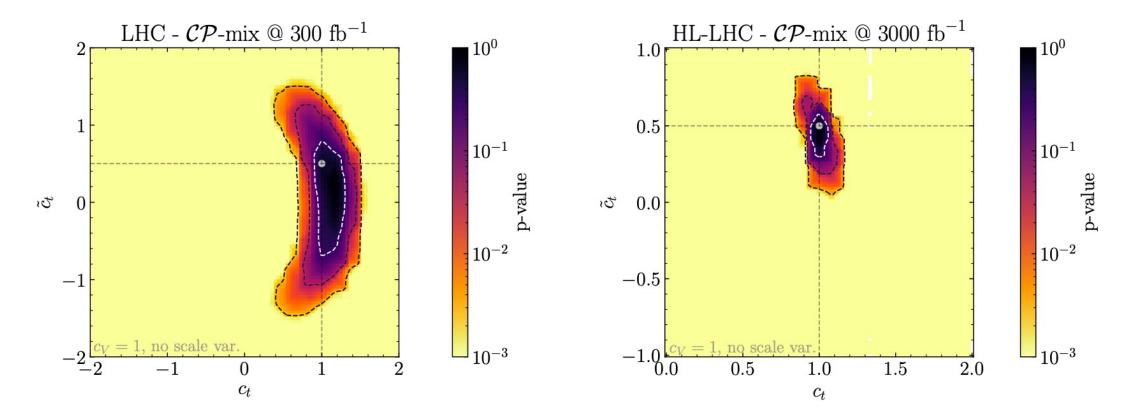
Interpretation in terms of CP-violating angle



Variation of c_V and renormalization scale



Limits in case of deviation from SM



• CP-mix:
$$c_t = 1$$
, $c_{\tilde{t}} = 0.5$, $c_V = 1$.

Fisher information for CP-mixed scenario

