

Constraining the CP character of the Higgs–top-quark interaction

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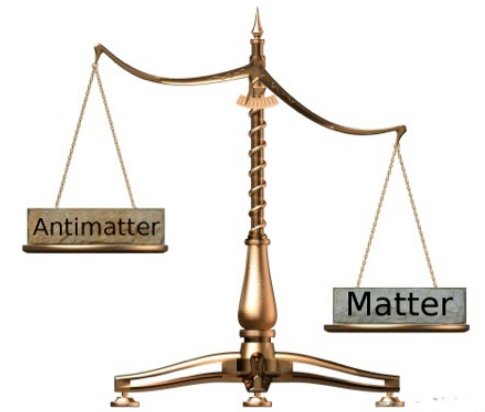


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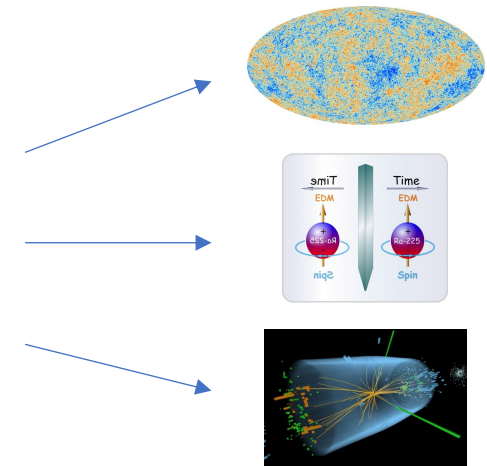
CP violation in the Higgs sector

- New sources of CP violation are necessary to explain the baryon asymmetry of the Universe.
- One possibility: CP violation in the Higgs sector.



Focus of this talk: Constraining CP violation in the top-Yukawa interaction at the LHC.

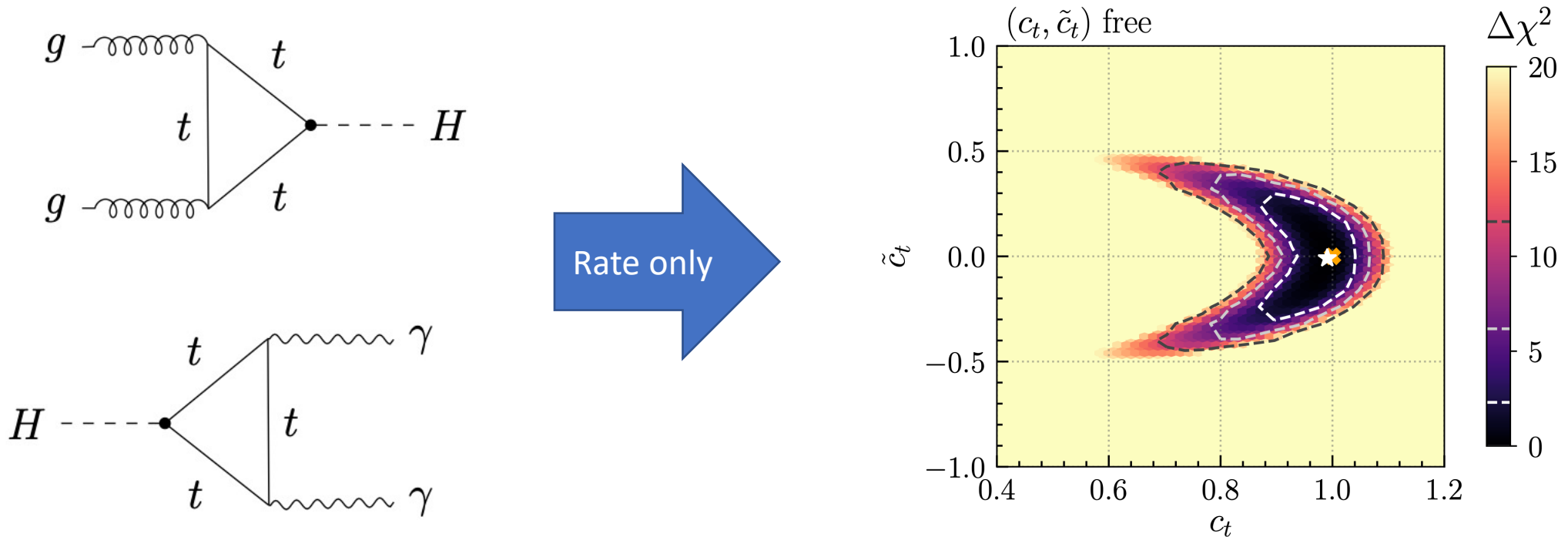
- CP violation in the Higgs sector can be constrained by
 - demanding significant contribution to the baryon asymmetry (BAU)
 - electric dipole measurements,
 - collider measurements.



Collider constraints @ loop level

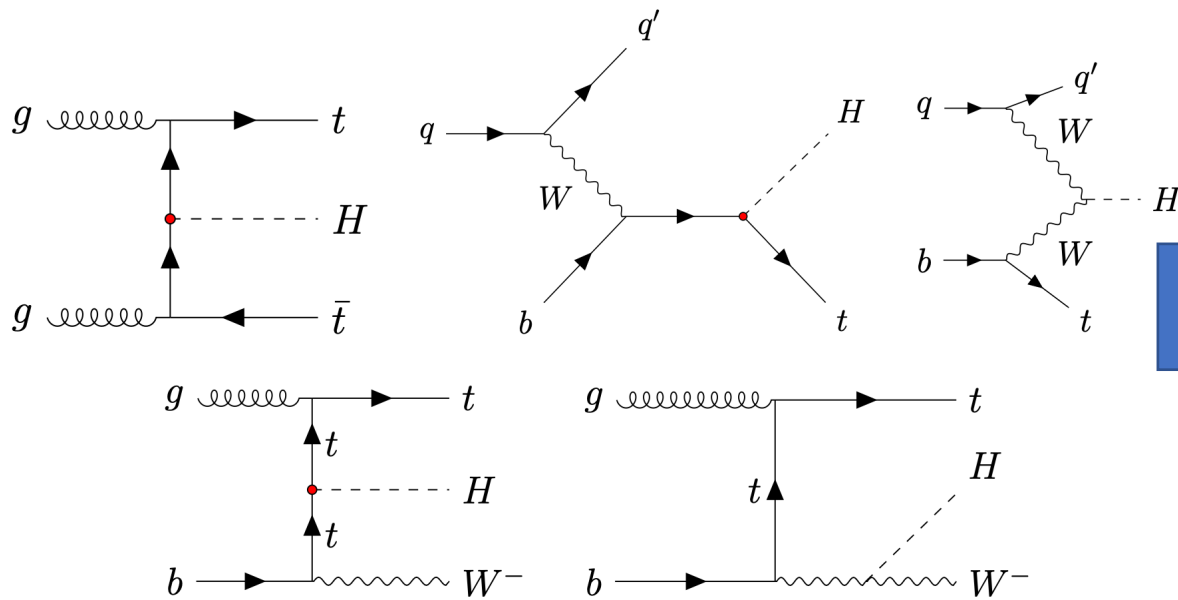
[HB et al., 2007.08542]

- Effective model: $\mathcal{L}_{\text{top-yuk}} = -\frac{y_t^{\text{SM}}}{\sqrt{2}} \bar{t} (c_t + i\gamma_5 \tilde{c}_t) t H.$
- Probe top-Yukawa coupling at the loop-level via $gg \rightarrow H, H \rightarrow \gamma\gamma$:

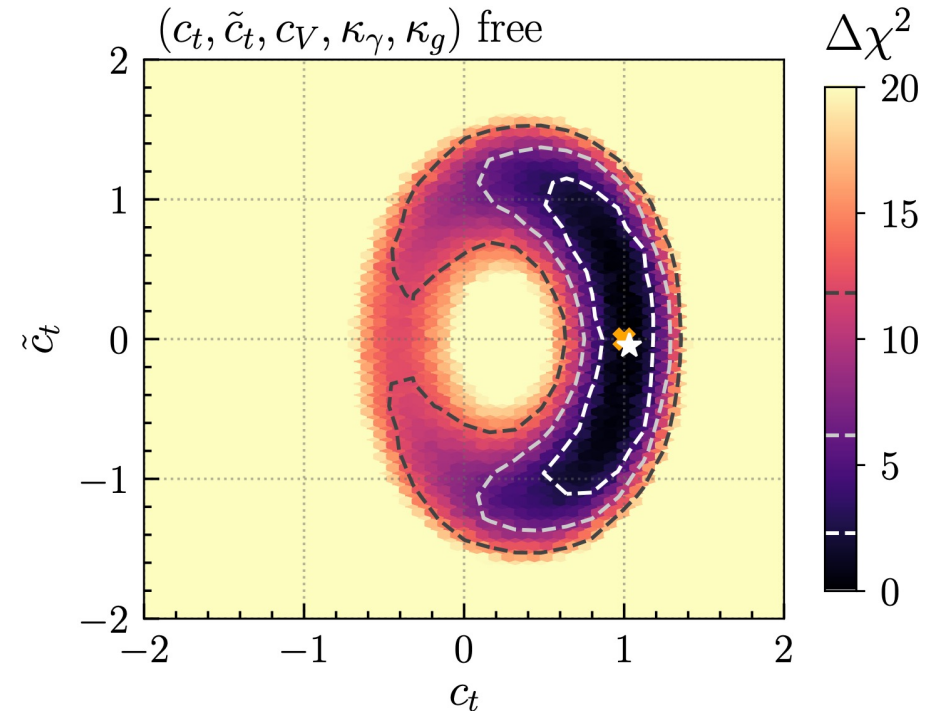


Collider constraints @ tree level

- Tree-level constraints: **top associated Higgs production**
- Direct access to top-Yukawa interaction \rightarrow less model dependence.
- Three sub channels contribute: $\bar{t}tH$, tH (or tHq), tWH .



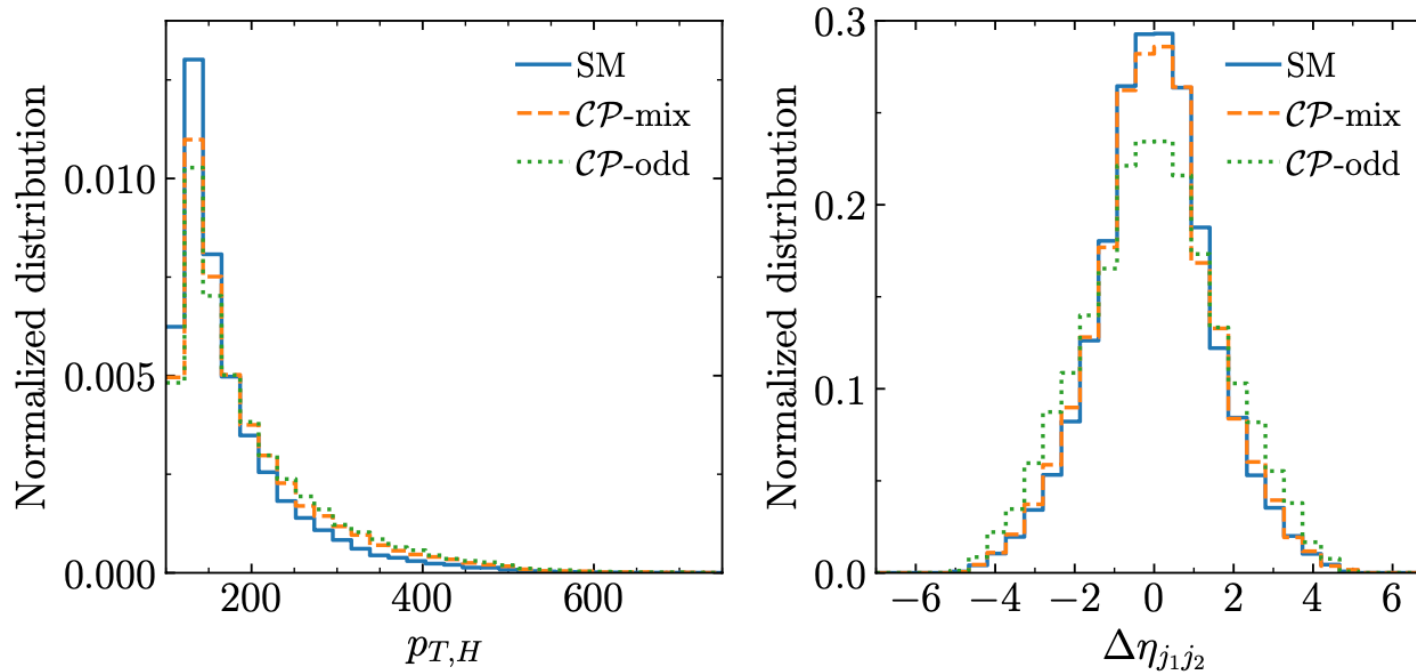
Rate only



Exploiting the kinematic information

[HB & Brass, 2110.10177]

Exemplary kinematic distributions for top-associated Higgs production

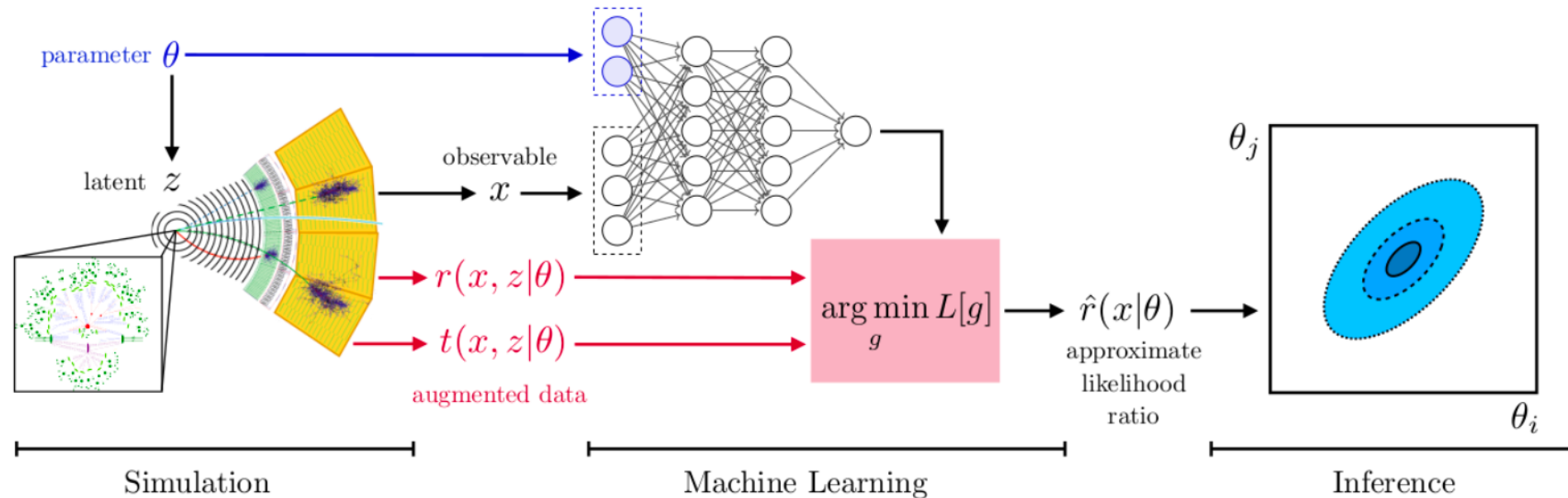


How to best exploit the full available information to constraint top-Yukawa interaction?

→ Focused on top-associated Higgs production with $H \rightarrow \gamma\gamma$ (demanding at least one lepton).

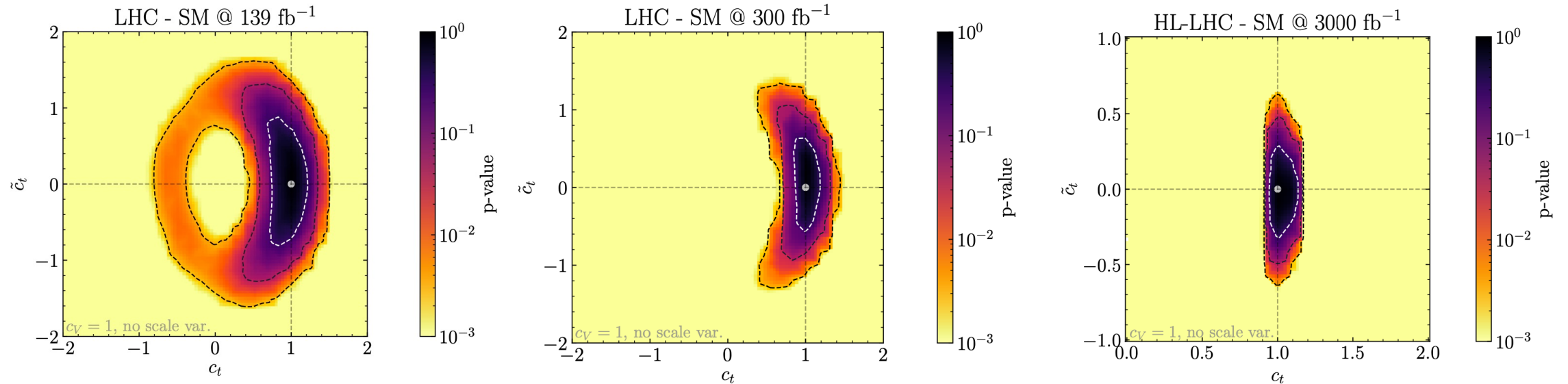
Machine-learning based inference

[Brehmer et al.,1906.01578,1805.12244,1805.00013,1805.00020,1808.00973]



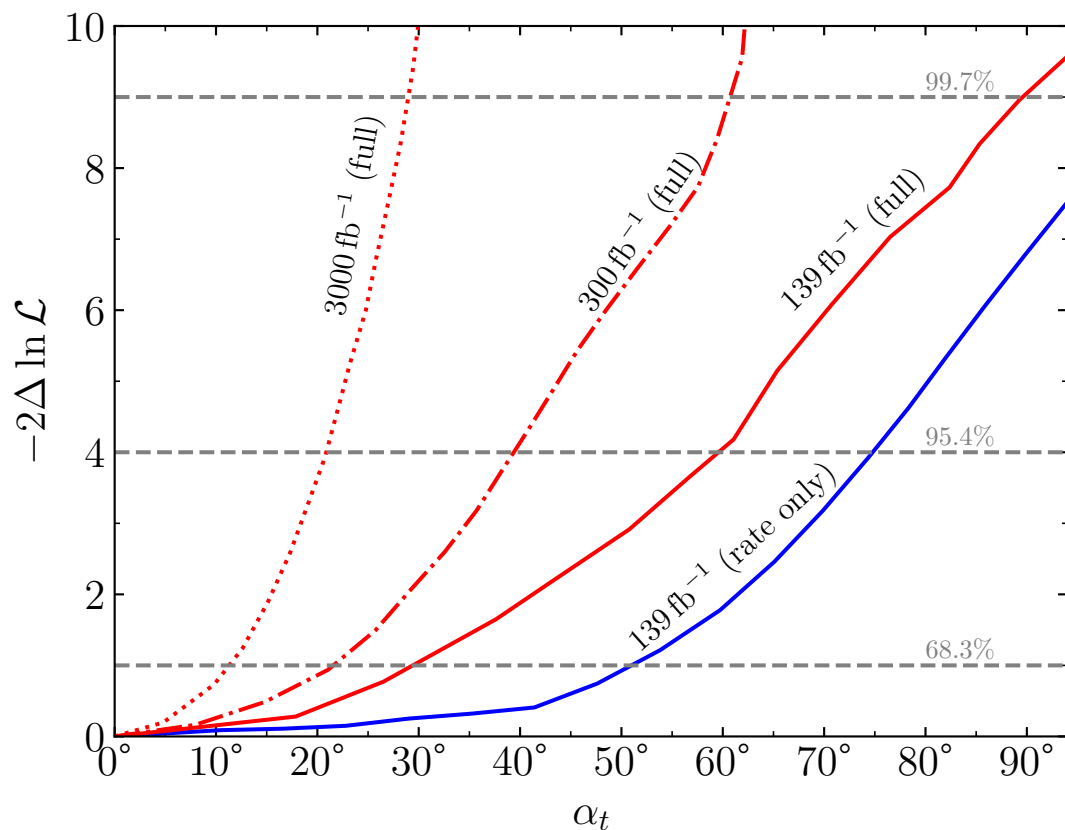
- Allows to extract the full available information (maximal sensitivity).
 - Use implementation in public code MadMiner [Brehmer,Kling,Espejo,Cranmer,1907.10621] designed to work with MadGraph + Pythia + Delphes.
 - Defined 47 observables as input for neural network.
- Averaged over ensemble of six neural networks to minimize ML uncertainty.

Expected limits at the (HL-)LHC



- Assumed here that Higgs–vector-boson coupling is SM-like ($c_V = 1$).
- Additional variation of c_V (and of the renormalization scale) only slightly weakens bounds.

Comparison of constraints on CP-violating phase



- CP-violating phase α_t :
$$\tan \alpha = \tilde{c}_t / c_t$$
- Exploiting full kinematic information significantly strengthen limits.
- Including full-hadronic channel and other Higgs decay channels will allow to further improve sensitivity.

Complementarity with eEDM and BAU [HB et al., 2202.11753]

Electron EDM

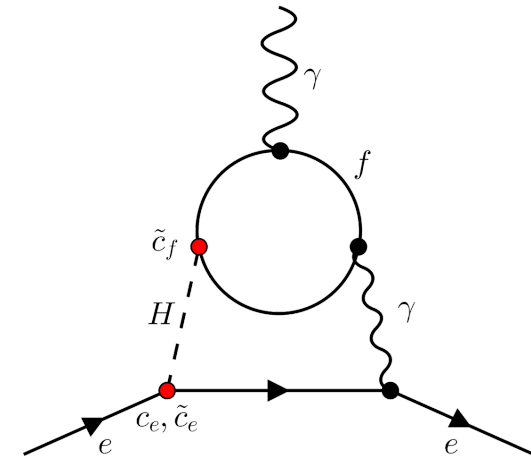
- Several EDMs are sensitive to CP violation in the Higgs sector.

- We consider only the electron EDM.

[Brod et al., '13, '15, '18, '22; Panico et al., '18; Altmannshofer et al., '20]

- $$\frac{d_e}{d_e^{\text{ACME}}} \simeq 870c_e\tilde{c}_t + \tilde{c}_e(610c_t - 1082.6c_V) + \dots$$

- Bounds strongly depend on assumptions about electron-Yukawa coupling.



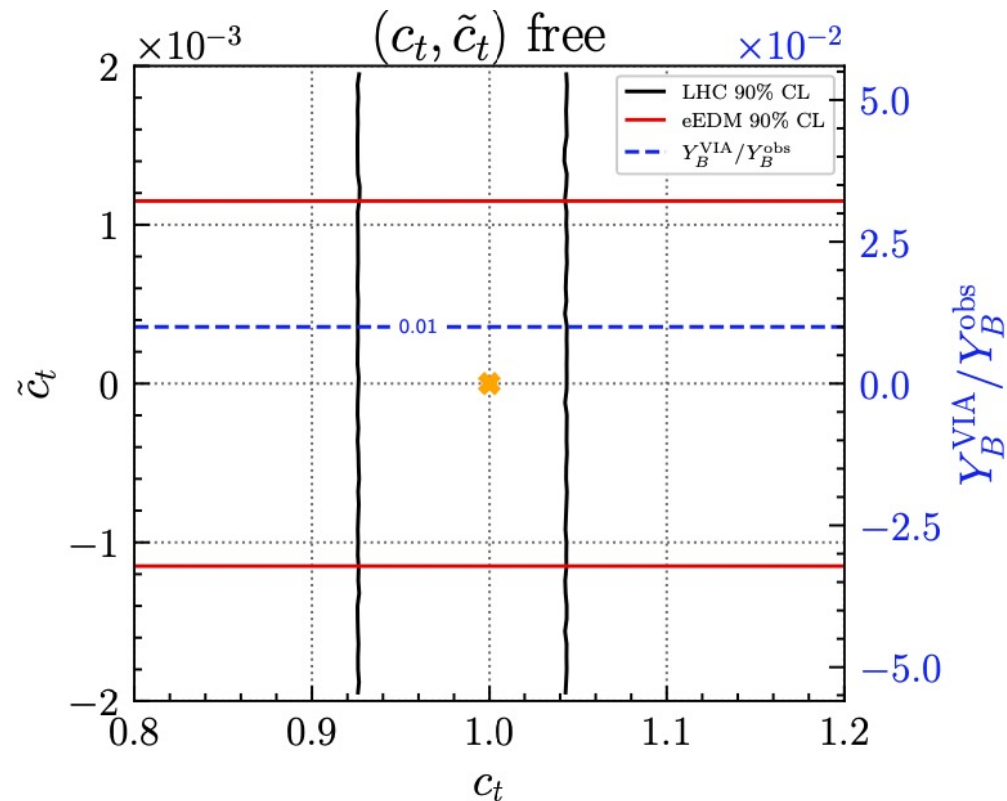
BAU

- Different techniques used in the literature to calculate BAU Y_B : vev-insertion approach (VIA) and WKB approximation.

[Huet&Nelson, '95; Carena et al., '96; Riotto, '97; Lee et al., '04; Joecy et al., '94; Kainulainen et al., '01, '02; Prokopec et al., '03, '04; Konstandin et al., '13, '14; Basler, '21]

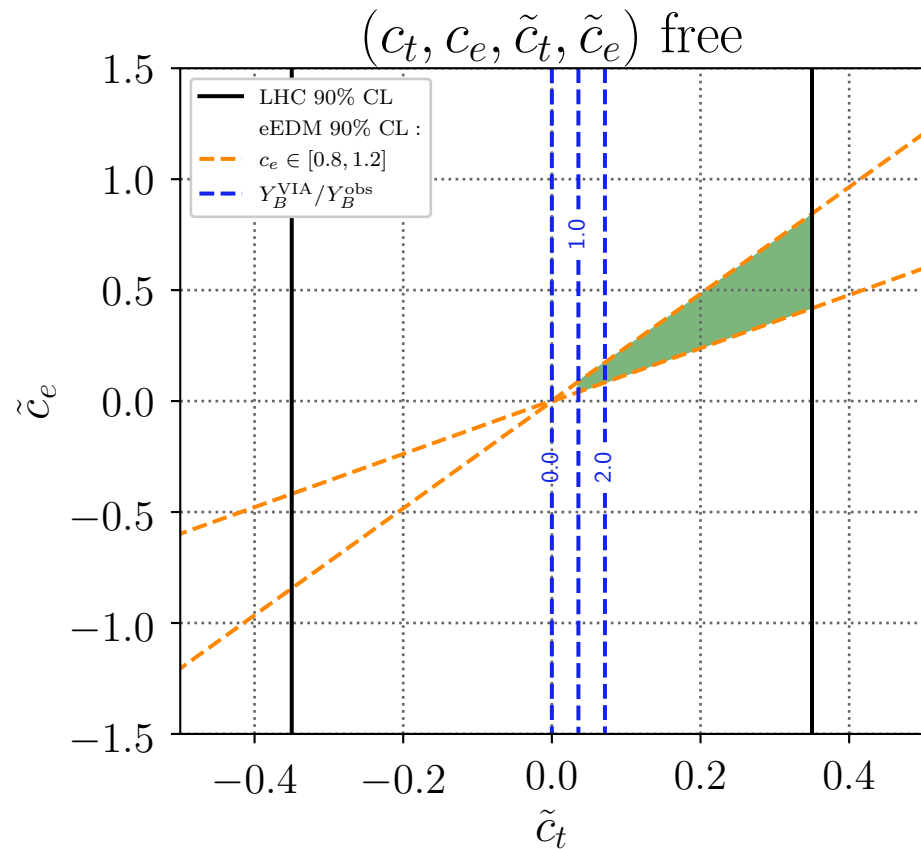
- VIA approach yields consistently higher results by orders of magnitude.
- We use VIA approach with bubble wall parameters close to optimal values for $Y_B \rightarrow Y_B$ values should be regarded as **upper bound** on what is theoretically achievable. [de Vries et al., '18; Fuchs et al., '20; Shapira, '21]

Constraints on top-Yukawa coupling



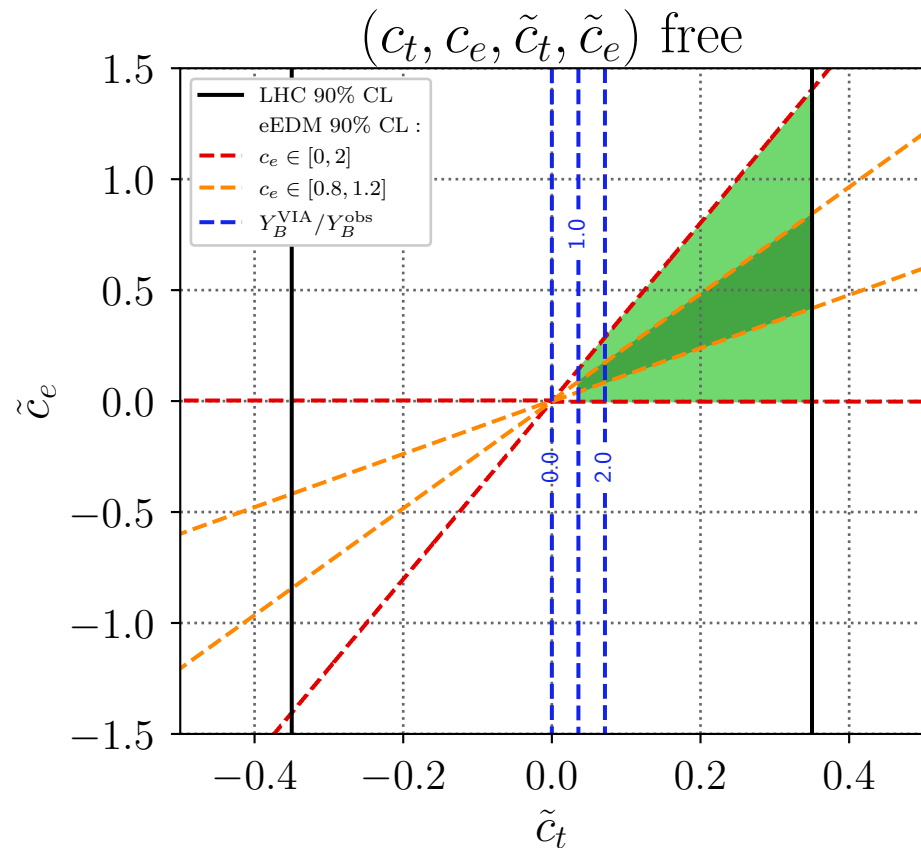
- EDM constraint strongly limits size of \tilde{c}_t
- Only tiny amount of BAU can be generated via CP-violating top-Yukawa coupling.
- However, strong dependence on electron-Yukawa coupling

Dependence on electron-Yukawa coupling



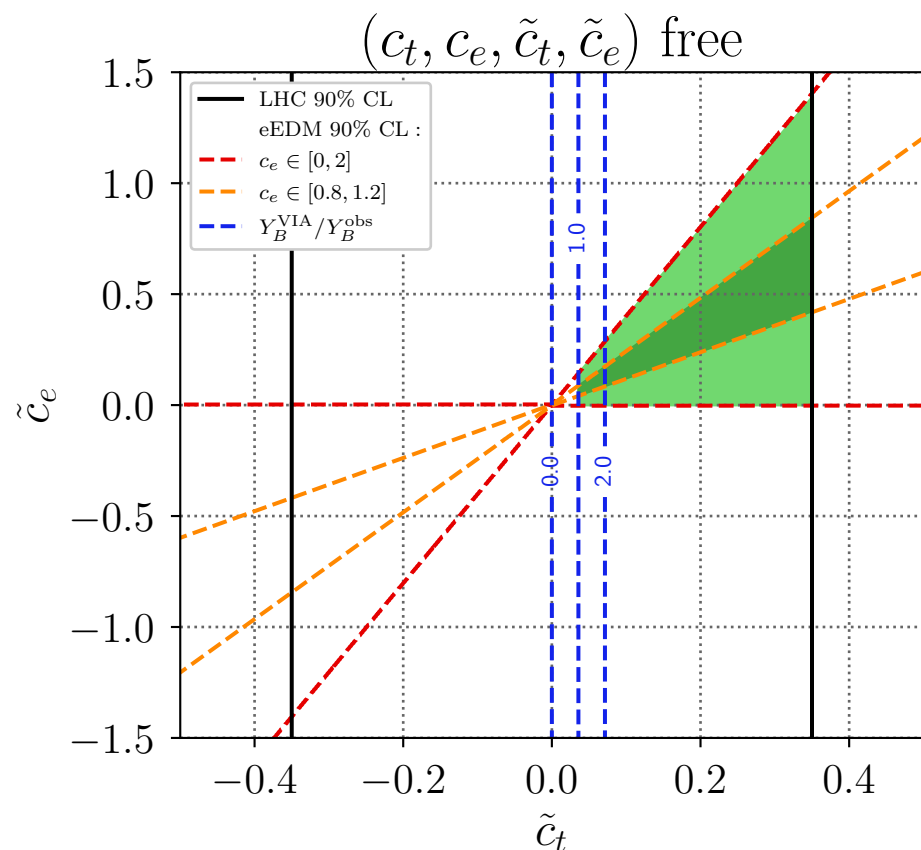
- Electron Yukawa-coupling only very weakly constrained ($g_e \leq 268$ at 95% CL).
- If c_e smaller, eEDM significantly weakened.
- Moreover, we can fine-tune CP-odd electron-Yukawa coupling such that $d_e < d_e^{\text{ACME}}$.
- Neutron EDM has similar dependence on first-generation quark-Yukawa couplings.

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➔ LHC bounds important since they do not depend on 1st gen. Yukawa couplings.

Conclusions

*Initial question: how well can we constrain
CP violation in the Higgs–top-quark interaction?*

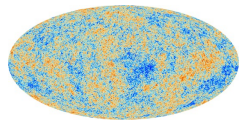
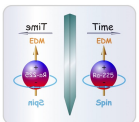
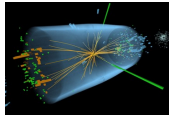
LHC constraints:

- $gg \rightarrow H$ and $H \rightarrow \gamma\gamma$ tightly constrain CP violation in the **top-Yukawa couplings** indirectly.
- Top-associated Higgs production is prime candidate to reduce model dependence.
- Strong constraints from top-associated Higgs production can be expected if full kinematic information is exploited.

EDM and baryogenesis constraints:

- EDM bounds put very strong bounds on a CP-violating top-Yukawa interaction.
- Only very small contribution to BAU realizable
- EDM interpretation, however, strongly depends on first generation Yukawa couplings.

Thanks for your attention!

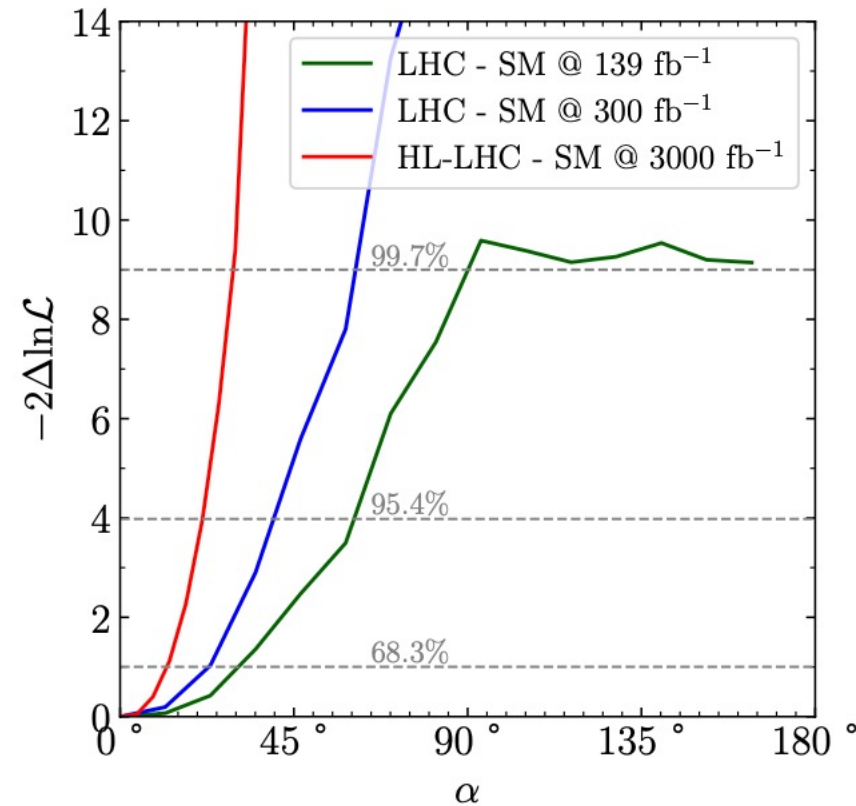


Appendix

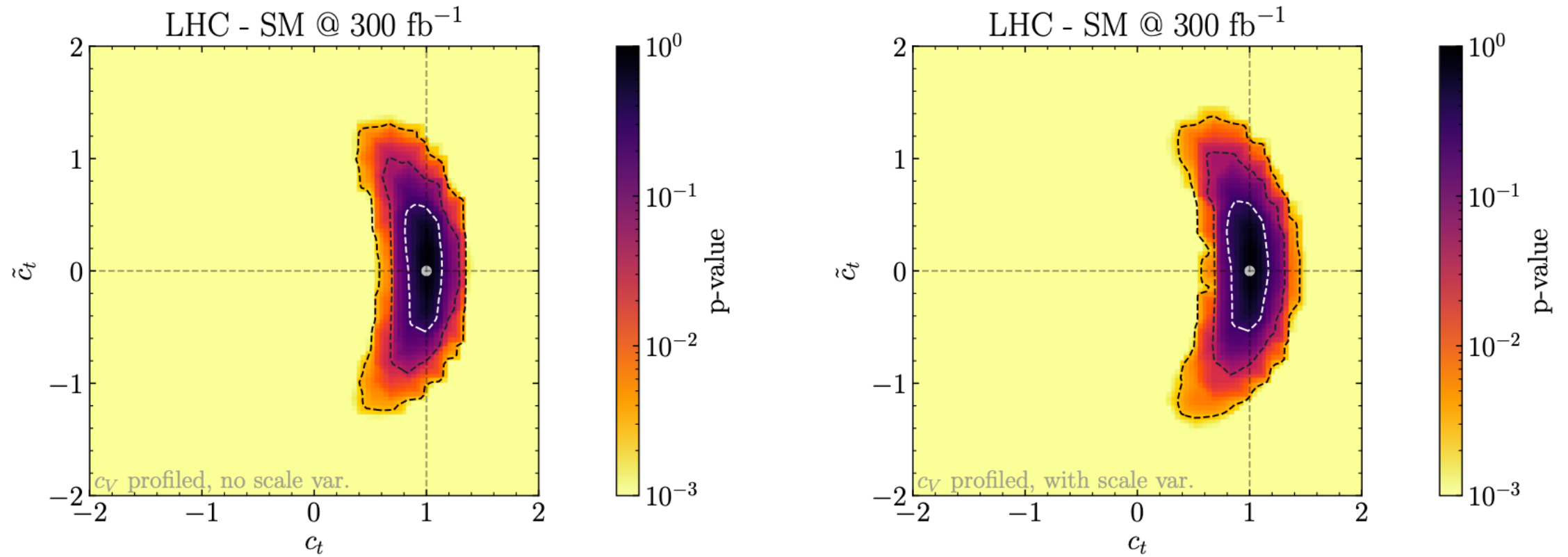
observable	condition
N_γ	≥ 2 (with $ \eta < 2.5$ and $p_T > 25$ GeV)
$(p_{T,1}^\gamma, p_{T,2}^\gamma)$	$\geq (35, 25)$ GeV
$m_{\gamma\gamma}$	$[105 - 160]$ GeV
$(p_{T,1}^\gamma/m_{\gamma\gamma}, p_{T,2}^\gamma/m_{\gamma\gamma})$	$\geq (0.35, 0.25)$
N_ℓ	≥ 1 (with $ \eta < 2.5$ and $p_T > 15$ GeV)
$m_{\ell\ell}$	$[80, 100]$ GeV vetoed if same flavour
N_{jet}	≥ 1 (with $ \eta < 2.5$ and $p_T > 25$ GeV)

Table 1: Summary of preselection cuts.

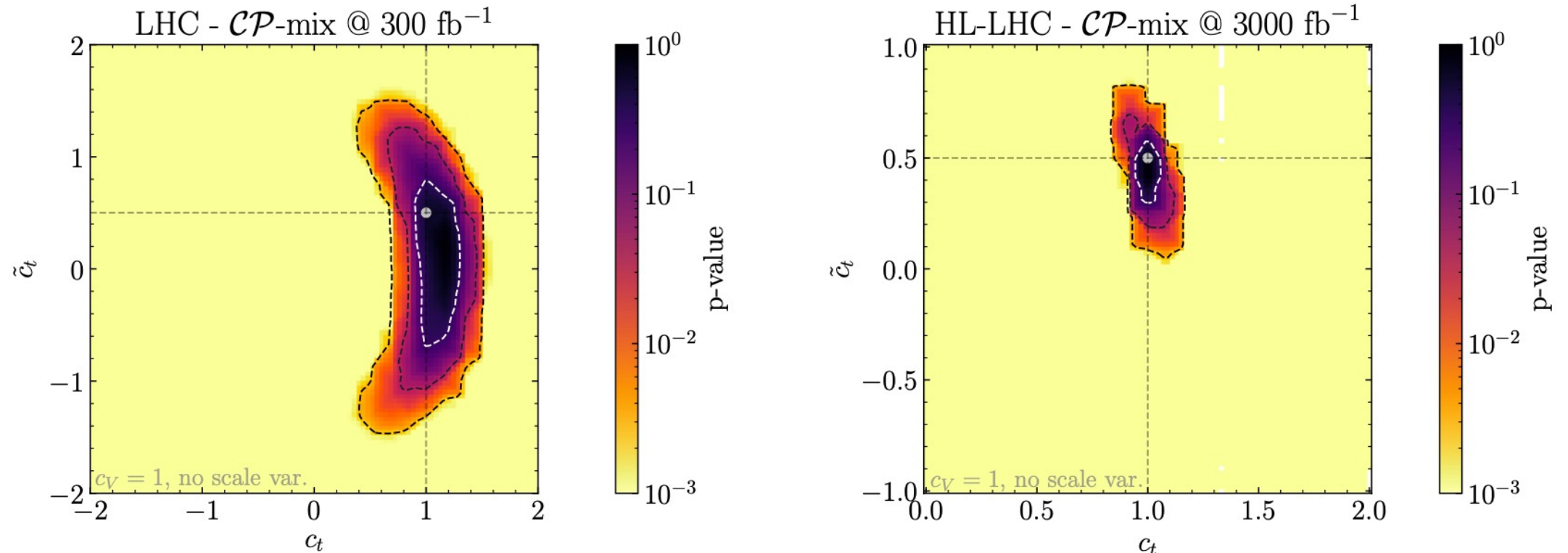
Interpretation in terms of CP-violating angle



Variation of c_V and renormalization scale

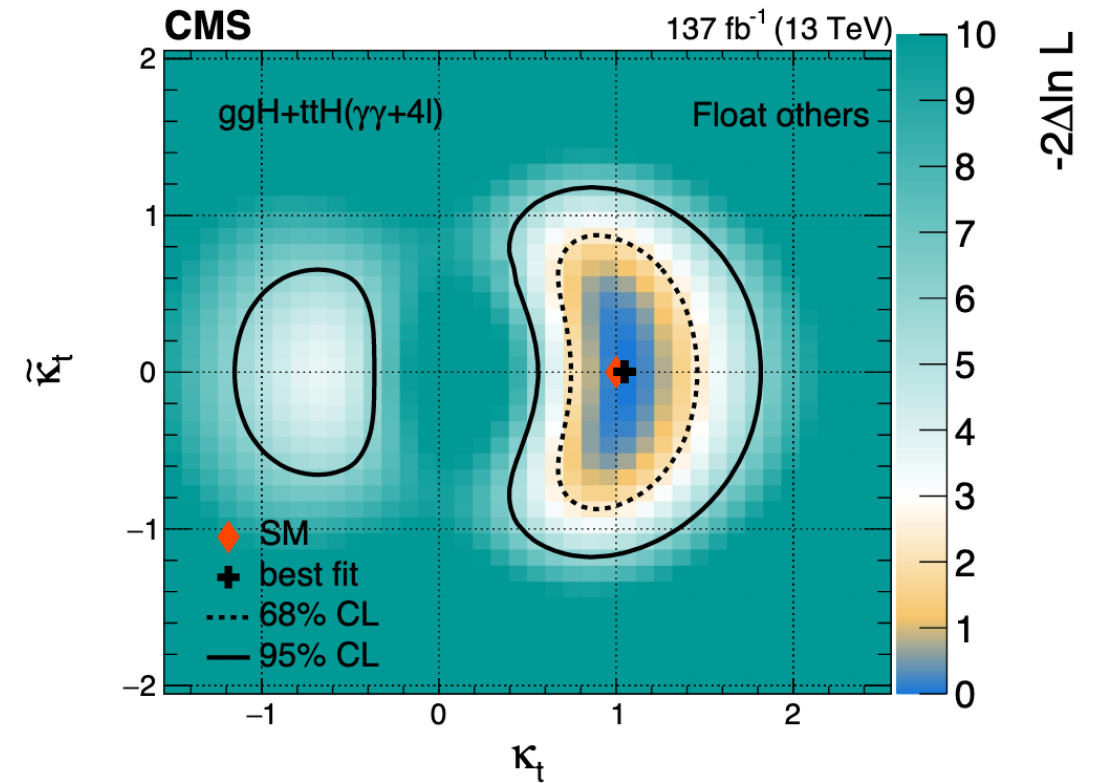
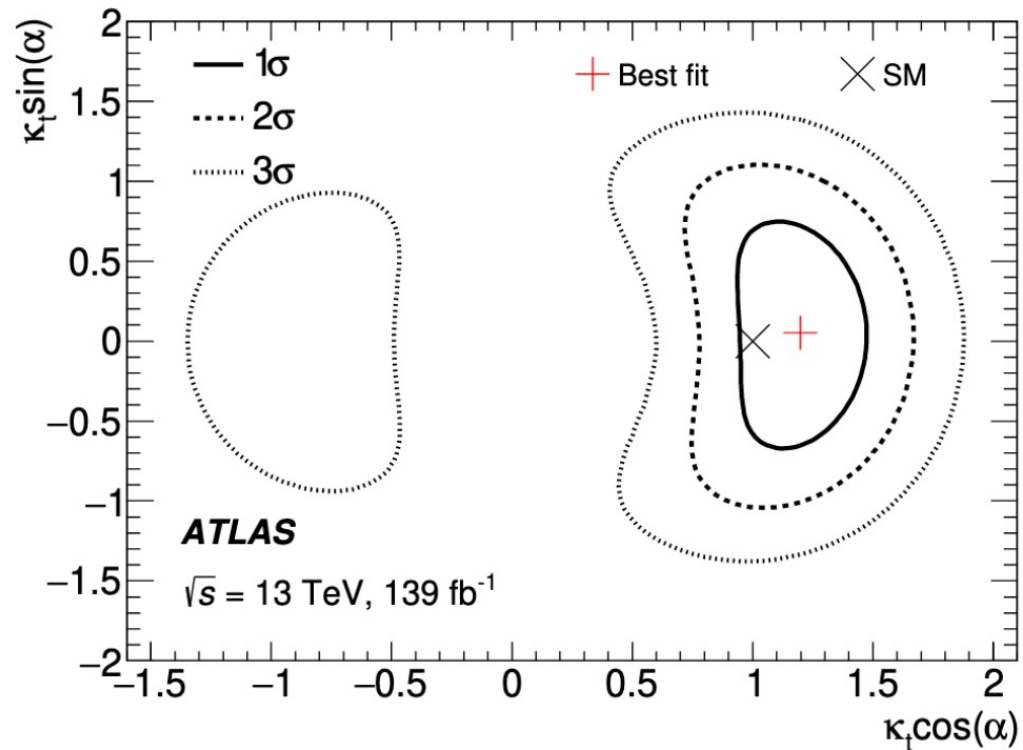


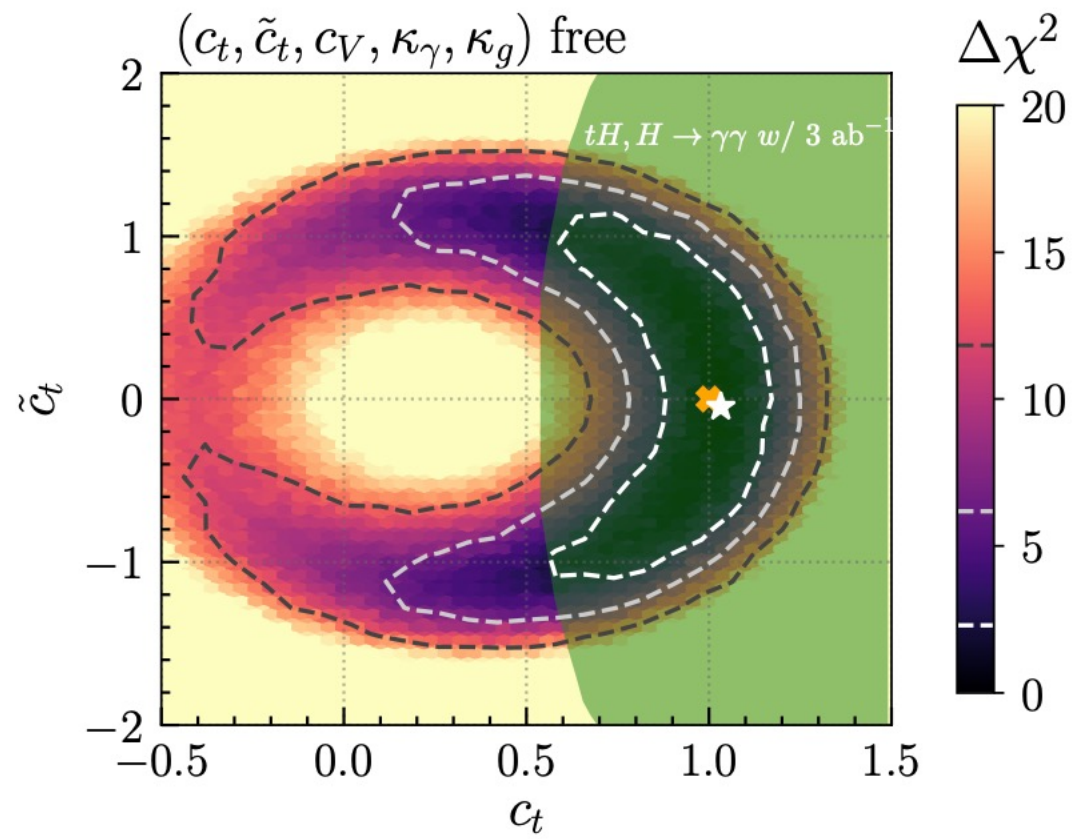
Limits in case of deviation from SM



- CP-mix: $c_t = 1, c_{\tilde{t}} = 0.5, c_V = 1.$

Experimental studies [ATLAS,2004.04545;CMS,2104.12152]





Which observables drive these constraints?

- Use Fisher matrix to evaluate information for different observables

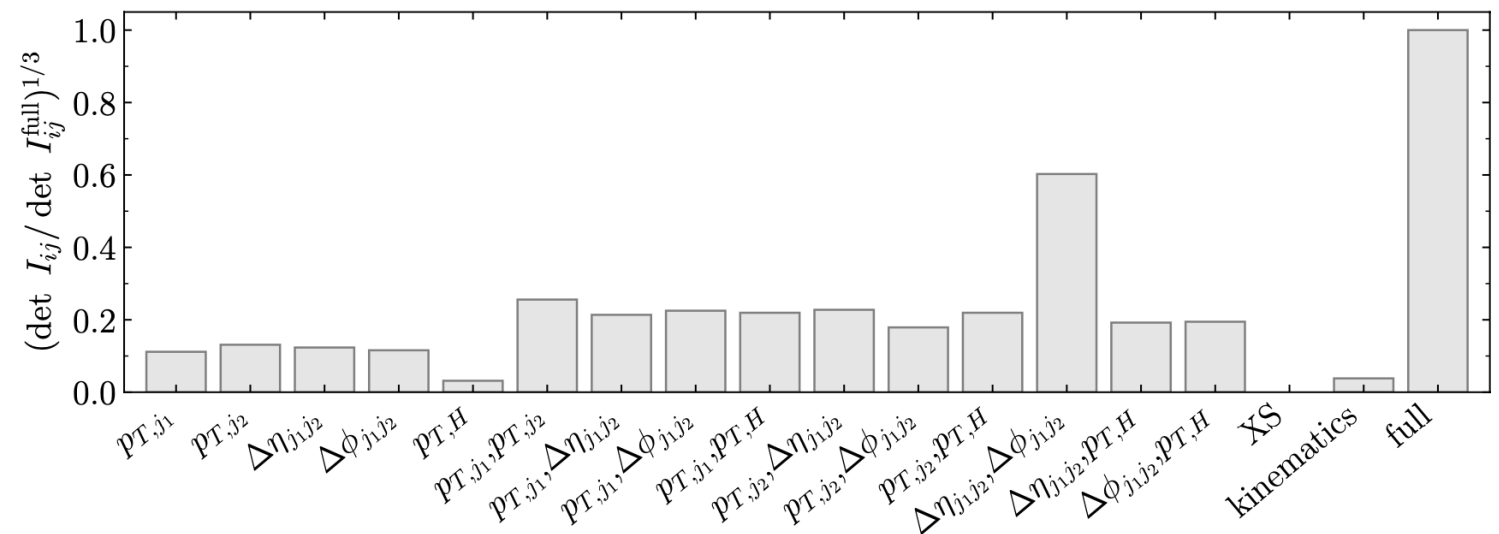
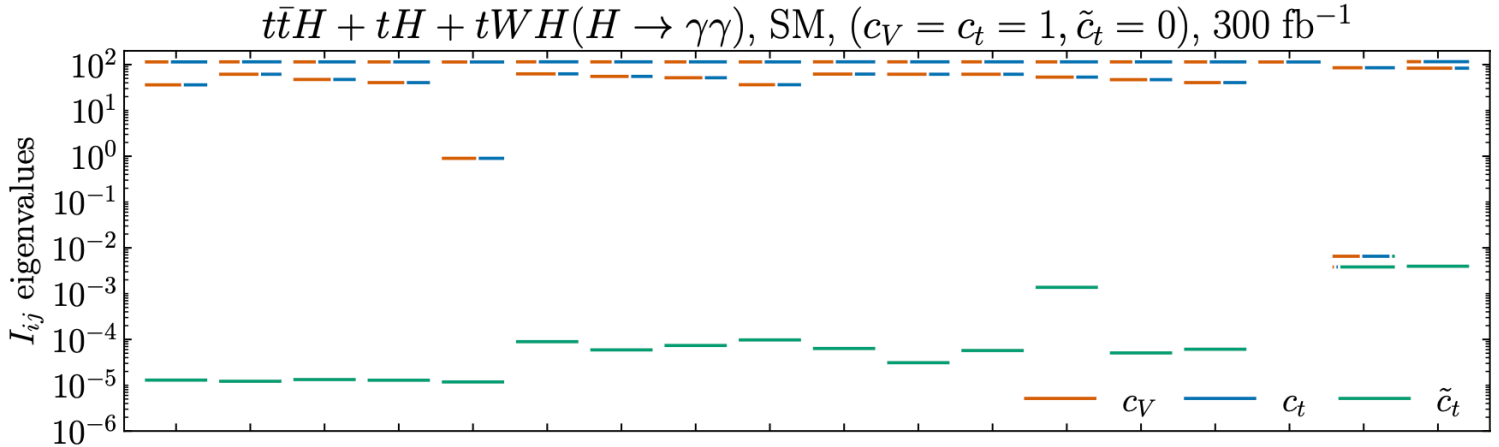
$$I_{ij}(\theta) = \mathbb{E} \left[\frac{\partial \log p_{\text{full}}(\{x\}|\theta)}{\partial \theta_i} \frac{\partial \log p_{\text{full}}(\{x\}|\theta)}{\partial \theta_j} \Big|_{\theta} \right], \quad \text{with} \quad \text{cov}(\hat{\theta}|\theta)_{ij} \geq I_{ij}^{-1}(\theta),$$

- E.g., for SM point we have

$$I_{ij}^{\text{full}}(\text{SM}) \simeq \begin{pmatrix} 91.4 & 13.7 & 0.1 \\ 13.7 & 108.2 & -0.1 \\ 0.1 & -0.1 & 0.004 \end{pmatrix}, \quad \text{with the parameter space spanned by} \quad \begin{pmatrix} c_V \\ c_t \\ c_{\tilde{t}} \end{pmatrix}$$

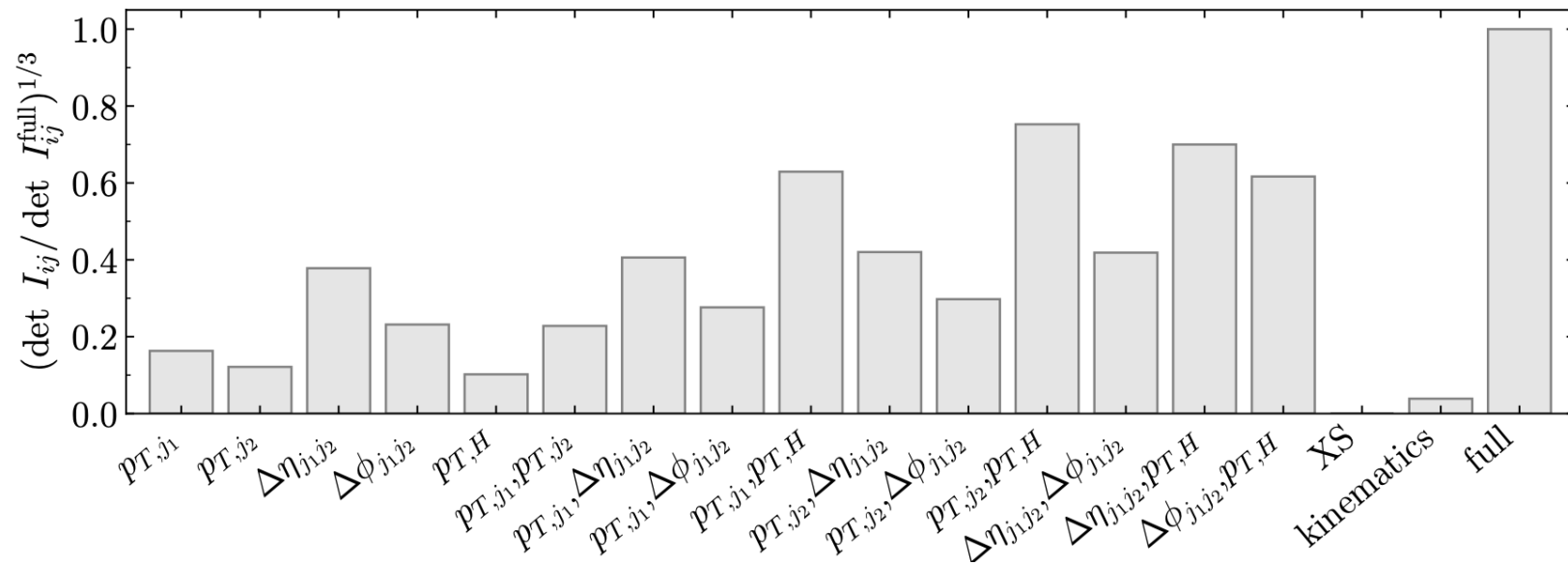
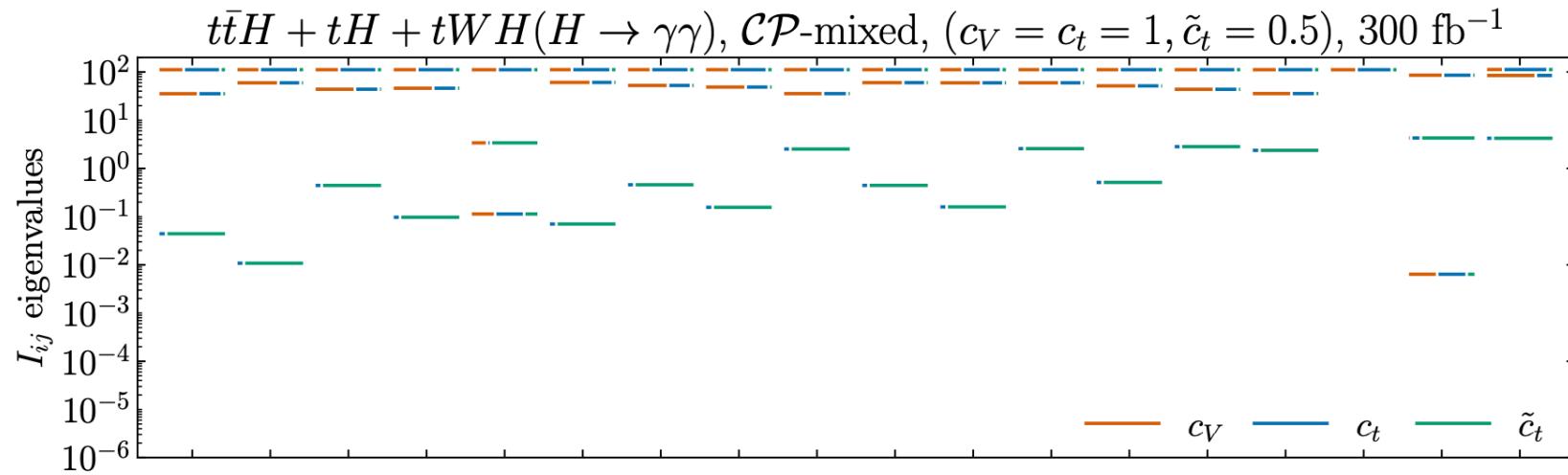
- Evaluate Fisher matrix for various 1D and 2D histograms, full likelihood, XS only, kinematics only.

Fisher information for SM scenario

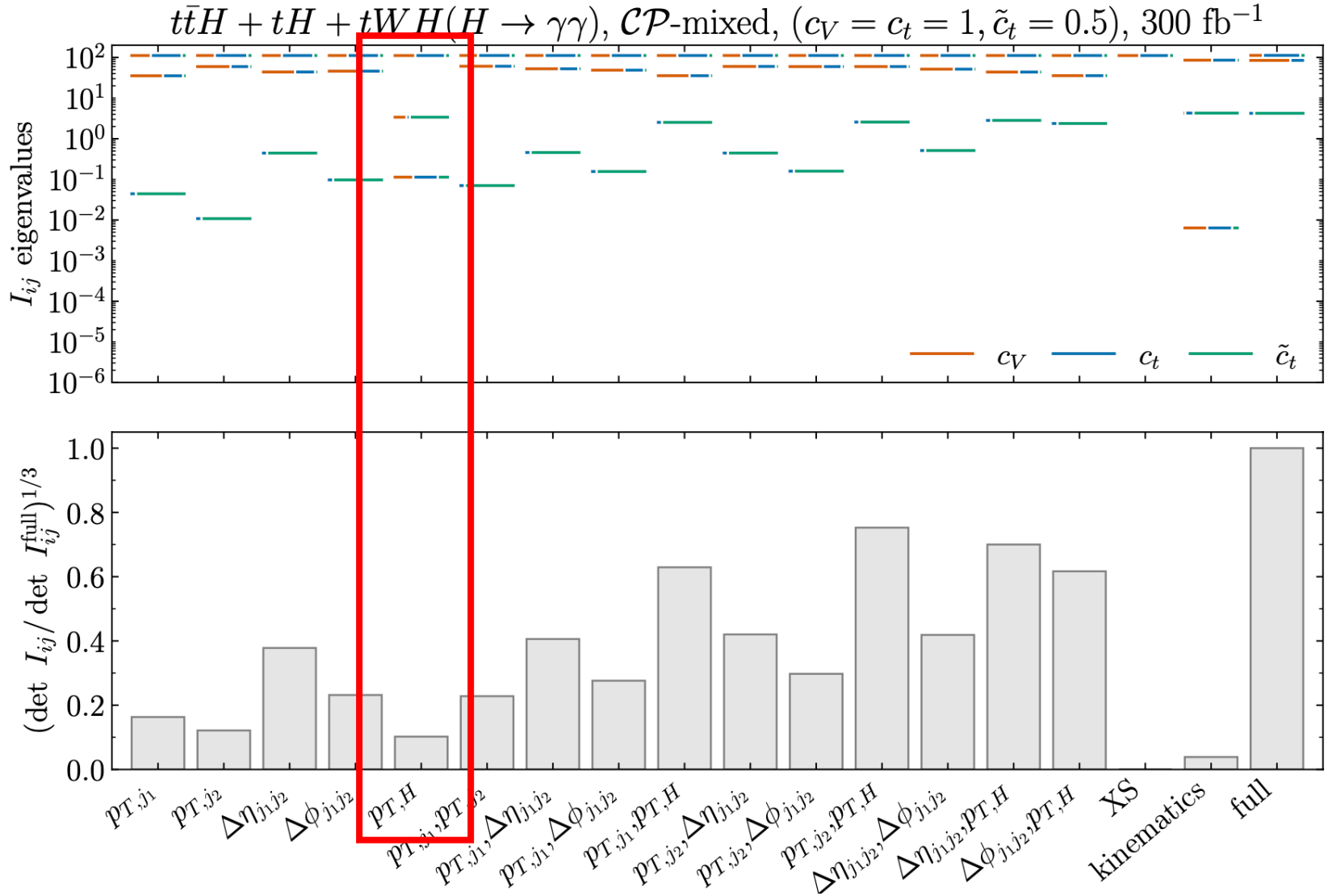


- $c_{\tilde{t}}$ not constrained by rate.
- Use of kinematic information mandatory.
- No single observable able to capture information about $c_{\tilde{t}}$.

Fisher information for CP-mixed scenario



Fisher information for CP-mixed scenario



- For CP-mixed scenario, Higgs p_T captures sizeable amount of information on $c_{\tilde{t}}$.
- ↓
- p_T binned STXS measurements useful to constrain CP violation in the top-Yukawa coupling.