

THERMODYNAMIC OF INTEGRABLE  
 $N=2$  THEORIES, SQUARED

BASED ON 2502.10356

WITH X. KERVYN & A. SFONDRINI

SUPERSTRINGS PROPAGATING IN  $AdS_3 \times S^3 \times T^4$  ARE  
CLASSICALLY INTEGRABLE [A. CAGNAZZO, K. ZAREMBO '12]

$$S = \int d\tau d\sigma \left[ -\frac{T}{2} \sqrt{-\gamma} \gamma^{ab} G_{\mu\nu} + \frac{K}{2\pi} \epsilon^{ab} B_{\mu\nu} \right] \partial_a X^\mu \partial_b X^\nu$$

+ FERMIONS

THERE IS A FAMILY OF SUCH THEORIES SPANNED BY

$T \geq 0$  STRING TENSION

$K = 0, 1, 2$  QUANTISED AMOUNT OF B FIELD

$$\left( \text{psu}(1|1)_L \oplus \text{psu}(1|1)_R \right)_{\text{c.e.}}^{\oplus 2}$$

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SUSY  
ALGEBRA

$$\{Q, S\} = \frac{1}{2}(E + M)$$

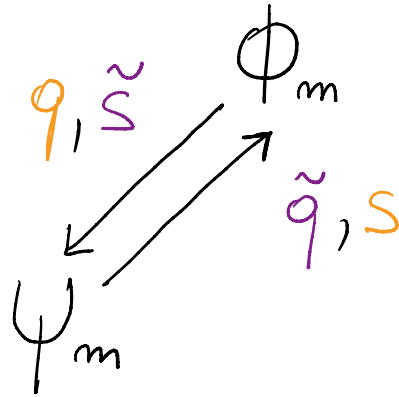
$$\{\tilde{Q}, \tilde{S}\} = \frac{1}{2}(E - M)$$

$$\{Q, \tilde{Q}\} = C \quad \{S, \tilde{S}\} = C^\dagger$$

CENTRAL  
EXTENSION

$\Phi_m$ : BOSON

$\Psi_m$ : FERMION



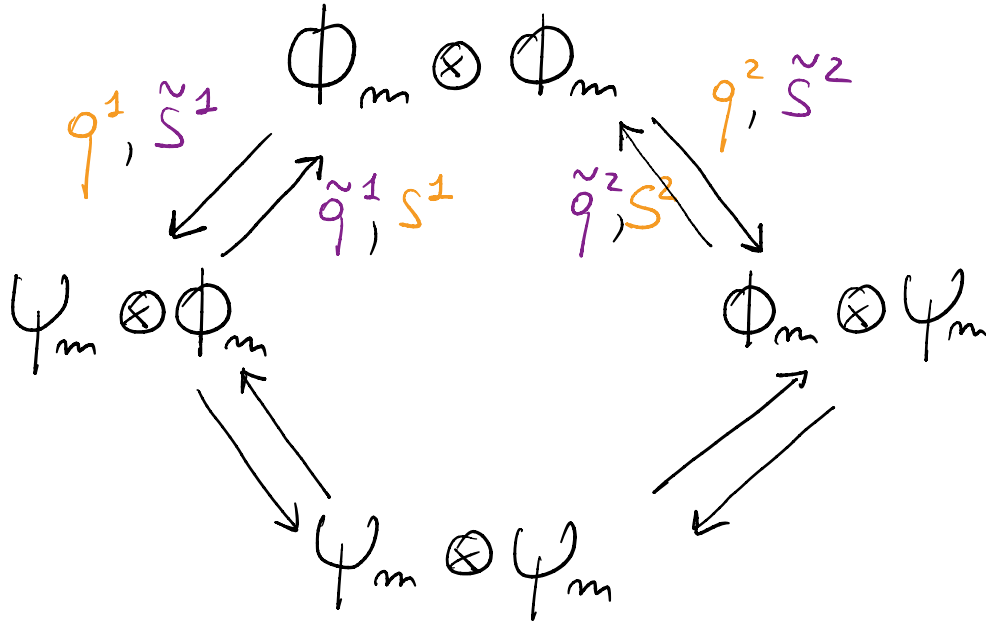
$$E_m(p) = \sqrt{\left(m + \frac{\kappa}{2\pi} p\right)^2 + 4h^2 \sin^2\left(\frac{p}{2}\right)}$$

$$m = 1, 2, \dots$$

$$\kappa \in \mathbb{N}$$

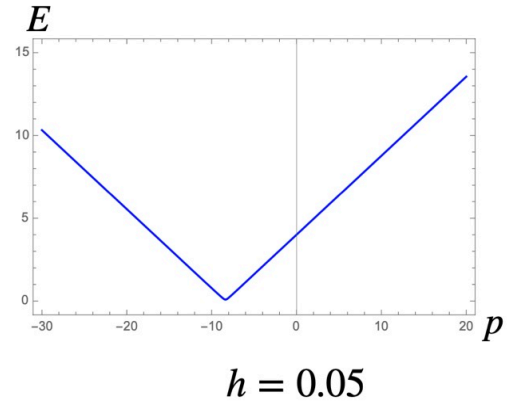
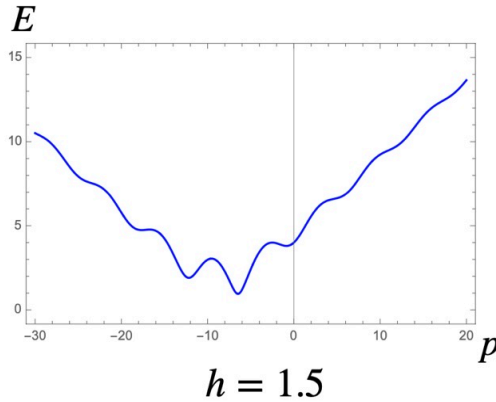
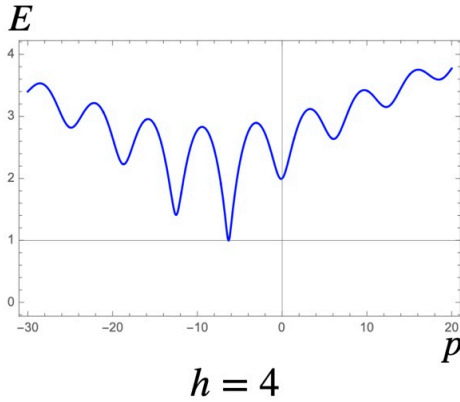
$$h \geq 0 \quad (\sim \text{'t Hooft coupling})$$

$\oplus 2$   
 $\Rightarrow$

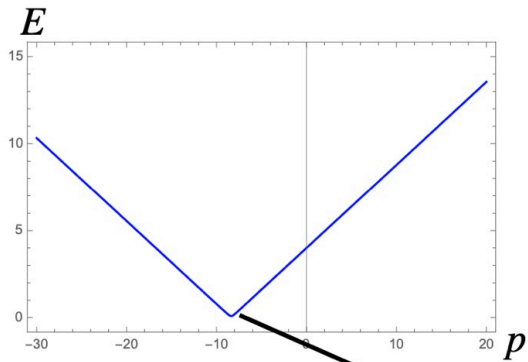


$$q^1 = q \otimes \mathbb{1} \quad , \quad q^2 = \mathbb{1} \otimes q, \quad \tilde{q}^1 = \tilde{q} \otimes \mathbb{1} \quad \dots$$

$$E_m(p) = \sqrt{\left(m + \frac{\kappa}{2\pi} p\right)^2 + 4h^2 \sin^2\left(\frac{p}{2}\right)}$$



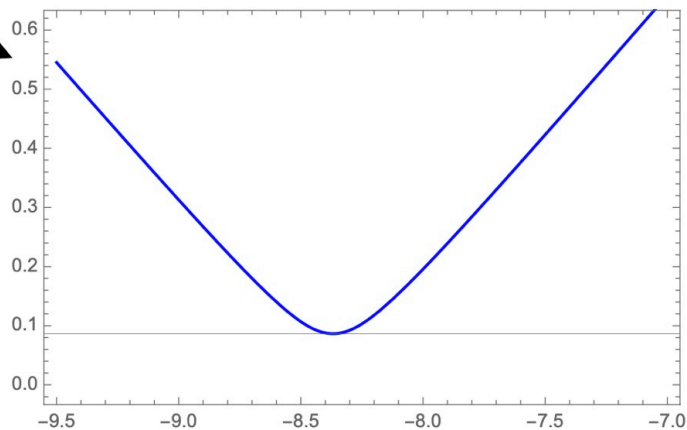
If  $h \ll 1 \Rightarrow p_{\text{MIN}} \approx -\frac{2\pi}{\kappa} m$



$$h = 0.05$$

Zoom in

$$p = -\frac{2\pi}{k}m + hq$$

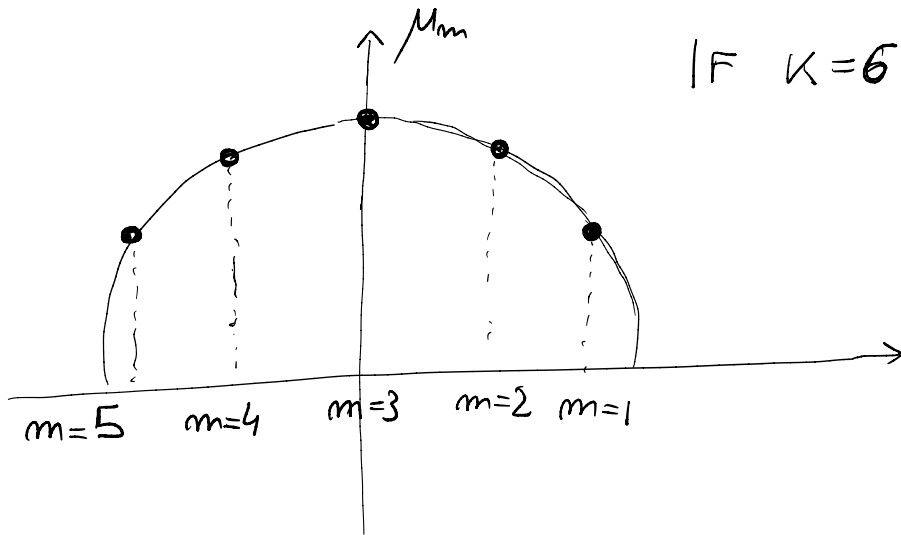


$$E_m(P) \rightarrow \mu_m \operatorname{ch} \theta$$

$\theta$ : RAPIDITY OF  
A RELATIVISTIC  
THEORY

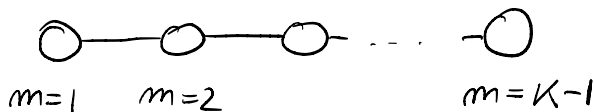
$$\mu_m = 2h \sin\left(\frac{\pi}{K} m\right)$$

$$m = 1, 2, \dots, K-1$$



SPECTRUM OF INTEGRABLE TODA

THEORIES OF  $A_{K-1}$  TYPE





# ZF ALGEBRA

$$\Phi_m = (1, 0) \quad \Psi_m = (0, 1)$$

$$S_{m_1, m_2}(\theta) = f_{m_1, m_2}(\theta) \begin{pmatrix} A(\theta) & 0 & 0 & 0 \\ 0 & C(\theta) & D(\theta) & 0 \\ 0 & B(\theta) & E(\theta) & 0 \\ 0 & 0 & 0 & F(\theta) \end{pmatrix}$$

DRESSING FACTOR

$$A(\theta) = 1, \quad B(\theta) = \frac{\text{sh} \left( \frac{\theta}{2} - \frac{i\pi}{2K} (m_1 - m_2) \right)}{\text{sh} \left( \frac{\theta}{2} + \frac{i\pi}{2K} (m_1 + m_2) \right)}, \quad \dots$$

$$f_{m_1, m_2}(\theta) = \frac{R\left(\theta - \frac{i\pi}{k}(m_1 + m_2)\right) R\left(\theta + \frac{i\pi}{k}(m_1 + m_2)\right)}{R\left(\theta - \frac{i\pi}{k}(m_1 - m_2)\right) R\left(\theta + \frac{i\pi}{k}(m_1 - m_2)\right)}$$

$$R(\theta) = \frac{\Gamma\left(1 - \frac{\theta}{2\pi i}\right)}{\Gamma\left(1 + \frac{\theta}{2\pi i}\right)}$$

$$\uparrow$$

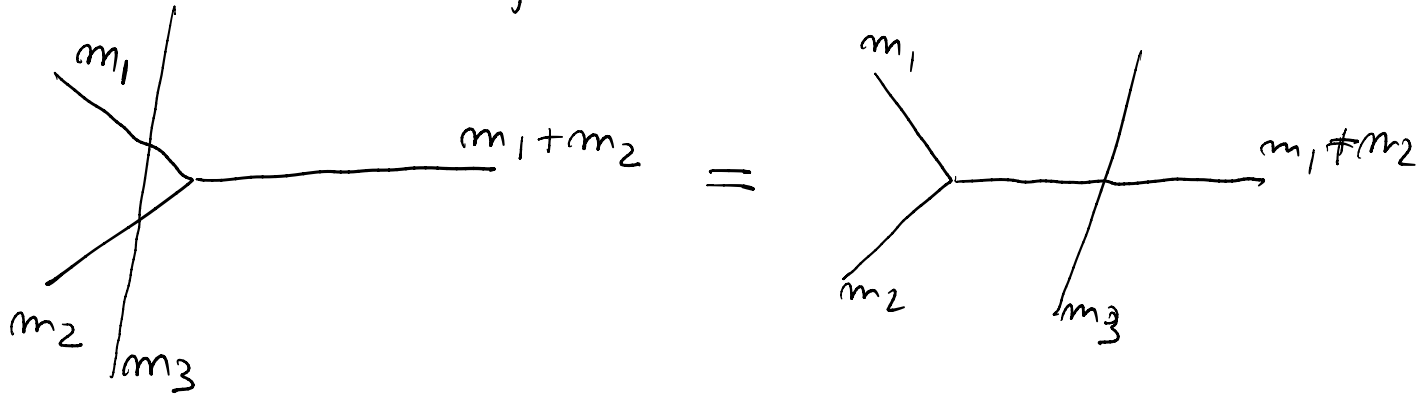
BARNES GAMMA FUNCTIONS

+ CDD

$f$  SATISFIES ALL THE PROPERTIES OF THE BOOTSTRAP.

•  $f_{m_1, m_2}(\theta) f_{m_2, m_1}(-\theta) = 1$

•  $f_{m_1, m_3}\left(\theta - \frac{i\pi}{k} m_2\right) f_{m_2, m_3}\left(\theta + \frac{i\pi}{k} m_1\right) = f_{m_1+m_2, m_3}(\theta)$



## WHAT WE DID IN THE PAPER

- PUT THE SYSTEM AT FINITE VOLUME AND DERIVE ITS TBA EQUATIONS AND  $Y$  SYSTEM.

$Y_p$        $p = 1, \dots, k-1.$       PHYSICAL  $Y$  FUNCTIONS

$Y_a$        $a = \uparrow, \downarrow.$       AUXILIARY  $Y$  FUNCTIONS

- SEARCH FOR CONSTANT SOLUTIONS TO  $Y$  SYSTEM  
( $\Rightarrow$  UV CENTRAL CHARGE)

IF SPECTRUM  $\phi_P, \psi_P$   $P=1, \dots, K-1$

$$Y_P^2 = \frac{(1 + Y_{P+1})(1 + Y_{P-1})}{\prod_{a=\uparrow, \downarrow} \left(1 + \frac{1}{Y_a}\right)^{1 \cdot \bar{I}_{Pa}}}$$

$$\bar{I}_{Pa} = \delta_{P1} \delta_{a\uparrow} + \delta_{P, K-1} \delta_{a\downarrow}$$

$$Y_{\uparrow}^2 = (1 + Y_1)^{-1} \left(1 + \frac{1}{Y_{\uparrow}}\right)^{-1} \left(\frac{1 + \frac{1}{Y_{\uparrow}}}{1 + \frac{1}{Y_{\downarrow}}}\right)^{\frac{1}{K}}$$

$$Y_{\downarrow}^2 = (1 + Y_{K-1})^{-1} \left(1 + \frac{1}{Y_{\downarrow}}\right)^{-1} \left(\frac{1 + \frac{1}{Y_{\downarrow}}}{1 + \frac{1}{Y_{\uparrow}}}\right)^{\frac{1}{K}}$$

THEN

$$Y_P = \frac{\Delta \text{im} \left( \frac{(2P+3)\pi}{2(k+1)} \right) \Delta \text{im} \left( \frac{(2P-1)\pi}{2(k+1)} \right)}{\Delta \text{im}^2 \left( \frac{\pi}{k+1} \right)}$$

$$Y_{\uparrow} = Y_{\downarrow} = \frac{\Delta \text{im} \left( \frac{\pi}{2(k+1)} \right)}{\Delta \text{im} \left( \frac{3\pi}{2(k+1)} \right)}$$

$$C = 3 \frac{(k-1)}{(k+1)}$$

$A_{k-1}$  TYPE  $N=2$  MINIMAL  
MODELS PERTURBED BY  
MOST RELEVANT OPERATOR

[P. FENDLEY, K. INTRILIGATOR '92]

# IF SPECTRUM $\Phi_P \otimes \Phi_P, \Phi_P \otimes \Psi_P, \Psi_P \otimes \Phi_P, \Psi_P \otimes \Psi_P$

$$Y_P^2 = \frac{(1 + Y_{P+1})(1 + Y_{P-1})}{\prod_{a=\uparrow, \downarrow} \left(1 + \frac{1}{Y_a}\right)^{2\bar{I}_{Pa}}}$$

$$\bar{I}_{Pa} = d_{P,1} d_{a\uparrow} + d_{P,K-1} d_{a\downarrow}$$

$$Y_{\uparrow}^2 = (1 + Y_1)^{-1} \left(1 + \frac{1}{Y_{\uparrow}}\right)^{-2} \left(\frac{1 + \frac{1}{Y_{\uparrow}}}{1 + \frac{1}{Y_{\downarrow}}}\right)^{\frac{2}{K}}$$

$$Y_{\downarrow}^2 = (1 + Y_{K-1})^{-1} \left(1 + \frac{1}{Y_{\downarrow}}\right)^{-2} \left(\frac{1 + \frac{1}{Y_{\downarrow}}}{1 + \frac{1}{Y_{\uparrow}}}\right)^{\frac{2}{K}}$$

THEN

$$Y_P = \frac{\sin\left(\frac{\pi}{K}(P-1)\right) \sin\left(\frac{\pi}{K}(P+1)\right)}{\sin^2\frac{\pi}{K}} \quad P=1, \dots, K-1$$

$$Y_{\uparrow} = Y_{\downarrow} = 0$$

$$C = 6 \left( 1 - \frac{1}{K} \right)$$

WHAT DOES THIS CORRESPOND  
TO?



# OUTLOOK AND REMARKS

• DOES THE TBA CONTAIN ANY INFORMATION ON THE THEORY BEFORE THE LIMIT ?

• FINITE VOLUME = DOUBLE WICK ROTATION

$$p \rightarrow i \tilde{E}$$

$$E \rightarrow i \tilde{p}$$

FOR A RELATIVISTIC THEORY:  $E^2 - p^2 = m^2 \rightarrow \tilde{E}^2 - \tilde{p}^2 = m^2$

$$\theta \rightarrow \theta + \frac{i\pi}{2}$$

- S MATRIX OF DIFFERENCE FORM

$$\Rightarrow S(\theta_1 - \theta_2) \longrightarrow S(\theta_1 - \theta_2)$$

- IN THE FULL MODEL THE S-MATRIX IS NOT INVARIANT.

DOES THE LIMIT COMMUTE WITH DOUBLE WICK ROTATION?

- CAN WE PERFORM THE SAME LIMIT ON SIMILAR MODELS AS

$$AdS_3 \times S^3 \times S^3 \times S^4 \quad ?$$

THANK you