Electroweak B+L Violation at High Energies.

Part 2: Rate Estimates and Signatures

Andreas Ringwald

BSM Physics Opportunities at 100 TeV CERN, Geneva, February 10-11, 2014





Recap

Cross-sections for exclusive B+L violation rapidly growing below

$$4\pi \frac{M_W}{\alpha_W} \simeq 30 \text{ TeV}$$

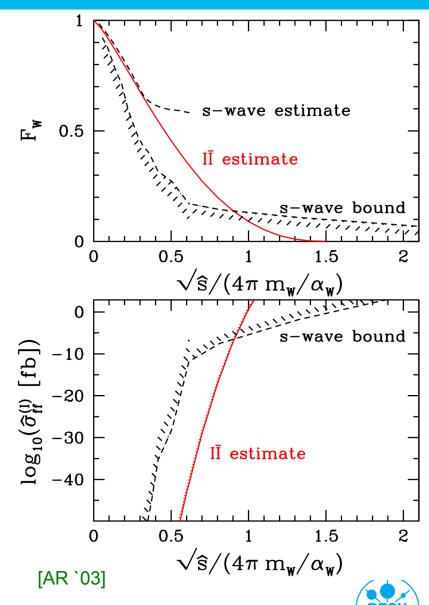
Total cross-section grows exponentially; dominated by multiple emission of EW gauge bosons

$$\hat{\sigma}_{ ext{ff}}^{(I_W)} ~pprox ~rac{1}{m_W^2} \left(rac{2\pi}{lpha_W}
ight)^{7/2} \, \mathrm{e}^{-rac{4\pi}{lpha_W} F_W(\epsilon)}$$

$$\epsilon \equiv \sqrt{\hat{s}}/(4\pi m_W/\alpha_W) \simeq \sqrt{\hat{s}}/(30 \text{ TeV})$$

$$F_W(\epsilon) = 1 - \frac{3^{4/3}}{2} \epsilon^{4/3} + \frac{3}{2} \epsilon^2 + \mathcal{O}(\epsilon^{8/3}).$$

- > For $\epsilon > 0.3 \div 0.75$ only estimates, educated guesses and bounds
- Need future hadron collider or search for analogous QCD processes



Recap

Cross-sections for exclusive B+L violation rapidly growing below

$$4\pi \frac{M_W}{\alpha_W} \simeq 30 \text{ TeV}$$

Total cross-section grows exponentially; dominated by multiple emission of EW gauge bosons

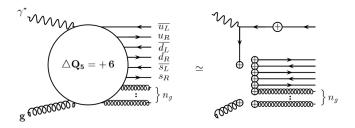
$$\hat{\sigma}_{
m ff}^{(I_W)} \;\; pprox \;\; rac{1}{m_W^2} \left(rac{2\pi}{lpha_W}
ight)^{7/2} \, {
m e}^{-rac{4\pi}{lpha_W} F_W(\epsilon)}$$

$$\epsilon \equiv \sqrt{\hat{s}}/(4\pi m_W/\alpha_W) \simeq \sqrt{\hat{s}}/(30 \text{ TeV})$$

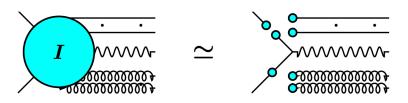
$$F_W(\epsilon) = 1 - \frac{3^{4/3}}{2} \epsilon^{4/3} + \frac{3}{2} \epsilon^2 + \mathcal{O}(\epsilon^{8/3}).$$

- > For $\epsilon > 0.3 \div 0.75$ only estimates, educated guesses and bounds
- Need future hadron collider or search for analogous QCD processes

- Small-size QCD-instanton induced processes in
 - DIS



virtual vector boson production





Instanton-Antiinstanton Valley Cross-Section Estimate

> Total cross-section via optical theorem [V.V. Khoze, AR `91; AR, F. Schrempp `98]

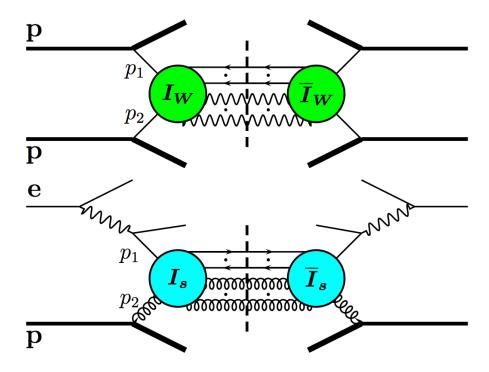
$$\hat{\sigma}_{\mathbf{p}_{1}\mathbf{p}_{2}}^{(I)} \sim \int d^{4}R \int_{0}^{\infty} d\rho \int_{0}^{\infty} d\overline{\rho} \, \mathbf{D}(\rho) \mathbf{D}(\overline{\rho}) \int dU e^{-\frac{4\pi}{\alpha g}\Omega\left(U,\frac{R^{2}}{\rho\overline{\rho}},\ldots\right)} \, e^{\mathrm{i}(p_{1}+p_{2})\cdot R - \sum_{i=1}^{2}\sqrt{-p_{i}^{2}} \, (\rho + \overline{\rho})}$$

- Ingredients:
 - Instanton-size distribution

$$D(\rho) \propto \mathrm{e}^{-2\pi/\alpha_g}$$

Instanton-Antiinstanton interaction

$$\Omega\left(U,R^2/(
hoar
ho),\ldots
ight)$$





Instanton-Antiinstanton Valley Cross-Section Estimate

Total cross-section via optical theorem

[Khoze,AR '91; AR, F. Schrempp '98]

$$\hat{\sigma}_{\mathbf{p}_{1}\mathbf{p}_{2}}^{(I)} \sim \int \mathrm{d}^{4}R \int_{0}^{\infty} \mathrm{d}\rho \int_{0}^{\infty} \mathrm{d}\overline{\rho} \, D(\rho)D(\overline{\rho}) \int \mathrm{d}U \mathrm{e}^{-\frac{4\pi}{\alpha g}\Omega\left(U,\frac{R^{2}}{\rho\overline{\rho}},\ldots\right)} \, \mathrm{e}^{\mathrm{i}(p_{1}+p_{2})\cdot R-\sum_{i=1}^{2}\sqrt{-p_{i}^{2}}(\rho+\overline{\rho})}$$

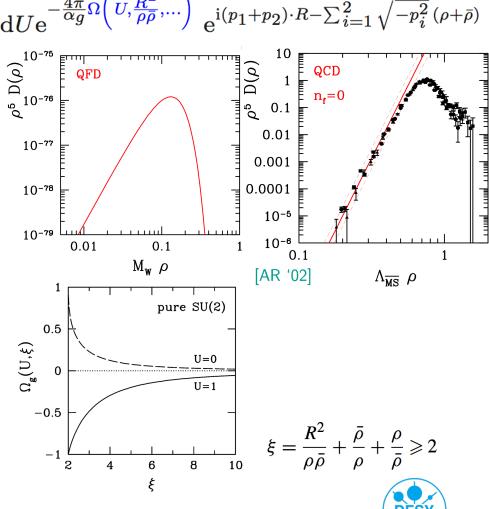
Ingredients:

Instanton-size distribution

$$D(\rho) \propto \mathrm{e}^{-2\pi/\alpha g}$$

Instanton-Antiinstanton interaction

$$\Omega\left(U,R^2/(\rho\bar{\rho}),\ldots\right)$$



Instanton-Antiinstanton Valley Cross-Section Estimate

Saddle point evaluation:

$$\hat{\sigma}^{(I)} \propto e^{-\Gamma_*} \equiv e^{-\frac{4\pi}{\alpha g} F_g(\epsilon)}$$

where

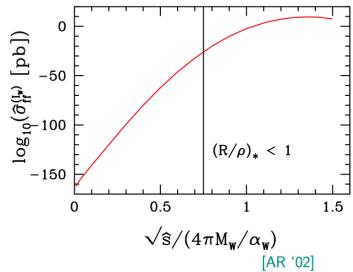
$$\epsilon \equiv \begin{cases} \sqrt{\hat{s}}/(4\pi M_W/\alpha_W) & (\text{QFD}) \\ \sqrt{\hat{s}}/Q' \equiv \sqrt{1/x'-1} & (\text{QCD}) \end{cases}$$

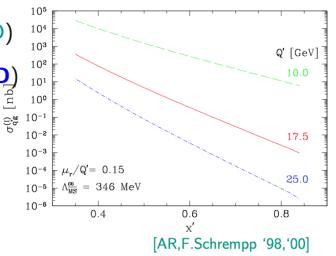
is a scaled cm energy and

$$F_g = 1 + \Omega_g(1, \xi_*) +$$

$$\begin{cases} -(\xi_* - 2) \frac{\partial}{\partial \xi_*} \Omega_g(1, \xi_*)_{|\xi_* = 2 + \left(\frac{R}{\rho}\right)_*^2} & \text{(QFD)} \\ 0 & \text{QCD} \end{cases}$$

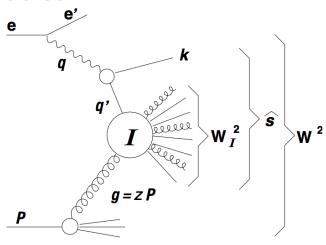
• Increasing $\epsilon \Rightarrow$ smaller $(R/\rho)_*$ probed \Rightarrow cross-section grow due to attractive nature of Ω_g in perturbative semiclassical regime







Kinematics:



Deep-inelastic scattering variables:

$$S = (e + P)^{2}$$

$$Q^{2} = -q^{2} = -(e - e')^{2}$$

$$x_{\mathrm{Bj}} = Q^{2} / (2P \cdot q)$$

$$y_{\mathrm{Bj}} = Q^{2} / (Sx_{\mathrm{Bj}})$$

$$W^{2} = (q + P)^{2} = Q^{2}(1/x_{\mathrm{Bj}} - 1)$$

$$\hat{s} = (q + g)^{2}$$

$$z = x_{\mathrm{Bj}} (1 + \hat{s}/Q^{2})$$

Variables of instanton-subprocess:

$$Q'^2 = -q'^2 = -(q - k)^2$$

 $x' = Q'^2 / (2 g \cdot q')$
 $W_I^2 = (q' + g)^2 = Q'^2 (1/x' - 1)$

"Fiducial" kinematical region from lattice constraints: [AR,F.Schrempp '99;'01]

$$\left(\rho^* \Lambda_{\overline{\text{MS}}}^{(0)} \lesssim 0.4, \frac{R^*}{\rho^*} \gtrsim 1.0\right) \Rightarrow \left(Q' / \Lambda_{\overline{\text{MS}}}^{(n_f)} \gtrsim 30.8, x' \gtrsim 0.35\right)$$



Event generator QCDINS 2.0:

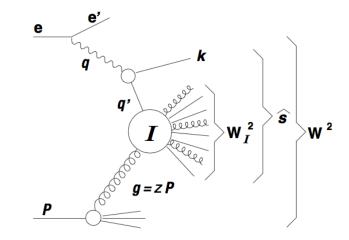
[Gibbs, AR, F. Schrempp '95; AR, F. Schrempp '00]

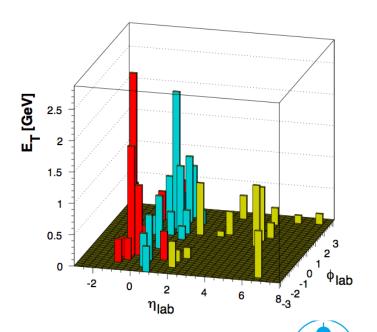
Hard subprocess:

- isotropic in q'g CM
- flavour democratic
- large parton multiplicity

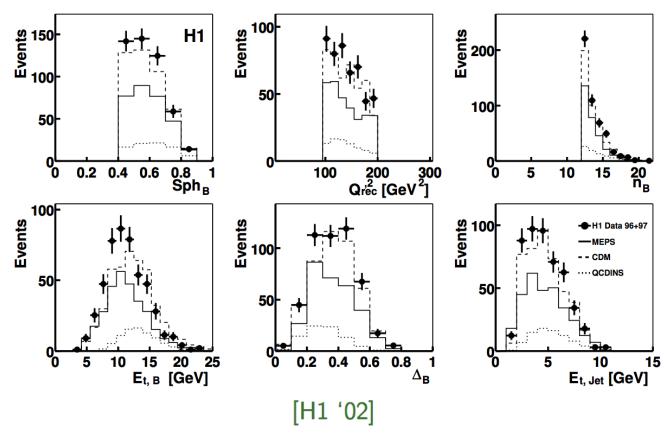
$$\langle n_q + n_g \rangle = 2 n_f - 1 + \mathcal{O}(1)/\alpha_s \gtrsim 8,$$

- Parton shower (HERWIG)
- Hadronization (HERWIG or JET-SET)



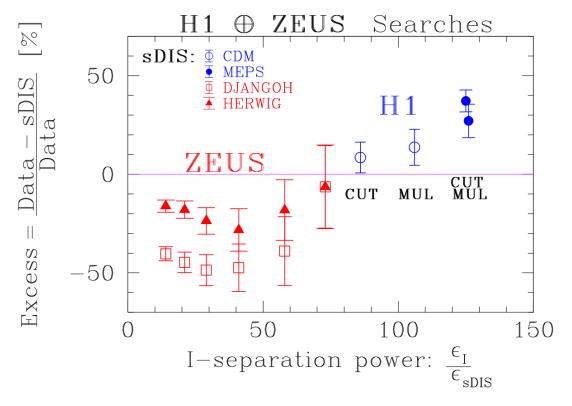


- Instanton-enriched samples by cuts on discriminating observables
- Large uncertainties in predictions of standard DIS processes





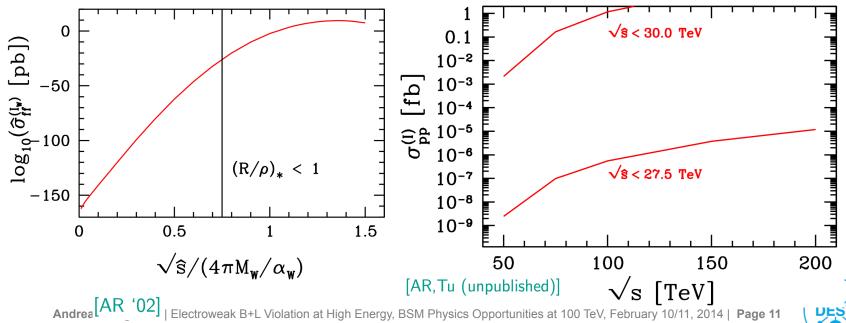
- Instanton-enriched samples by cuts on discriminating observables
- Large uncertainties in predictions of standard DIS processes
- H1/ZEUS "excess" increases with separation power (ratio of efficiencies)





QFD-Instanton Induced Processes at Future Hadron Coll.

- H1/ZEUS limits on QCD-instanton induced processes suggest:
 - Instanton-antiinstanton estimate reliable, as long as $(R/\rho)_* \geq 1$
 - For $(R/\rho)_* < 0.5 \div 1$, rapid growth, as implied by Ω , stops.
- Implications for QFD-instantons:
 - $(R/\rho)_*$ < $0.5 \div 1$ corresponds to $\epsilon < 0.75 \div 1.15$, $\sqrt{\hat{s}} < 22 \div 35$ TeV
 - At these energies, cross-section estimate reaches observable values



QFD-Instanton Induced Processes at Future Hadron Coll.

Phenomenology of QFD-instantons

[AR,F.Schrempp,Wetterich '91; Gibbs,AR,Webber,Zadrozny '94]

- No background from perturbative Standard Model processes by requiring
 - ≥ 4 identified charged e's or μ 's
 - $-E_T \geq \text{several TeV}$
- Event generator HERBVI:

[Gibbs, Webber '95]

- B-violation cannot be established
- L-violation verifiable: measure

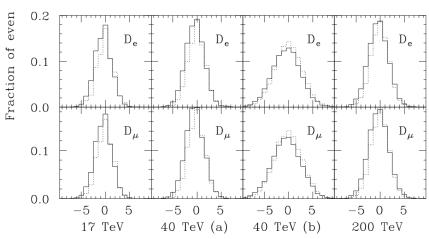
$$D_{\ell} = N_{\ell^{-}} - N_{\ell^{+}};$$

need $\sim 10^3$ events

[Gibbs, AR, Webber, Zadrozny '94]

Simulations performed			
Energy (TeV)		n_B estimate	$\sqrt{\hat{s}_0}$ (TeV)
17		$1/lpha_W$	5
40	(a)	$1/lpha_W$	18
40	(b)	LOME	18
200		$1/lpha_W$	18

[Gibbs, AR, Webber, Zadrozny '94]





Summary

- Electroweak B+L violation central building block of our understanding of baryogenesis in big bang cosmology
- Characteristic scale of B+L violation

$$4\pi \frac{M_W}{\alpha_W} \simeq 30 \text{ TeV}$$

- Cross-sections for B+L violating processes
 - exponentiall small below this scale
 - exponentially growing below this scale
 - may reach observable values near this scale
- New computational methods needed
- Unique opportunity for a 100 TeV collider!
- In the meantime, search for small-size QCD-instanton induced processes

