A new 3D grid method for accurate electronic structure calculation of polyatomic molecules: The Voronoi-cell finite difference method

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Abstract

We introduce a new computational method on unstructured grids in the threedimensional (3D) spaces to investigate the electronic structure of polyatomic molecules. The Voronoi-cell finite difference (VFD) method realizes a simple discrete Laplacian operator on unstructured grids based on Voronoi cells and their natural neighbors. The feature of unstructured grids enables us to choose intuitive pictures for an optimal molecular grid system. The new VFD method achieves highly adaptability by the Voronoi-cell diagram and yet simplicity by the finite difference scheme. It has no limitation in local refinement of grids in the vicinity of nuclear positions and provides an explicit expression at each grid without any integration. This method augmented by unstructured molecular grids is suitable for solving the Schrödinger equation with the realistic 3D Coulomb potentials regardless of symmetry of molecules. For numerical examples, we test accuracies for electronic structures of one-electron polyatomic systems: linear H_2^+ and triangular H_3^{++} . We also extend VFD to the density functional theory (DFT) for many-electron polyatomic molecules.

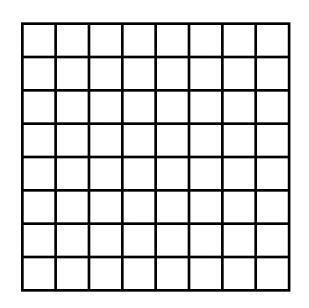


Introduction

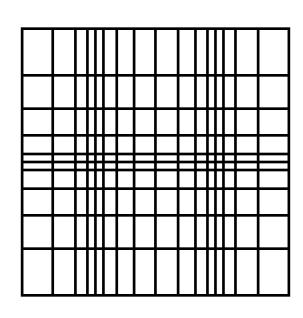
- Voronoi-cell FD (VFD): based on Voronoi diagram and its natural neighbors
- Unstructured grids: optimal molecular grids
- High adaptability like FE: no limitation on local grid refinement
- Simplicity like FD: a simple and explicit matrix form of the discrete Laplacian operator
- Direct solutions of the Schrödinger / Poisson equations on unstructured grids
- No integration for constructing the Hamiltonian matrix
- 3D realistic Coulomb potential



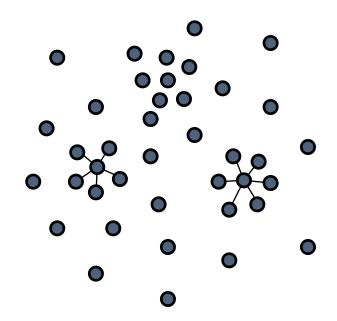
Unstructured grids



uniformly structured



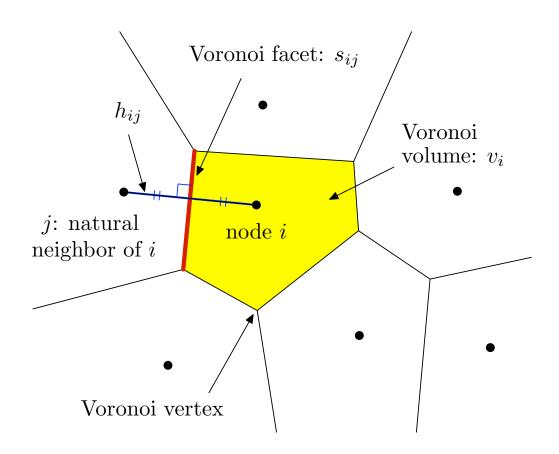
non-uniformly structured



randomly unstructured



Voronoi diagram



Voronoi cell: uniquely defined on unstructured grids in n-dimension

$$T_i = \{ \mathbf{x} \in \mathbb{R}^n : d(\mathbf{x}, \mathbf{x}_i) \le d(\mathbf{x}, \mathbf{x}_j) \text{ for } \forall j \ne i \}$$



Voronoi-cell FD

Differentiation
$$\nabla^2 \varphi = \lim_{\int_V d\tau \to 0} \frac{\int_S \nabla \varphi \cdot \mathbf{n} \, d\sigma}{\int_V d\tau} \text{ (Gauss's theorem)}$$

- i) volume and surface integrals are computed by Voronoi volumes and Voronoi facet areas
- ii) a simple difference scheme is used for directional derivatives

Discrete Laplacian:
$$(\nabla^2 \varphi)_i = \frac{1}{v_i} \sum_{j=1}^{natural neighbors} \frac{\varphi_j - \varphi_i}{h_{ij}} s_{ij}$$

Integration

Simple nodal quadrature using Voronoi volumes

$$\int_{V} f(\mathbf{x}) d\tau \approx \sum_{i} f(\mathbf{x}_{i}) v_{i}$$



PDE solver by VFD

Schrödinger equation
$$\left(-\frac{1}{2} \mathbf{\nabla}^2 + U(\mathbf{x}) \right) \psi(\mathbf{x}) = E \psi(\mathbf{x})$$

$$\Rightarrow (\tilde{\mathbf{T}} + \mathbf{U}) \tilde{\mathbf{C}} = \tilde{\mathbf{E}} \tilde{\mathbf{C}}$$

$$ilde{T}_{ij} = egin{cases} -rac{1}{2\sqrt{v_iv_j}}rac{s_{ij}}{h_{ij}} & (i
eq j), & ext{After symmetrization} \ rac{1}{2v_i}\sum_{k}^{ ext{neighbors}}rac{s_{ik}}{h_{ik}} & (i = j) \ U_{ij} = \delta_{ij}U(\mathbf{x}_i) & ext{Potential is given by a diagonal noise} \end{cases}$$

Potential is given by a diagonal matrix without any integration

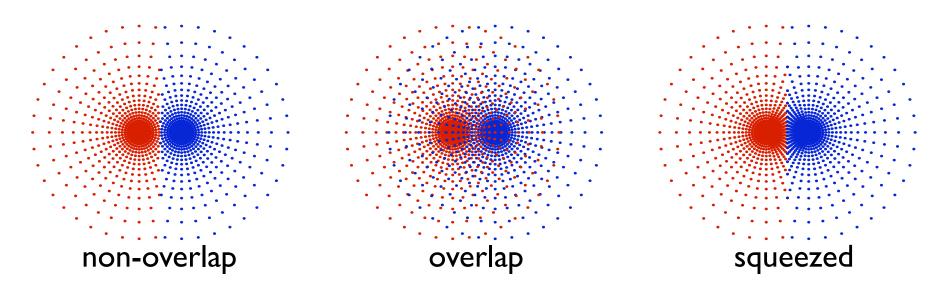
Poisson equation

$$\nabla^2 v_{\rm h}(\mathbf{x}) = -4\pi \rho(\mathbf{x}) \Rightarrow \mathbf{v}_{\rm h} = -4\pi \mathbf{L}^{-1} \boldsymbol{\rho}$$

Easy to incorporate the boundary condition within L



Molecular grids



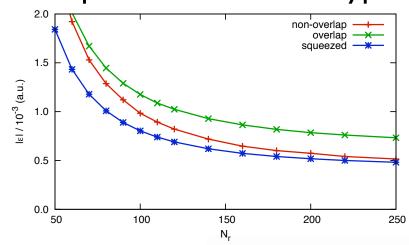
Molecular grids are intuitively constructed by combination of spherical atomic grids in 3D

- radial part:

$$r(x) = L\frac{1+x}{1-x} \quad (-1 < x < 1)$$

- angular part: Lebedev grids

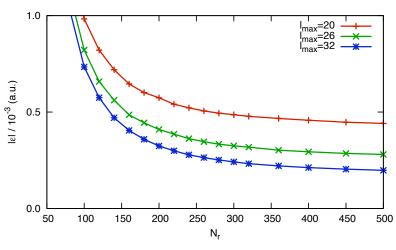
Comparison of different types

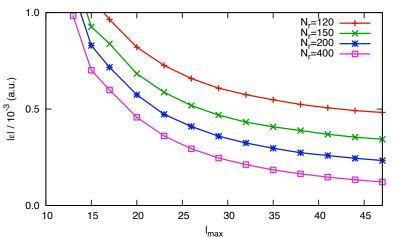




Accuracy assessment

Errors of H_2^+ ground energies by varying N_r and I_{max}





Errors of H₂⁺ bound energies

Symmetry	Exact	LCAO-GTO		VFD	
		$\varepsilon(a.u.)$	$ \varepsilon (\%)$	$\varepsilon(\mathrm{a.u.})$	$ \varepsilon (\%)$
$1\sigma_g$	$-1.102\ 634\ 21$	1.41×10^{-6}	0.000	-1.22×10^{-4}	0.011
$1\sigma_u$	$-0.667\ 534\ 39$	1.69×10^{-6}	0.000	-1.55×10^{-4}	0.023
$1\pi_u$ (2)	$-0.428\ 771\ 82$	1.09×10^{-4}	0.026	-7.92×10^{-5}	0.018
$2\sigma_g$	$-0.360\ 864\ 88$	5.29×10^{-5}	0.015	-6.65×10^{-5}	0.018
$2\sigma_u$	$-0.255\ 413\ 17$	2.49×10^{-4}	0.097	-7.14×10^{-5}	0.028
$3\sigma_g$	$-0.235\ 777\ 63$	3.46×10^{-3}	1.468	-1.20×10^{-4}	0.051
$1\pi_g$ (2)	-0.226 699 63	4.23×10^{-3}	1.867	-1.81×10^{-4}	0.080

Basis set: aug-cc-pV6Z / Grid parameters: N_r =400, L=1, l_{max} =47 / H_2^+ : R=2.0 a.u.

Errors of triangular H₃⁺⁺ bound energies

Symmetry	FE	LCAO-GTO	VFD	Difference
$1a_1'$	$-1.909\ 570\ 99$	$-1.909\ 569$	$-1.909\ 787$	-2.18×10^{-4}
1e'(2)		$-1.138\ 578$	-1.138979	-4.02×10^{-4}
$1a_2''$		-0.869699	$-0.870\ 008$	-3.10×10^{-4}
$2a_1'$		-0.704~969	$-0.705\ 131$	-1.62×10^{-4}
2e'(2)		-0.534978	$-0.535\ 372$	-3.93×10^{-4}
3e'(2)		$-0.484\ 387$	$-0.485\ 081$	-6.94×10^{-4}

Basis set: aug-cc-pV6Z / Grid parameters: N_r =400, L=1, l_{max} =41 / H_3^{++} : R=1.68 a.u.

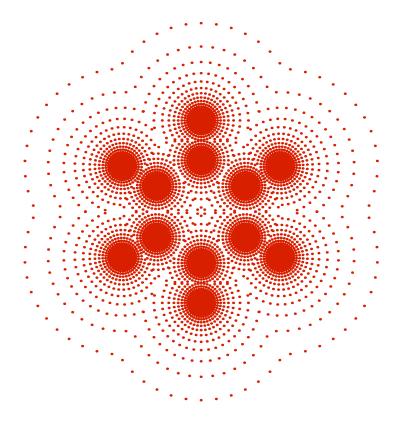


DFT results

LDA energies (in a.u.) of benzene

Symmetry	LCAO-GTO	VFD	Difference
$1a_{1g}$	-9.791	-9.797	-0.006
$1e_{1u}(2)$	-9.791	-9.797	-0.006
$1e_{2g}(2)$	-9.790	-9.797	-0.007
$1b_{2u}$	-9.790	-9.796	-0.006
$2a_{1g}$	-0.778	-0.775	0.003
$2e_{1u}(2)$	-0.676	-0.673	0.003
$2e_{2g}(2)$	-0.545	-0.543	0.002
$3a_{1g}$	-0.478	-0.477	0.001
$2b_{2u}$	-0.411	-0.410	0.001
$1b_{1u}$	-0.407	-0.404	0.003
$3e_{1u}(2)$	-0.379	-0.377	0.002
$1a_{2u}$	-0.341	-0.340	0.001
$3e_{2g}(2)$	-0.305	-0.303	0.002
$1e_{1g}(2)$	-0.240	-0.240	0.000
$E_{ m total}$	-230.177	-230.211	-0.034

Basis set: 6-311++G(3df,3pd) / Grid parameters: N_r =200, L=0.5, l_{max} =26



2D sketch of *unstructured* molecular grids for benzene



Conclusion

- VFD is a new numerical grid method based on Voronoi diagram and applied for electronic structure calculations of polyatomic molecules.
- VFD allows us to employ intuitive pictures for unstructured molecular grids.
- With realistic Coulomb potential, eigenvalues of I-e systems are solved within ~10-4 a.u. accuracy. DFT calculations show ~0.01% total energy difference from large-basis-set LCAO results.
- VFD is extensible to time-dependent problems due to ease of the potential matrix and applicable to problems demanding highly adaptive refinement.



References

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