



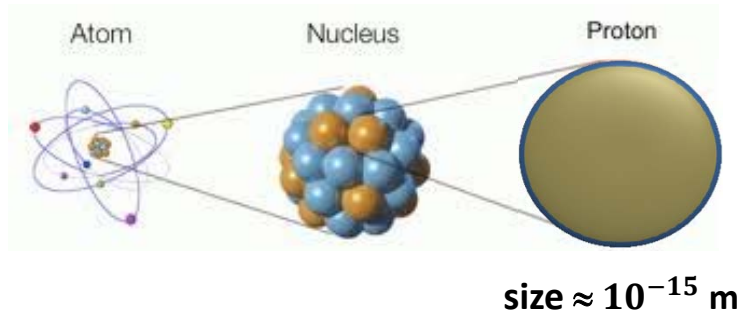
Transverse Momentum Distributions: an experimental update

Luciano L. Pappalardo

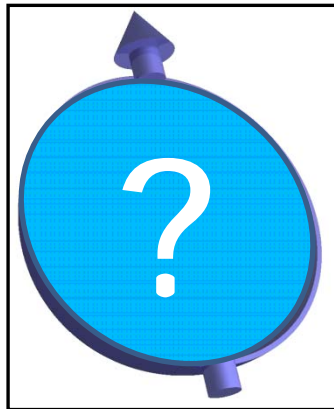
INFN & University of Ferrara

QNP2012 - Palaiseau (France) - April 16-20 2012

Looking deeply into the proton

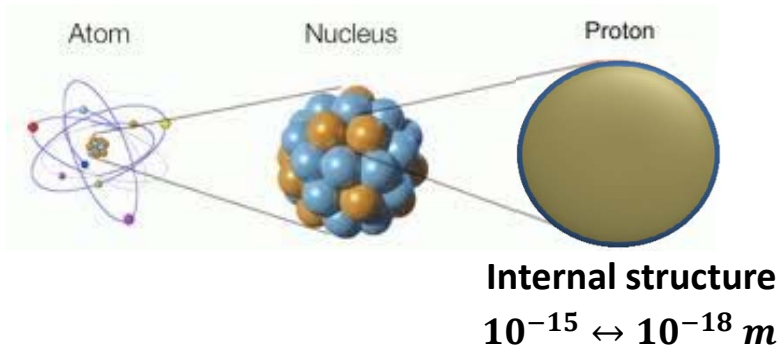


< 1960

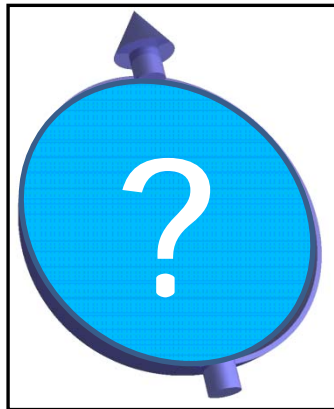


- Elementary particle?

Looking deeply into the proton

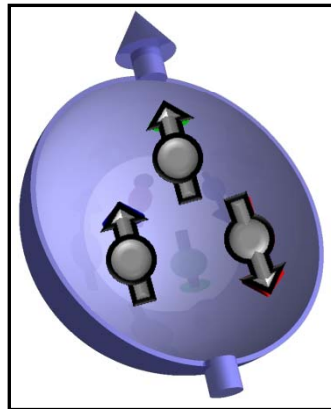


< 1960



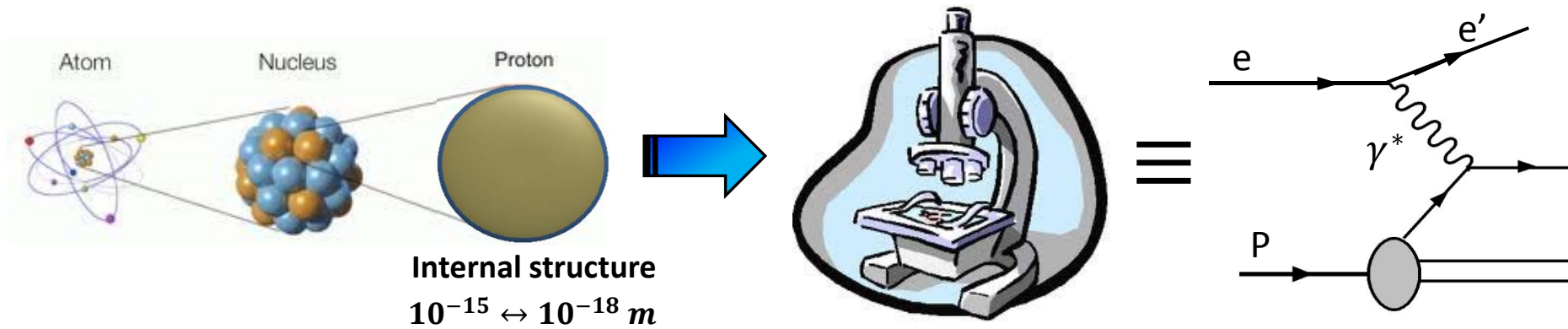
- Elementary particle?

1964

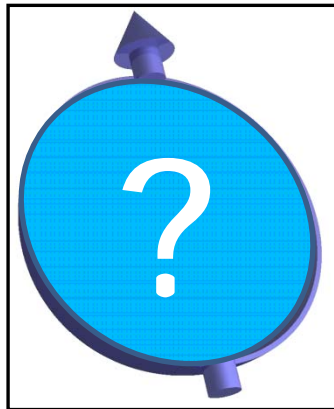


- Quark hypothesis
(Gell-Mann - Zweig)

Looking deeply into the proton

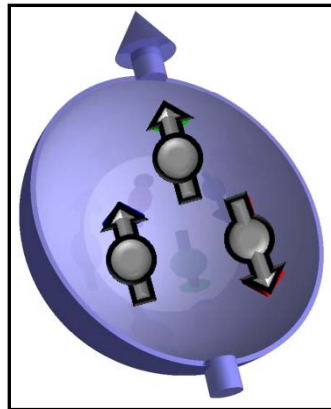


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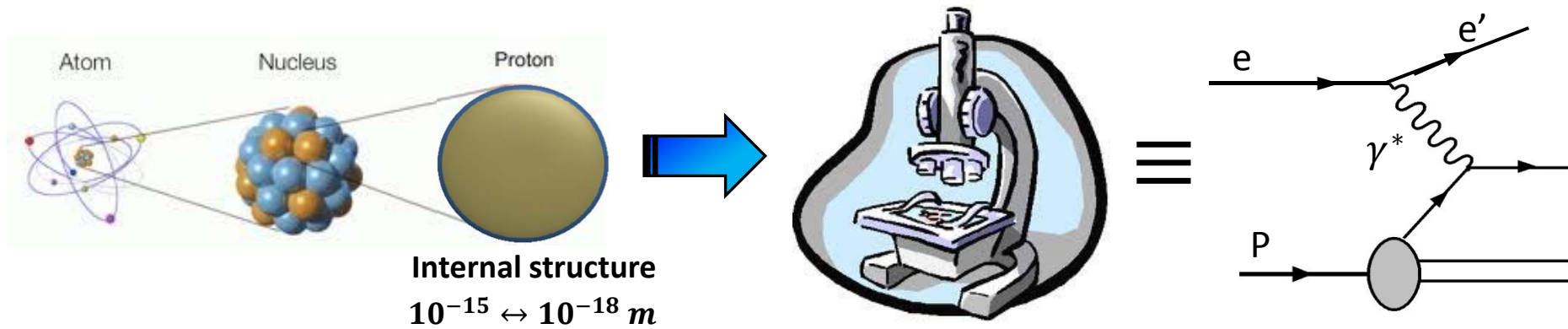
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1964-1969

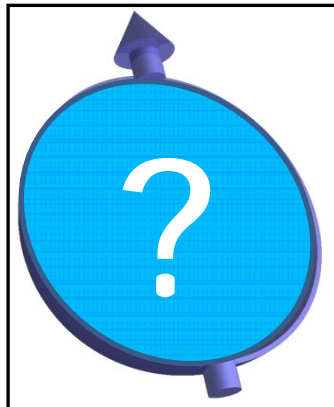


- Quark hypothesis
(Gell-Mann - Zweig)
- Scaling at SLAC ('69)
- Parton Model
(Faynman, Bjorken)

Looking deeply into the proton

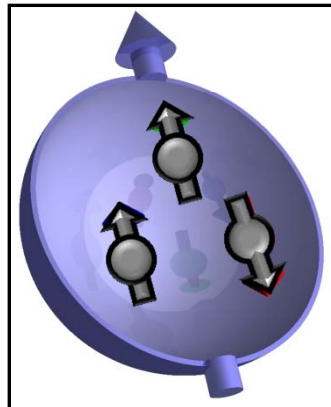


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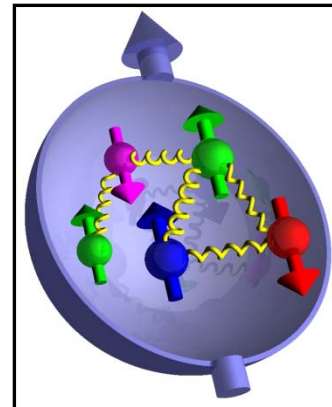
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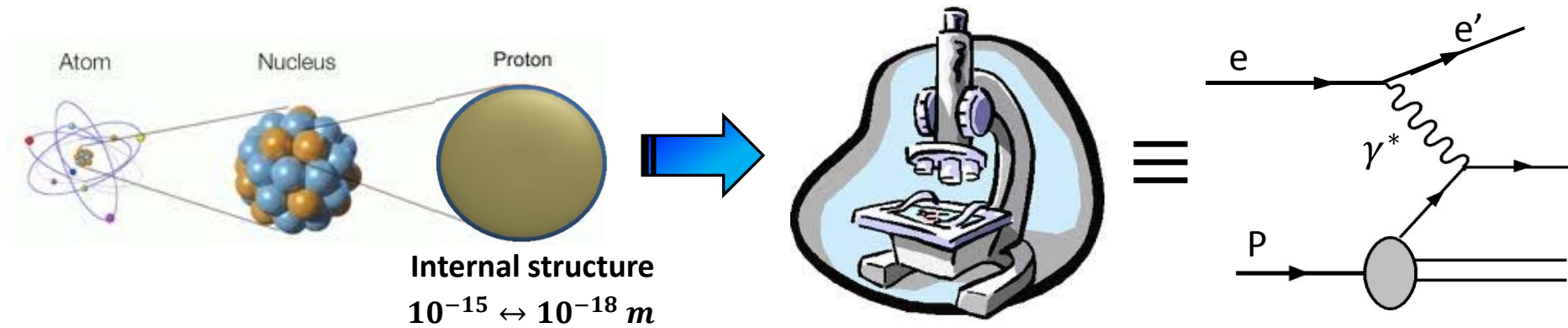
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1972-1973

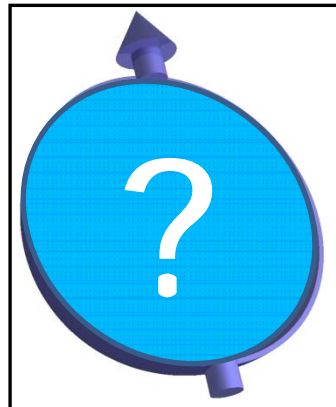


- QCD Lagrangian
 - colors, sea quarks,
 gluons
 - discovery of gluons
 (PETRA '73)

Looking deeply into the proton

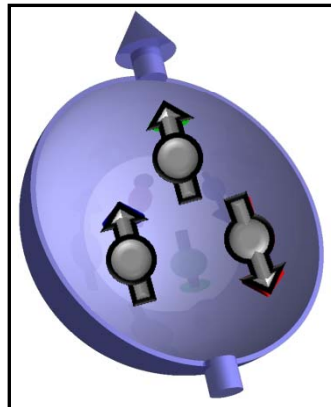


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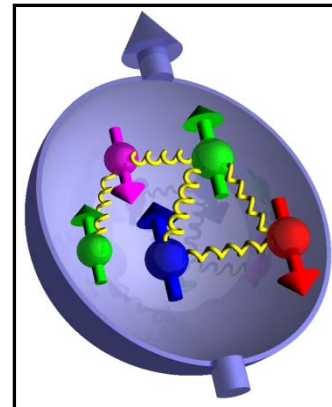
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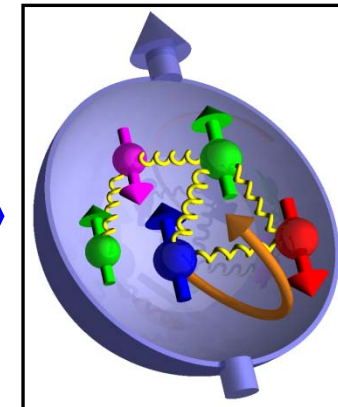
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- QCD Lagrangian
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> 1988



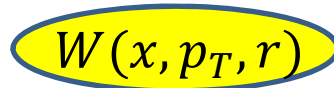
- EMC experiment
 - the spin crisis
 $\frac{1}{2} = \Delta\Sigma + \Delta G + L_{q,g}$
 - quest for ΔG & $L_{q,g}$

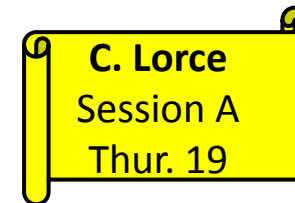
What is the final goal?

Understand the **full phase-space distribution of the partons:**

- Where are they located? $\rightarrow x, y, z \equiv r$
 - How do they move? $\rightarrow p_x, p_y, p_z \equiv x, p_T$
- $W(x, p_T, r)$ **Wigner function**

...but $\Delta x \Delta p \geq \frac{\hbar}{2}$ \rightarrow no simultaneous knowledge of momentum and position!


$$W(x, p_T, r)$$



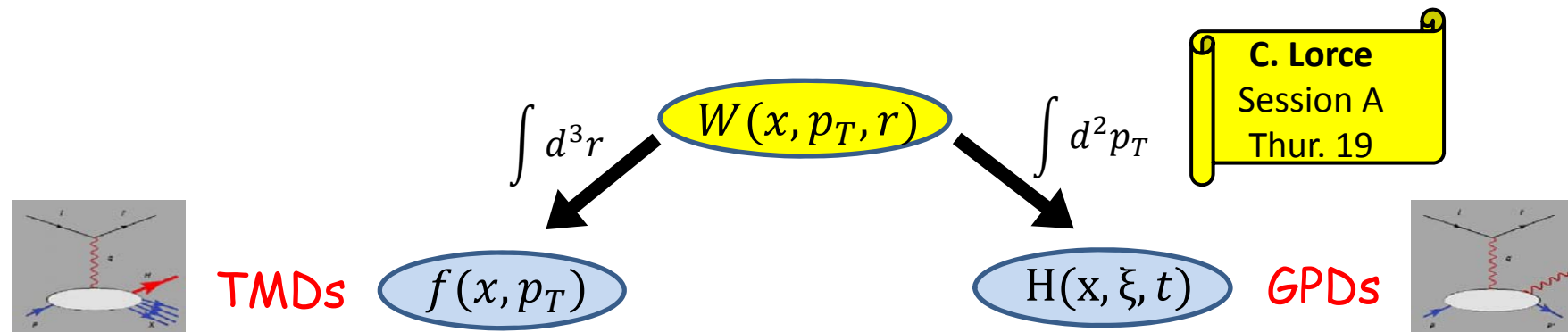
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\rightarrow can only measure quantities integrated over momentum and/or position



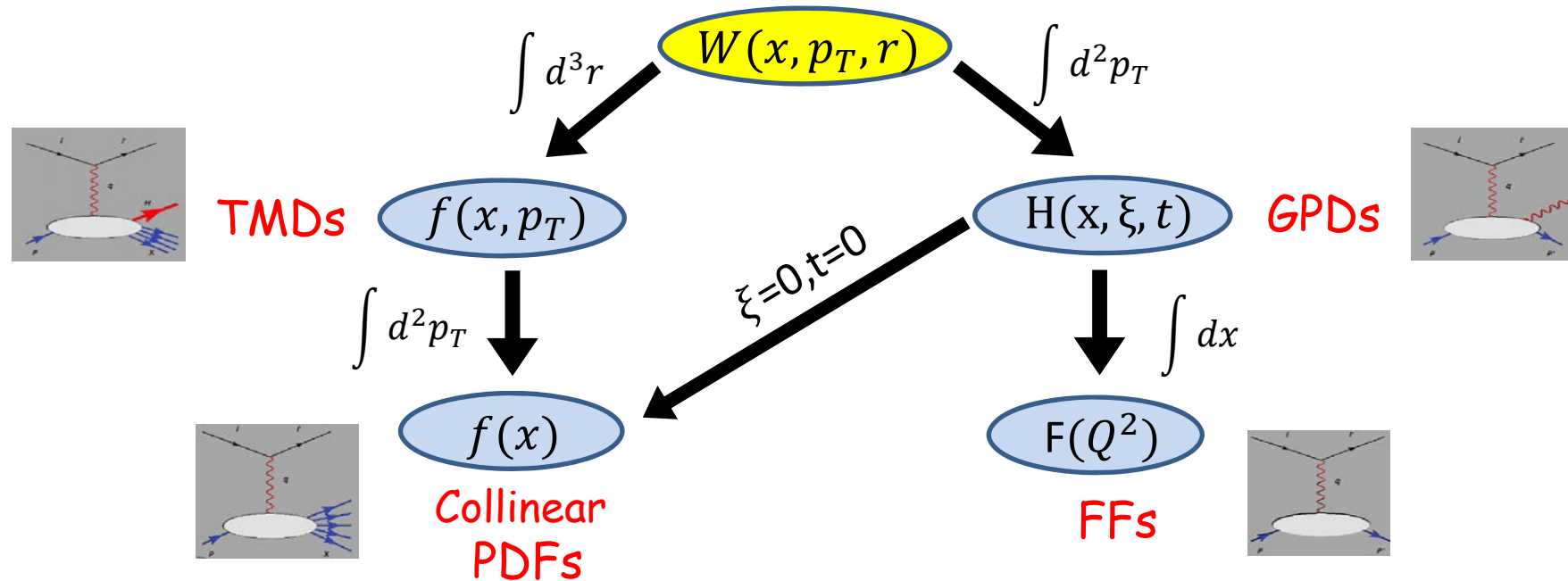
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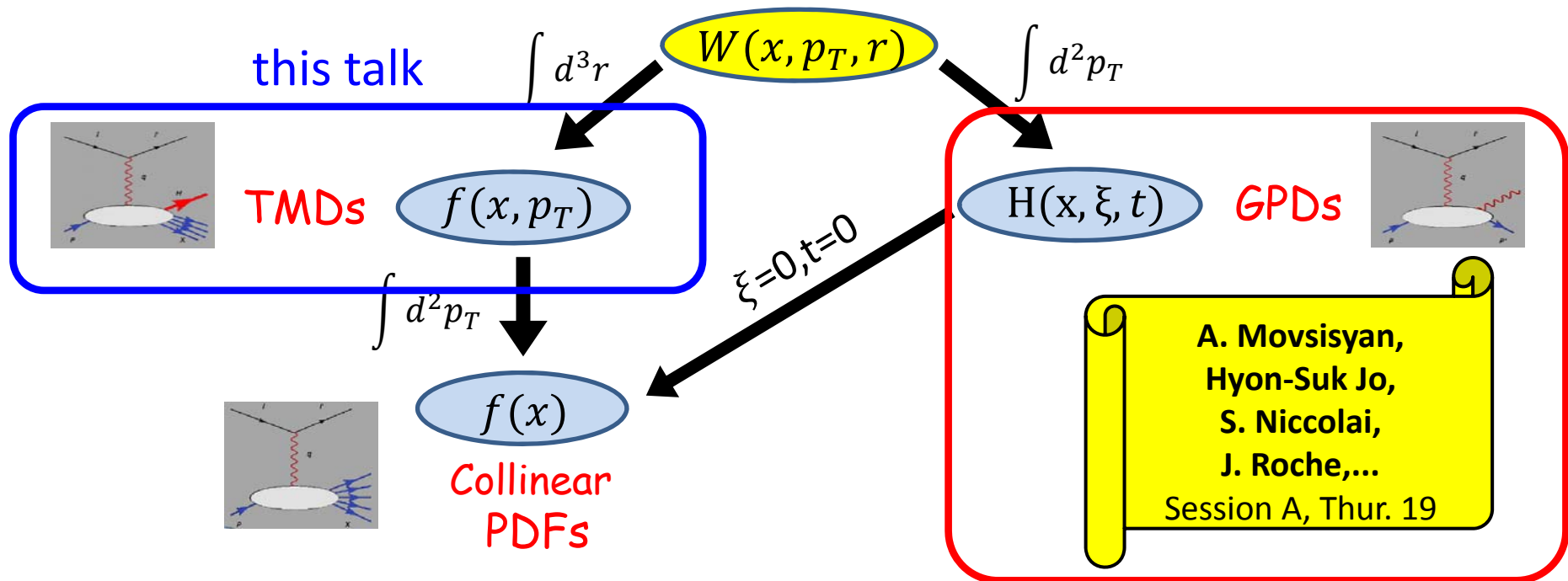


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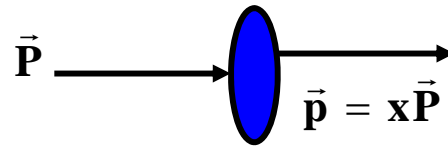
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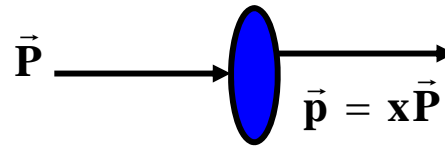
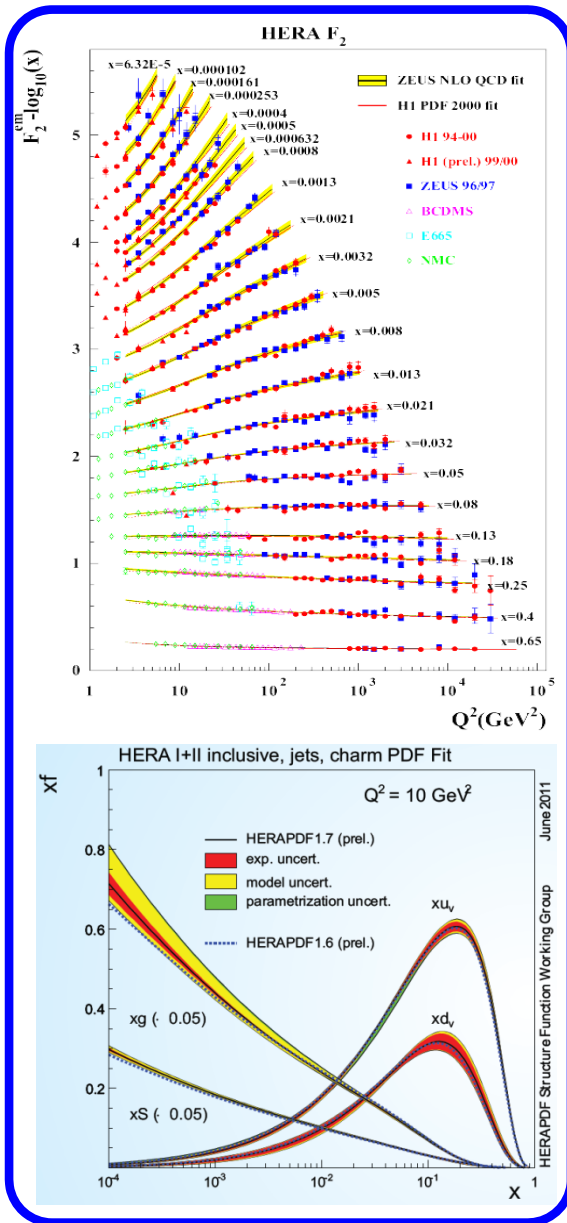


The nucleon collinear structure



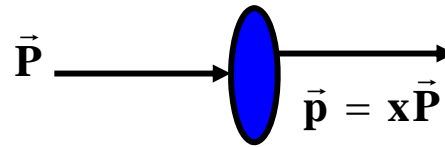
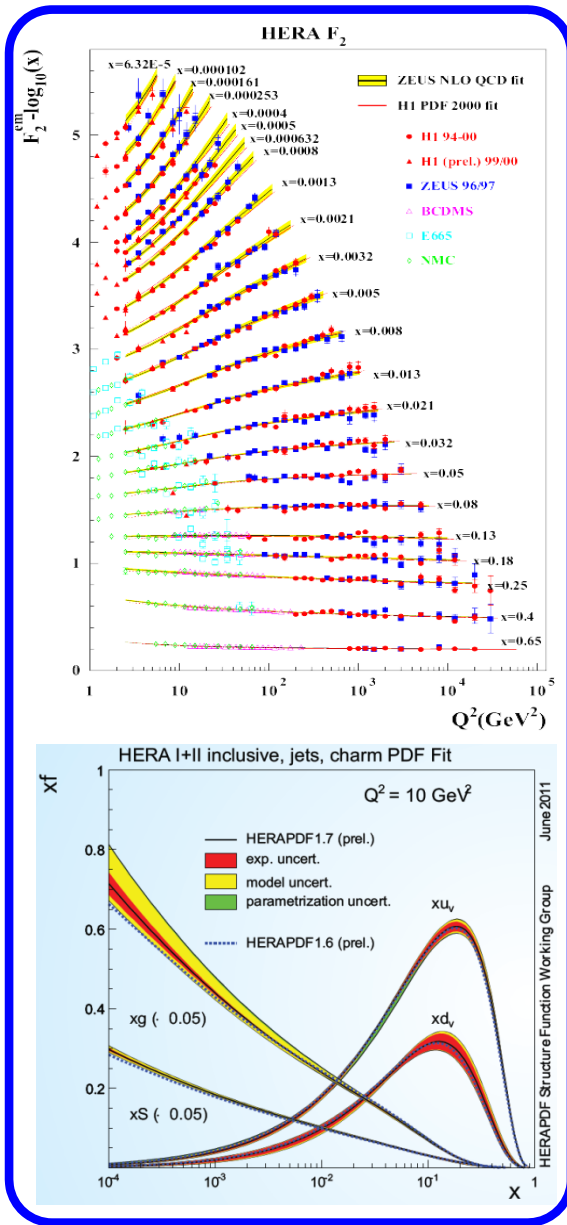
		quark		
		U	L	T
n u c l e o n	U	f_1		
	L		g_1 -	
	T			h_1 - \longrightarrow

The nucleon collinear structure

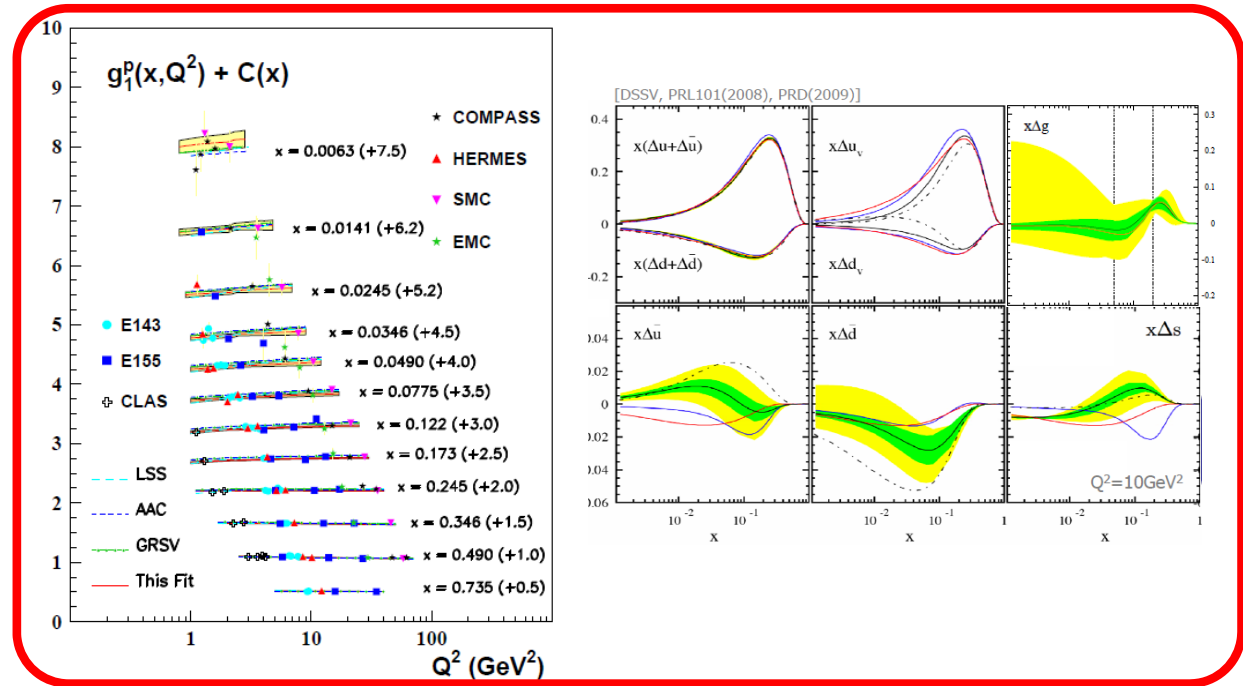


		quark		
		U	L	T
n u c l e o n	U	f_1		
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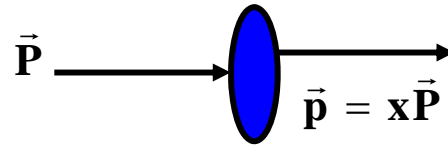
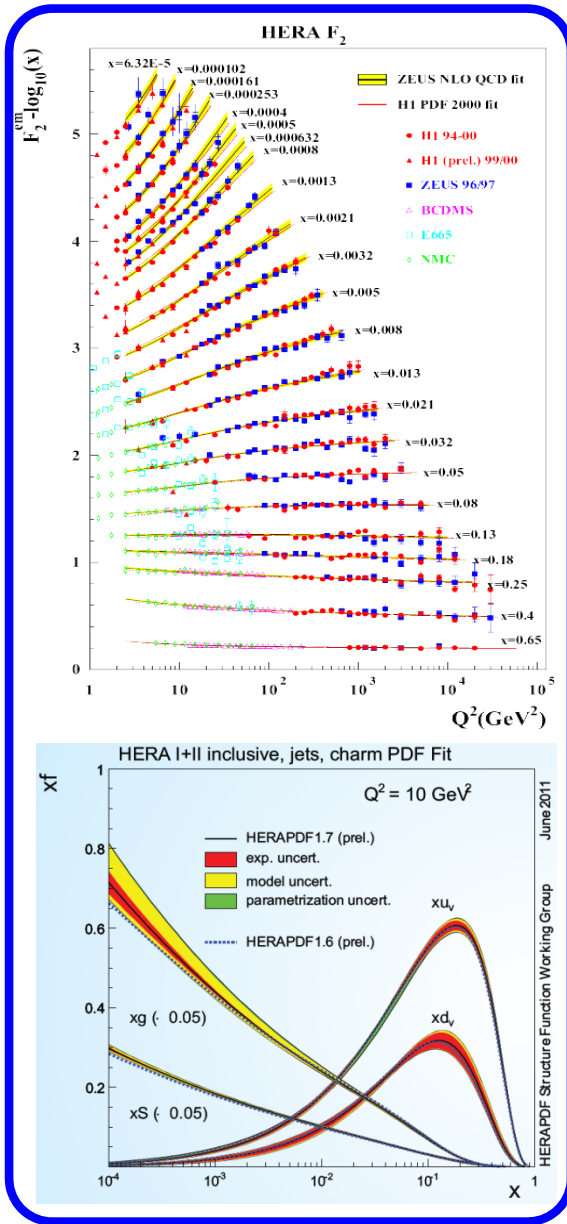
The nucleon collinear structure



		quark		
		U	L	T
nucleon	U	f_1 ○		
	L		g_1 ○ - ○	
	T			h_1 ○ - ○



The nucleon collinear structure

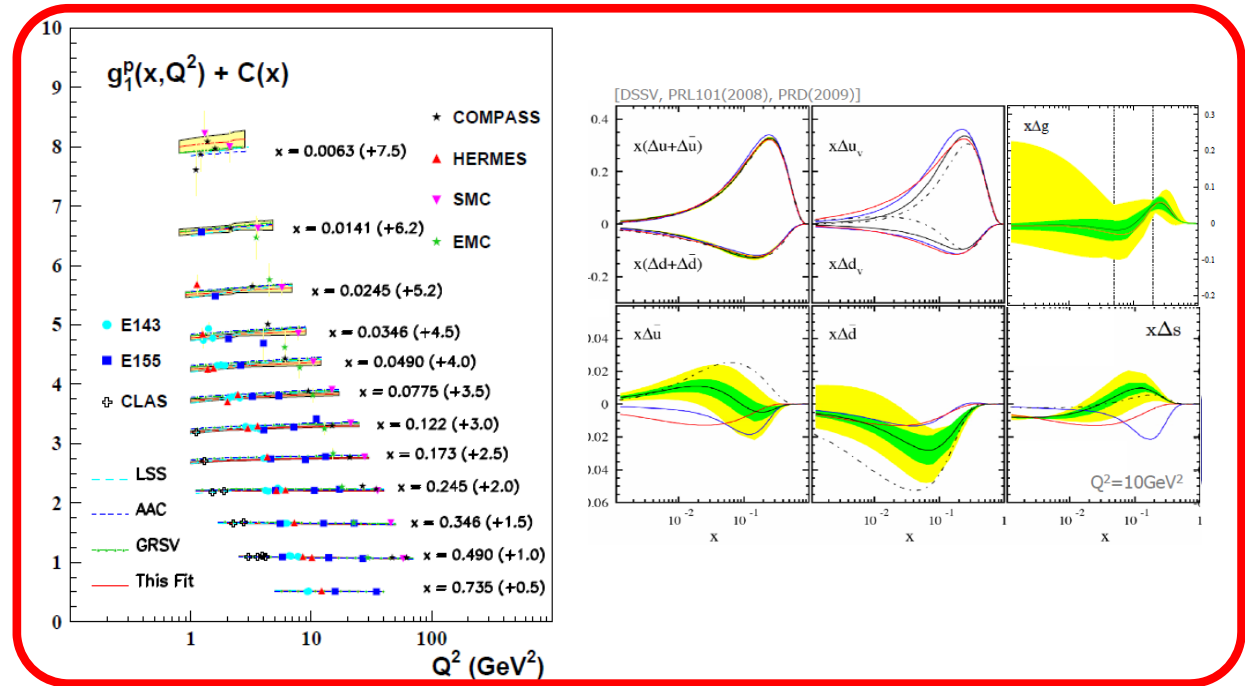


		quark		
		U	L	T
nucleon	U	f_1 (circle with dot)		
	L		g_1 (circle with dot and minus)	
	T			h_1 (circle with dot and dash)

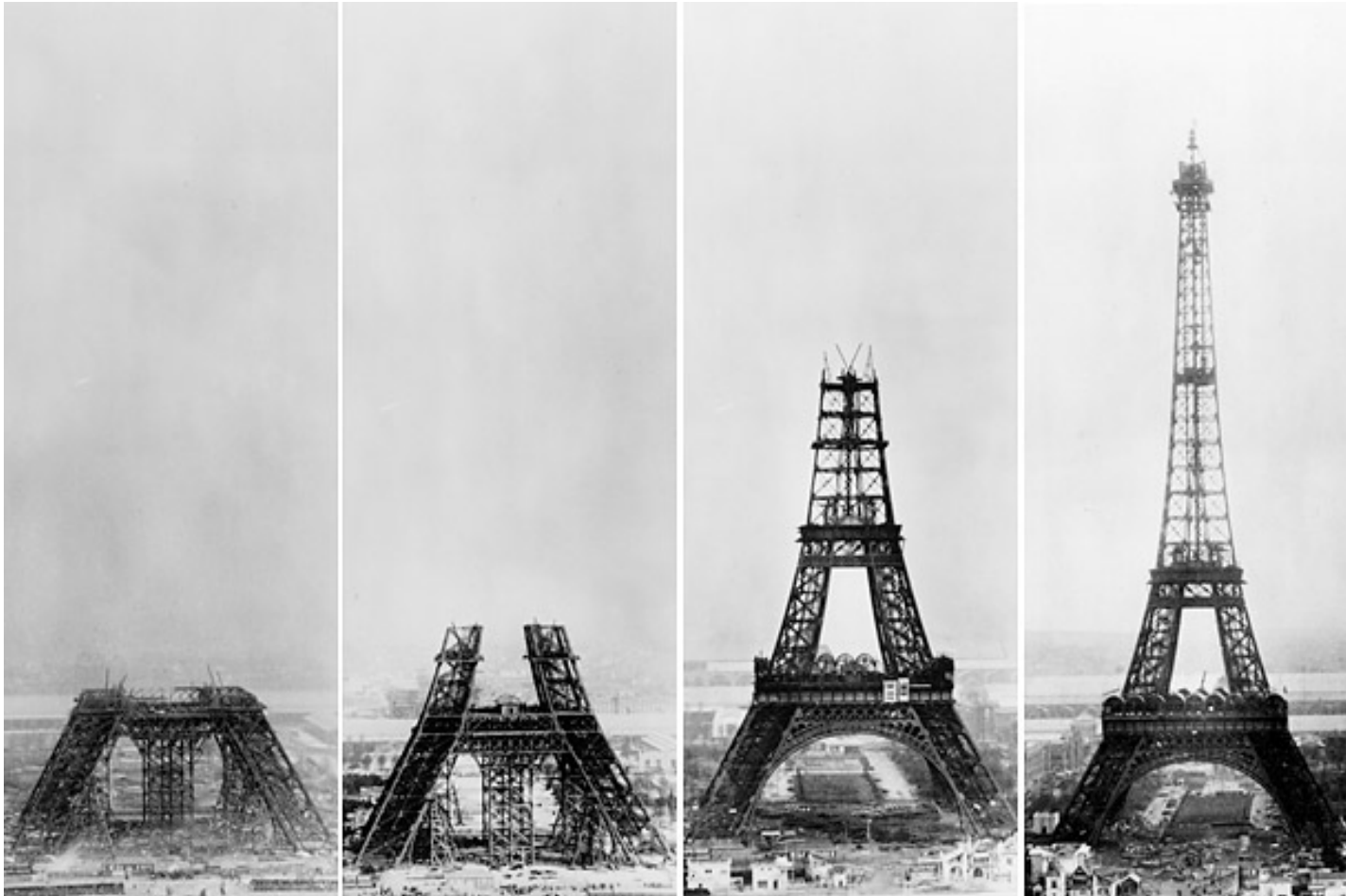
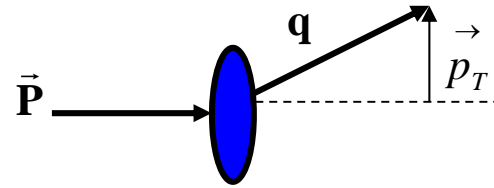
Left: Nucleon with a red quark and a red gluon, with a red octagon labeled 'STOP'.

Right: Parton distribution function (PDF) diagram showing a nucleon P with a quark q and a gluon g , and a parton distribution function DF with a quark q and a gluon g .

- chiral-odd!
- differs from g_1 due to relativistic effects
- no mix with gluons in spin- $\frac{1}{2}$ the nucleon
- valence object



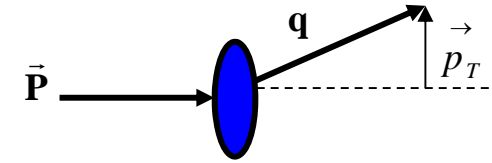
...now let's go transverse!



The non-collinear structure of the nucleon

		quark		
		U	L	T
n u c l e o n	U	f_1		h_1^+
	L		g_1	h_{1L}^+
	T	f_{1T}^+	g_{1T}^+	h_{1T}^+

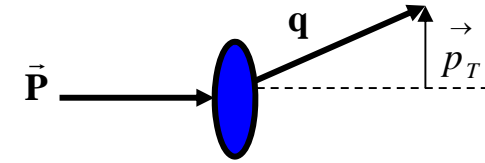
momentum helicity Boer-Mulders transversity pretzelosity
Sivers worm-gears



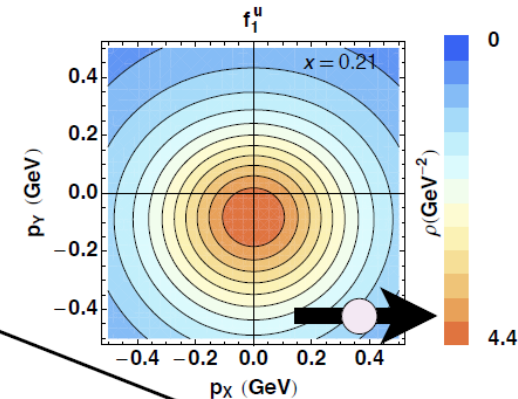
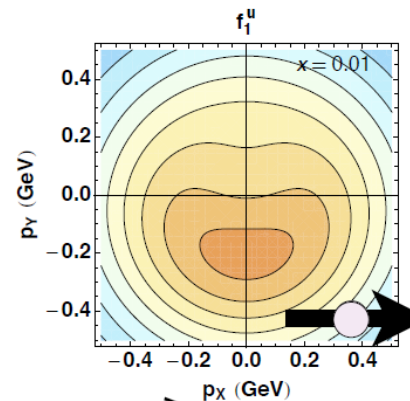
- TMDs depend on p_T
- Vanish when integrated over p_T
- Describe correlations between p_T and quark or nucleon spin (**spin-orbit correlations**)

The non-collinear structure of the nucleon

		quark		
		U	L	T
nucleon	U	f_1		h_1^\perp
	L		g_1	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}^\perp	h_{1T}^\perp



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- Provide a **3-dim picture** of the nucleon in momentum space (**nucleon tomography**)

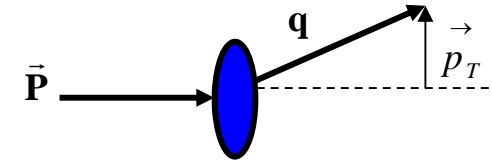


Based on model calculation
A.B., Conti, Guagnelli, Radici, arXiv:1003.1328

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momentum, helicity, Boer-Mulders, transversity, pretzelosity, Sivers, worm-gears

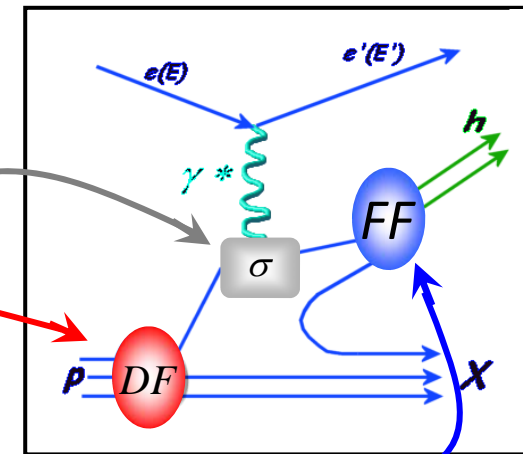


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- Provide a **3-dim picture** of the nucleon in momentum space (**nucleon tomography**)
- Mostly studied in **SIDIS**

		quark		
		U	L	T
hadron	U	D_1		H_1^\perp
	L		G_{1L}	H_{1L}^\perp
	T	D_{1T}^\perp	G_{1T}^\perp	H_{1T}^\perp

Factorization \rightarrow
(key result of QCD!)

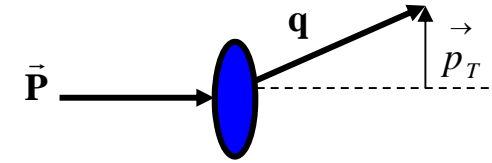
$$\sigma^{ep \rightarrow ehX} = \sum_q \text{DF} \otimes \sigma^{eq \rightarrow eq} \otimes \text{FF}$$



The non-collinear structure of the nucleon

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		U	L	T
nucleon	U	f_1		h_1^\perp
	L		g_1	h_{1L}^\perp
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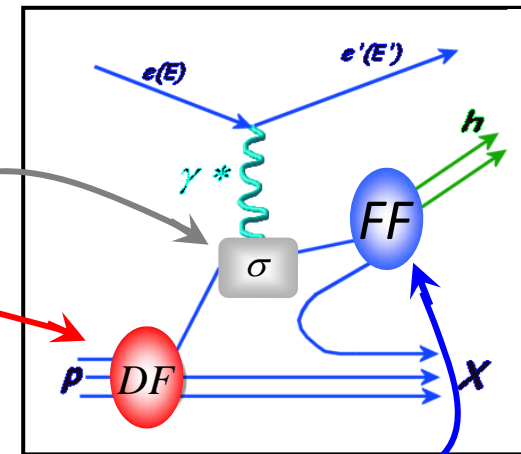


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		quark		
		U	L	T
hadron	U	D_1		H_1^\perp
	L		G_{1L}	H_{1L}^\perp
	T	D_{1T}^\perp	G_{1T}^\perp	H_{1T}^\perp

unpol. FF
chiral-even

Collins FF
chiral-odd



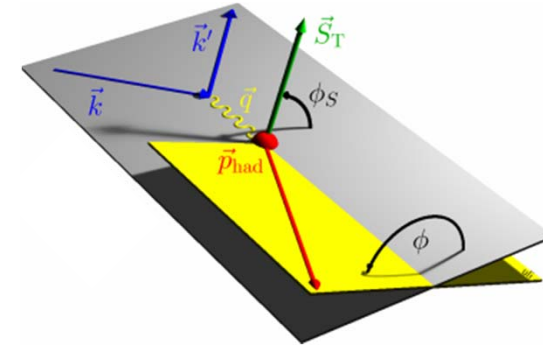
Factorization →
(key result of QCD!)

$$\sigma^{ep \rightarrow ehX} = \sum_q \text{DF} \otimes \sigma^{eq \rightarrow eq} \otimes \text{FF}$$

The SIDIS cross-section

$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2 y^2}{xyQ^2 2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right)$$

$$\left\{ \begin{aligned} & \left[F_{UU,T} + \epsilon F_{UU,L} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} \right] \\ & + \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right] \\ & + S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right] \\ & + S_L \lambda_l \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right] \\ & + S_T \left[\sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi-\phi_S)} + \epsilon F_{UT,L}^{\sin(\phi-\phi_S)} \right) \right. \\ & \quad + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi+\phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi-\phi_S)} \\ & \quad + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \\ & \quad \left. + \sqrt{2\epsilon(1-\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi-\phi_S)} \right] \\ & + S_T \lambda_l \left[\sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi-\phi_S)} \right. \\ & \quad + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \\ & \quad \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi-\phi_S)} \right] \end{aligned} \right\}$$



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$$\left\{ \begin{array}{l} F_{UU,T} + \epsilon F_{UU,L} \\ + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} \end{array} \right.$$

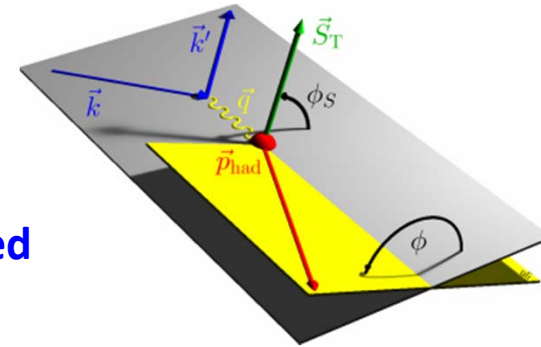
unpolarized

$$+ \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right]$$

beam polarization

$$+ S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right]$$

$$+ S_L \lambda_l \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right]$$



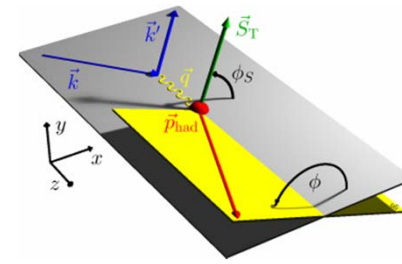
target polarization

$$+ S_T \left\{ \begin{array}{l} \sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi - \phi_S)} \right) \\ + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi - \phi_S)} \\ + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \\ + \sqrt{2\epsilon(1+\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi - \phi_S)} \end{array} \right.$$

beam and target polarization

$$+ S_T \lambda_l \left\{ \begin{array}{l} \sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} \\ + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \\ + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi - \phi_S)} \end{array} \right\}$$

Transversity



$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2 y^2}{xyQ^2 2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x} \right)$$

$$\left\{ \begin{aligned} & \left[F_{UU,T} + \epsilon F_{UU,L} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} \right] \end{aligned} \right.$$

$$+ \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right]$$

$$+ S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right]$$

$$+ S_L \lambda_l \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right]$$

$$+ S_T \left[\begin{aligned} & \sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi - \phi_S)} \right) \\ & + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi - \phi_S)} \\ & + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \\ & + \sqrt{2\epsilon(1-\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi - \phi_S)} \end{aligned} \right]$$

$$+ S_T \lambda_l \left\{ \begin{aligned} & \sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi - \phi_S)} \end{aligned} \right\}$$

Describes probability to find transversely polarized quarks in a transversely polarized nucleon

Distribution Functions

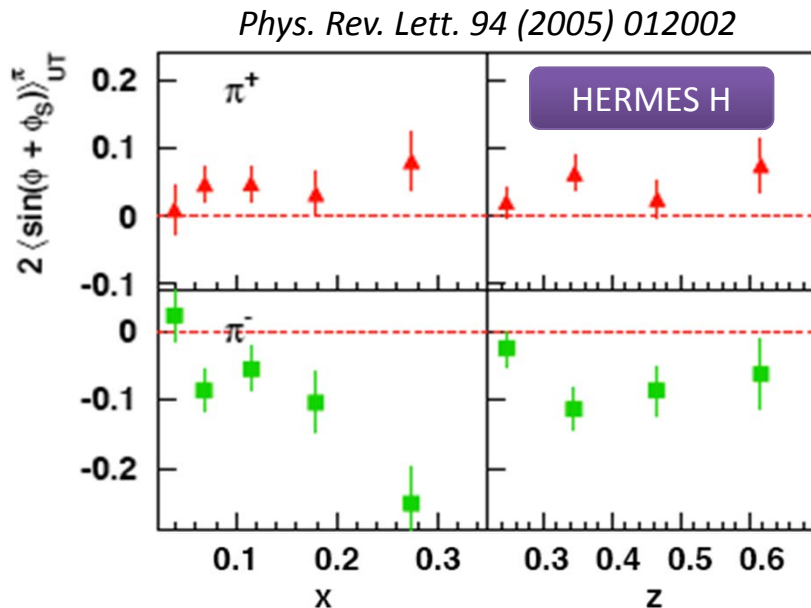
		quark		
		U	L	T
nucleon	U	f_1		h_1^\perp
	L		g_1	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}^\perp	h_1 h_{1T}^\perp

Fragmentation Functions

		quark		
		U	L	T
h	U	D_1		H_1^\perp

Transversity

$$A_{UT}^{\sin(\varphi + \varphi_S)} \propto \frac{\sum_q e_q^2 h_1^q(x, p_T^2) \otimes H_1^{q,\perp}(z, k_T^2)}{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^q(z, k_T^2)}$$



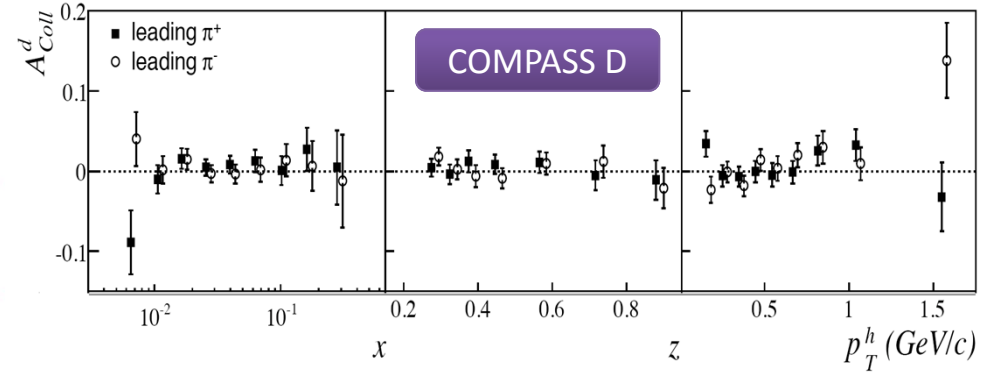
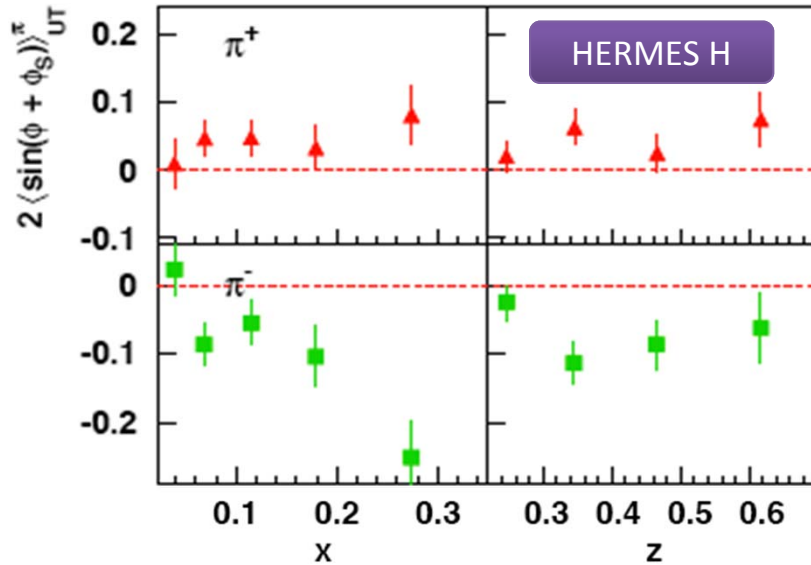
- **First evidence by HERMES (05)**
- Transv. pol. H target
- Limited statistics (2002-2003)
- **Non-zero Collins amplitudes**
- **Non-zero transversity & Collins FF**

$$H_1^{\perp, \text{unfav}}(z) \approx -H_1^{\perp, \text{fav}}(z)$$

Transversity

$$A_{UT}^{\sin(\phi + \phi_S)} \propto \frac{\sum_q e_q^2 h_1^q(x, p_T^2) \otimes H_1^{q,\perp}(z, k_T^2)}{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^q(z, k_T^2)}$$

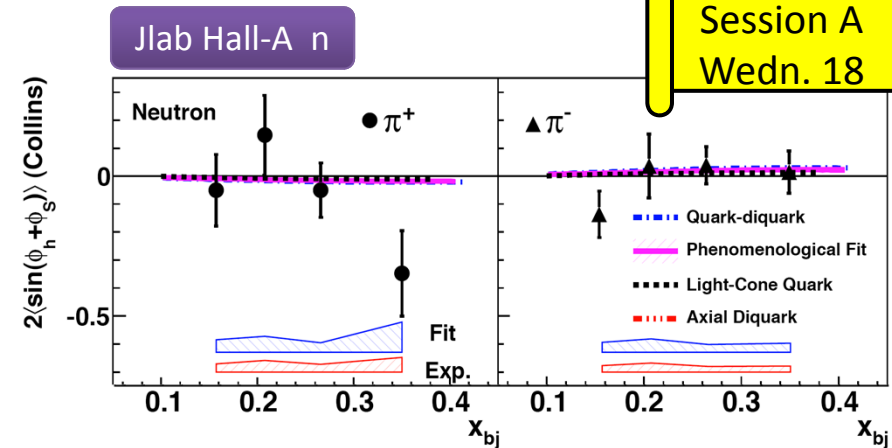
Phys. Rev. Lett. 94 (2005) 012002



- First evidence by HERMES (05)
- Transv. pol. H target
- Limited statistics (2002-2003)
- Non-zero Collins amplitudes
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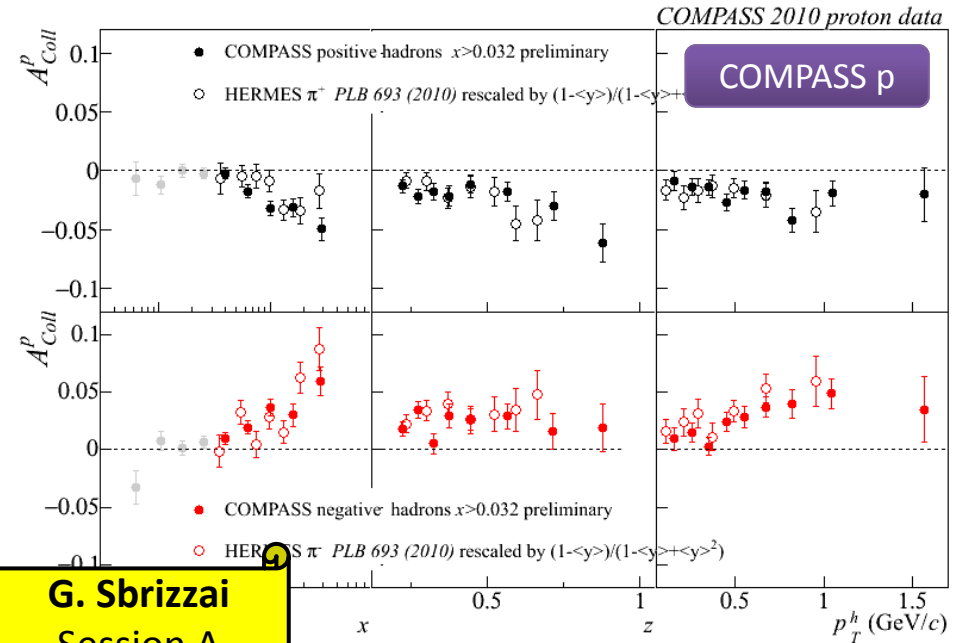
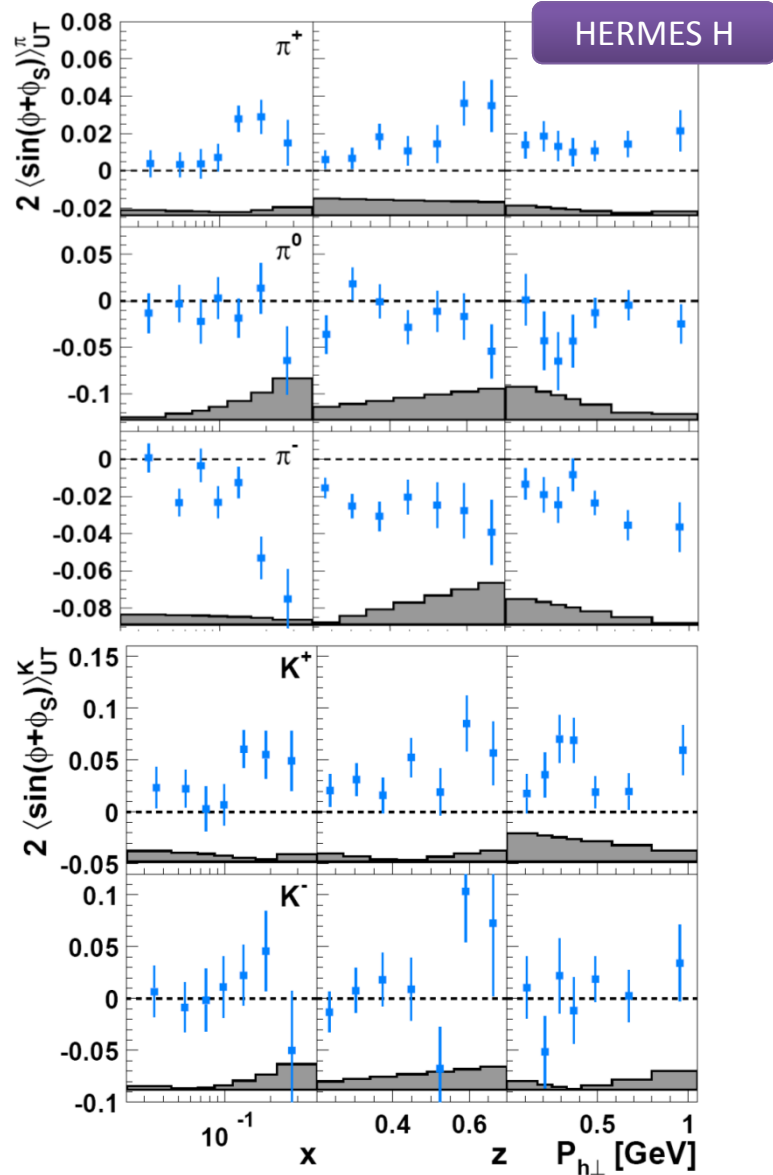
$$H_1^{\perp, unfav}(z) \approx -H_1^{\perp, fav}(z)$$

- Consistent with COMPASS zero amplitude on Deuterium

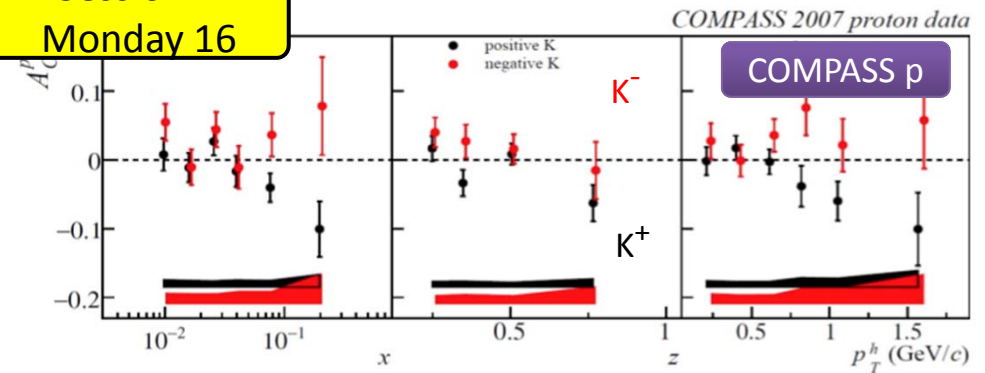


E. Cisbani
Session A
Wedn. 18

Transversity



G. Sbrizzai
Session A
Monday 16

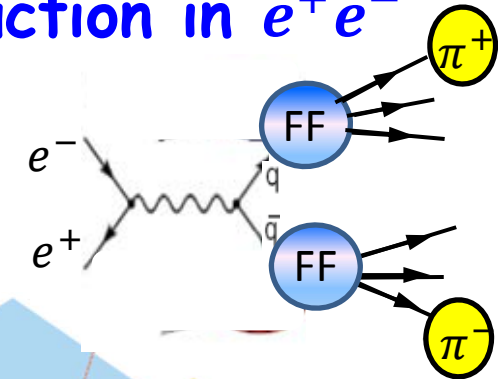


COMPASS π results on p consistent with HERMES!

- K^- amplitudes are not in agreement
- statistics for kaons relatively poor
- **Need new data (e.g. from JLab @12Gev)**

Collins FF from inclusive hadron-pair production in e^+e^-

Describes correlation between quark polarization and observed hadron transverse momentum

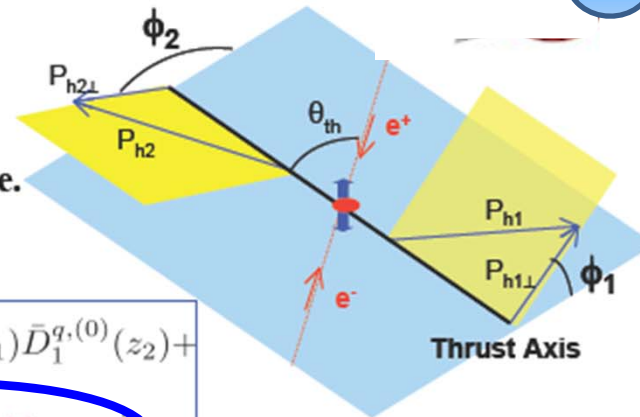


$\phi_1 + \phi_2$ or Thrust RF

θ : angle between the e^+e^- axis and the thrust axis;
 $\phi_{1,2}$: azimuthal angles between $\mathbf{P}_{h1(h2)}$ and the scattering plane.

All quantities in e^+e^- center of mass

$$\frac{d\sigma(e^+e^- \rightarrow h_1 h_2 X)}{d\Omega dz_1 dz_2 d\phi_1 d\phi_2} = \sum_{q,\bar{q}} \frac{3\alpha^2 e_q^2}{Q^2} \frac{z_1^2 z_2^2}{4} \left[(1 + \cos^2\theta) D_1^{q,(0)}(z_1) \bar{D}_1^{q,(0)}(z_2) + \sin^2(\theta) \cos(\phi_1 + \phi_2) H_1^{\perp,(1),q}(z_1) \bar{H}_1^{\perp,(1),q}(z_2) \right]$$

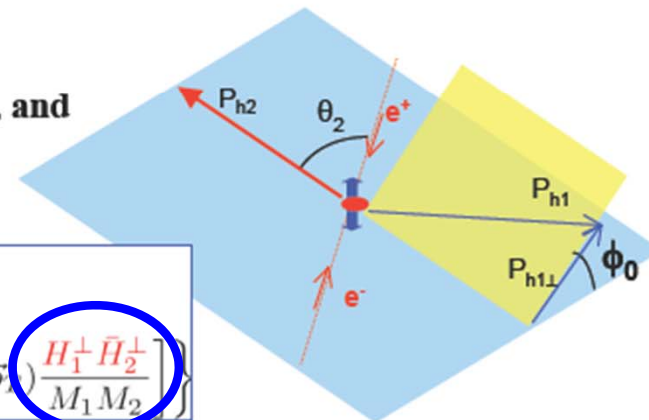


$2\phi_0$ or P_{h2} RF

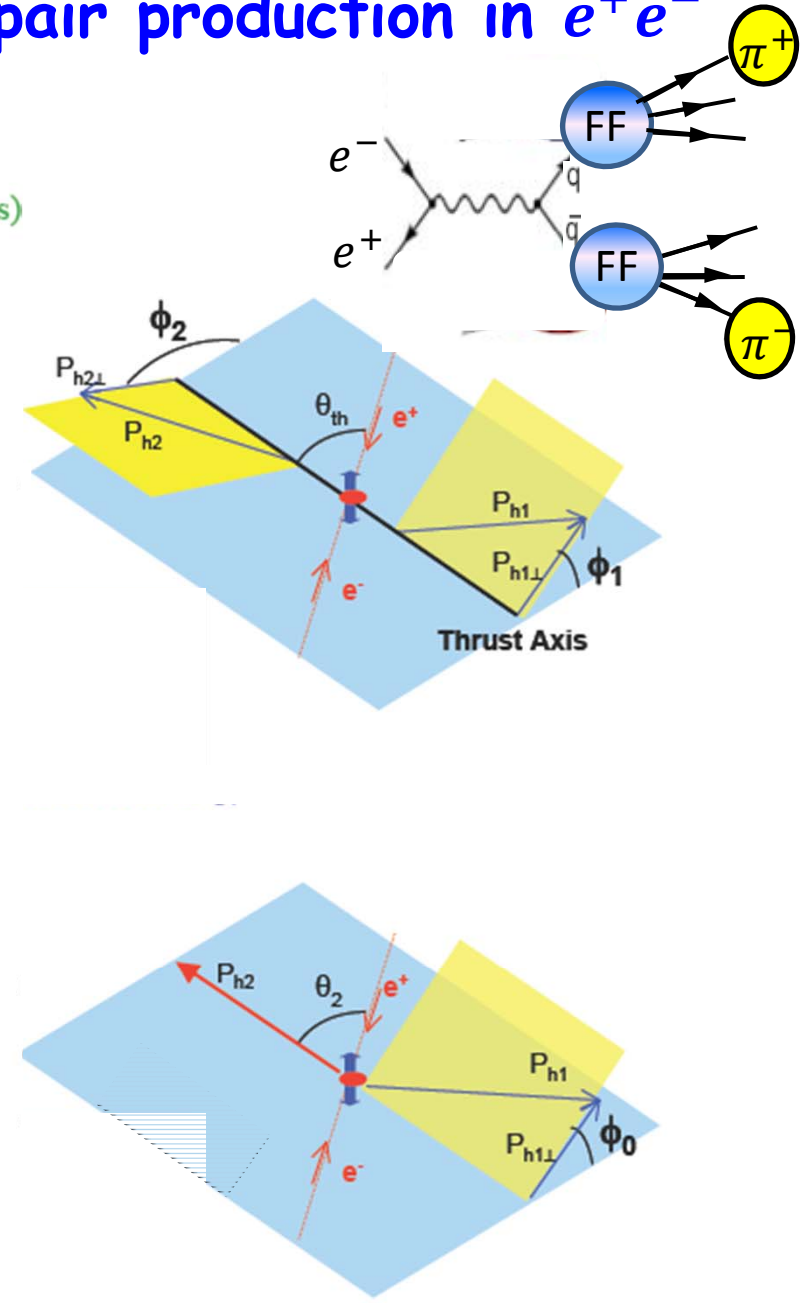
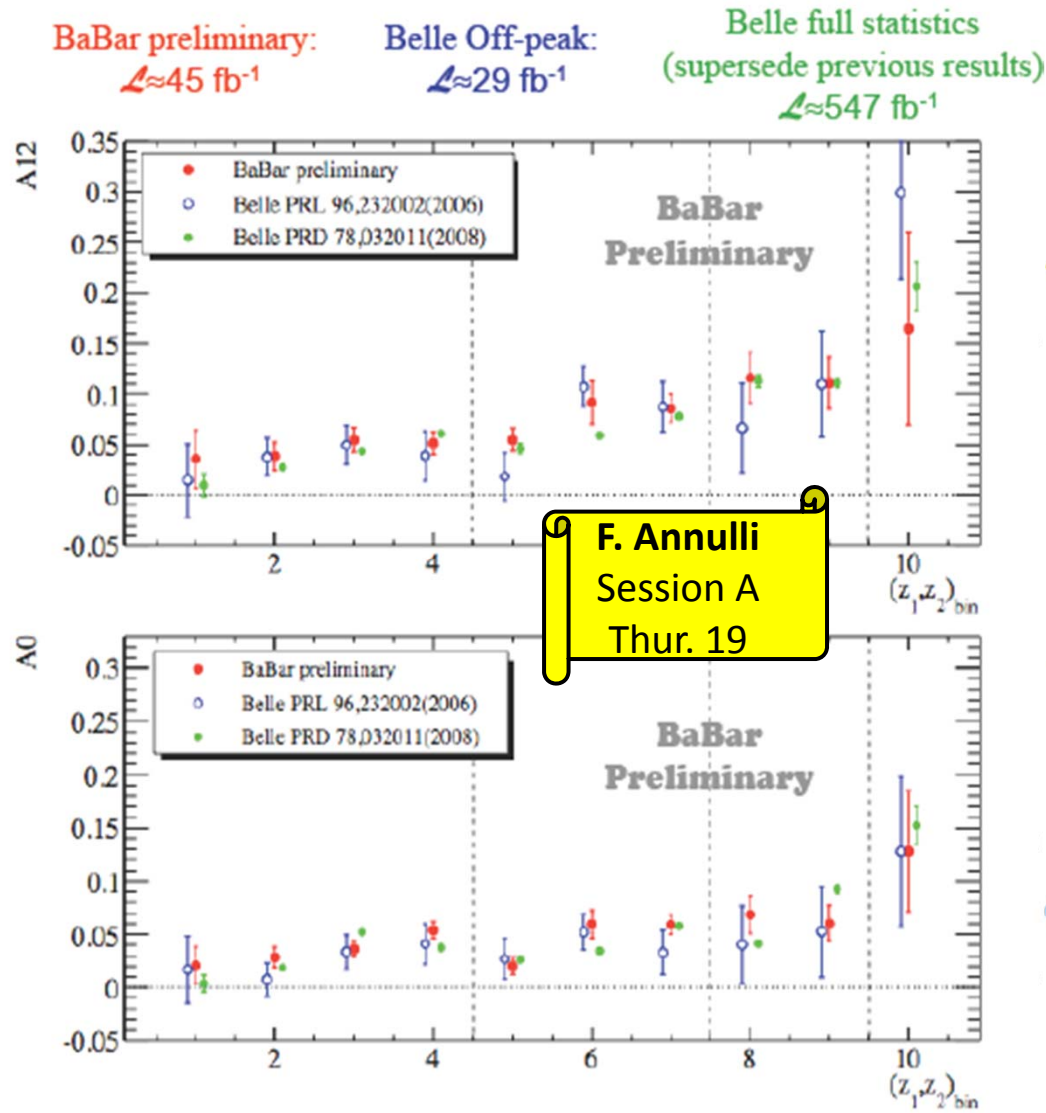
θ_2 : angle between the e^+e^- axis and \mathbf{P}_{h2} ;
 ϕ_0 : angle between the plane spanned by \mathbf{P}_{h2} and the e^+e^- axis, and the direction of \mathbf{P}_{h1} perpendicular to \mathbf{P}_{h2} .

All quantities in e^+e^- center of mass

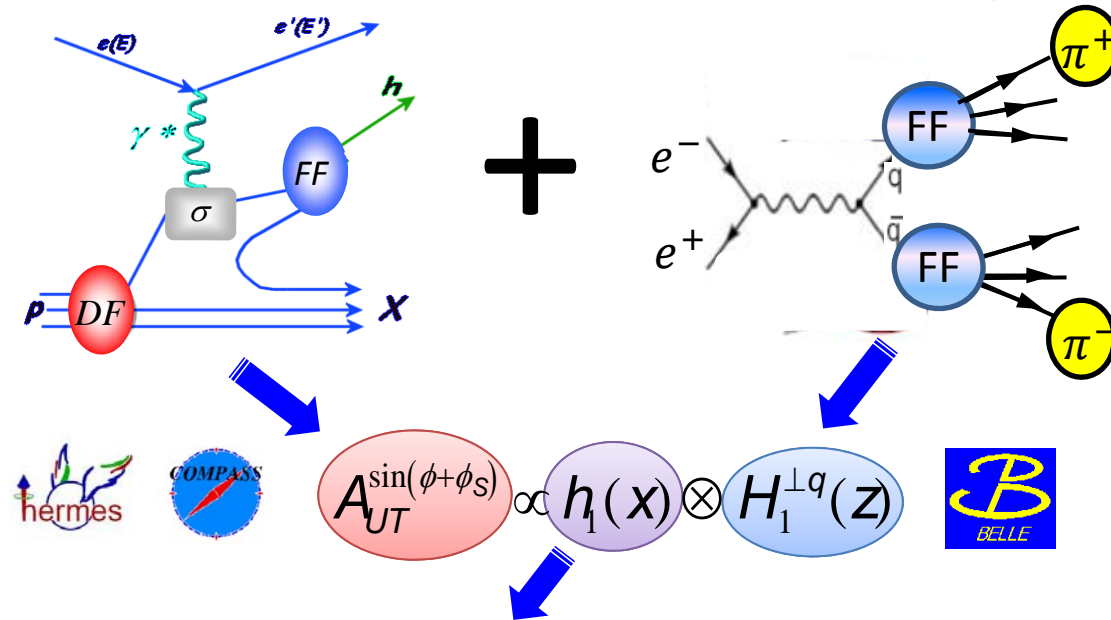
$$\frac{d\sigma(e^+e^- \rightarrow h_1 h_2 X)}{d\Omega dz_1 dz_2 d^2\vec{q}_T} = \frac{3\alpha^2}{Q^2} z_1^2 z_2^2 \left\{ A(y) \mathcal{F}[D_1 \bar{D}_2] + B(y) \cos(2\phi_0) \mathcal{F} \left[(2\hat{h} \cdot \vec{k}_T \hat{h} \cdot \vec{p}_T - \vec{k}_T \cdot \vec{p}_T) \frac{H_1^\perp \bar{H}_2^\perp}{M_1 M_2} \right] \right\}$$



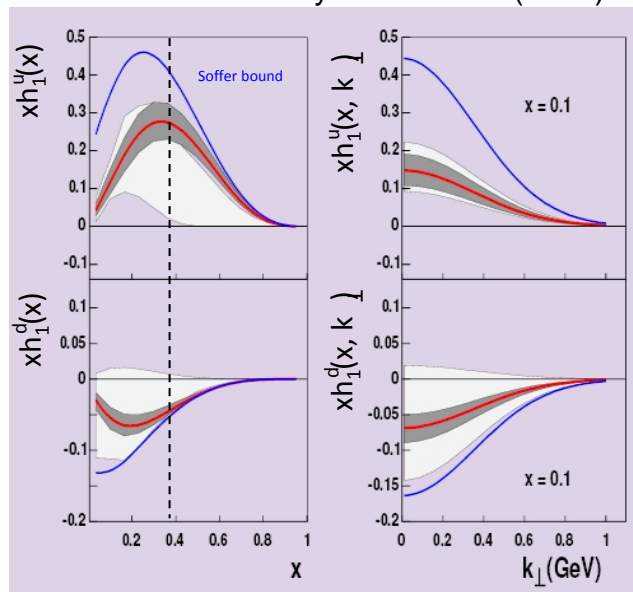
Collins FF from inclusive hadron-pair production in e^+e^-



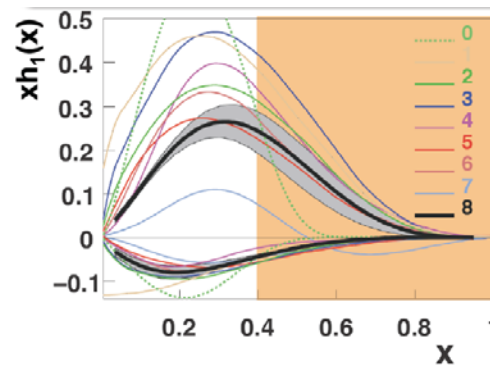
First extraction of Transversity



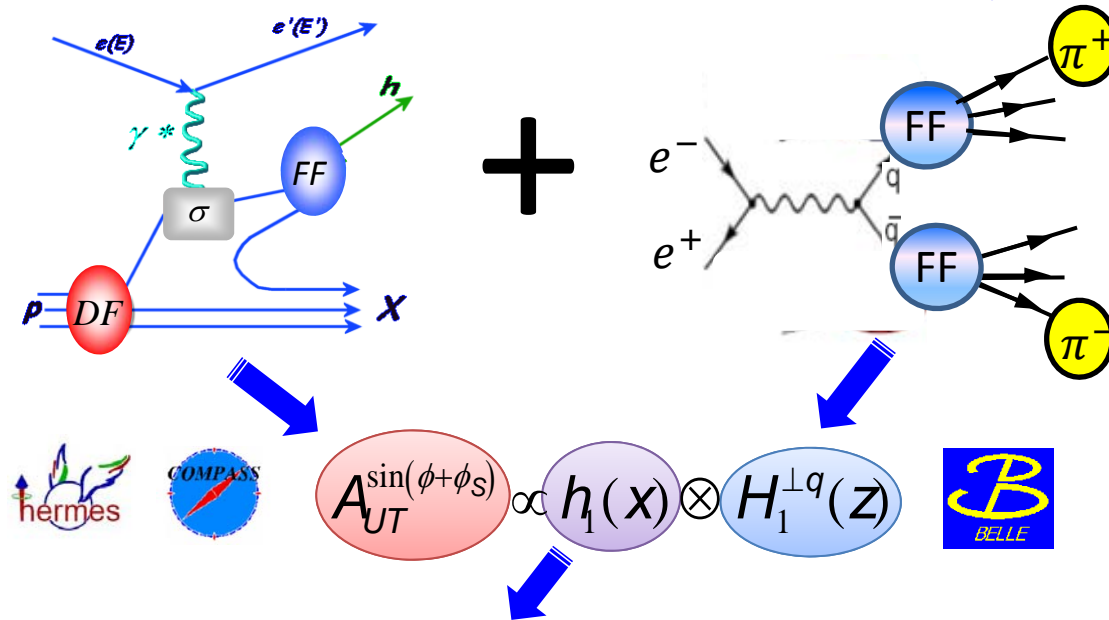
Anselmino et al. Phys. Rev. D 75 (2007)



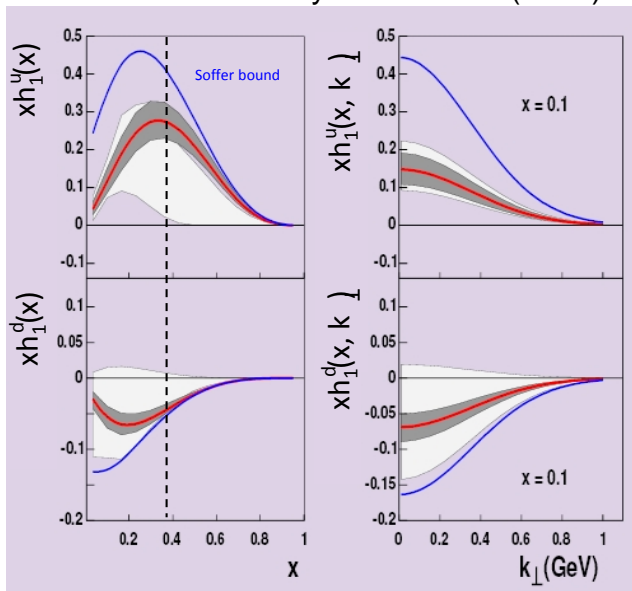
0. chiral color-dielectric model [Barone et al., PLB 390 (97)]
1. Soffer bound [Soffer et al., PRD 65 (02)]
2. $h_1 = g_1$ [Korotov et al., EPJC 18 (01)]
3. Chiral quark-soliton model [Schweitzer et al., PRD 64 (01)]
4. chiral-quark soloiton model [Wakamatsu, PLB 509 (01)]
5. light-cone constituent quark model [Pasquini et al., PRD 72 (05)]
6. quark-diquark model [Cloet, Bentz, Thomas, PLB 659 (08)]
7. quark-diquark model [Bacchetta, Conti, Radici, PRD 78 (08)]
8. Parametrization [Anselmino et al., arXiv 0807.0173]



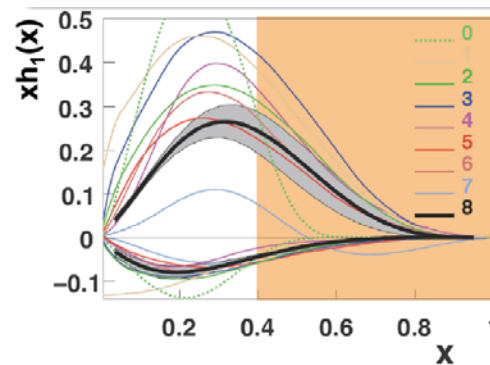
First extraction of Transversity



Anselmino et al. Phys. Rev. D 75 (2007)

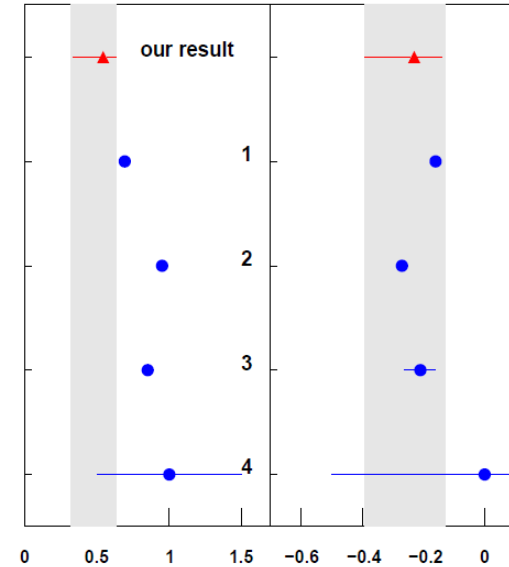


0. chiral color-dielectric model [Barone et al., PLB 390 (97)]
1. Softer bound [Soffer et al., PRD 65 (02)]
2. $h_1 = g_1$ [Korotov et al., EPJC 18 (01)]
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7. quark-diquark model [Bacchetta, Conti, Radici, PRD 78 (08)]
8. Parametrization [Anselmino et al., arXiv 0807.0173]



Tensor charge

$$\delta q = \int_0^1 dx [h_1^q(x) - \bar{h}_1^q(x)]$$



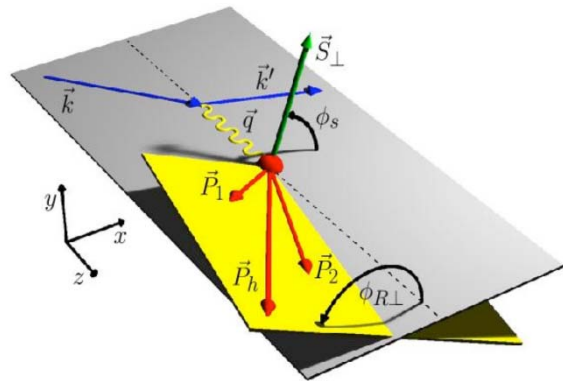
- 1: Quark-diquark model
- 2: Chiral quark soliton model
- 3: Lattice QCD
- 4: QCD sum rules

$$\delta u = 0.54_{-0.22}^{+0.09} \quad \delta d = -0.23_{-0.16}^{+0.09}$$

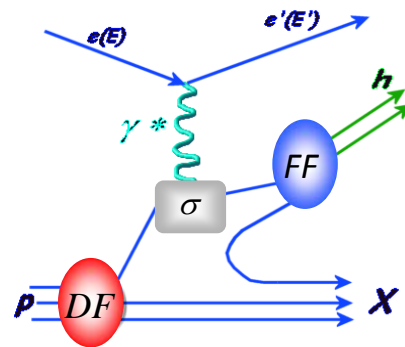
M. Anselmino et al hep-ph:0812.4366

**Need data in valence region
(x > 0.4) -> JLab @ 12 GeV**

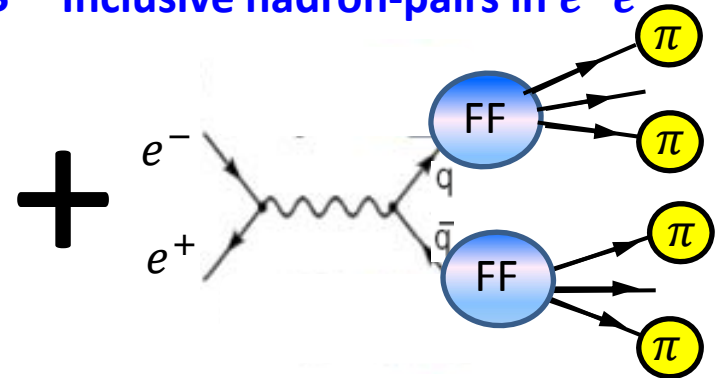
Transversity



2h Semi-Inclusive DIS



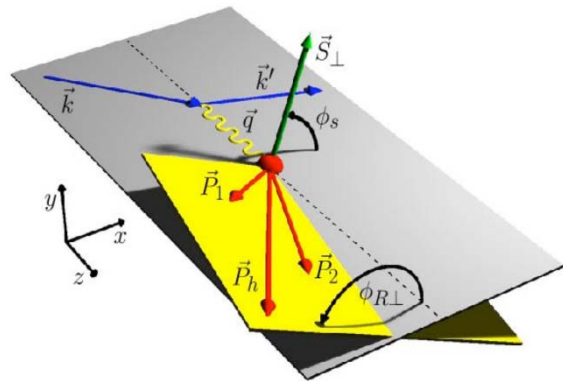
Inclusive hadron-pairs in e^+e^-



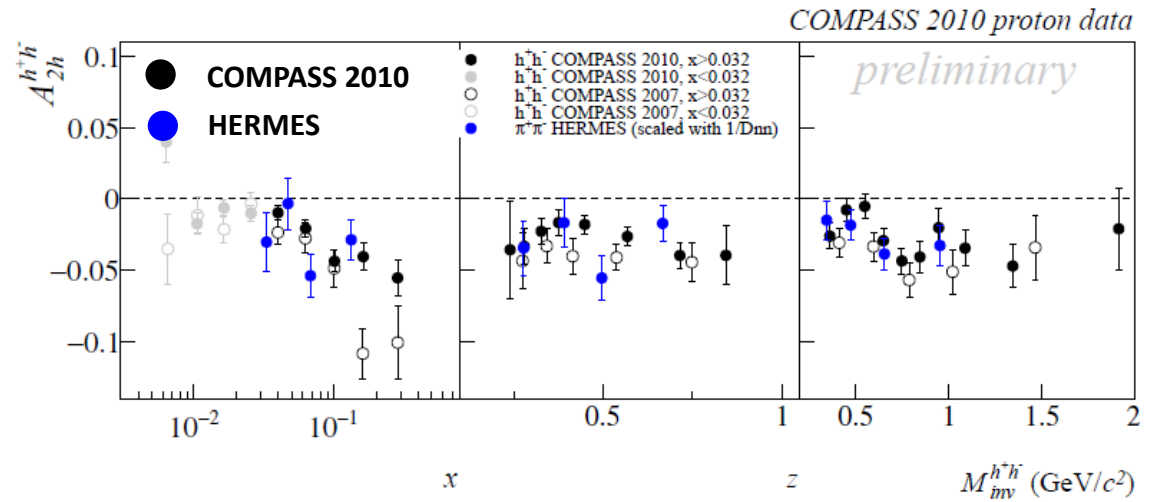
$$A_{UT}^{\sin(\phi_R + \phi_S)\sin\theta} \propto \frac{\sum_q e_q^2 h_1(x, Q^2) H_1^\triangleleft(z, M_h^2, Q^2)}{\sum_q e_q^2 f_1(x, Q^2) D_1^\triangleleft(z, M_h^2, Q^2)}$$

- Survives integration over transverse momentum
- Collinear factorization (simple product)
- DGLAP evolution

Transversity

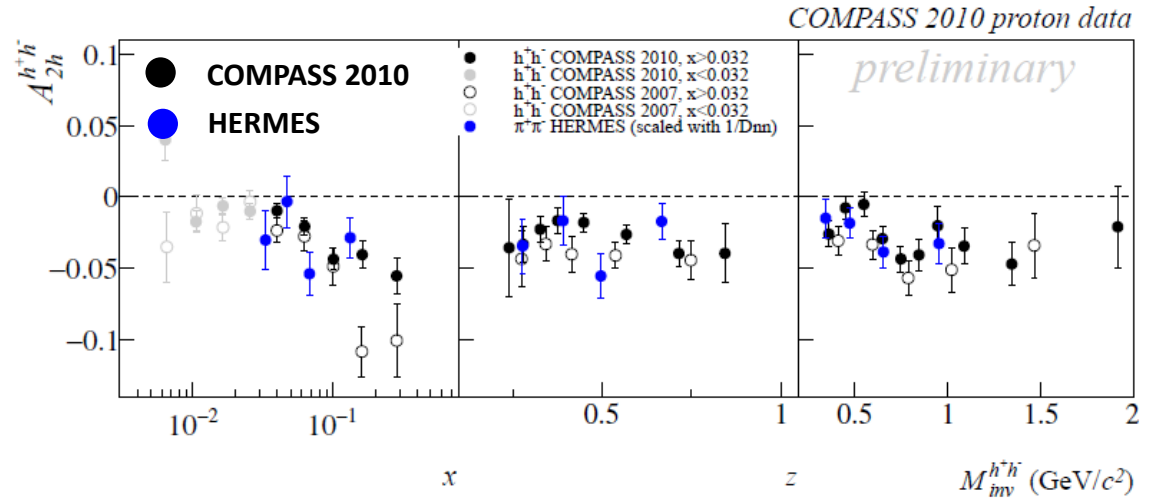
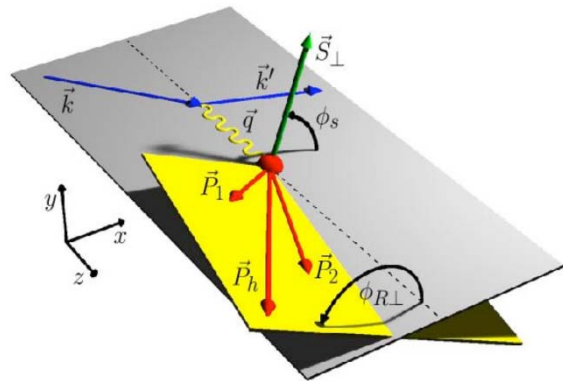


$$A_{UT}^{\sin(\phi_R+\phi_S)\sin\theta} \propto \frac{\sum_q e_q^2 h_1(x, Q^2) H_1^\triangleleft(z, M_h^2, Q^2)}{\sum_q e_q^2 f_1(x, Q^2) D_1^\triangleleft(z, M_h^2, Q^2)}$$



- Survives integration over transverse momentum
- Collinear factorization (simple product)
- DGLAP evolution
- **Large asymmetries @ HERMES & COMPASS**

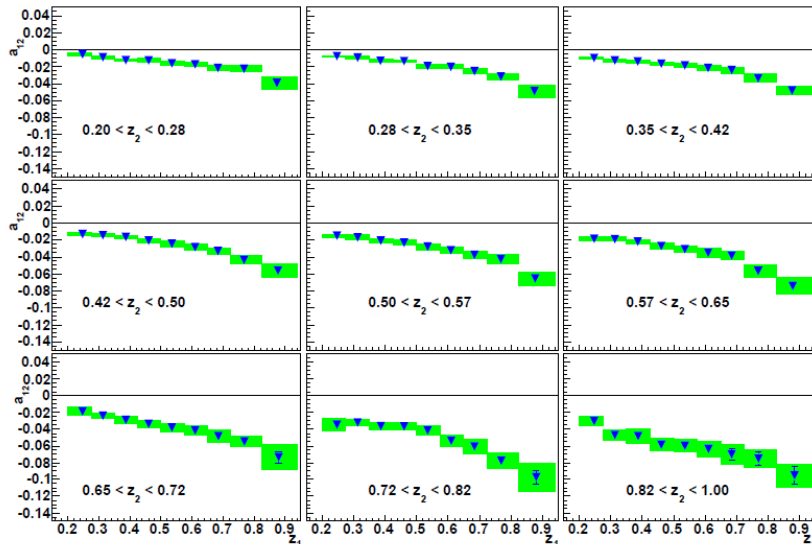
Transversity



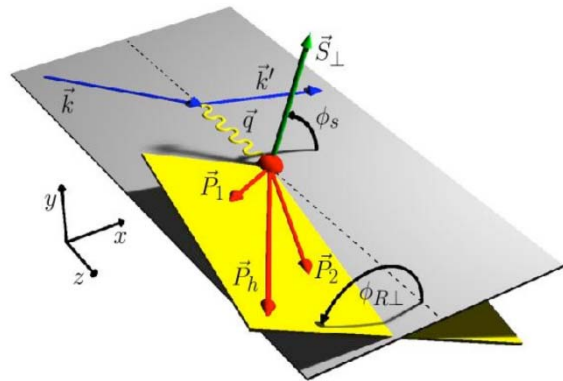
$$A_{UT}^{\sin(\phi_R + \phi_S)\sin\theta} \propto \frac{\sum_q e_q^2 h_1(x, Q^2) H_1^\triangleleft(z, M_h^2, Q^2)}{\sum_q e_q^2 f_1(x, Q^2) D_1^\triangleleft(z, M_h^2, Q^2)}$$

- Survives integration over transverse momentum
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- **Large asymmetries @ HERMES & COMPASS**
- **H_1^\triangleleft chiral-odd, measured at BELLE**

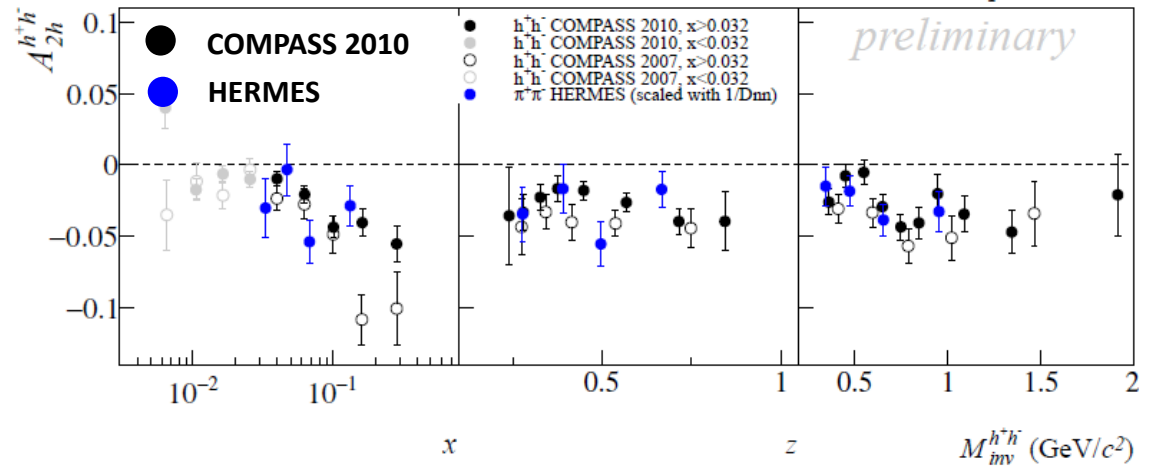
Phys.Rev.Lett.107:072004,2011



Transversity



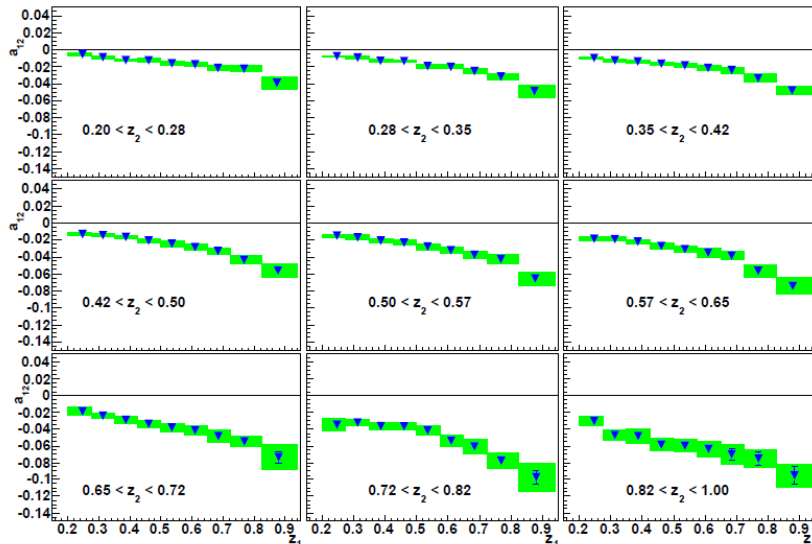
COMPASS 2010 proton data



$$A_{UT}^{\sin(\phi_R + \phi_S)\sin\theta} \propto \frac{\sum_q e_q^2 h_1(x, Q^2) H_1^\triangleleft(z, M_h^2, Q^2)}{\sum_q e_q^2 f_1(x, Q^2) D_1^\triangleleft(z, M_h^2, Q^2)}$$

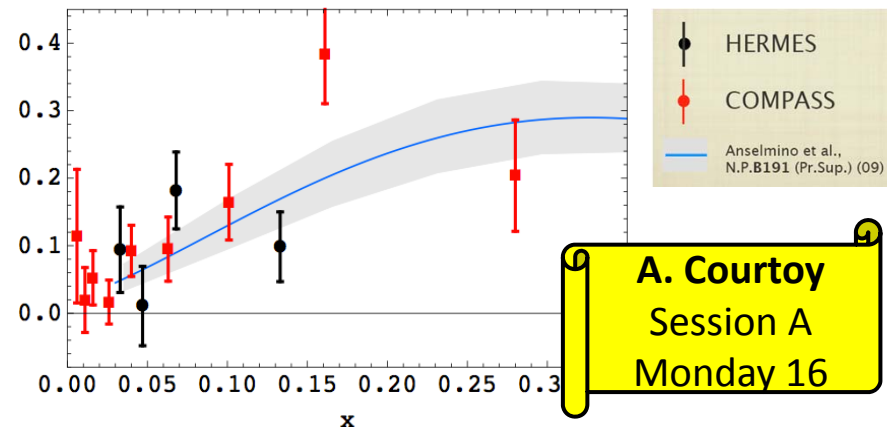
- Survives integration over transverse momentum
- Collinear factorization (simple product)
- DGLAP evolution
- **Large asymmetries @ HERMES & COMPASS**
- **H_1^\triangleleft chiral-odd, measured at BELLE**
- **Independent extraction of transversity!**

Phys.Rev.Lett.107:072004,2011



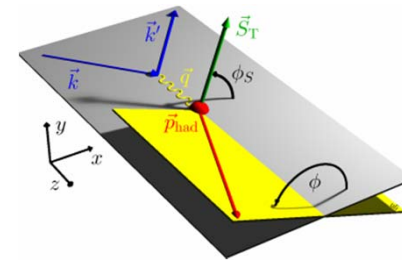
$$x h_1^u(x) - \frac{x}{4} h_1^d(x)$$

Bacchetta et al., PRL 107 (11)



A. Courtoy
Session A
Monday 16

Sivers function



$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2 y^2}{xyQ^2 2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x} \right)$$

$$\left\{ \begin{aligned} & \left[F_{UU,T} + \epsilon F_{UU,L} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} \right] \end{aligned} \right.$$

$$+ \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right]$$

$$+ S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right]$$

$$+ S_L \lambda_l \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right]$$

$$+ S_T \left[\sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi - \phi_S)} \right) \right. \\ \left. + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi - \phi_S)} \right. \\ \left. + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \right. \\ \left. + \sqrt{2\epsilon(1+\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi - \phi_S)} \right]$$

$$+ S_T \lambda_l \left\{ \begin{aligned} & \left[\sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi - \phi_S)} \right] \end{aligned} \right\}$$

Describes correlation between quark transverse momentum and nucleon transverse polarization

Distribution Functions

		quark		
		U	L	T
nucleon	U	f_1		h_1^\perp
	L		g_1	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}^\perp	h_{1T}^\perp

Fragmentation Functions

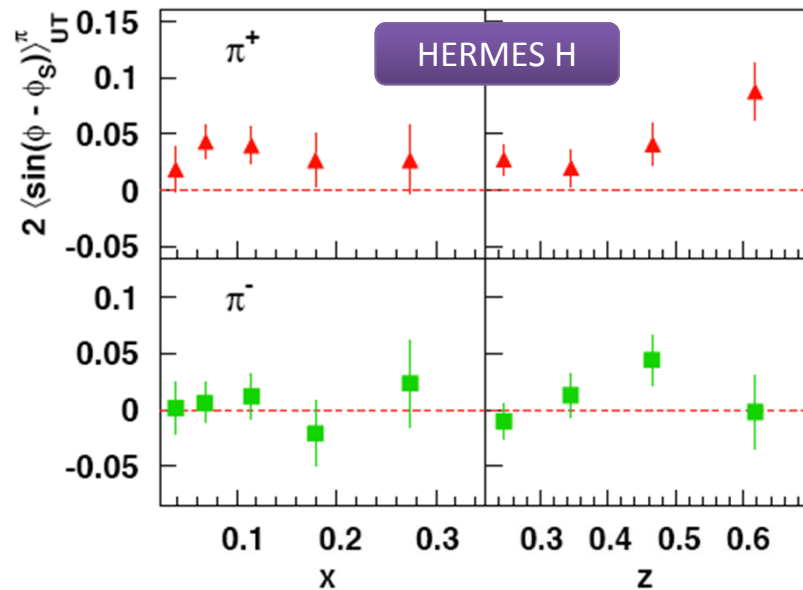
		quark		
		U	L	T
h	U	D_1		H_1^\perp

Sivers function

$$A_{UT}^{\sin(\varphi + \varphi_S)} \propto - \frac{\sum_q e_q^2 f_{1T}^{\perp,q}(x, p_T^2) \otimes D_1^q(z, k_T^2)}{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^q(z, k_T^2)}$$

First evidence by HERMES (2005)

Phys. Rev. Lett. 94 (2005) 012002

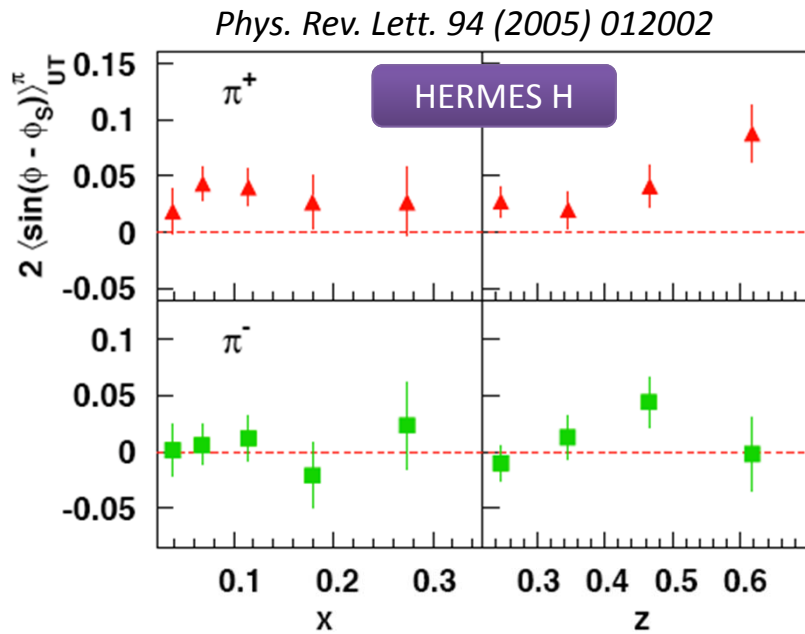


- Transv. pol. H target
- Limited statistics (2002-2003)
- **Non-zero Sivers amplitudes for π^+**
- **Non-zero Sivers function !!**

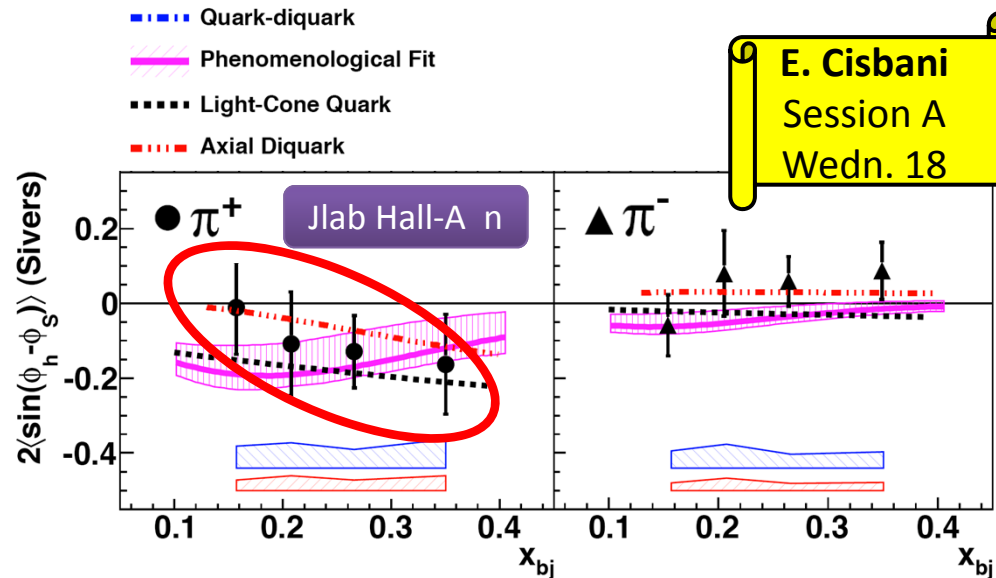
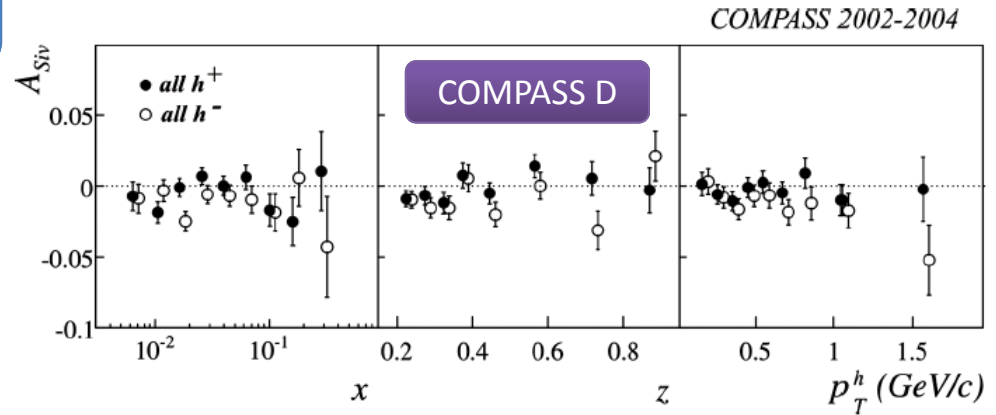
Sivers function

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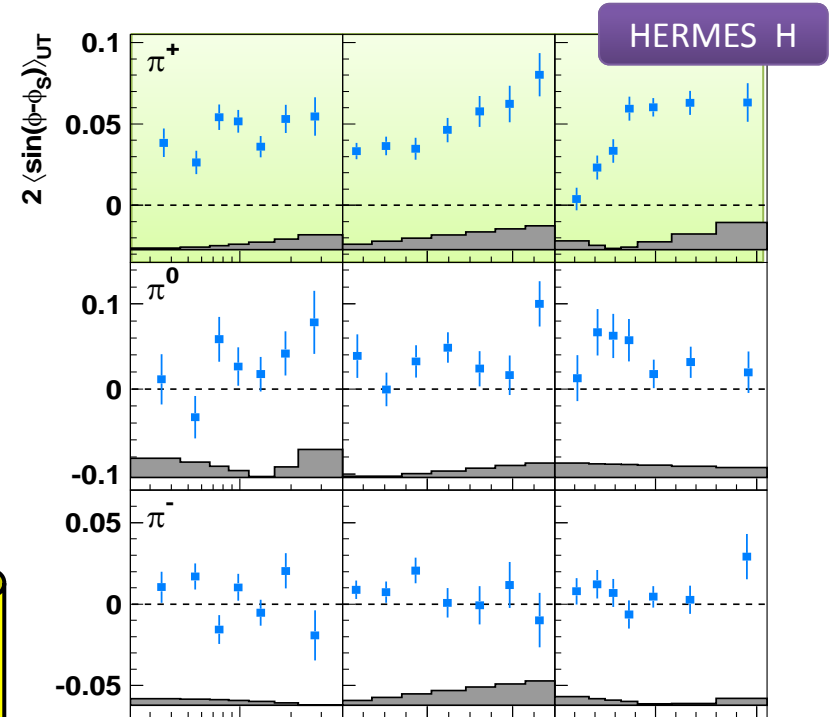
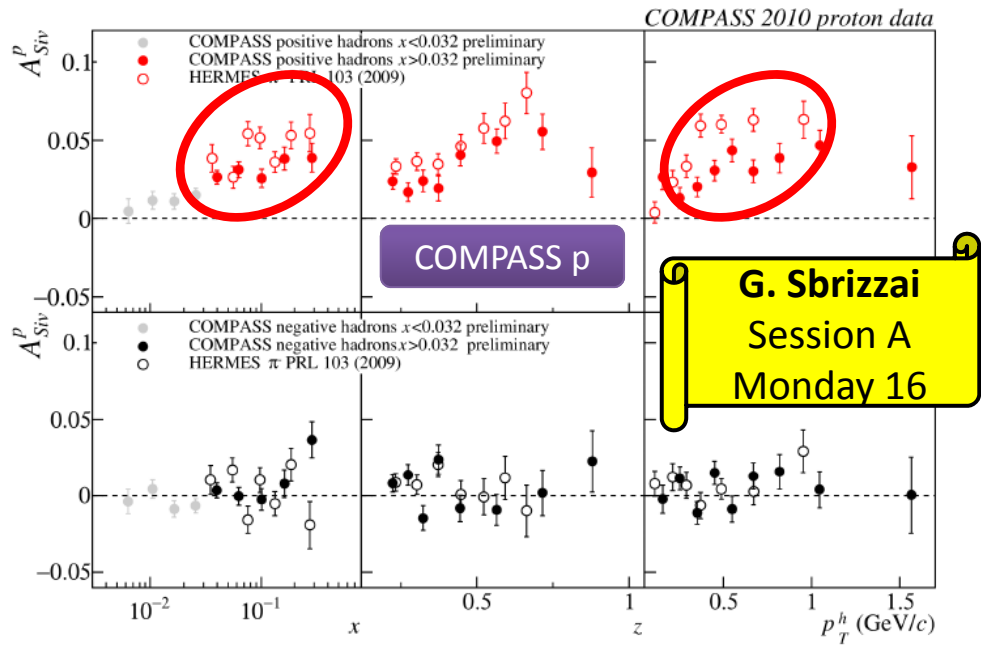


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- **Non-zero Sivers function !!**



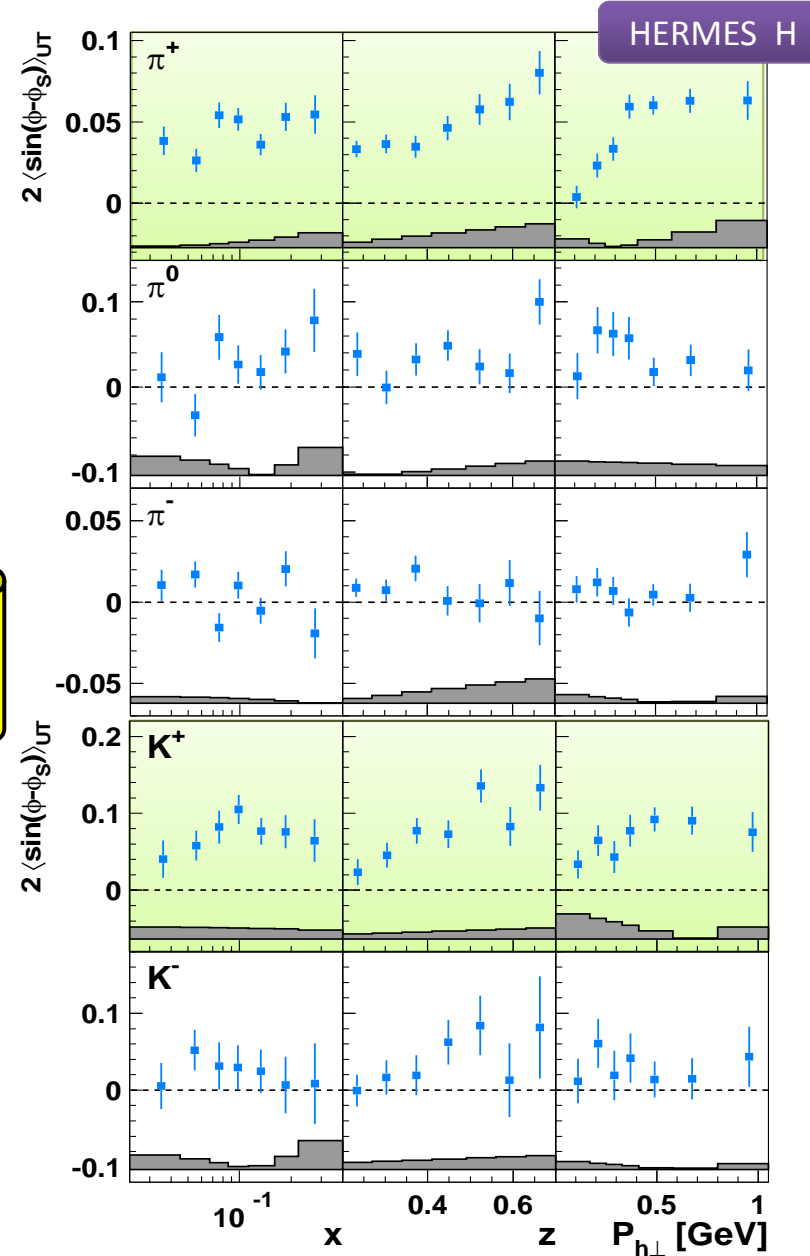
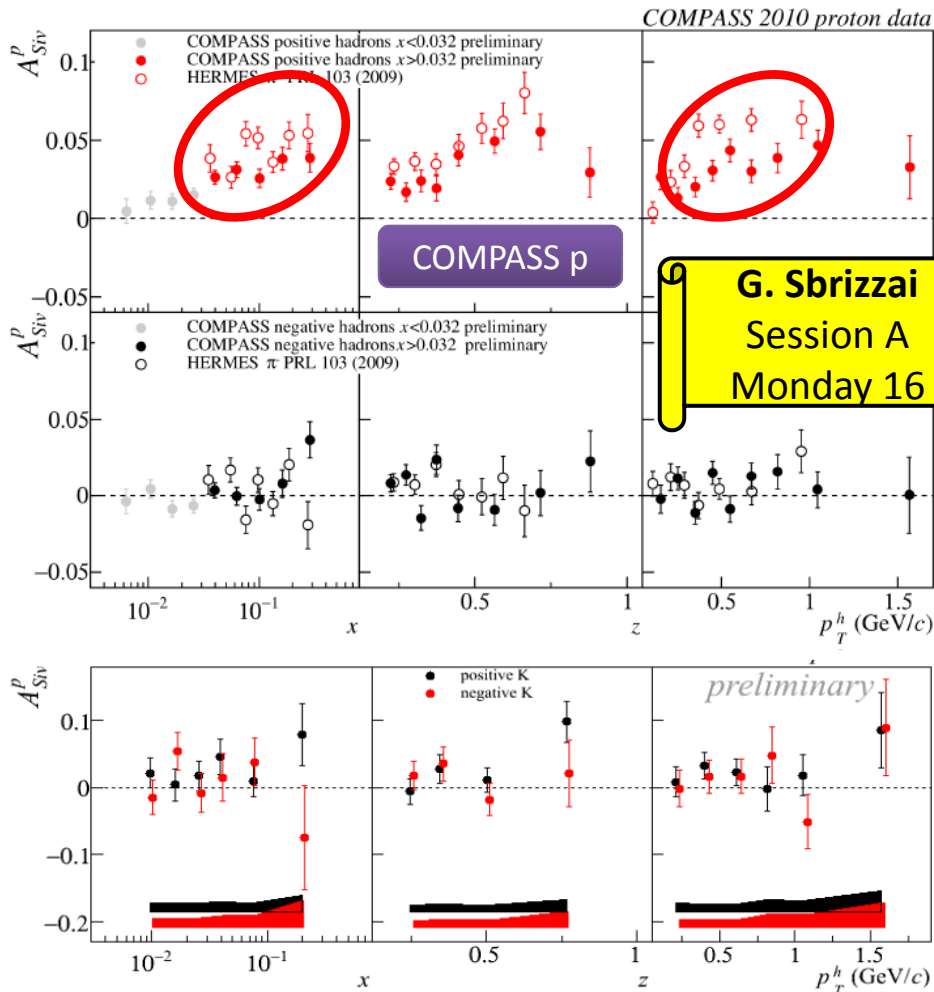
Sivers function

- COMPASS amplitudes smaller than HERMES
- New studies: difference reduces substantially in the low y region ($0.05 < y < 0.1$)
- TMD evolution? (Anselmino)



Sivers function

- COMPASS amplitudes smaller than HERMES
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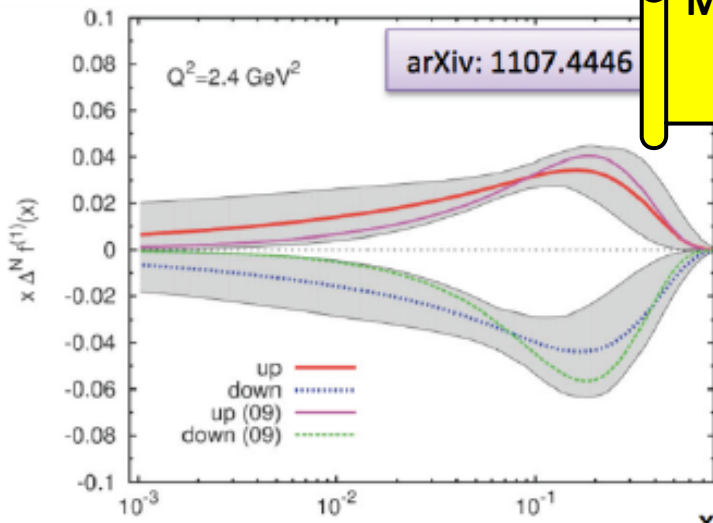
$\pi^+ \equiv |u\bar{d}\rangle$, $K^+ \equiv |u\bar{s}\rangle \rightarrow$ role of sea quarks?

Extracting the Sivers function

Torino

$$A_{UT}^{\sin(\Phi-\Phi_s)} \propto f_{1T}^\perp(x) \otimes D_1^q(z)$$

Fit to Compass d , Hermes p
+ unpol FF



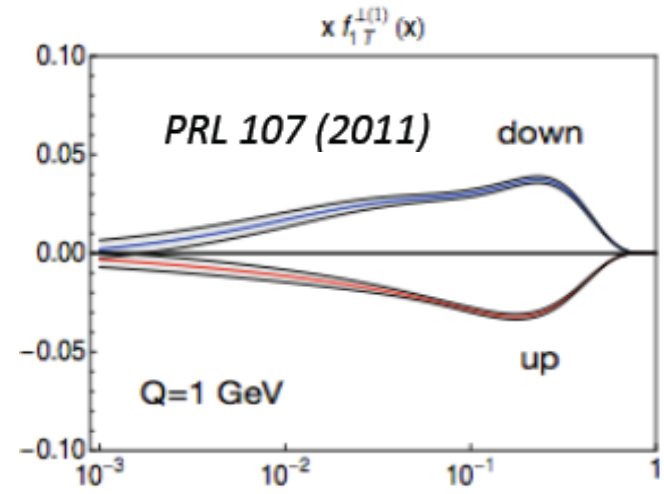
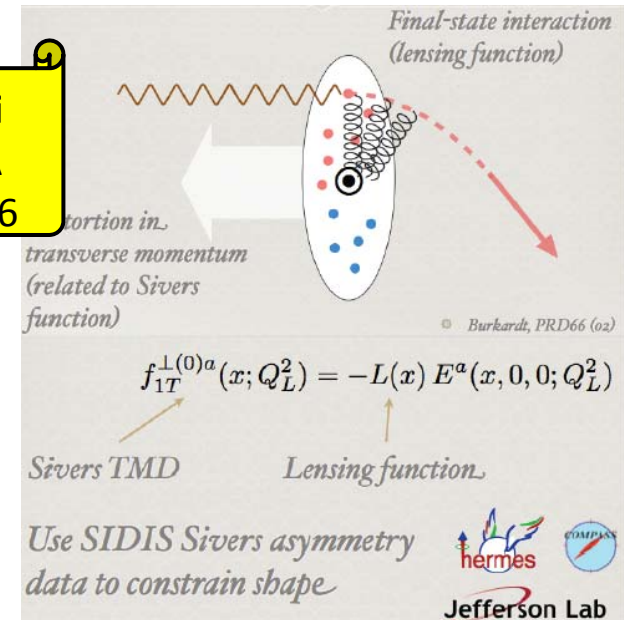
New approach: TMD evolution!

arXiv: 1107.4446

M. Anselmino
Session A
Monday 16

M. Radici
Session A
Monday 16

Pavia

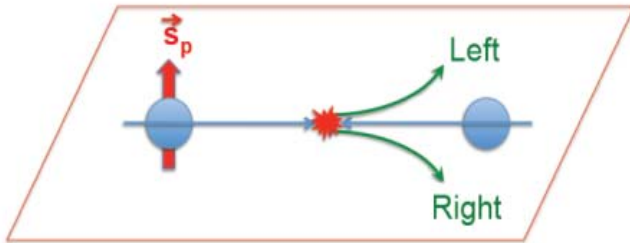


Also put constraints on GPD E_q
and quark OAM!!

Extraction of Sivers
through Bessel weighting
avoids convolution!

L. Gamberg
Session A
Monday 16

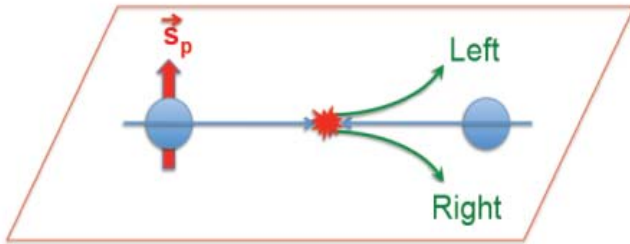
Transversity & Sivers in pp scattering



$$A_N \equiv \frac{\Delta\sigma(l, \vec{s})}{\sigma(l)} = \frac{\sigma(l, \vec{s}) - \sigma(l, -\vec{s})}{\sigma(l, \vec{s}) + \sigma(l, -\vec{s})}$$

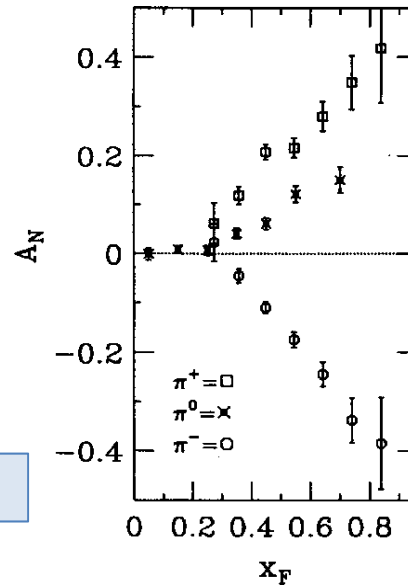
Naive (collinear) pQCD predicts $A_N \approx 0$

Transversity & Sivers in pp scattering



$$A_N \equiv \frac{\Delta\sigma(l, \vec{s})}{\sigma(l)} = \frac{\sigma(l, \vec{s}) - \sigma(l, -\vec{s})}{\sigma(l, \vec{s}) + \sigma(l, -\vec{s})}$$

Naive (collinear) pQCD predicts $A_N \approx 0$



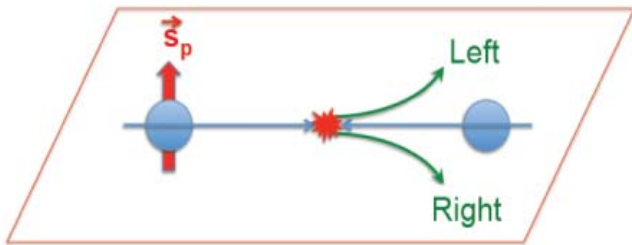
In '90s E-704 exp. @ Fermilab reported large A_N asymmetries (up to 40%!!)

$$x_F \sim \langle z \rangle P_{jet} / P_L$$

Possible explanations:

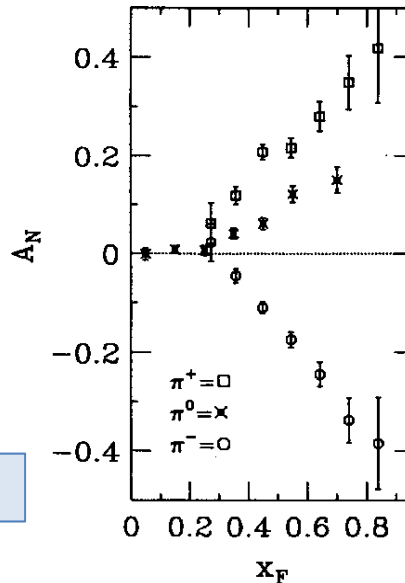
- **Collins effect** (Transv. x Collins FF)
- **Sivers effect** (orbital motion of quarks)
- **Twist-3 effects**
- Combination of above

Transversity & Sivers in pp scattering



$$A_N \equiv \frac{\Delta\sigma(l, \vec{s})}{\sigma(l)} = \frac{\sigma(l, \vec{s}) - \sigma(l, -\vec{s})}{\sigma(l, \vec{s}) + \sigma(l, -\vec{s})}$$

Naive (collinear) pQCD predicts $A_N \approx 0$

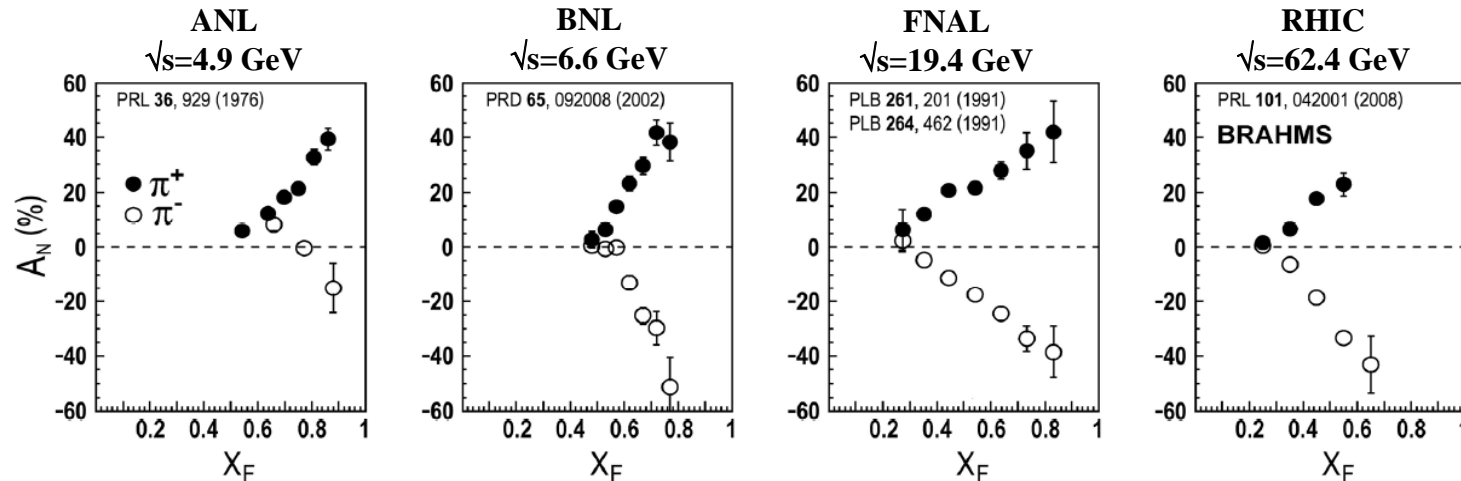


In '90s E-704 exp. @ Fermilab reported large A_N asymmetries (up to 40%!!)

$$x_F \sim \langle z \rangle P_{jet} / P_L$$

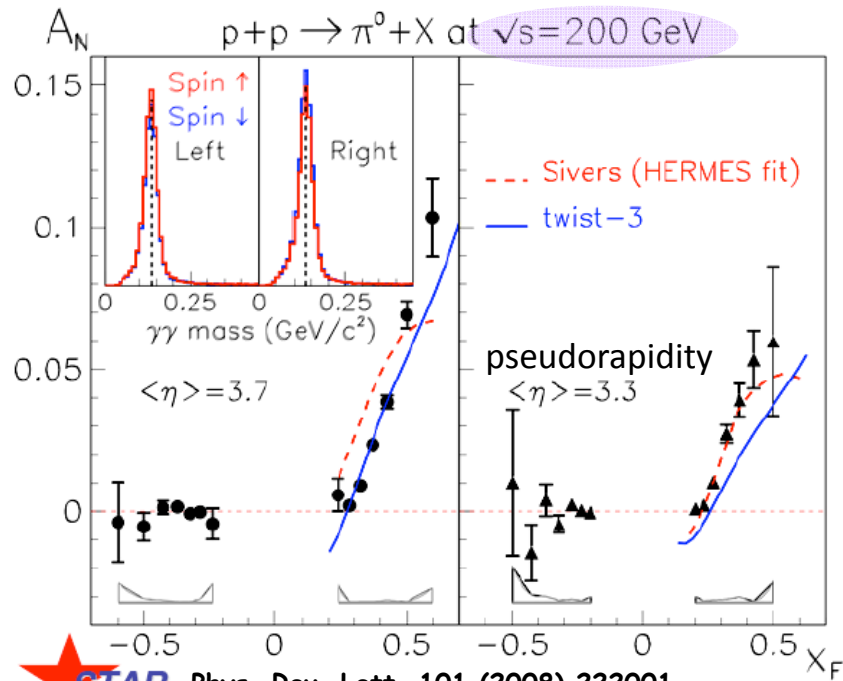
Possible explanations:

- **Collins effect** (Transv. x Collins FF)
- **Sivers effect** (orbital motion of quarks)
- **Twist-3 effects**
- Combination of above

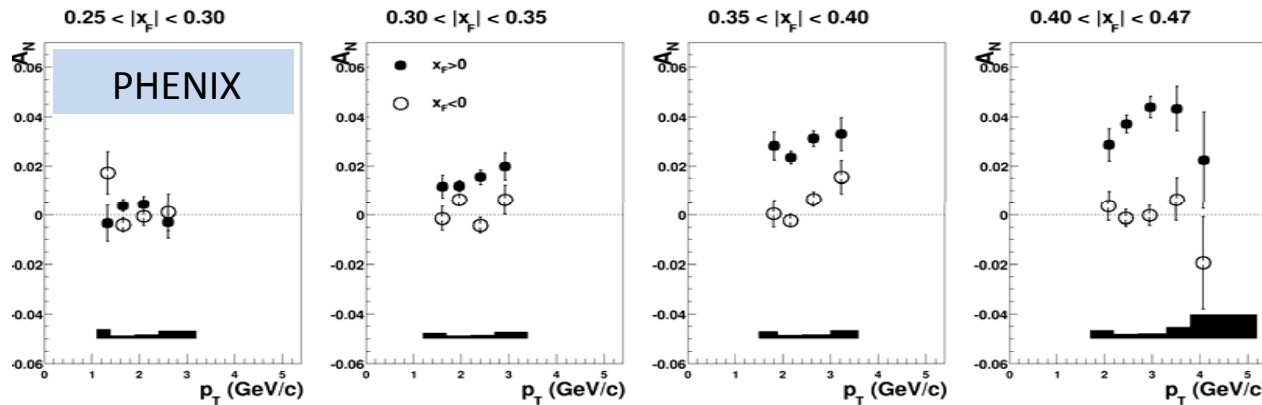
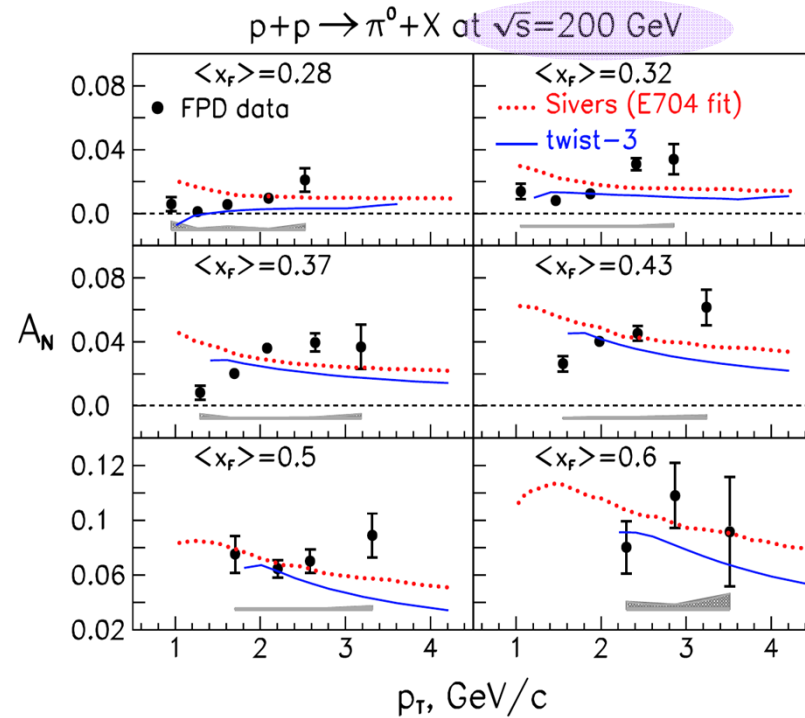


- Reproduced by various exps. over 35 years!
- **persistent with energy!**

Transversity & Sivers in pp scattering: RHIC



STAR Phys. Rev. Lett. 101 (2008) 222001



- Large asymmetries measured by STAR, PHENIX and BRAHMS
- No strong dependence on \sqrt{s} from 19.4 to 200 GeV
- Admixture of Transversity, Sivers and Twist-3 effects?

Boer-Mulders function: DY

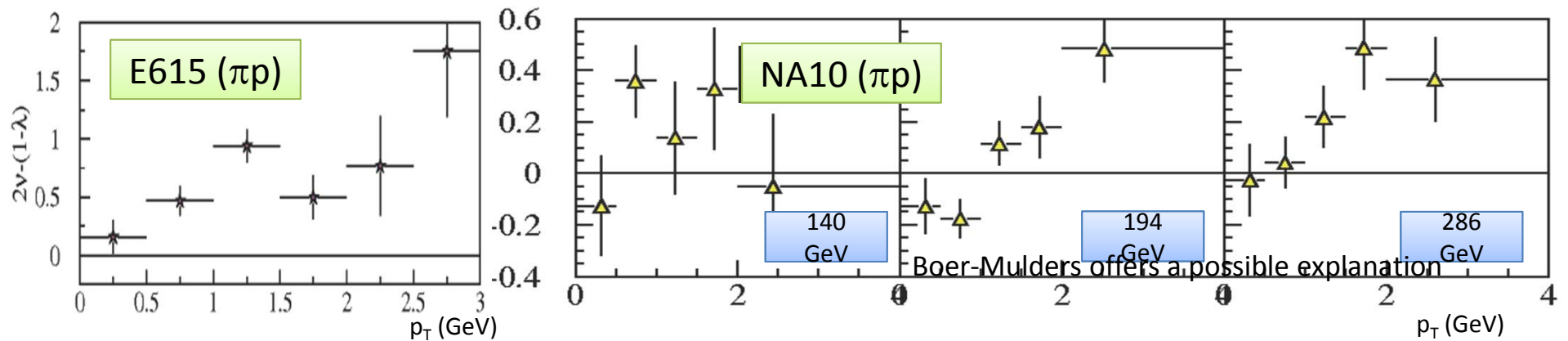
$$\frac{d\sigma^{hp \rightarrow eeX}}{d\Omega} \propto 1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi$$

$$(1 - \lambda) - 2\nu = 0 \quad \text{Lam-Tung rel.}$$

Boer-Mulders function: DY

$$\frac{d\sigma^{hp \rightarrow eeX}}{d\Omega} \propto 1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi$$

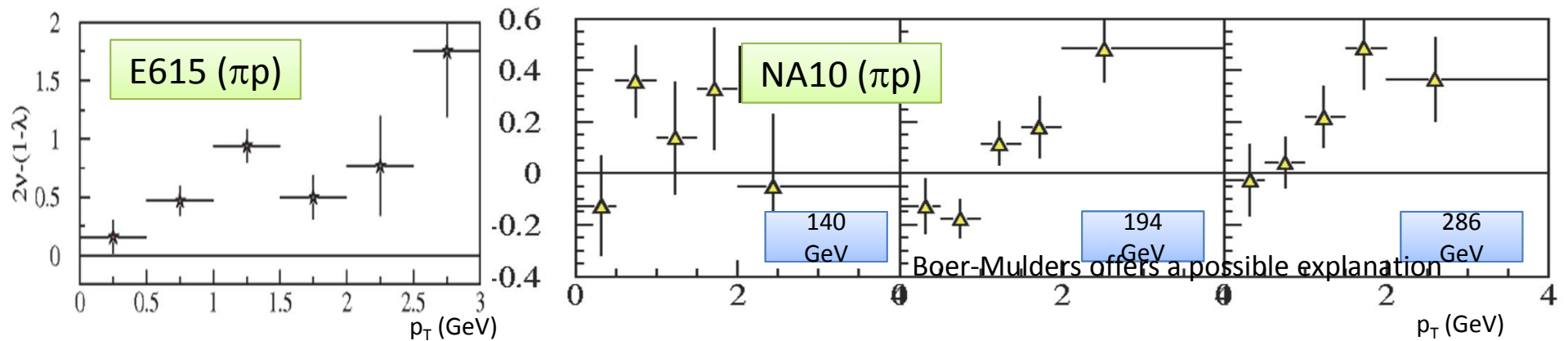
$(1 - \lambda) - 2\nu = 0$ **Lam-Tung rel.**
violated in pion-induced DY!



Boer-Mulders function: DY

$$\frac{d\sigma^{hp \rightarrow eex}}{d\Omega} \propto 1 + \lambda \cos^2 \theta + \mu \sin 2\theta \cos \phi + \frac{\nu}{2} \sin^2 \theta \cos 2\phi$$

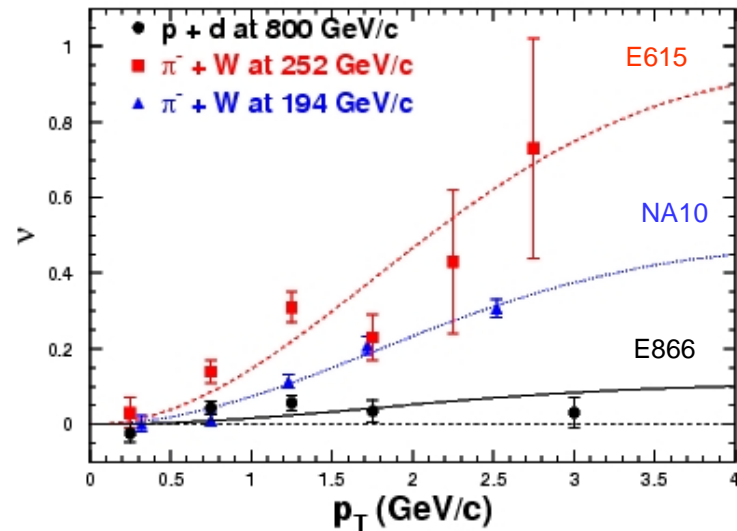
$(1 - \lambda) - 2\nu = 0$ Lam-Tung rel.
violated in pion-induced DY!



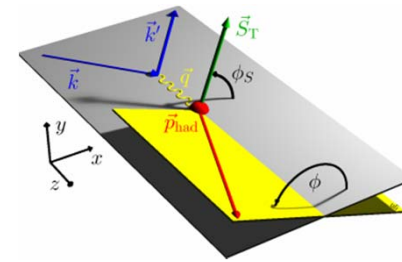
Naive parton model:
 $d\sigma/d\Omega \propto 1 + \cos^2 \vartheta \Rightarrow \nu \approx 0$

Boer-Mulders effect offers an explanation:

$$\nu \approx h_{1q}^\perp \times h_{1\bar{q}}^\perp$$



Boer-Mulders function: SIDIS



$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2 y^2}{xyQ^2 2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right)$$

$$\left\{ \begin{aligned} & [F_{UU,T} + \epsilon F_{UU,L} \\ & + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)}] \end{aligned} \right.$$

$$+ \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right]$$

$$+ S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right]$$

$$+ S_L \lambda_l \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right]$$

$$+ S_T \left[\begin{aligned} & \sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi - \phi_S)} \right) \\ & + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi - \phi_S)} \\ & + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \\ & + \sqrt{2\epsilon(1+\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi - \phi_S)} \end{aligned} \right]$$

$$+ S_T \lambda_l \left\{ \begin{aligned} & \sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi - \phi_S)} \end{aligned} \right\}$$

Describes correlation between quark transverse momentum and transverse spin in unpolarized nucleon

Distribution Functions

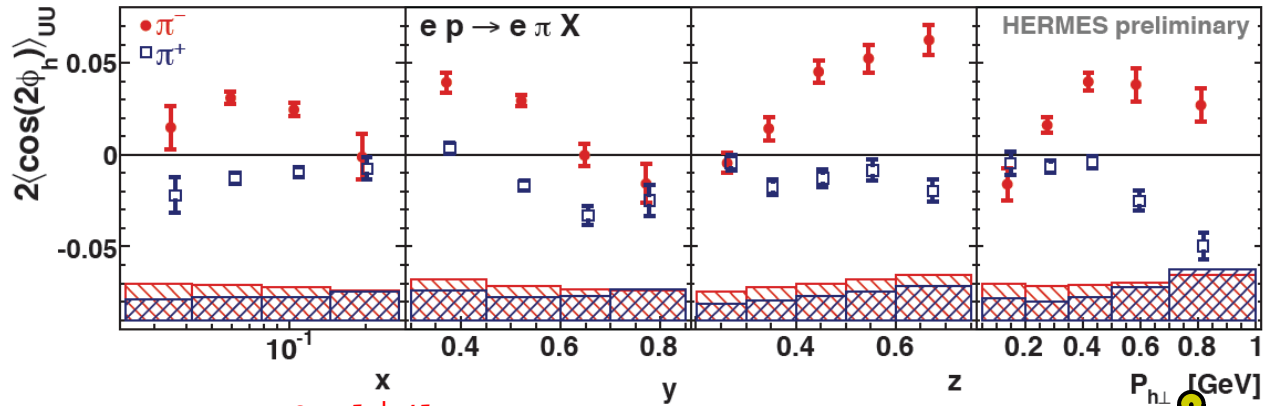
		quark		
		U	L	T
nucleon	U	f_1		h_T^\perp
	L		g_1	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}^\perp	h_1 h_{1T}^\perp

Fragmentation Functions

		quark		
		U	L	T
h	U	D_1		H_1^\perp

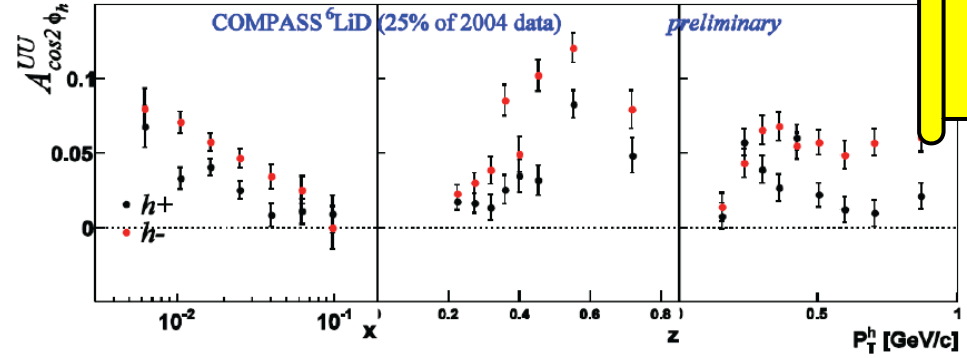
Boer-Mulders function: SIDIS ($\cos 2\phi$)

$$\sigma_{UU}^{\cos(2\phi)} \propto h_1^\perp \otimes H_1^\perp + [f_1 \otimes D_1 + \dots]/Q^2$$



- Opposite signs for π^+ / π^-
- Similar results on H & D target
→ same sign for u and d?
- K^+ / K^- amplitudes are larger and of same sign

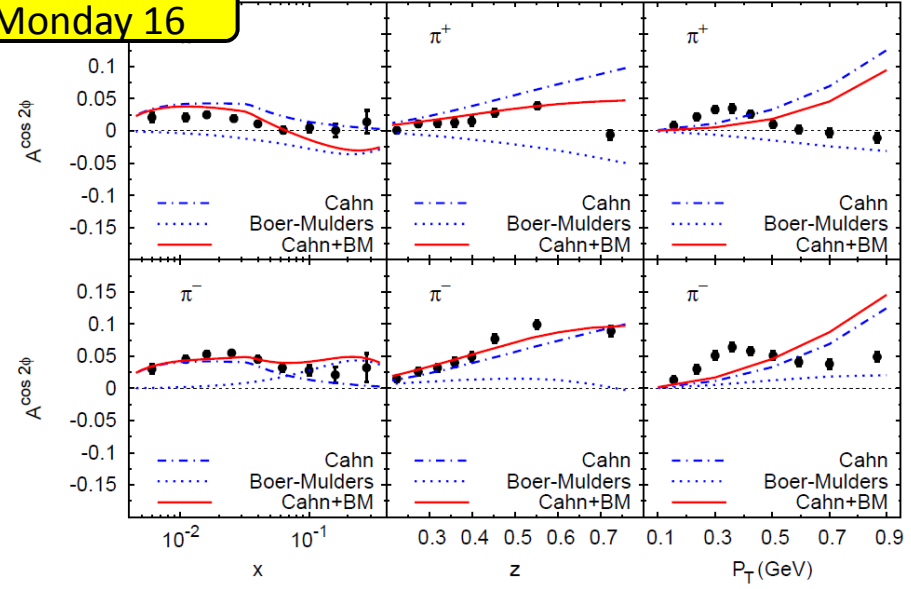
- Same sign for h^+ / h^-



G. Sbrizzai
Session A
Monday 16

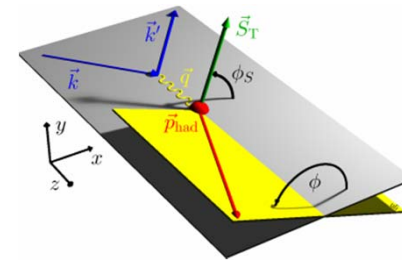
COMPASS D

arXiv: 0912.5194



- Amplitudes are significant
- Clear evidence of BM effect
- Strong dependence on kinematic variables
- Issue on data consistency for h^+
- New 2-dim extraction from COMPASS!
- Discrepancy with theory: uncert. on Chan + HT

Worm-gear h_{1L}^\perp



$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2 y^2}{xyQ^2 2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x} \right)$$

$$\left\{ \begin{aligned} & \left[F_{UU,T} + \epsilon F_{UU,L} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} \right] \\ & + \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right] \\ & + S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right] \\ & + S_L \lambda_l \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right] \\ & + S_T \left[\sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi-\phi_S)} + \epsilon F_{UT,L}^{\sin(\phi-\phi_S)} \right) \right. \\ & \left. + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi+\phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi-\phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi-\phi_S)} \right] \\ & + S_T \lambda_l \left[\sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi-\phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi-\phi_S)} \right] \end{aligned} \right\}$$

Describes the probability to find transversely polarized quarks in a longitudinally polarized nucleon

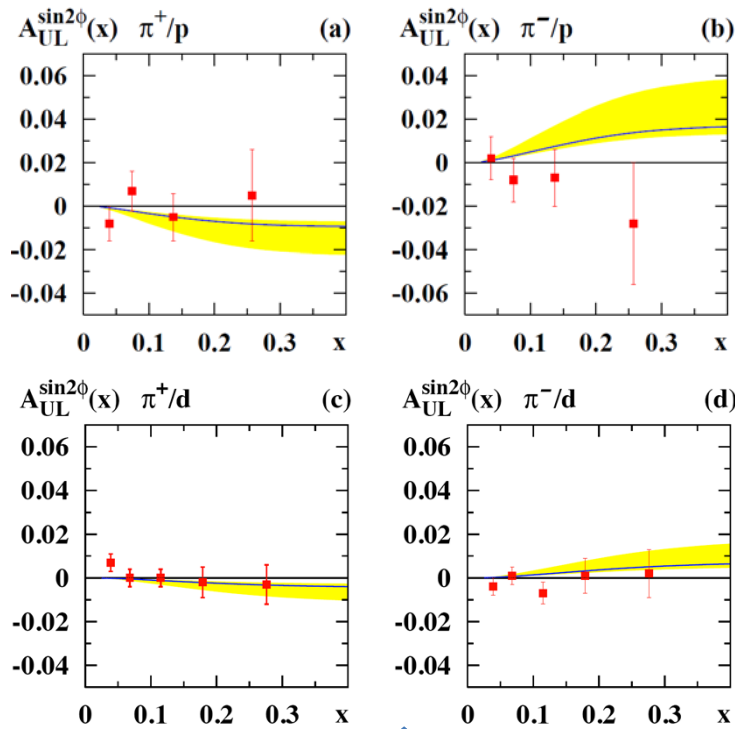
Distribution Functions

		quark		
		U	L	T
nucleon	U	f_1		h_1^\perp
	L		g_1	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}^\perp	h_{1T}^\perp

Fragmentation Functions

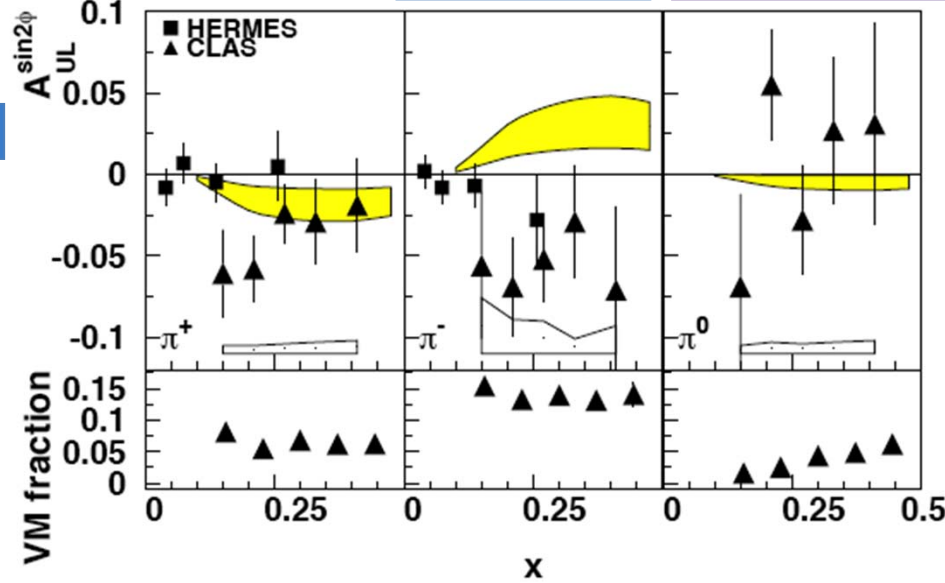
		quark		
		U	L	T
h	U	D_1		H_1^\perp

Worm-gear h_{1L}^\perp

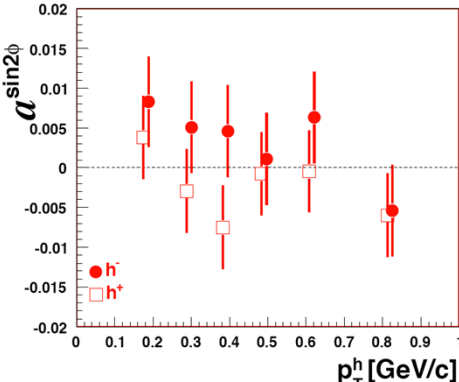
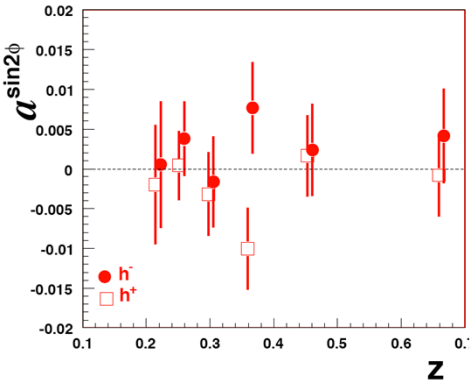
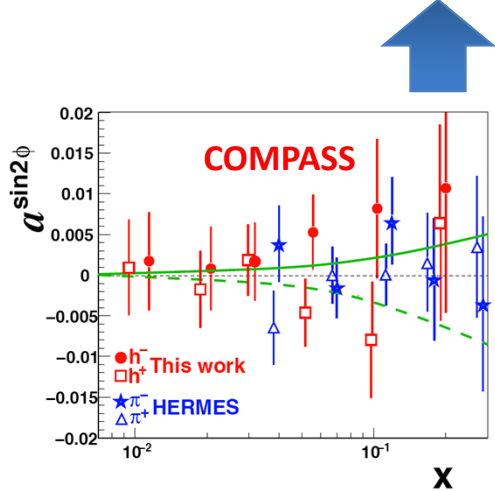


arXiv: 0902.0689

Proton arXiv: 1007.1562



- Consistent with zero @ HERMES & COMPASS
- Significant amplitude @ CLAS!



Deuteron

arXiv: 1003.4549

G. Sbrizzai
 Session A
 Monday 16

Worm-gear g_{1T}^\perp

$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2 y^2}{xyQ^2 2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x} \right)$$

$$\left\{ \begin{aligned} & \left[F_{UU,T} + \epsilon F_{UU,L} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} \right] \\ & + \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right] \\ & + S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right] \\ & + S_L \lambda_l \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right] \\ & + S_T \left[\sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi-\phi_S)} + \epsilon F_{UT,L}^{\sin(\phi-\phi_S)} \right) \right. \\ & \left. + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi+\phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi-\phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi-\phi_S)} \right] \end{aligned} \right.$$

$$\left. \left\{ \begin{aligned} & + S_T \lambda_l \left[\sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi-\phi_S)} \right. \right. \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi-\phi_S)} \right] \end{aligned} \right\}$$

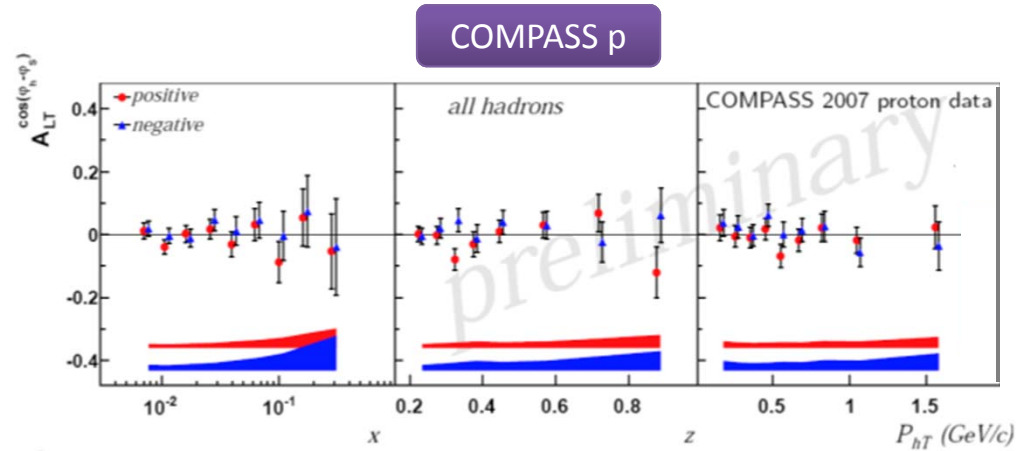
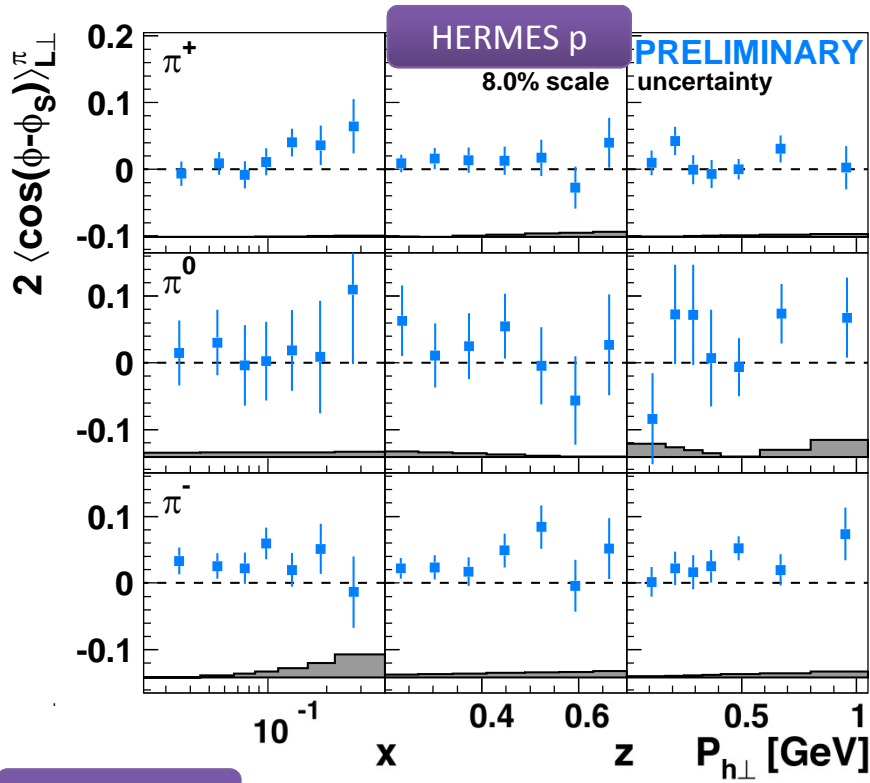
Describes the probability to find longitudinally polarized quarks in a transversely polarized nucleon!

- The only TMD that is both **chiral-even** and **naïve-T-even**
- requires interference between wave funct. components that differ by 1 unit of OAM

Distribution Functions				
		quark		
		U	L	T
nucleon	U	f_1		h_1^\perp
	L		g_1	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}^\perp	h_1 h_{1T}^\perp

Fragmentation Functions				
		quark		
		U	L	T
h	U	D_1		H_1^\perp

Worm-gear g_{1T}^\perp



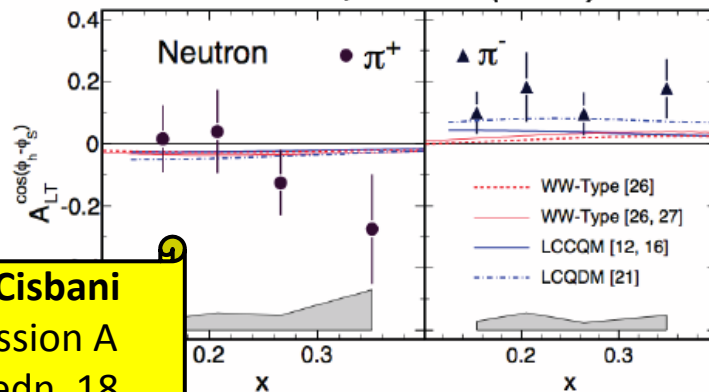
Statistics not enough to investigate relations supported by many theoretical models:

$$g_{1T}^q = -h_{1L}^{\perp q} \quad (\text{supported by Lattice QCD and first data})$$

$$g_{1T}^{q(1)}(x) \approx x \int_x^1 \frac{dy}{y} g_1^q(y) \quad (\text{Wandura-Wilczek type approximation})$$

Jlab Hall-A n

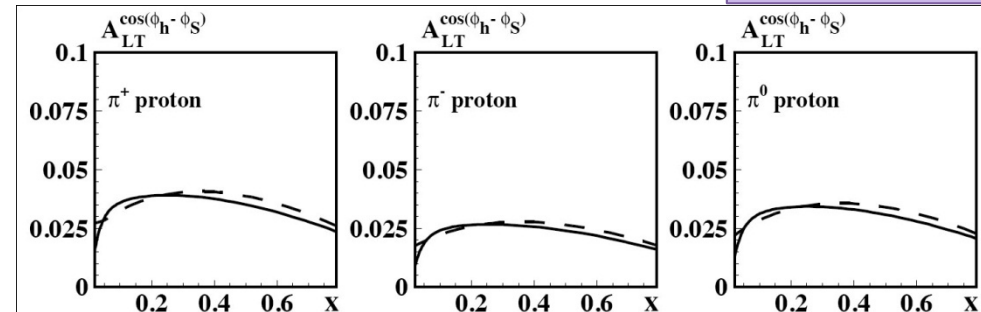
Jlab n, PRL108(2012)



E. Cisbani
Session A
Wedn. 18

From constituent quark model:

arXiv: 0903.1271



Pretzelosity

$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi dP_{h\perp}^2} = \frac{\alpha^2 y^2}{xyQ^2 2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right)$$

$$\left\{ \begin{aligned} & \left[F_{UU,T} + \epsilon F_{UU,L} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} \right] \\ & + \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right] \\ & + S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right] \\ & + S_L \lambda_l \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right] \\ & + S_T \left[\sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi - \phi_S)} \right) \right. \\ & \left. + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi - \phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi - \phi_S)} \right] \\ & + S_T \lambda_l \left[\sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \right. \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi - \phi_S)} \right] \end{aligned} \right\}$$

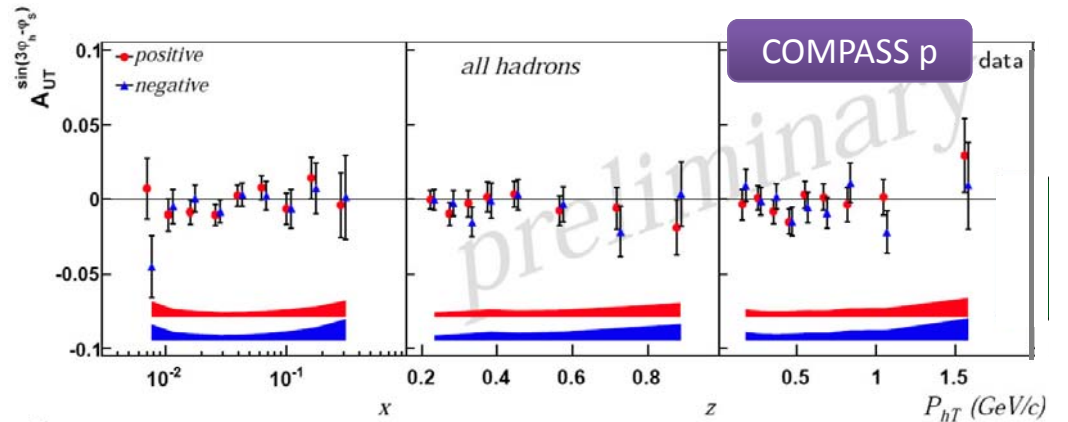
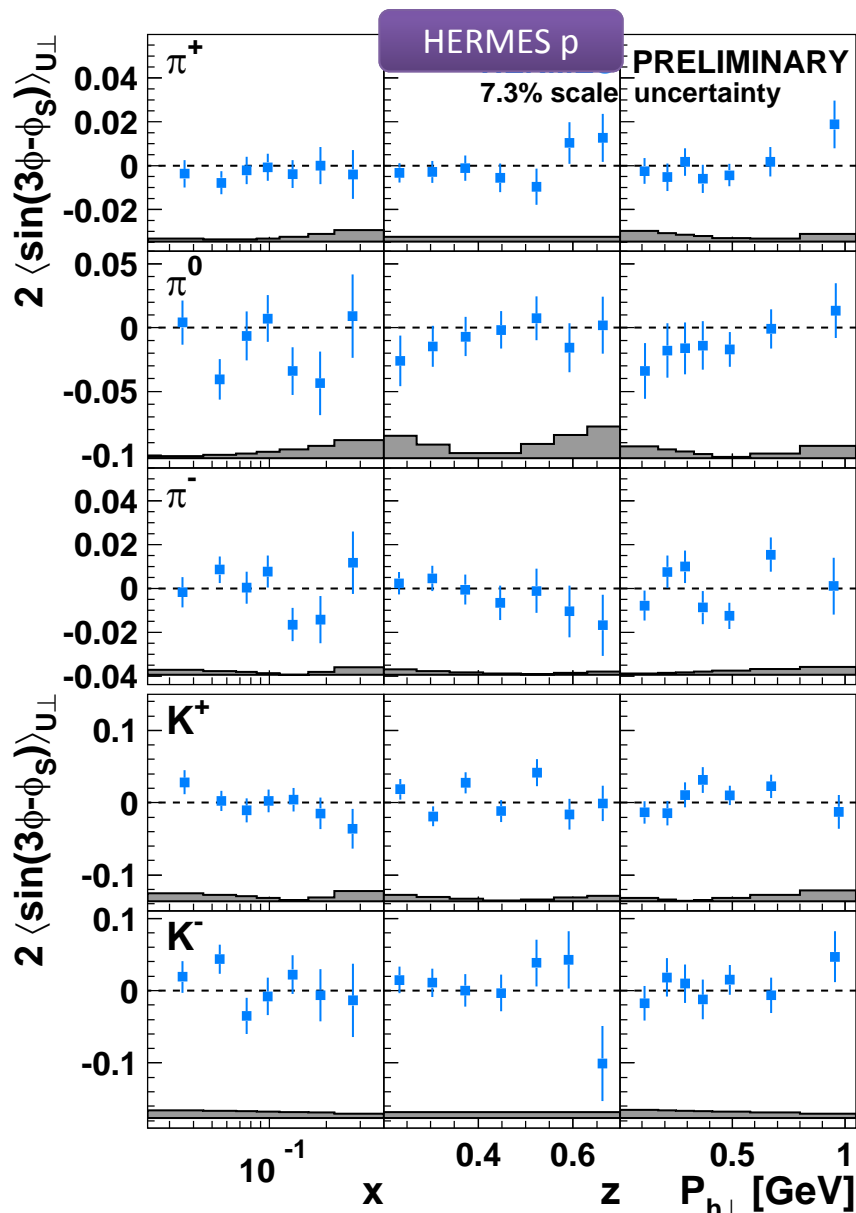
Describes correlation between quark transverse momentum and transverse spin in a transversely pol. nucleon

- Sensitive to the D-wave component
- Sensitive to **non-spherical shape** of the nucleon
- Kinematically suppressed w.r.t. Collins and Sivers by a factor $P_{h\perp}^2$

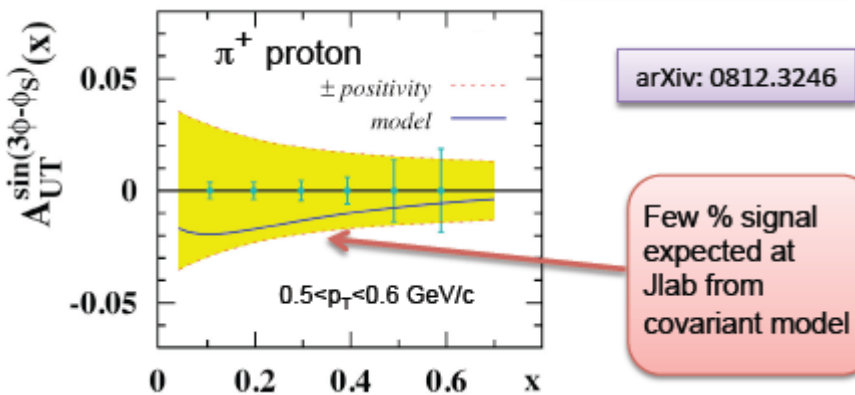
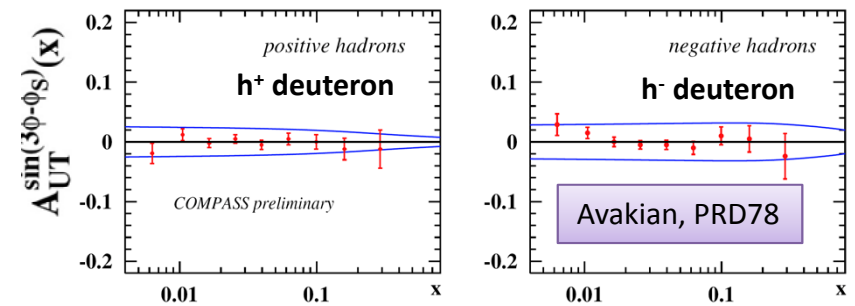
Distribution Functions				
		quark		
		U	L	T
nucleon	U	f_1		h_1^\perp
	L		g_1	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}^\perp	h_{1T}^\perp

Fragmentation Functions				
		quark		
		U	L	T
h	U	D_1		H_1^\perp

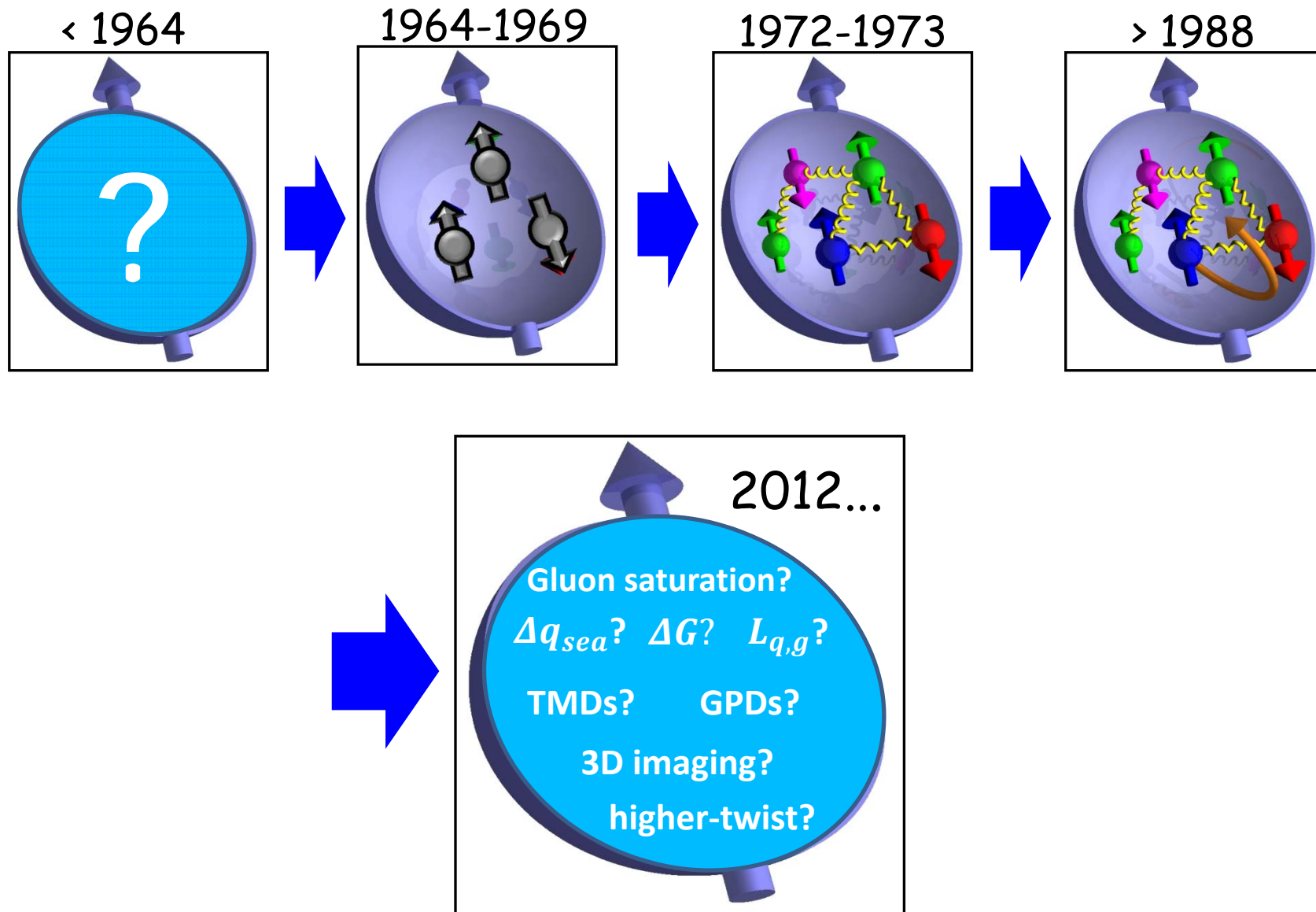
Pretzelocity



- Proton and deuteron data consistent with zero
 - Statistical power of existing data is not enough to observe significant signals



Looking deeply into the proton



The treasure map!

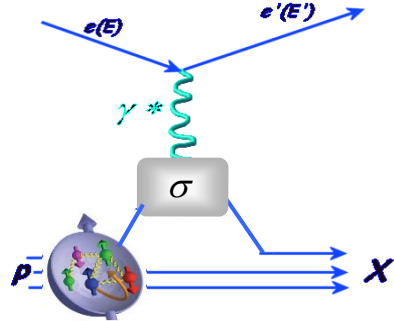


Back-up

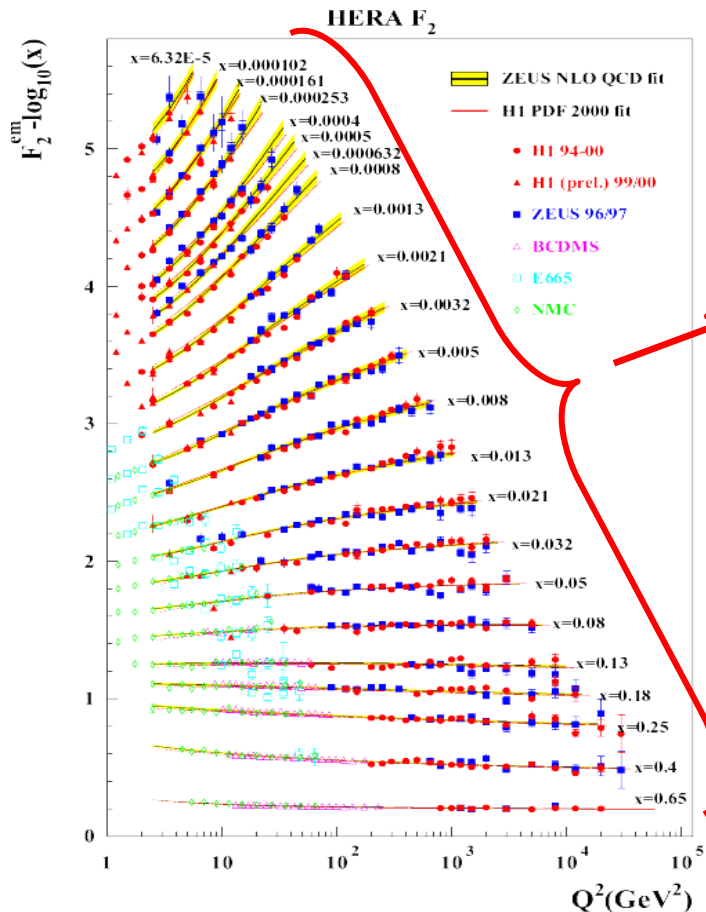
The nucleon collinear structure: momentum

	U	L	T
U	f_1		
L			
T			

$$f_1(x) = q^+(x) + q^-(x)$$

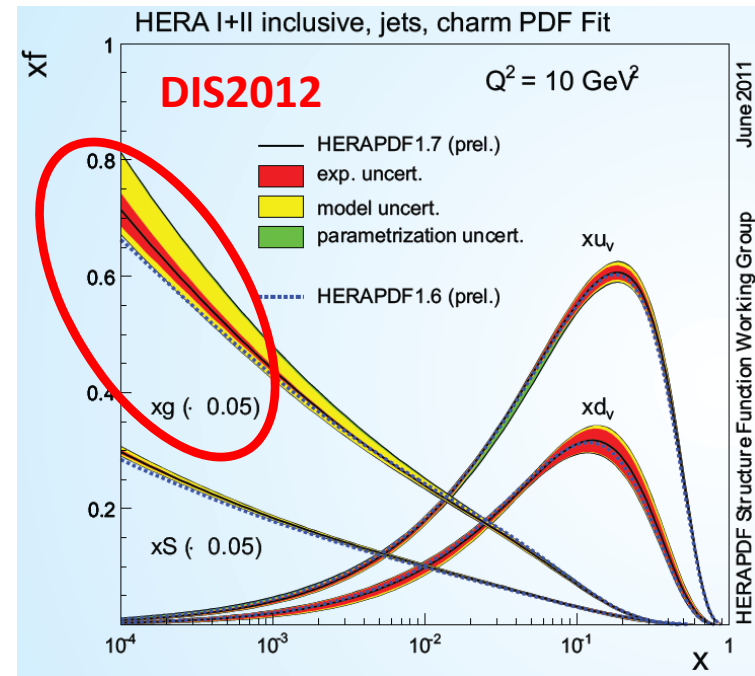


- 40 years of **inclusive DIS** experiments
- fixed target exp. & HERA collider
- **very precise data!**
- wide kinematic range in x and Q^2
- prediction of Q^2 dependence → triumph of QCD!



4 orders of magnitude in $x!$

5 orders of magnitude in $Q^2!$

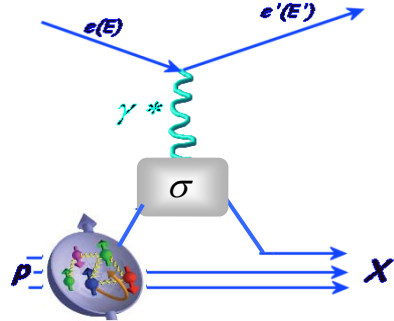


very good knowledge of longitudinal momentum structure of the nucleon!
 ...but gluons should saturate!

The nucleon collinear structure: helicity

	U	L	T
U			
L		g_{1L}	
T			

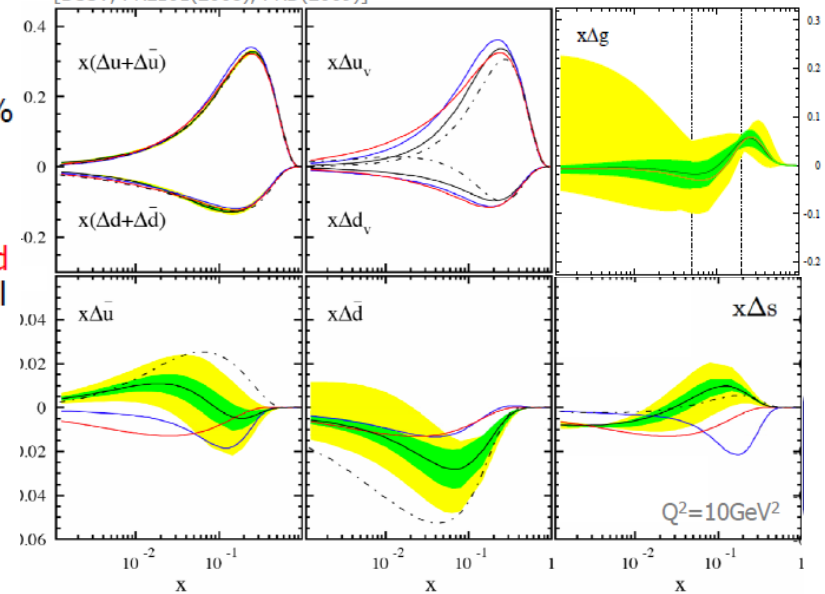
$g_1(x) = q^+(x) - q^-(x)$



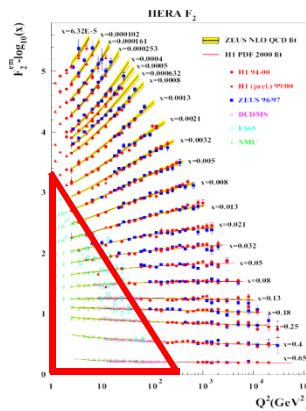
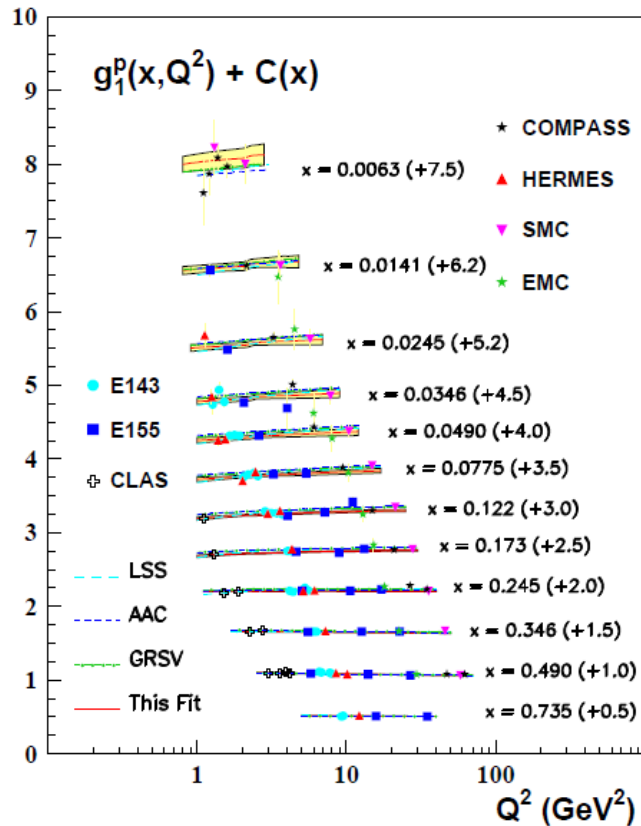
- **inclusive DIS** (+ semi-inclusive → flavor tagging)
- Fixed target experiments
- polarized beams and target → challenging
- relatively limited kinematic coverage in x and Q^2

global NLO fit of helicity distributions

[DSSV, PRL101(2008), PRD(2009)]



good knowledge of helicity distrib. of valence quarks
Sea quarks and gluon helicity still poorly constrained...



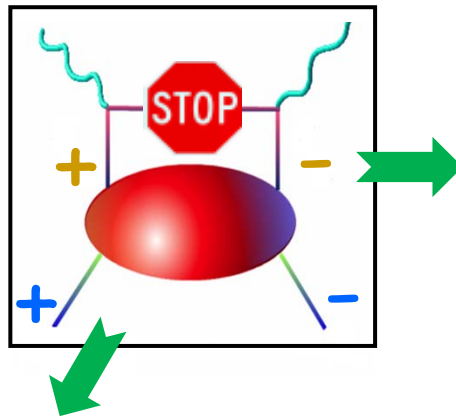
- DSSV
- DSSV+1
- DSSV 2%
- DNS
- GRSV std
- - - GRSV val

The nucleon collinear structure: transversity

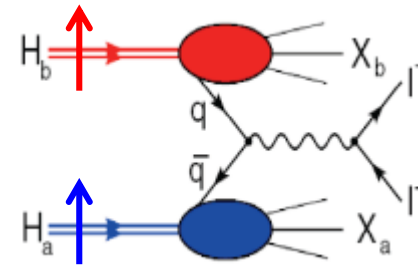
	U	L	T
U			
L			
T			h_1

$$h_1(x) = q^\uparrow(x) - q^\downarrow(x)$$

Chiral-odd!!!



Double-polarized DY



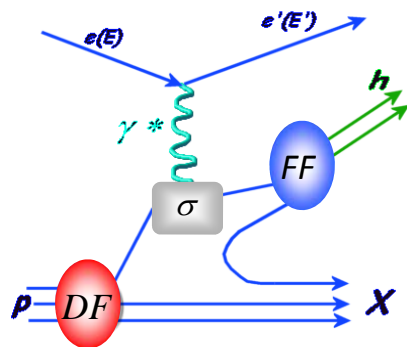
Golden process!

...but challenging:

- very small X-section
- huge background
- polarizing \bar{p} beams

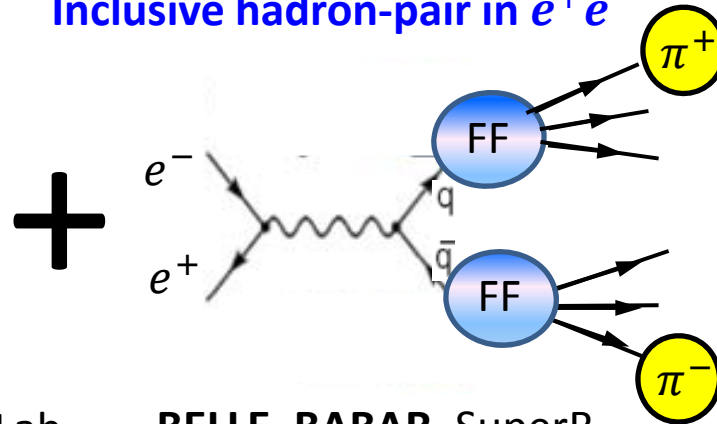
Future exp. RHIC, CERN, FAIR

Semi-Inclusive DIS



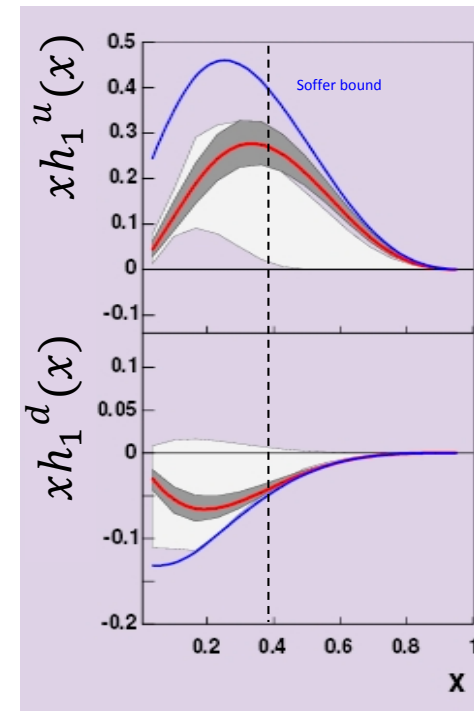
HERMES, COMPASS, JLab,...

Inclusive hadron-pair in e^+e^-

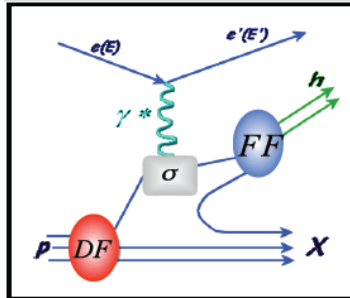


BELLE, BABAR, SuperB,...

- Very recent: first evidence (2005), first extraction (2007)!
- limited coverage in x (< 0.4)
- more data in valence region (JLab) to constrain tensor charge
- sea quark transversity completely unconstrained

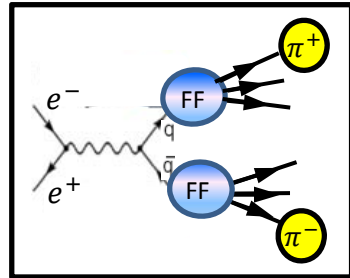


Reactions and experiments



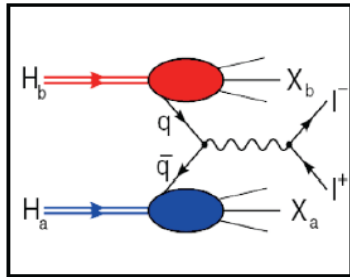
SIDIS: rich phenomenology, the most explored so far

SIDIS
$$\sigma^{ep \rightarrow ehX} = \sum_q \text{DF} \otimes \sigma^{eq \rightarrow eq} \otimes \text{FF}$$



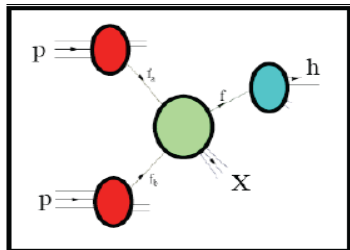
e⁺e⁻: B-factories as powerful fragmentation laboratories

e⁺e⁻
$$\sigma^{ee \rightarrow hhX} = \sum_q \sigma^{qq \rightarrow ee} \otimes \text{FF} \otimes \text{FF}$$



DY: challenging for experiments (only unpolarized so far)

DY
$$\sigma^{pp \rightarrow eeX} = \sum_q \text{DF} \otimes \text{DF} \otimes \sigma^{qq \rightarrow ee}$$



Hadron reactions: challenging for theory (ISI + FSI)

pp
$$\sigma^{pp \rightarrow hX} = \sum_q \text{DF} \otimes \text{DF} \otimes \sigma^{qq \rightarrow qq} \otimes \text{FF}$$



Boer-Mulders function: SIDIS

$$\frac{d^5\sigma}{dx dy dz d\phi_h dP_{h\perp}} \propto \left\{ F_{UU,T} + \varepsilon F_{UU,L} + 2\sqrt{\varepsilon(1+\varepsilon)} \cos\phi_h F_{UU}^{\cos\phi_h} + \varepsilon \cos 2\phi_h F_{UU}^{\cos 2\phi_h} \right\}$$

Twist-2:

$$d\sigma_{UU}^{\cos 2\phi} \propto \cos 2\phi \cdot \sum_q e_q^2 I \left[\frac{2(\hat{P}_{h\perp} \cdot \vec{k}_T)(\hat{P}_{h\perp} \cdot \vec{p}_T) - \vec{k}_T \cdot \vec{p}_T}{MM_h} h_1^\perp H_1^{\perp q} \right]$$

Boer & Mulders PRD 57 (1998)

Boer-Mulders effect

Pure kinematic effect due to transverse momentum of partons in the nucleon

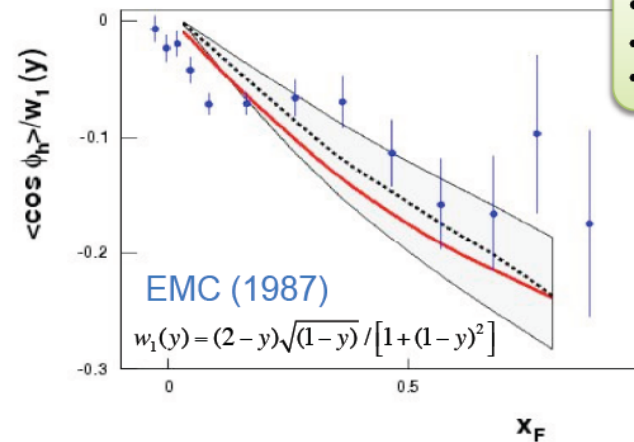
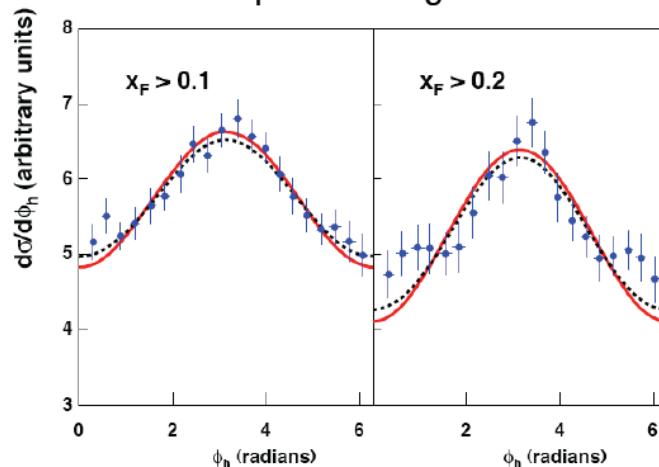
Cahn effect

Cahn PLB 78 (1978)

Twist-3:

$$d\sigma_{UU}^{\cos\phi} \propto \cos\phi \cdot \sum_q e_q^2 \frac{2M}{Q} I \left[-\frac{(\hat{P}_{h\perp} \cdot \vec{p}_T)}{M_h} x h_1^\perp H_1^{\perp q} - \frac{(\hat{P}_{h\perp} \cdot \vec{k}_T)}{M} x f_1 D_1 + \dots \right]$$

Till 2008: qualitative agreement with Cahn expectations

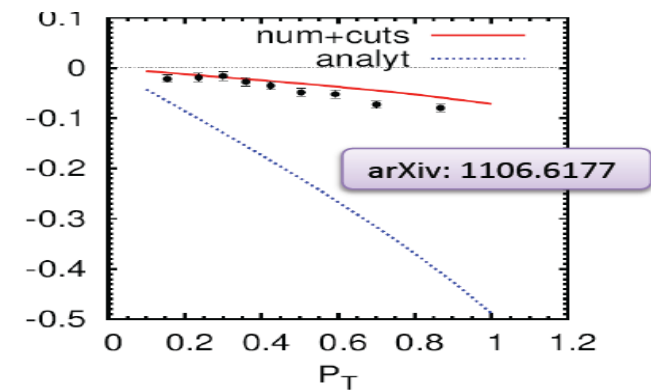
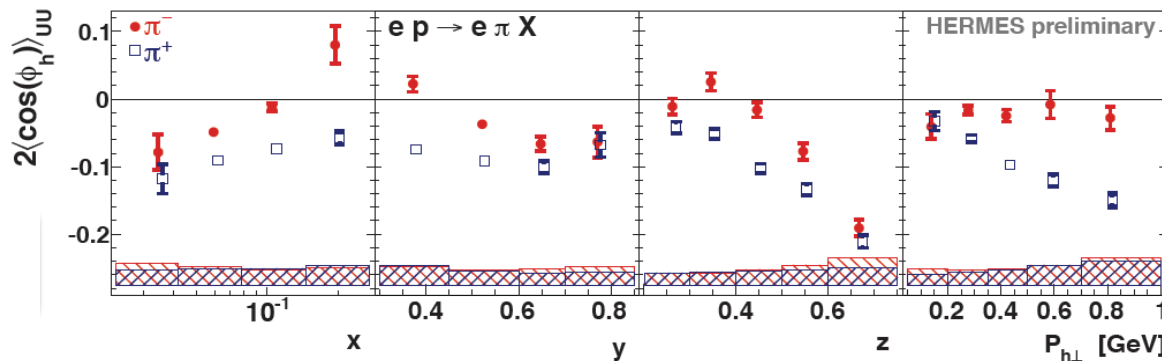
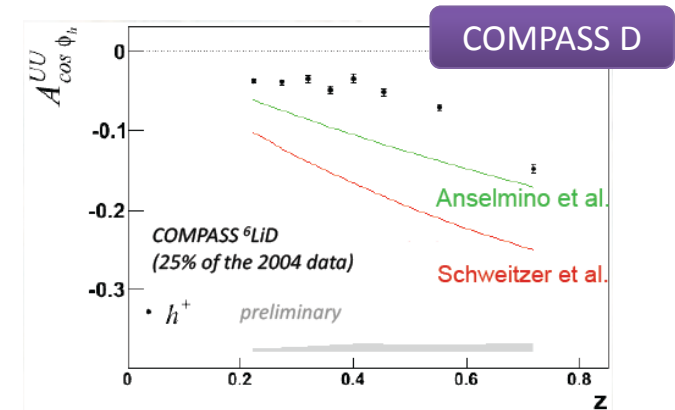
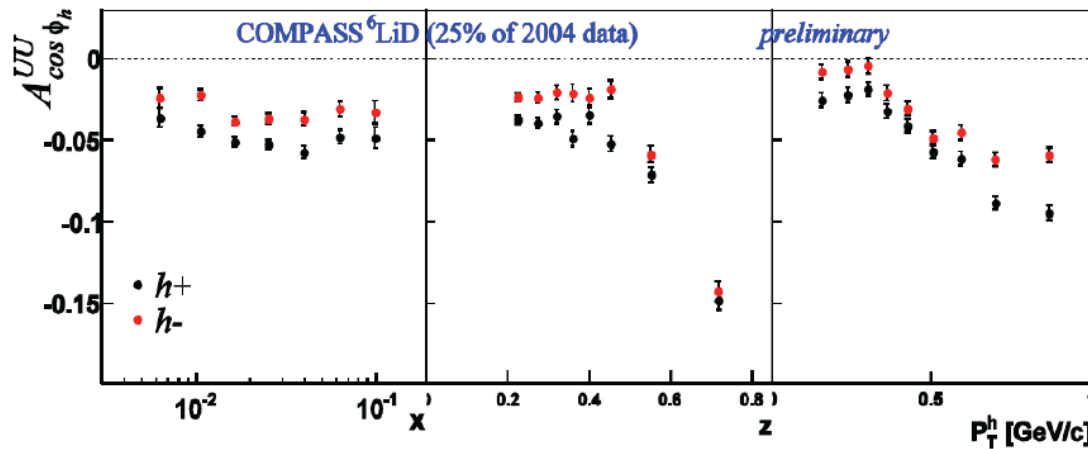


- No hadron identification
- No charge separation
- Poor statistics for cos2φ

Boer-Mulders function: SIDIS ($\cos\phi$)

$$\sigma_{UU}^{\cos(\phi)} \propto [f_1 \otimes D_1 + h_1^\perp \otimes H_1^\perp + \dots] / Q$$

- No dependence on hadron charge is expected
- Difference between h^+ / h^- due to Boer-Mulders term
- Predictions overestimate data
- BM or twist-3?



Unpolarized TMDs

$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2 y^2}{xyQ^2 2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right)$$

$$\left\{ \begin{aligned} & [F_{UU,T} + \epsilon F_{UU,L} \\ & + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)}] \end{aligned} \right.$$

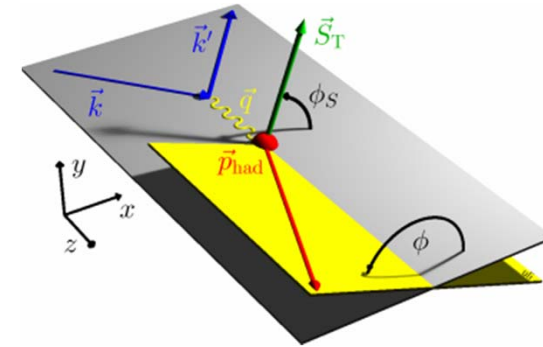
$$+ \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right]$$

$$+ S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right]$$

$$+ S_L \lambda_l \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right]$$

$$+ S_T \left[\begin{aligned} & \sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi - \phi_S)} \right) \\ & + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi - \phi_S)} \\ & + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \\ & + \sqrt{2\epsilon(1-\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi - \phi_S)} \end{aligned} \right]$$

$$+ S_T \lambda_l \left\{ \begin{aligned} & \sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi - \phi_S)} \end{aligned} \right\}$$



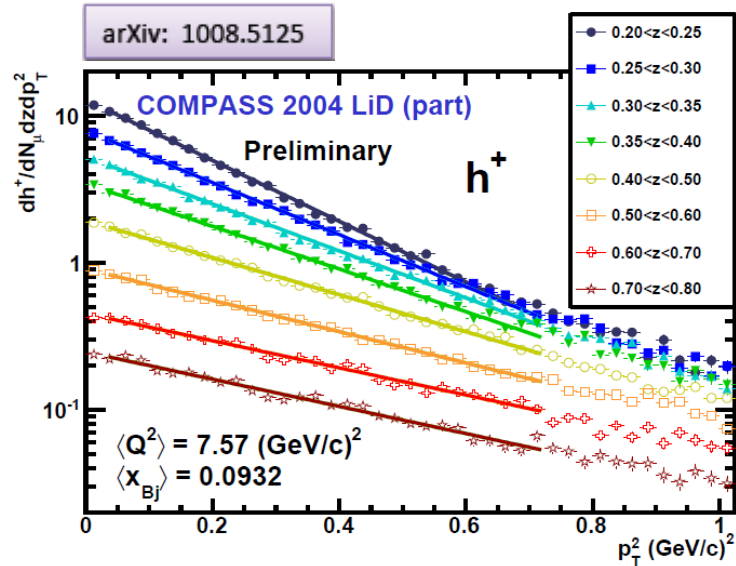
Distribution Functions

		quark		
		U	L	T
nucleon	U	f_1		h_1^\perp
	L		g_1	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}^\perp	h_{1T}^\perp

Fragmentation Functions

		quark		
		U	L	T
h	U	D_1		H_1^\perp

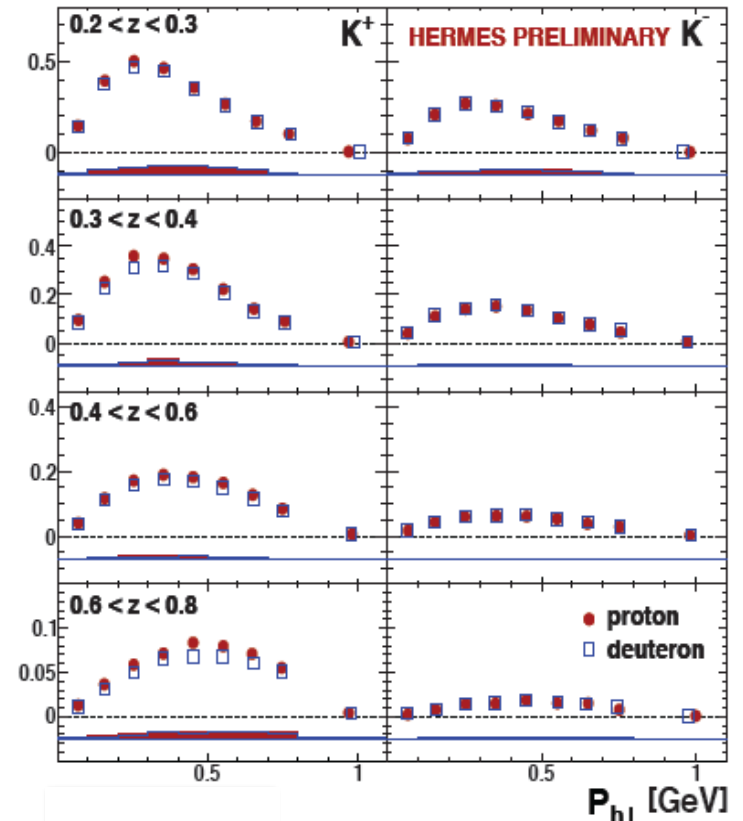
Unpolarized TMDs: $P_{h\perp}$ -unintegrated distributions



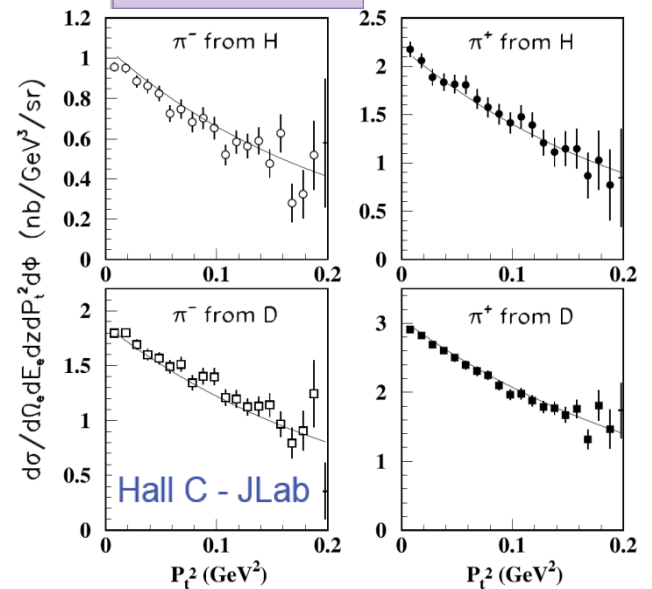
$$M_N^h = \frac{1}{N_N^{DIS}(Q^2)} \frac{dN_N^h(z, Q^2)}{dz} = \frac{\sum_q e_q^2 \int dx f_{1q}(x, Q^2) D_{1q}^h(z, Q^2)}{\sum_q e_q^2 \int dx f_{1q}(x, Q^2)}$$

disentangling the z and $P_{h\perp}$ dependencies allows to access intrinsic transverse momenta p_T and k_T :
 $\langle P_{h\perp}^2 \rangle = z^2 \langle k_T^2 \rangle + \langle p_T^2 \rangle \rightarrow 3D \text{ analysis } (x, z, P_{h\perp})$

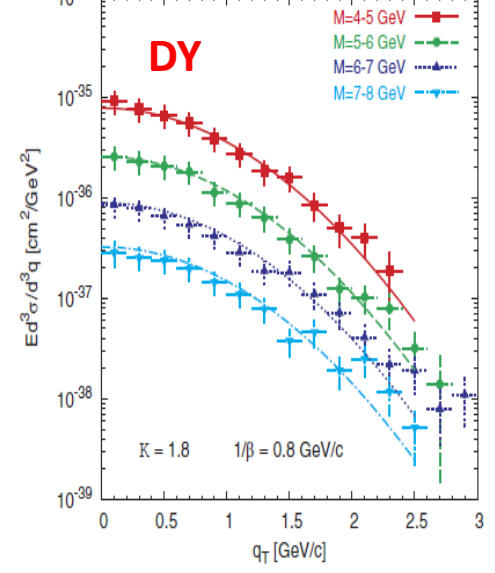
Hadron multiplicities



arXiv: 0709.3020

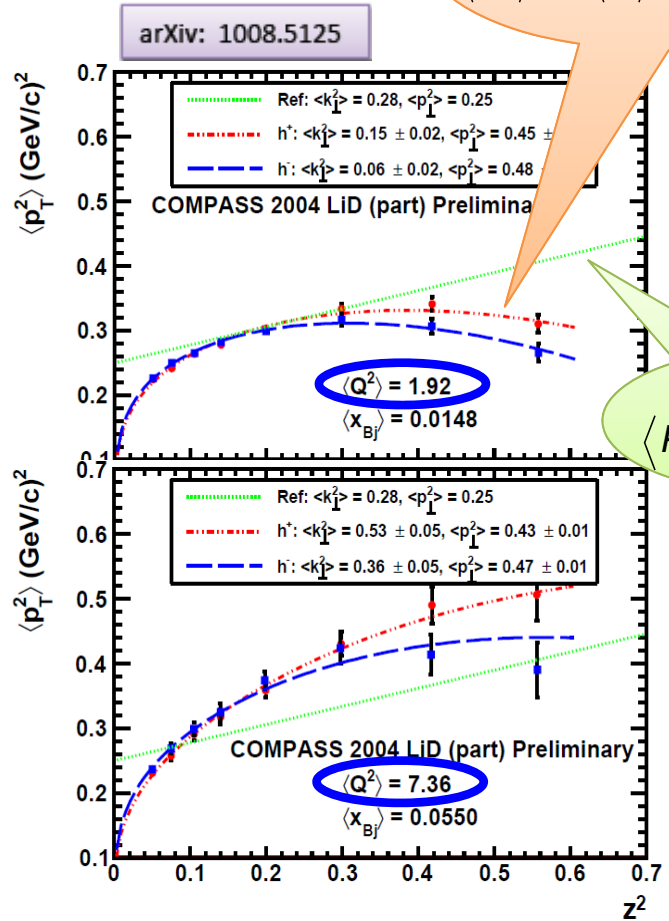


D'Alesio, Murgia, PRD70 (04)
 data E288, E605, R



Unpolarized TMDs: energy dependence & evolution

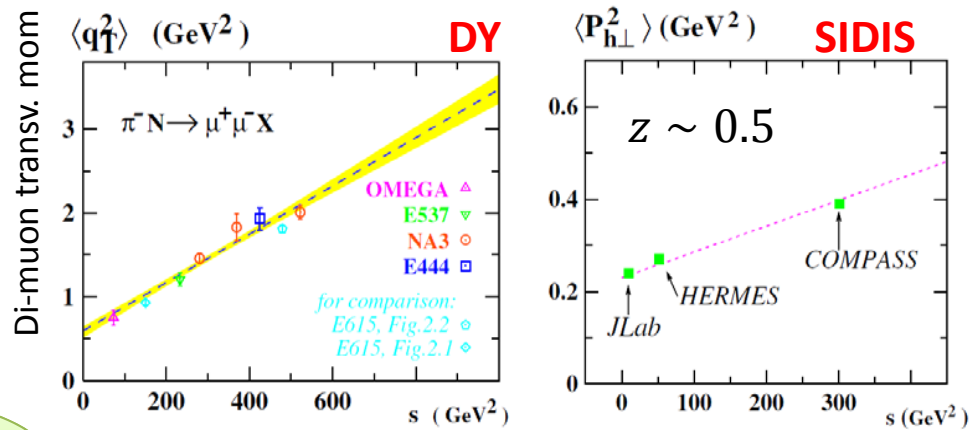
Is p_T independent of z ?
 $\langle P_{h\perp}^2 \rangle = z^2 \langle k_T^2 \rangle + z^\alpha (1-z)^\beta \langle p_T^2 \rangle$



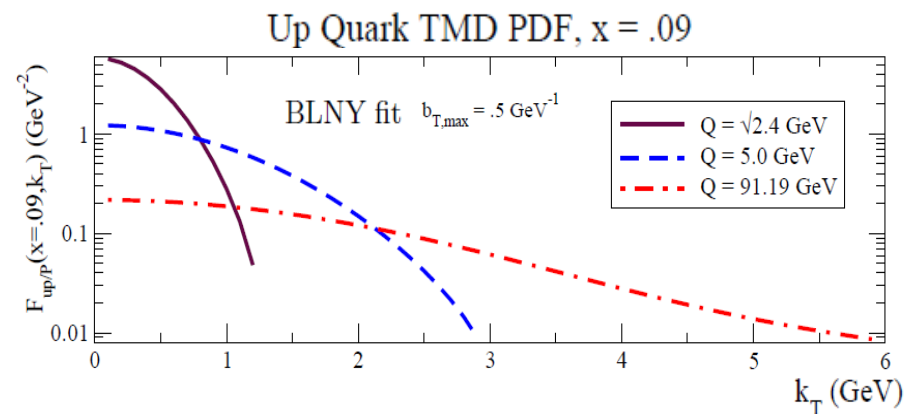
Gaussian ansatz
 $\langle P_{h\perp}^2 \rangle = z^2 \langle k_T^2 \rangle + \langle p_T^2 \rangle$

- Hint of z -dependence
- Hint of flavour dependence
- Q^2 -dependence

P. Schweitzer, T. Teckentrup, A. Metz, PRD81(10)

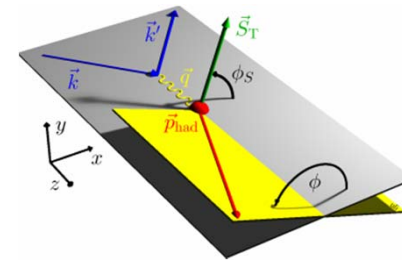


Transverse-momentum broadening in SIDIS & DY
 Available data suggest that intrinsic p_T in $f_1(x, p_T)$ are compatible in SIDIS and DY \rightarrow **universality!**



S.M. Aybat & T.C. Rogers arXiv: 1101.5057v2
 "QCD evolution has a strong quantitative effect on TMDs"

Helicity TMDs



$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2 y^2}{xyQ^2 2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x}\right)$$

$$\left\{ \begin{aligned} & \left[F_{UU,T} + \epsilon F_{UU,L} \right. \\ & \left. + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} \right] \end{aligned} \right.$$

$$+ \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right]$$

$$+ S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right]$$

$$+ S_L \lambda_l \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right]$$

$$+ S_T \left[\begin{aligned} & \sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi - \phi_S)} \right) \\ & + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi - \phi_S)} \\ & + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \\ & + \sqrt{2\epsilon(1-\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi - \phi_S)} \end{aligned} \right]$$

$$+ S_T \lambda_l \left\{ \begin{aligned} & \left[\sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} \right. \\ & + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \\ & \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi - \phi_S)} \right] \end{aligned} \right\}$$

Describes probability to find longitudinally polarized quarks in a longitudinally polarized nucleon

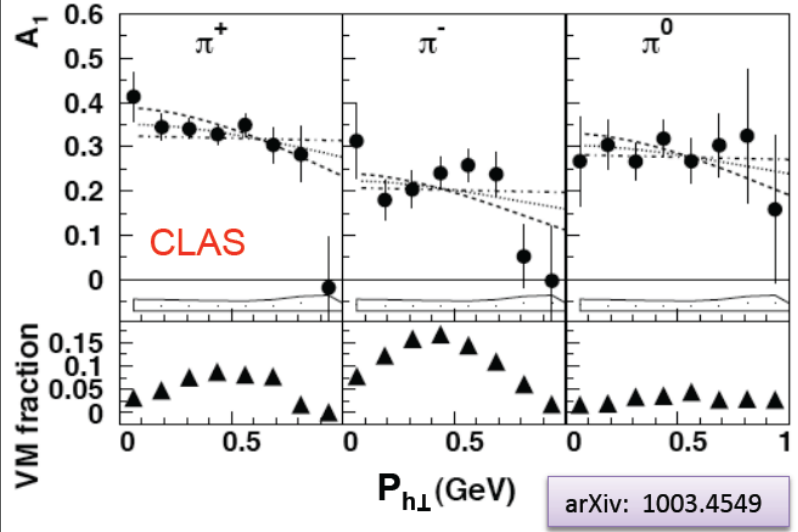
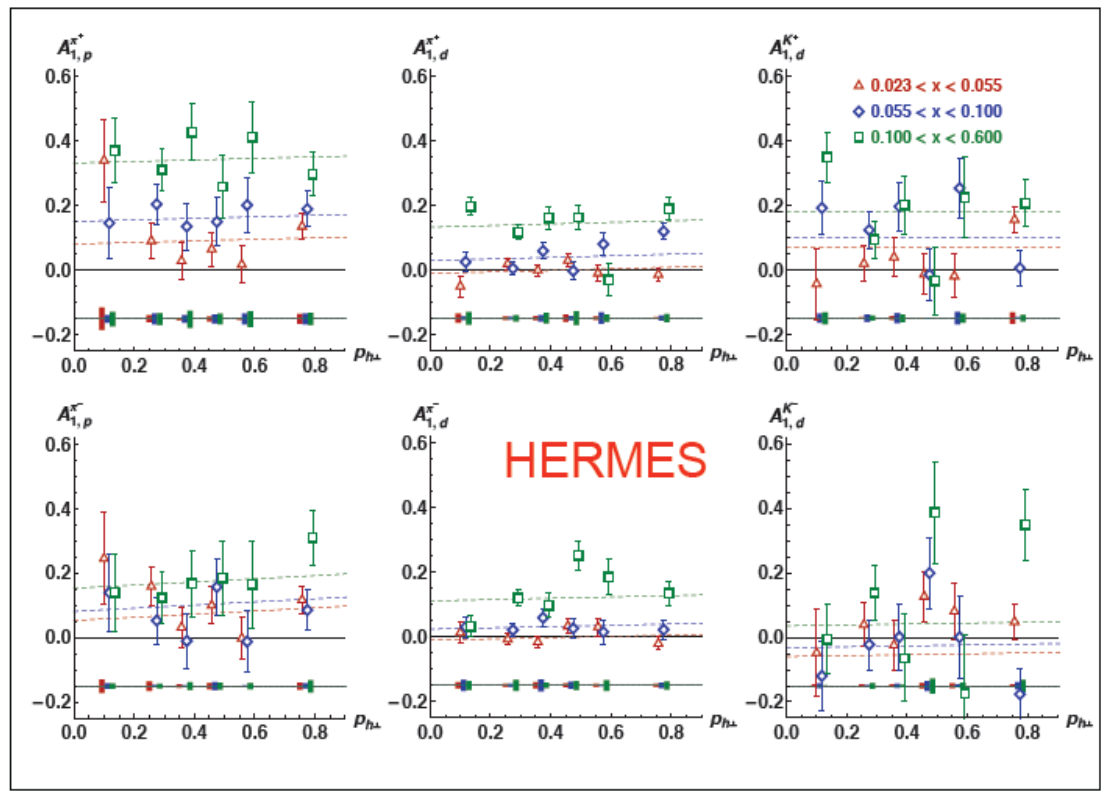
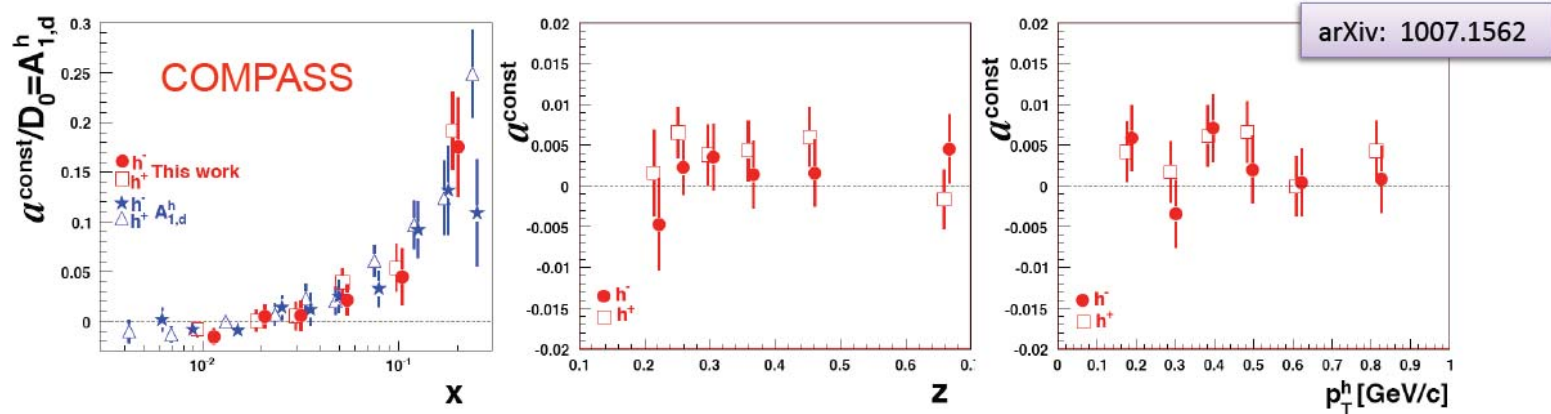
Distribution Functions

		quark		
		U	L	T
nucleon	U	f_1		h_1^\perp
	L		g_1	h_{1T}^\perp
	T	f_{1T}^\perp	g_{1T}^\perp	h_1

Fragmentation Functions

		quark		
		U	L	T
h	U	D_1		H_1^\perp

Helicity TMDs: $P_{h\perp}$ -unintegrated A_{LL} DSAs



Only a weak dependence on $P_{h\perp}$ is observed

Higher-twist

$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2 y^2}{xyQ^2 2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x} \right)$$

$$\left\{ \begin{array}{l} F_{UU,T} + \epsilon F_{UU,L} \\ + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} \end{array} \right]$$

$$+ \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right]$$

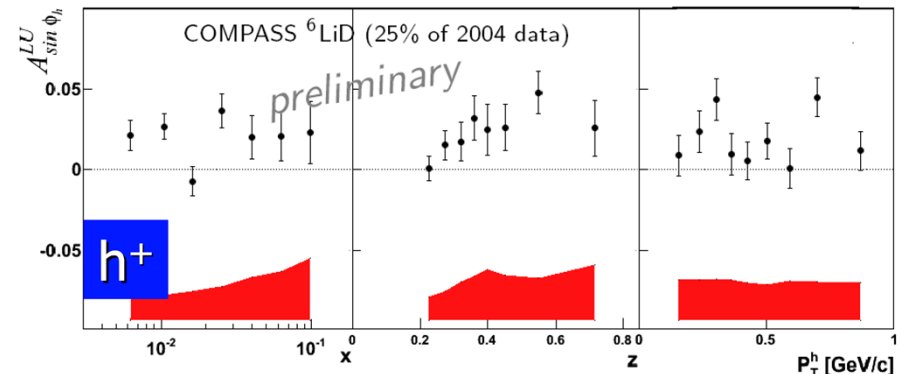
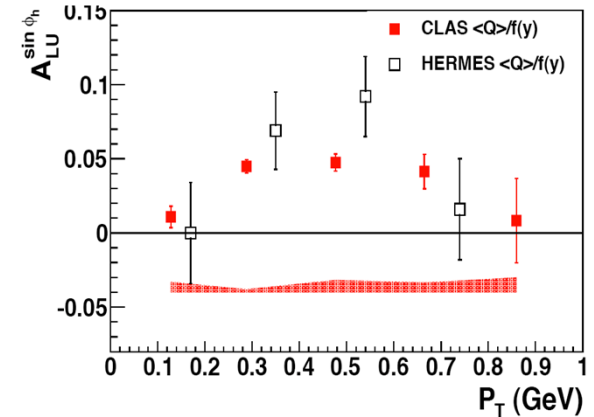
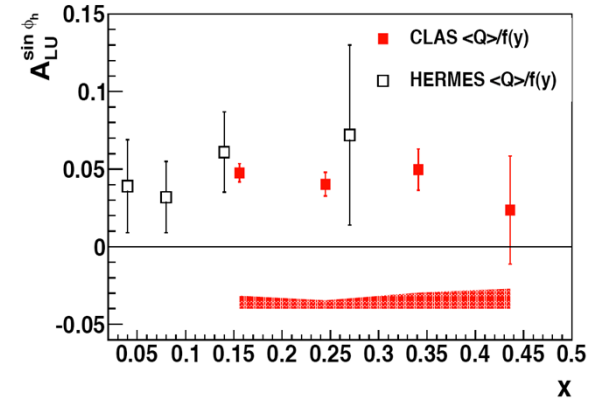
$$+ S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right]$$

$$+ S_L \lambda_l \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right]$$

$$+ S_T \left[\sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi - \phi_S)} \right) \right. \\ \left. + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi + \phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi - \phi_S)} \right. \\ \left. + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \right. \\ \left. + \sqrt{2\epsilon(1-\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi - \phi_S)} \right]$$

$$+ S_T \lambda_l \left[\sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi - \phi_S)} \right. \\ \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \right. \\ \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi - \phi_S)} \right]$$

$$\sigma_{LU}^{\sin(\phi)} \propto [e \otimes H_1^\perp + g^\perp \otimes D_1 + \dots] / Q$$

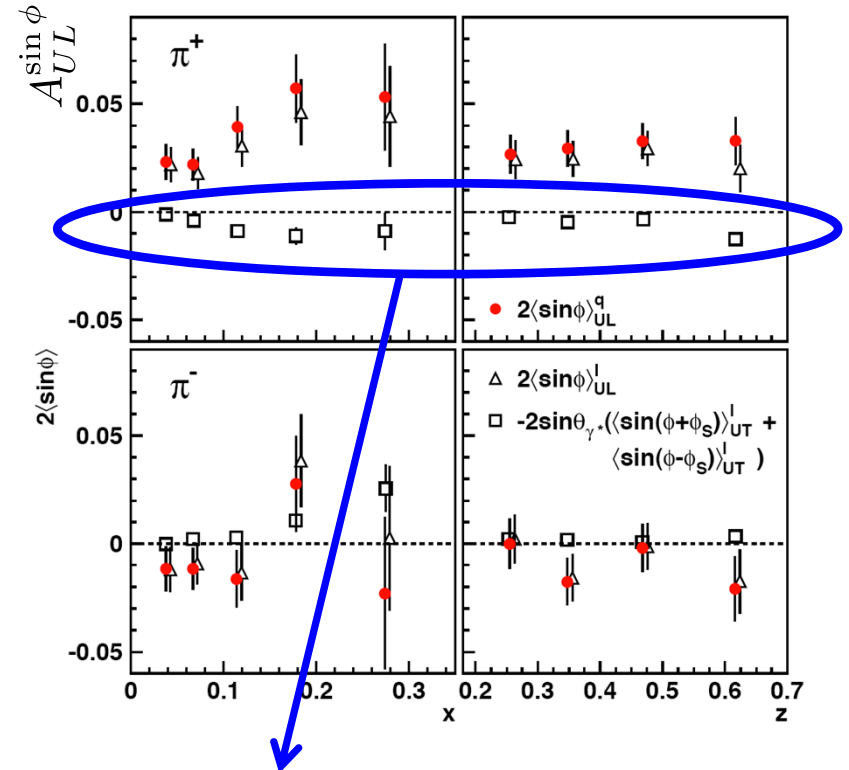


Higher-twist

$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2 y^2}{xyQ^2 2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x} \right)$$

$$\left\{ \begin{array}{l} \left[F_{UU,T} + \epsilon F_{UU,L} \right. \\ \left. + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)} \right] \\ + \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)} \sin(\phi) F_{LU}^{\sin(\phi)} \right] \\ + S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon \sin(2\phi) F_{UL}^{\sin(2\phi)} \right] \\ + S_L \lambda_l \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi) F_{LL}^{\cos(\phi)} \right] \\ + S_T \left[\sin(\phi - \phi_S) \left(F_{UT,T}^{\sin(\phi-\phi_S)} + \epsilon F_{UT,L}^{\sin(\phi-\phi_S)} \right) \right. \\ \left. + \epsilon \sin(\phi + \phi_S) F_{UT}^{\sin(\phi+\phi_S)} + \epsilon \sin(3\phi - \phi_S) F_{UT}^{\sin(3\phi-\phi_S)} \right. \\ \left. + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} \right. \\ \left. + \sqrt{2\epsilon(1+\epsilon)} \sin(2\phi - \phi_S) F_{UT}^{\sin(2\phi-\phi_S)} \right] \\ + S_T \lambda_l \left[\sqrt{1-\epsilon^2} \cos(\phi - \phi_S) F_{LT}^{\cos(\phi-\phi_S)} \right. \\ \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \right. \\ \left. + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi - \phi_S) F_{LT}^{\cos(2\phi-\phi_S)} \right] \end{array} \right\}$$

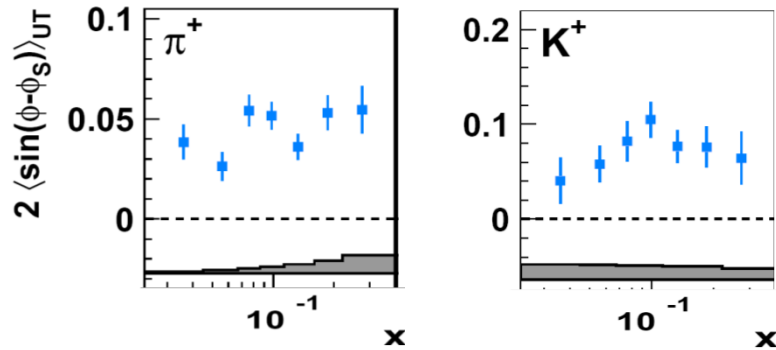
$$\sigma_{UL}^{\sin(\phi)} \propto [h_L \otimes H_1^\perp + f_L^\perp \otimes D_1 + \dots] / Q$$



Collins and Sivers contributions are small
 → Measured asymmetry is a genuine higher-twist effect

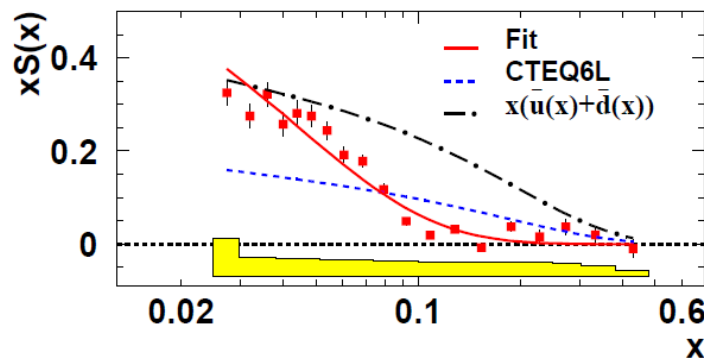
Sivers kaons amplitudes: open questions

π^+/K^+ production dominated by u-quarks, but:

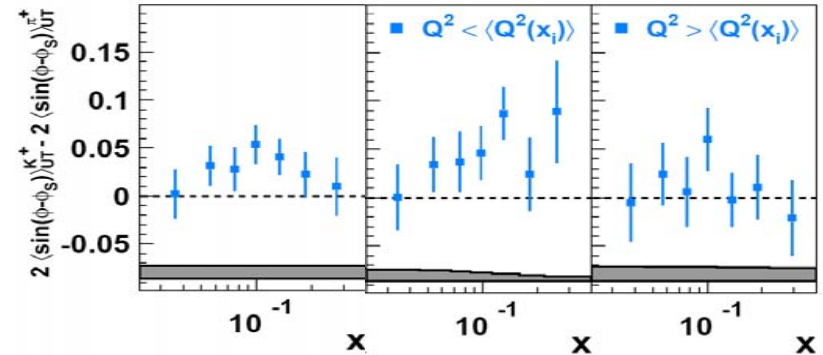


$$\pi^+ \equiv |u\bar{d}\rangle, K^+ \equiv |u\bar{s}\rangle \rightarrow$$

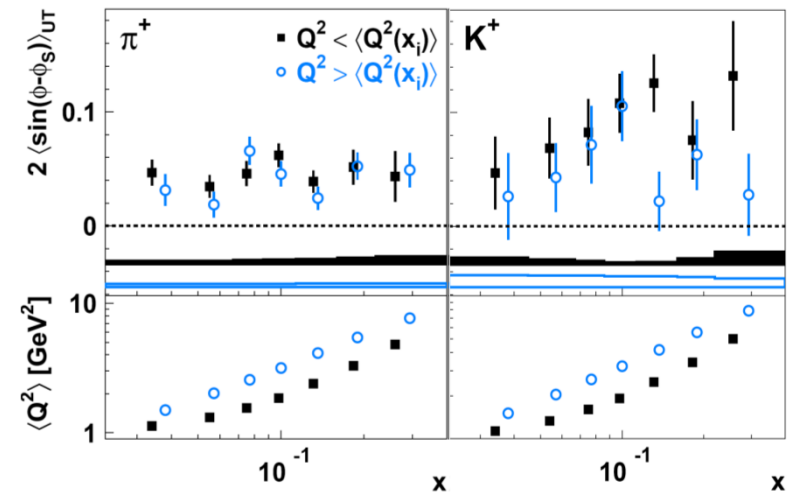
different role of various sea quarks ?



impact of different k_T dependence of FFs in the convolution integral

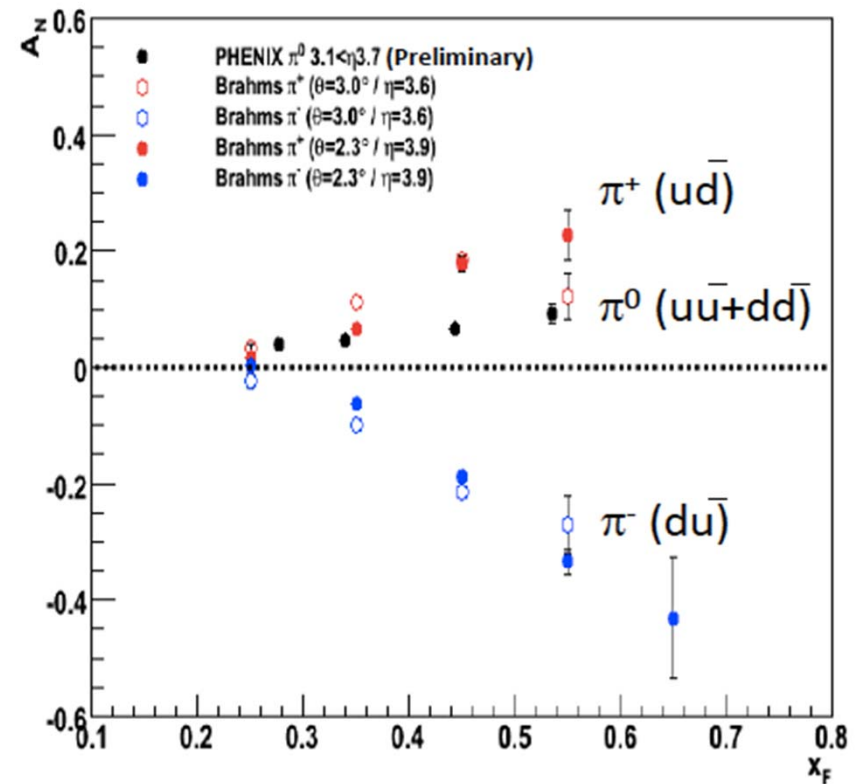
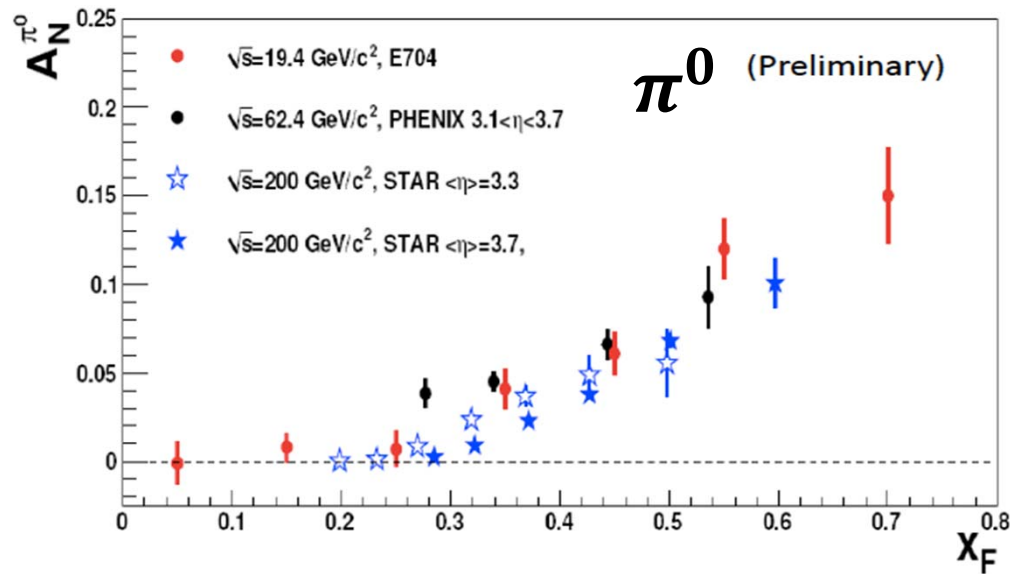


only in low- Q^2 region significant
(90% C.L.) deviation is observed



Each x -bin divided into two Q^2 bins
Higher-twist contrib. for kaons?

Transversity & Sivers in pp scattering: RHIC



- Large asymmetries measured by STAR, PHENIX and BRAHMS
- No strong dependence on \sqrt{s} from 19.4 to 200 GeV
- Spread of data probably due to different acceptance in pseudorapidity and/or p_T
- Could be due to admixture of Transversity, Sivers and Twist-3 effects