

Transverse target Single-Spin Asymmetry (SSA) in inclusive electroproduction of charged pions and kaons [1].

Deutsches Elektronen-Synchrotron (DESY)

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[1] HERMES collaboration, Phys. Lett. B 728 (2014) 183



Outline

>> Introduction

- ★ Inclusive electro-production of pions and kaons.
- ★ Transverse Single-Spin Asymmetry.

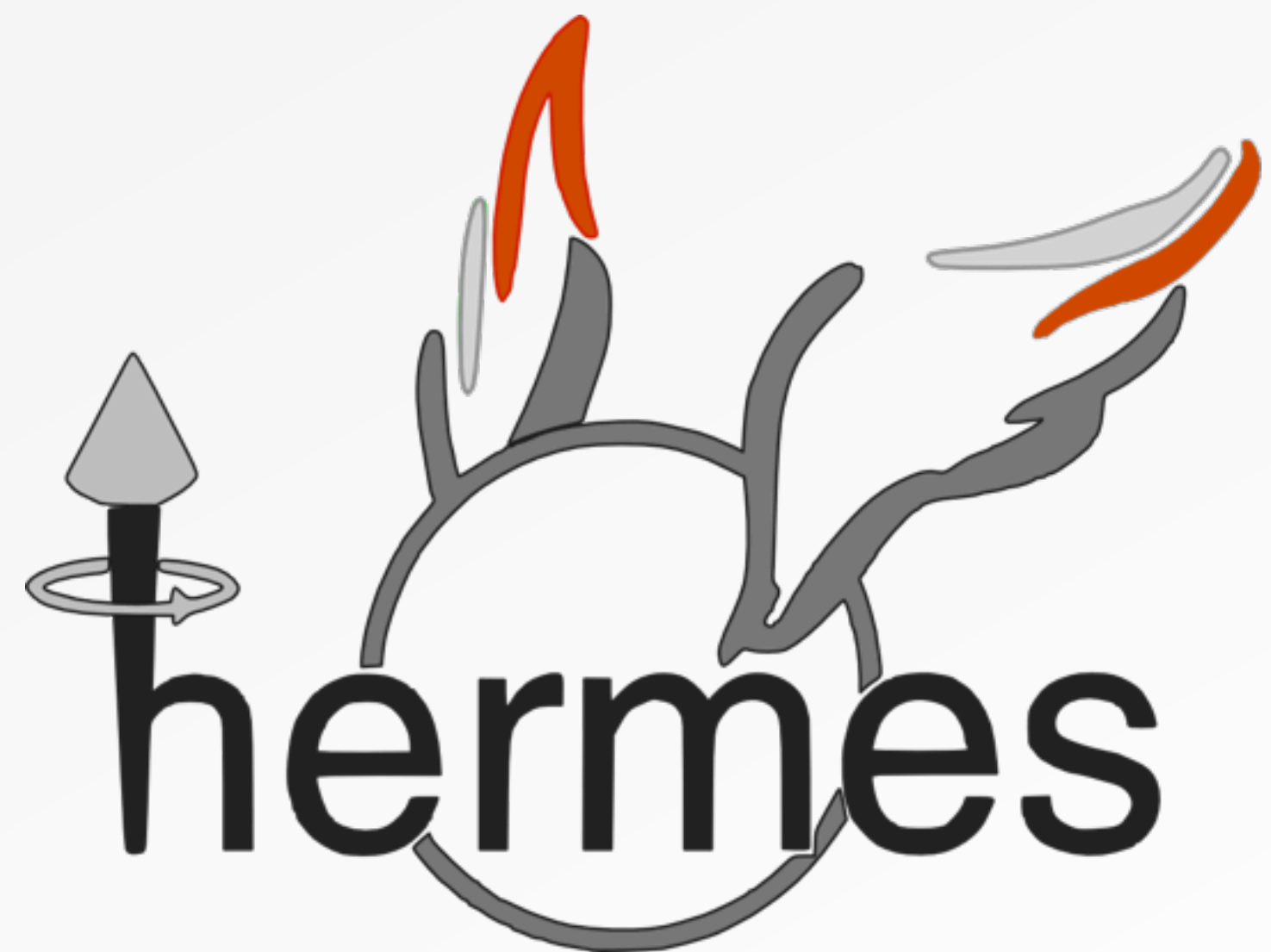
>> The measurement

- ★ The HERMES experiment.
- ★ Event selection.
- ★ Yields and kinematics.
- ★ Extraction of the amplitudes.

>> The Asymmetries

- ★ 1-dimensional projections.
- ★ 2-dimensional binning.
- ★ Systematics.
- ★ Categorized Asymmetries

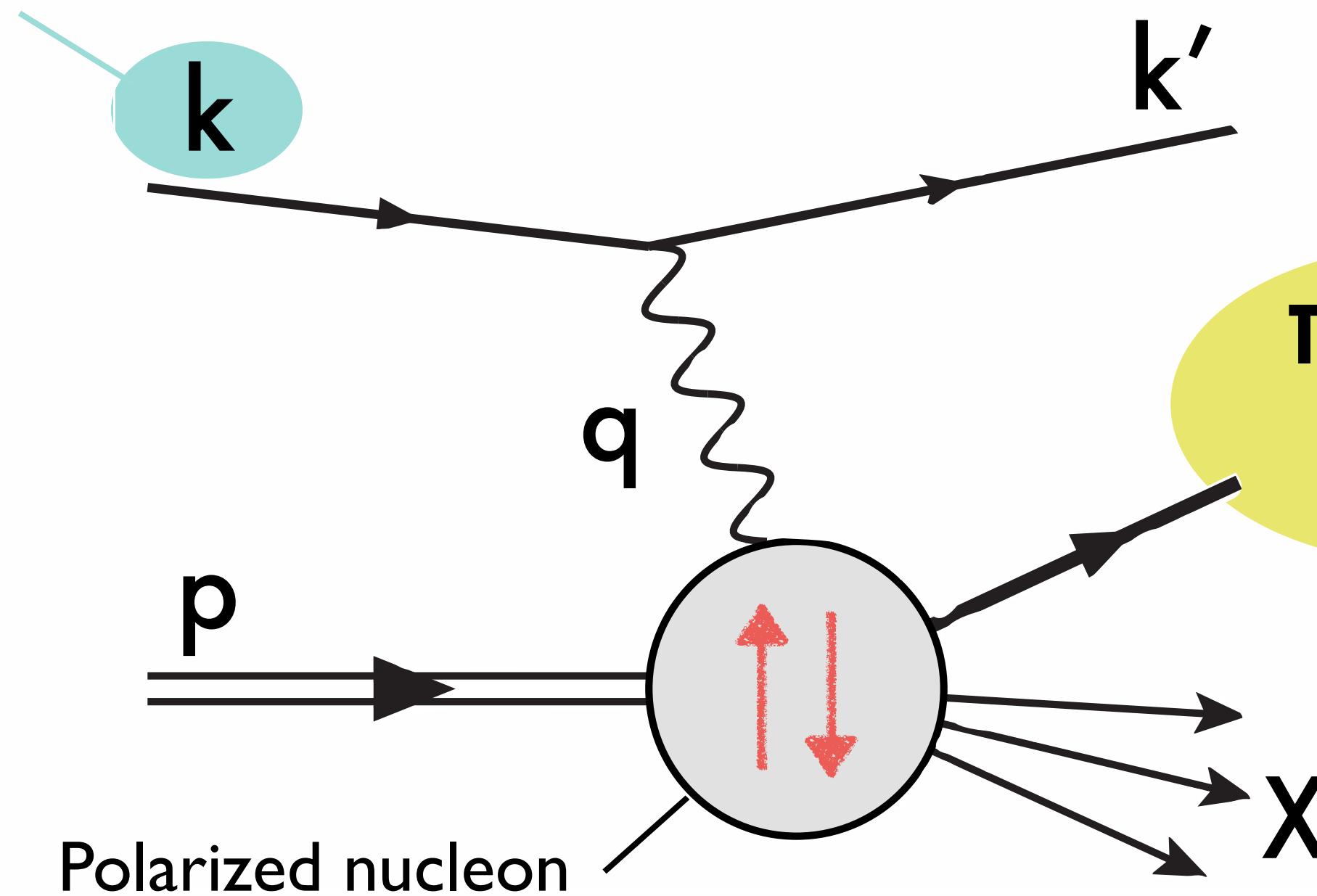
>> Summary and Outlook



Inclusive Electro-production

of charged π 's and K 's off a polarized nucleon

Unpolarized electron beam



Only one of these mesons is measured in the final state

Spin-dependent cross section

$$d\sigma_{UT} \propto \vec{S}_T \cdot (\hat{P}_h \times \hat{k})$$

$$lN^\uparrow \rightarrow hX$$

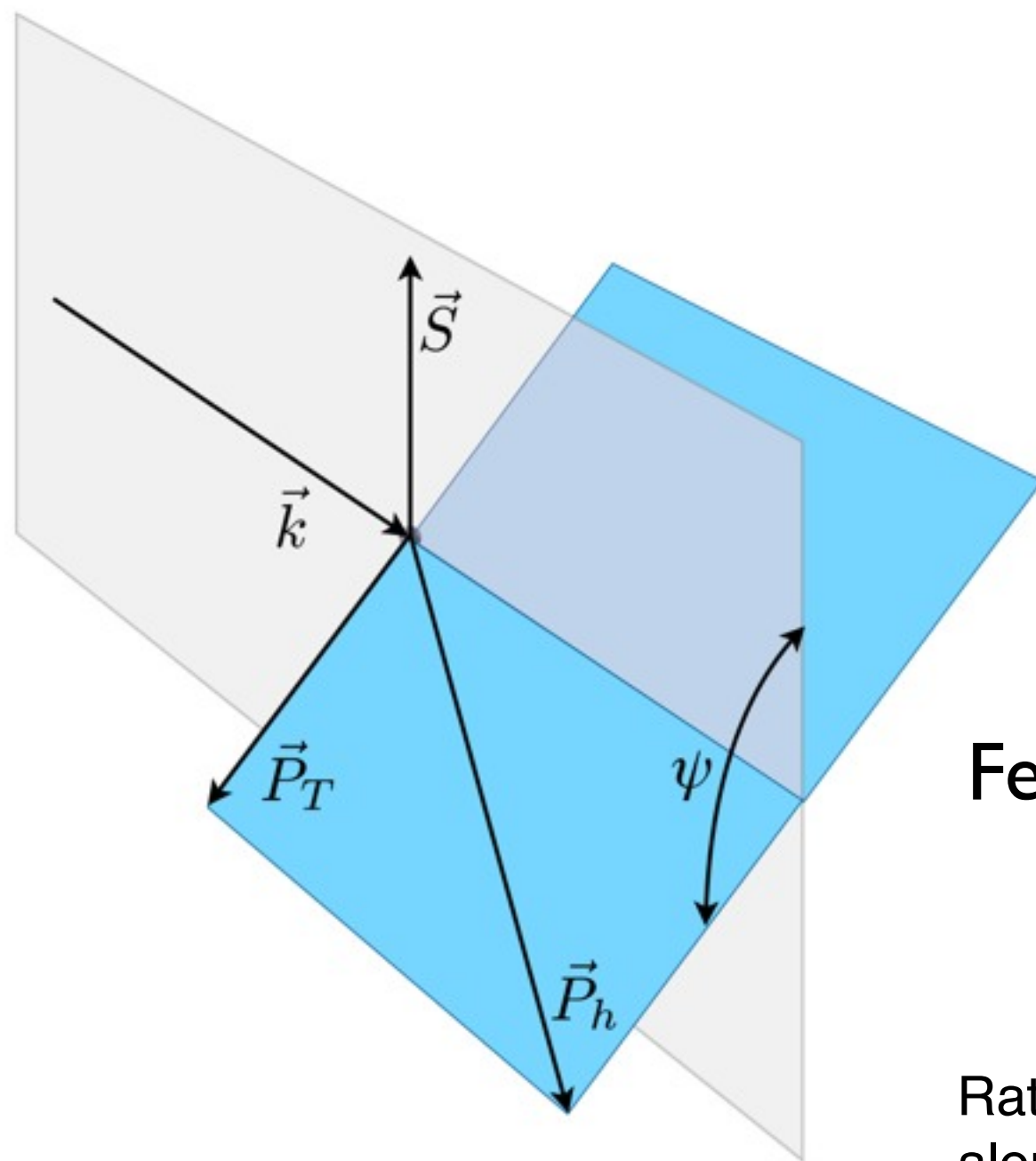
Transverse Single-Spin Asymmetry (SSA)

in respect to the polarization axis of the nucleon

Azimuthal asymmetry

$$A(x_F, P_T, \psi) = \frac{d\sigma_{UT}(x_F, P_T, \psi)}{d\sigma_{UU}(x_F, P_T)} = A_{UT}^{\sin\psi}(x_F, P_T) S_T \sin\psi$$

Asymmetry amplitude



Independent variables

$$\{x_F, P_T, \psi\}$$

Feynman x

$$x_F = 2P_z^{CM} / \sqrt{s}$$

Ratio of the longitudinal hadron momentum P_L along the beam direction to its maximum possible value

Left-Right Asymmetry

$$A_N(x_F, P_T) = \frac{d\sigma_L - d\sigma_R}{d\sigma_L + d\sigma_R} = -\frac{2}{\pi} A_{UT}^{\sin\psi}(x_F, P_T)$$

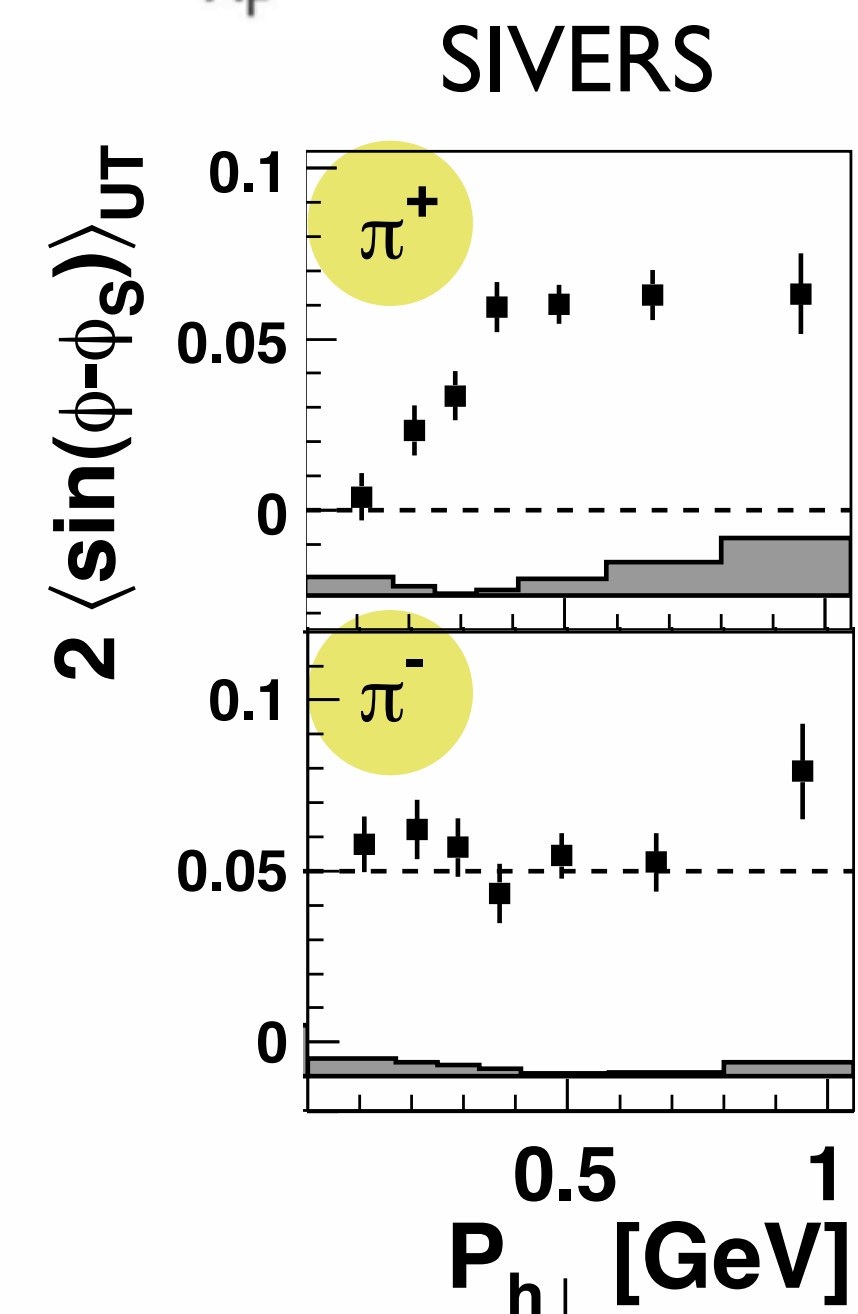
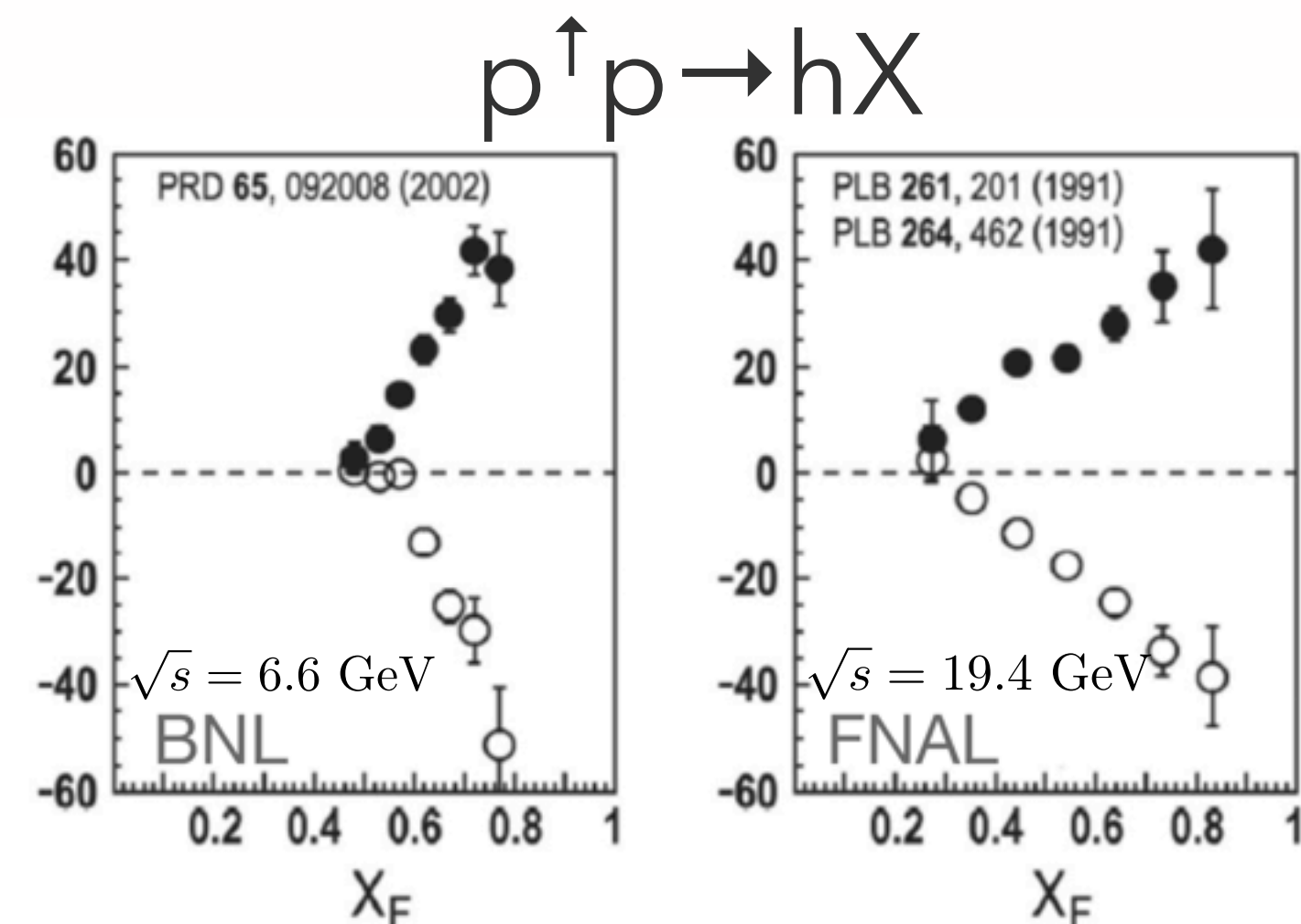
Motivation

>> Experimental:

- ★ Large values of A_N were observed in $p^\uparrow p \rightarrow hX$ reactions, but never measured in $IN^\uparrow \rightarrow hX$.
- ★ Measurement of **Sivers**^[2] and **Collins**^[3] effects on SIDIS: $IN^\uparrow \rightarrow l'hX$.
- ★ Lots of $IN^\uparrow \rightarrow hX$ data in HERMES!
- ★ Involves only one hadron.

>> Theoretical puzzle. Two existing approaches:

- ★ TMD's and fragmentation functions^[4] ($Q^2 \gg \Lambda_{\text{QCD}}^2$). Sivers and Collins effects.
- ★ Twist-3 parton correlation functions^[5] ($P_T \gg \Lambda_{\text{QCD}}$).



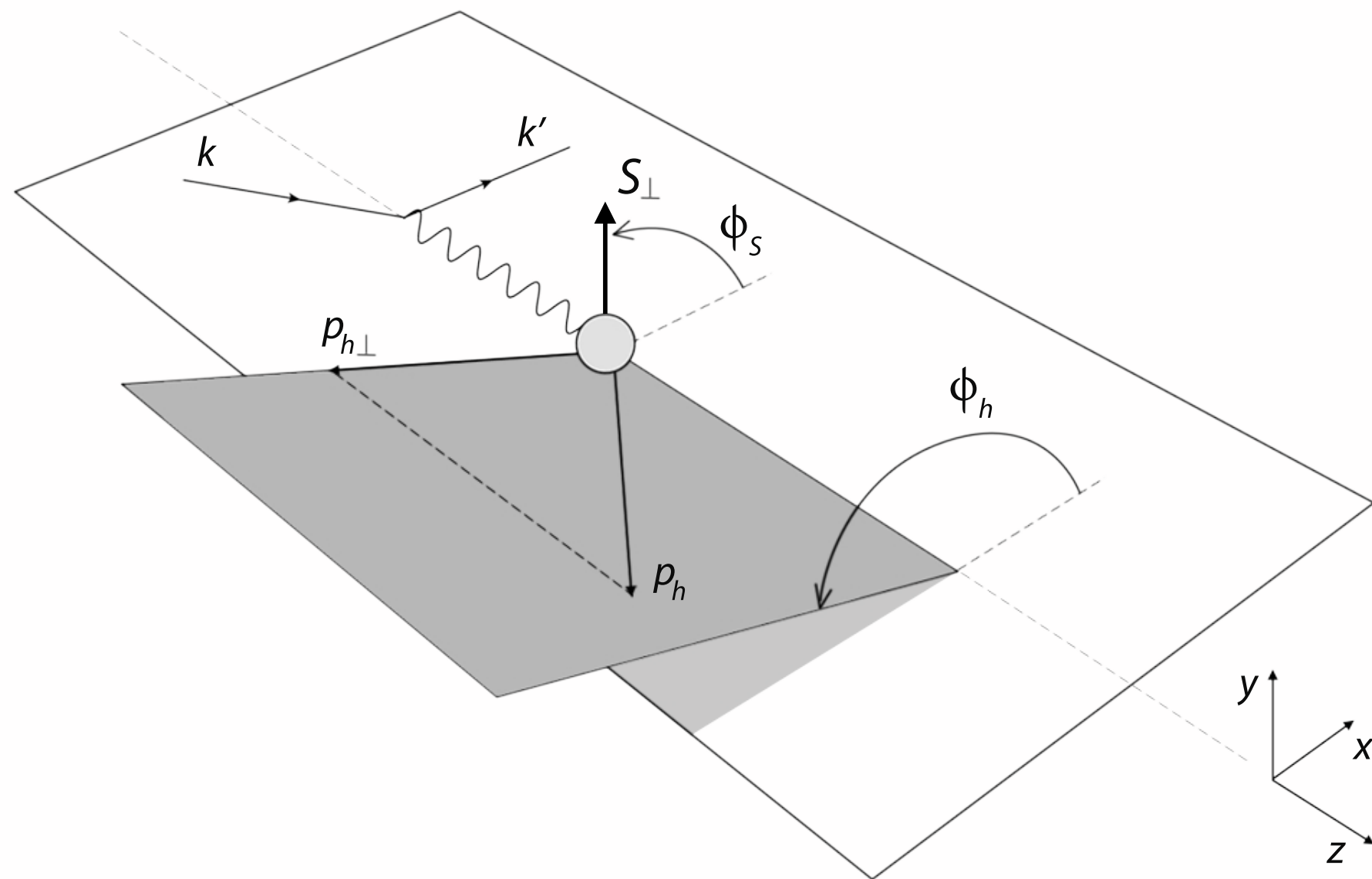
[2] *Phys. Rev. Lett.* **103**, 152002 (2009).

[3] *Phys. Lett. B* **693** (2010) 11-16.

[4] *Phys. Lett. B* **362** (1995) 164–172.

[5] *Phys. Rev. D* **59** (1999) 014004.

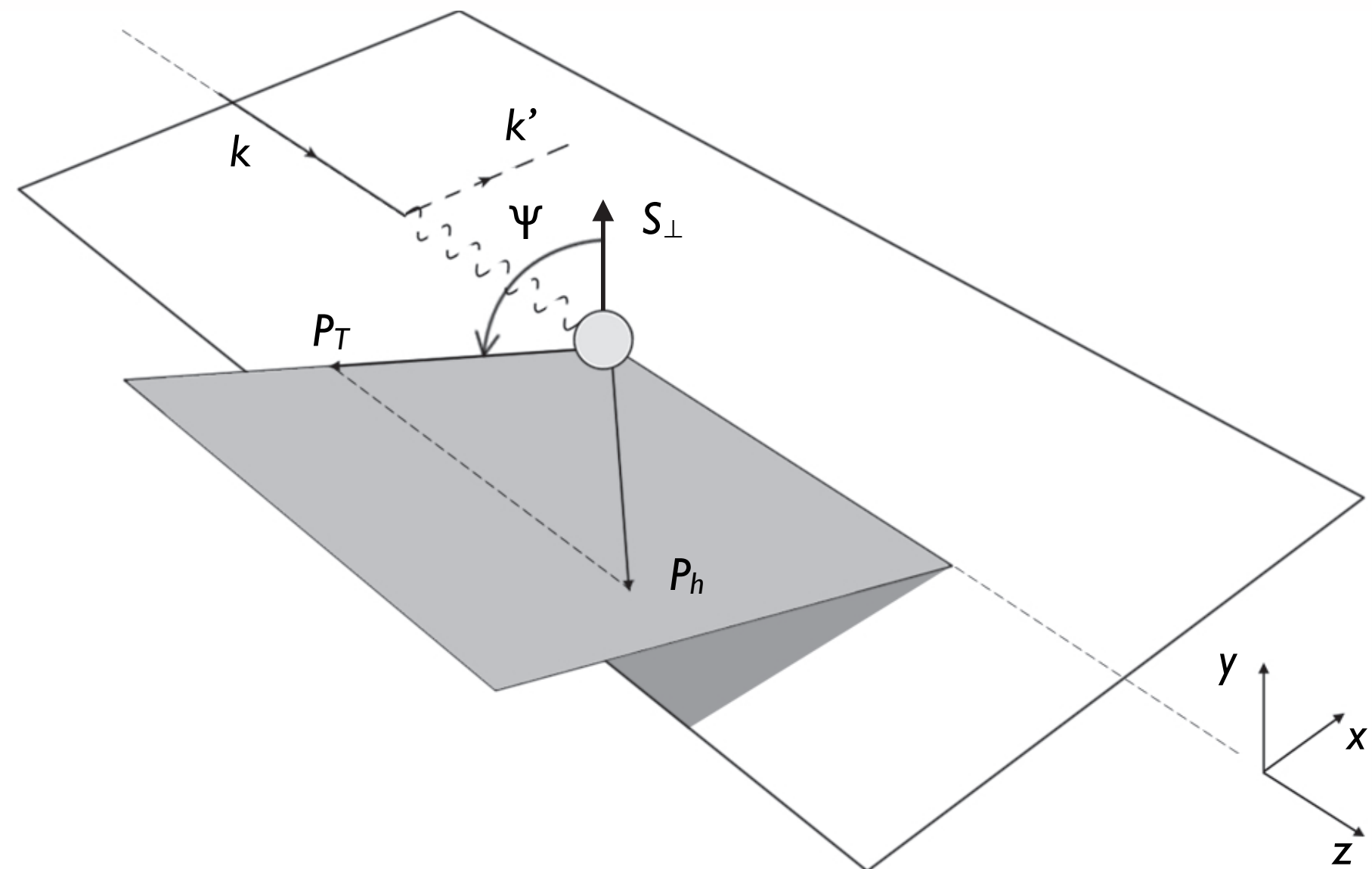
Similarity between Sivers and inclusive SSA



semi-inclusive DIS

$$d\sigma_{UT}^{ep \rightarrow ehX} \sim |\mathbf{S}_\perp| \left[\sin(\phi_h - \phi_s) f_{1T}^\perp \otimes D_1 + \sin(\phi_h + \phi_s) h_1 \otimes H_1^\perp + \dots \right]$$

↗ **SIVERS**
↘ **COLLINS**



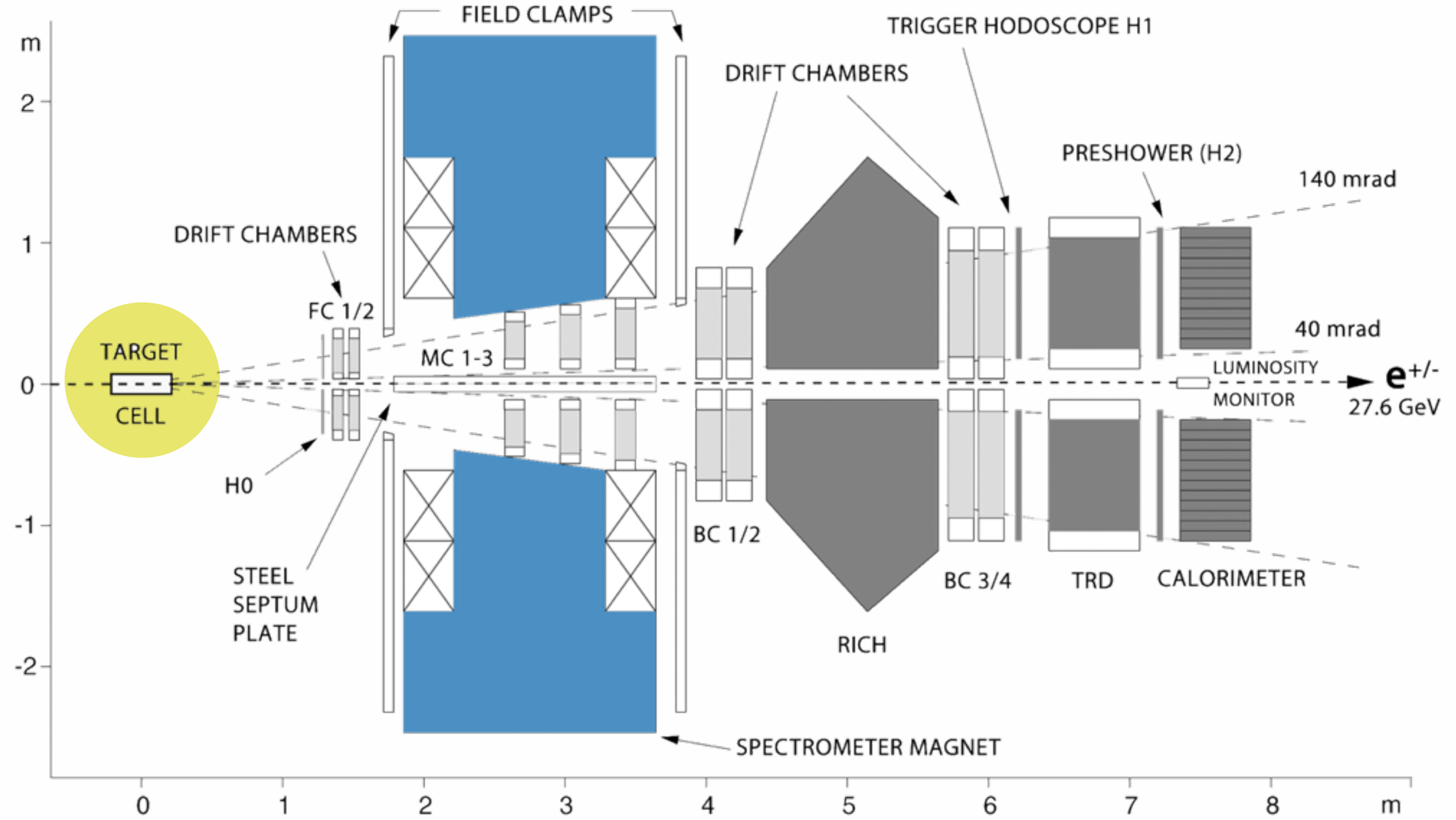
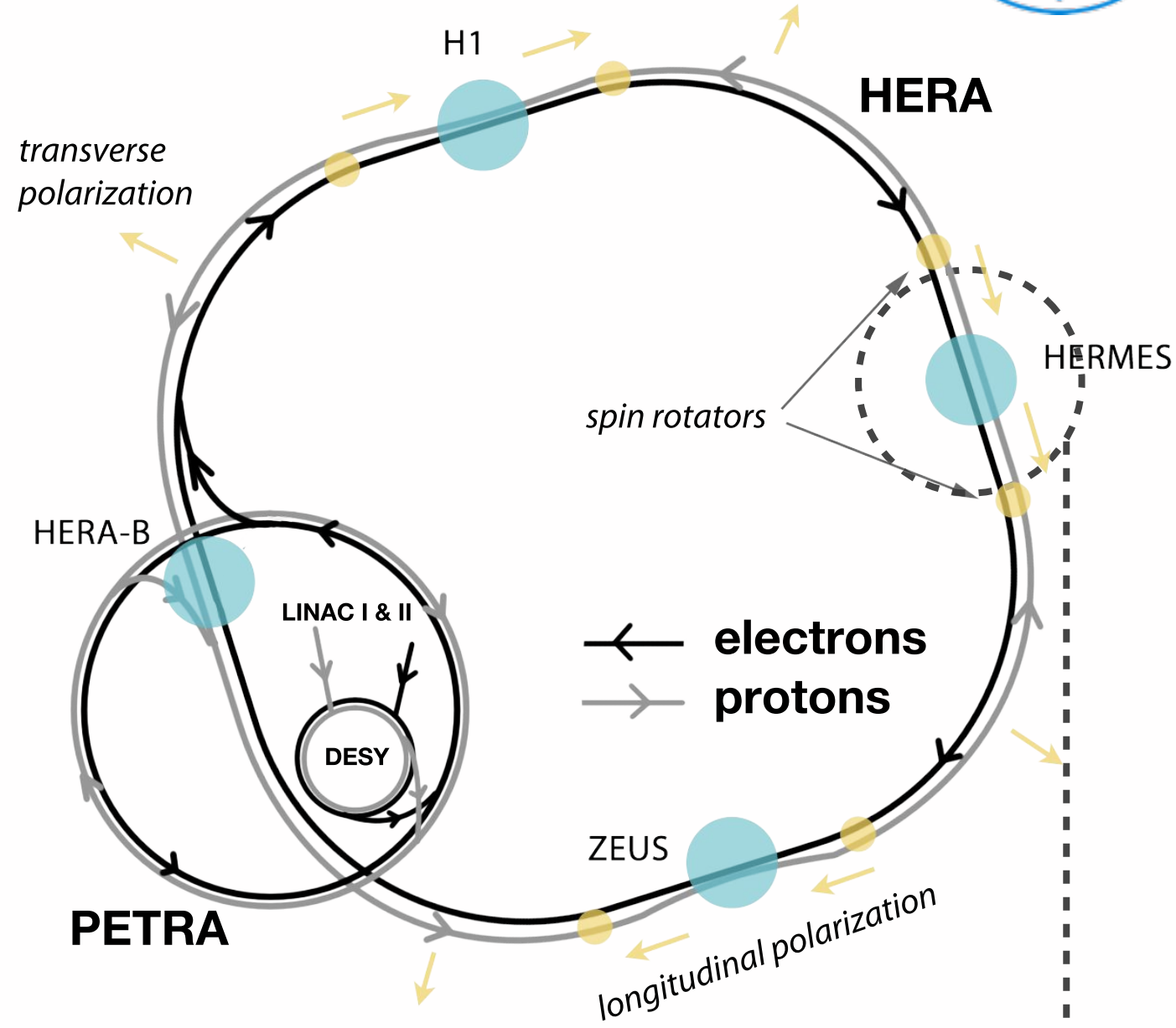
Inclusive hadron production

$$d\sigma_{UT}^{ep \rightarrow hX} \propto A_{UT}^{\sin \psi} S_T \sin \psi$$

↘ **Inclusive SSA amplitude**



HERA accelerator at



- ★ Lepton e^+/e^- beam at 27.6 GeV.
- ★ Transversely polarized H target. Polarization state reversed every 1-3 min.

Event Selection

- ★ Integrated luminosity \sim **146 pb⁻¹**.
- ★ Selected events with at least 1 hadron track in acceptance. (<0.1% lepton contamination).
- ★ Hadrons selected in **2-15 GeV** energy range.
 - >> Further PID using a dual-radiator ring-imaging Cherenkov detector (RICH): RICH unfolding.
- ★ Trigger by signal coincidence of different sub-detectors and energy deposition (>1.4 GeV) in the calorimeter:
 - >> \sim 100 % efficient for electrons.
 - >> Finite efficiency for single hadrons calculated as a function of h-type, momentum and detector position.
 - >> Event-wise determination of average efficiencies.
- ★ Current polarization degree of the target
 $\langle S_{\perp} \rangle = 0.713 \pm 0.063$

For every track/event:

PID weight

$$w_{id} = P^{-1}(\text{id}^{\text{true}} | \text{id}^{\text{tag}})$$

Event weight (trig. eff.)

$$\langle \epsilon \rangle = 1 - \prod_{\text{event}} (1 - \epsilon_w)$$

or 1 if a detected lepton

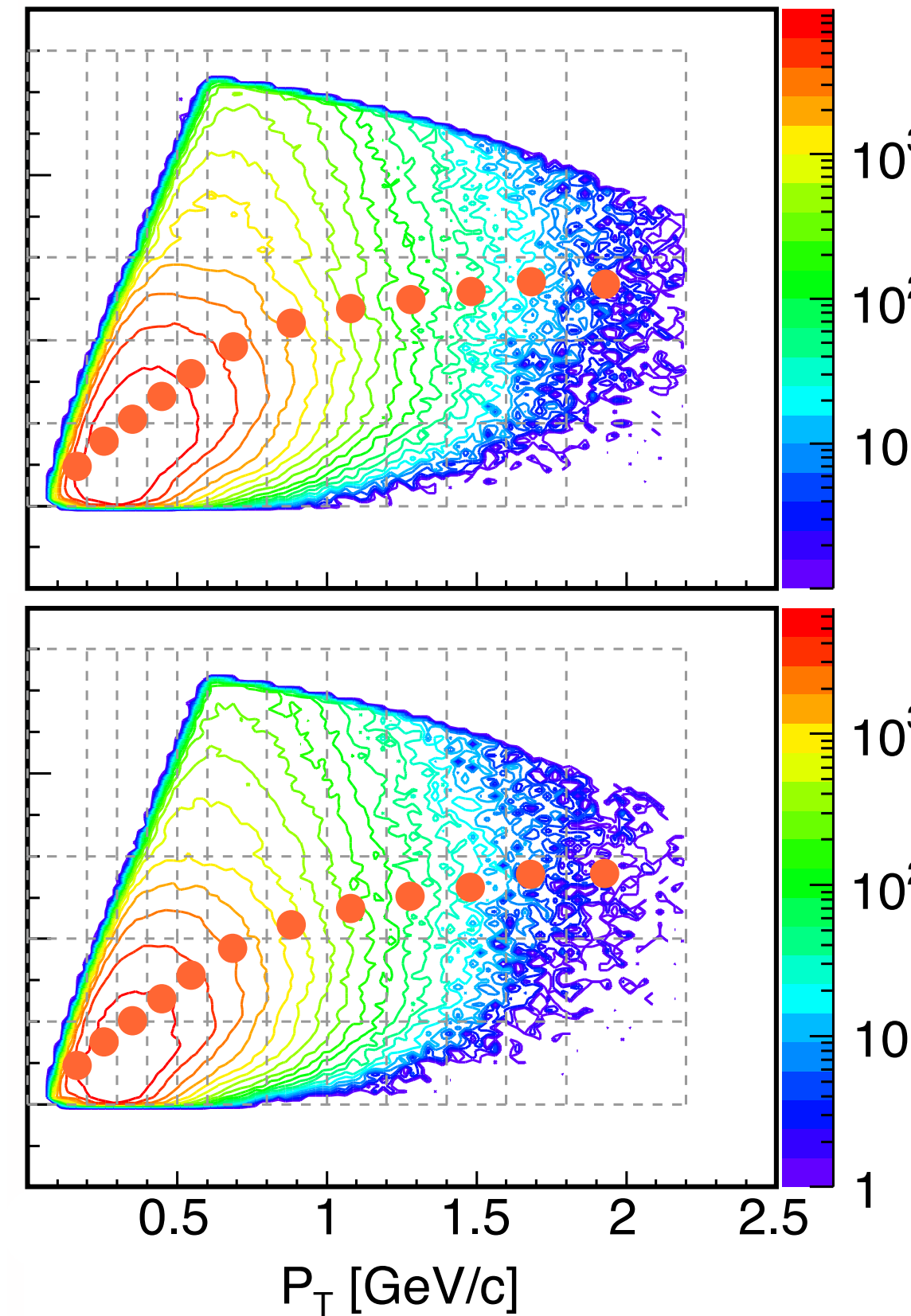
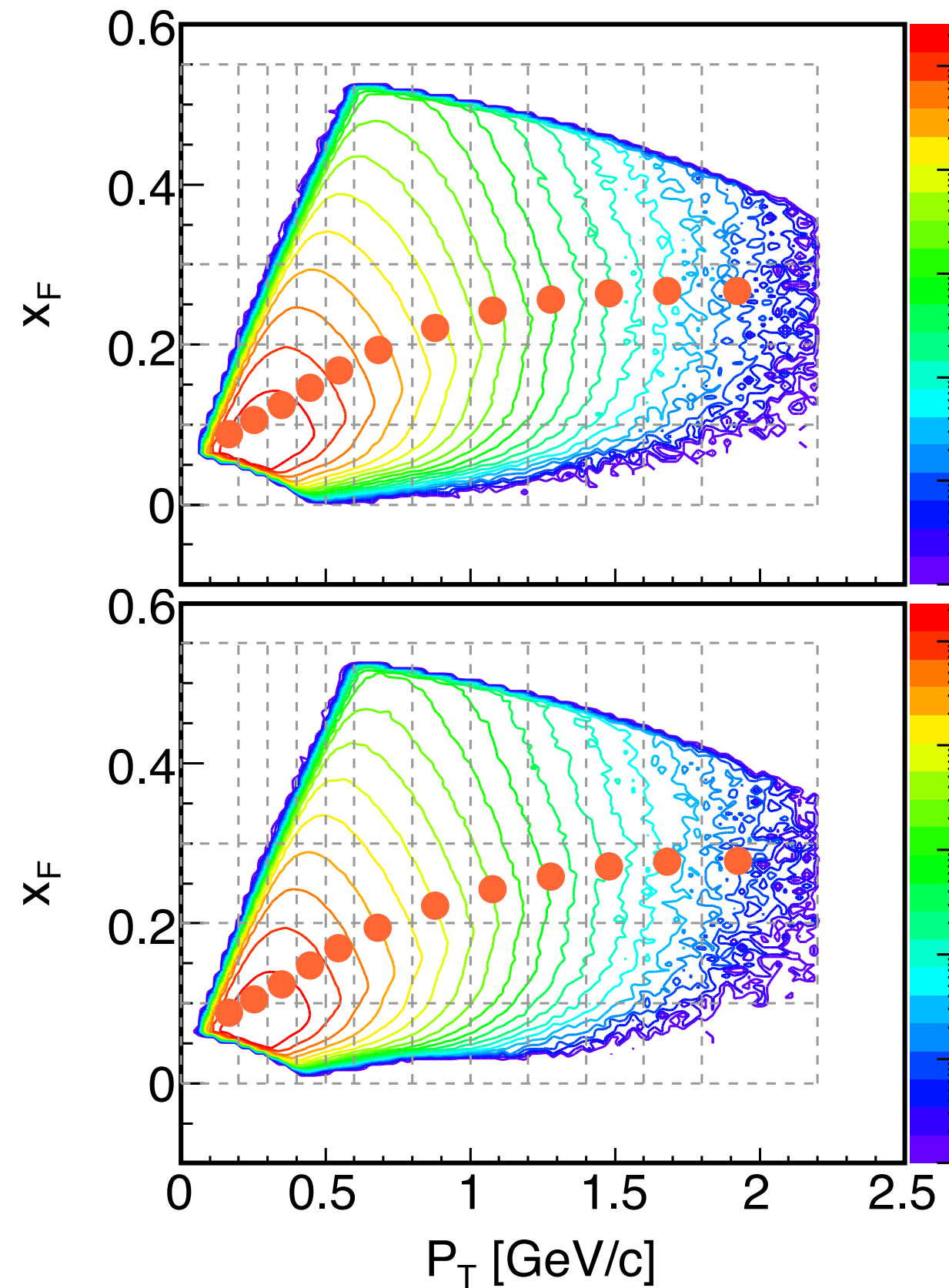
Final weight

$$W_{\text{event}} = \frac{w_{id}}{\langle \epsilon \rangle}$$

Kinematics and Yields

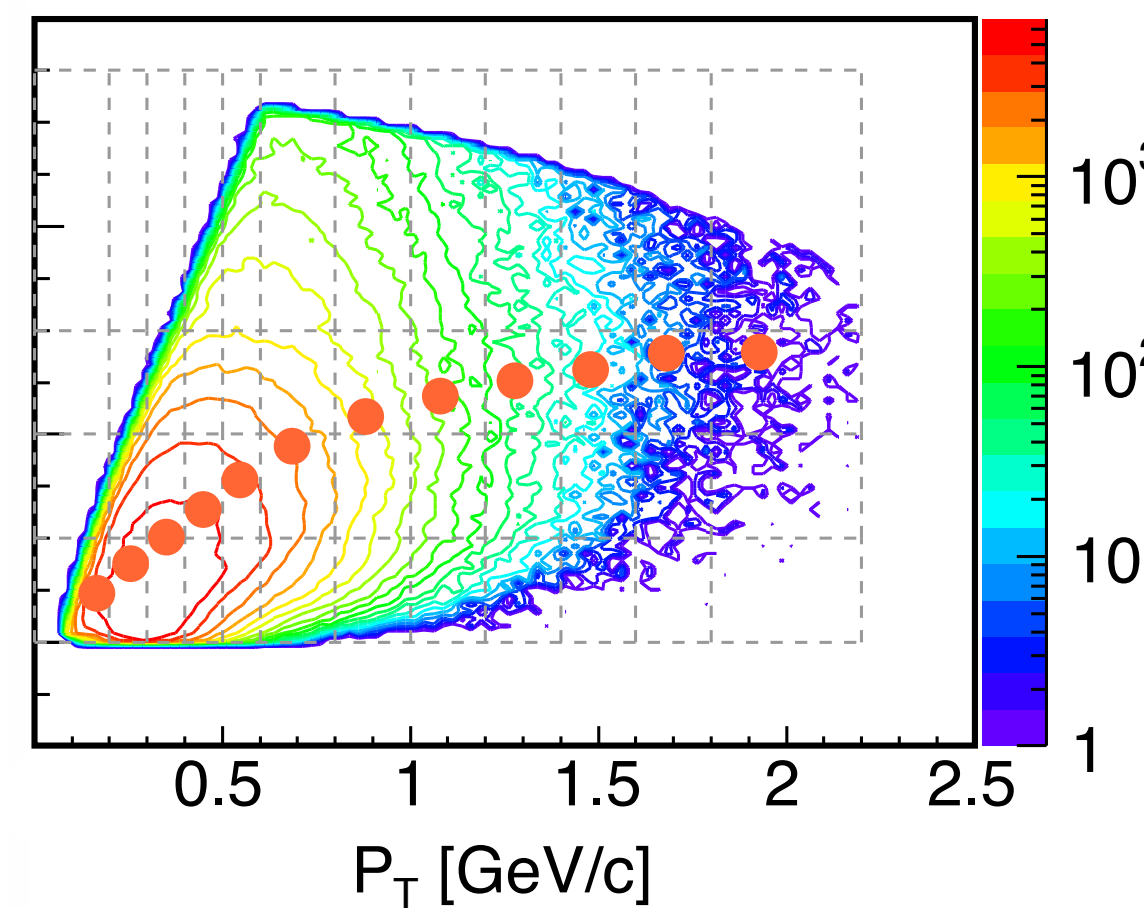
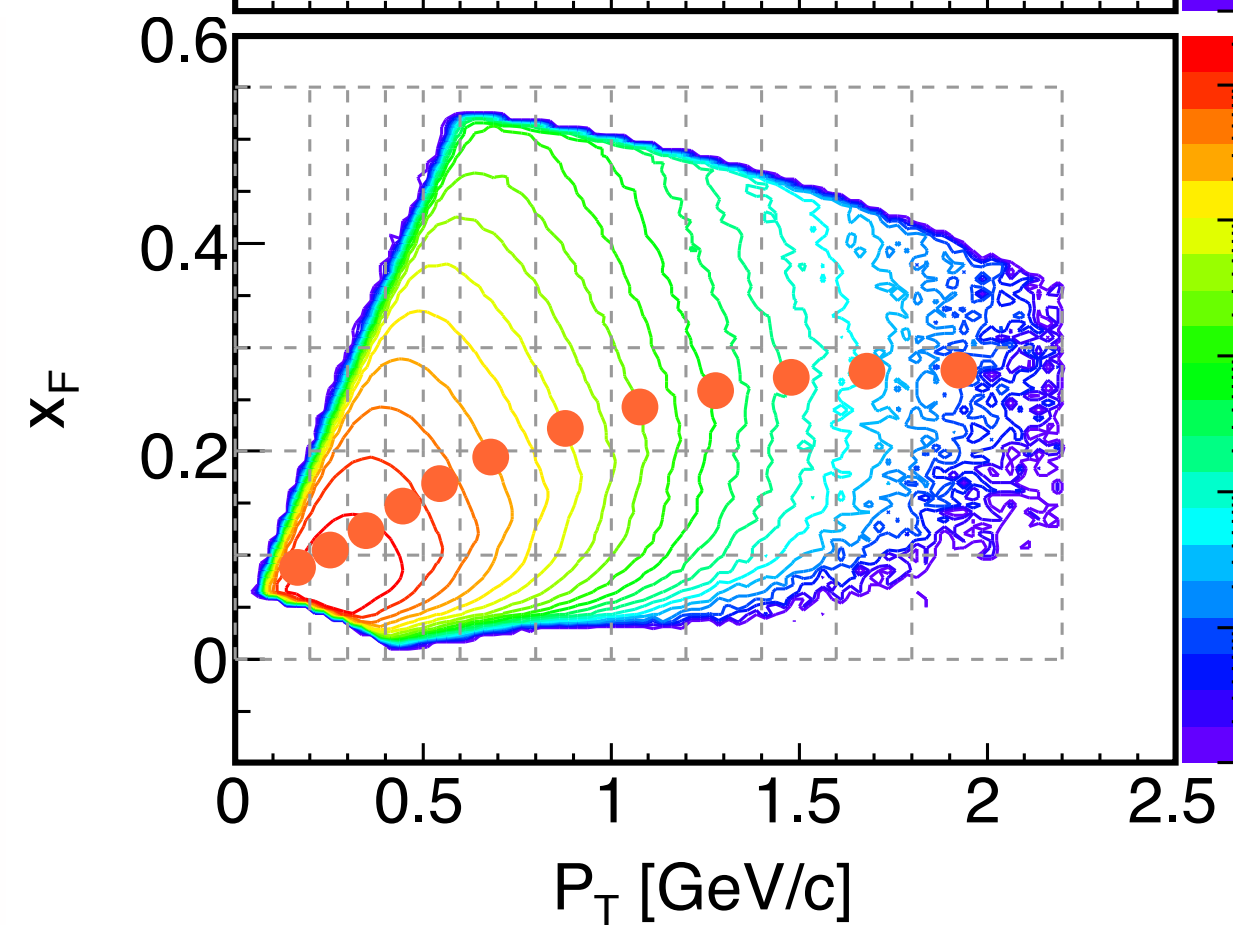
Corrected by trig. eff.

π^+
172 M



K^+
14.5 M

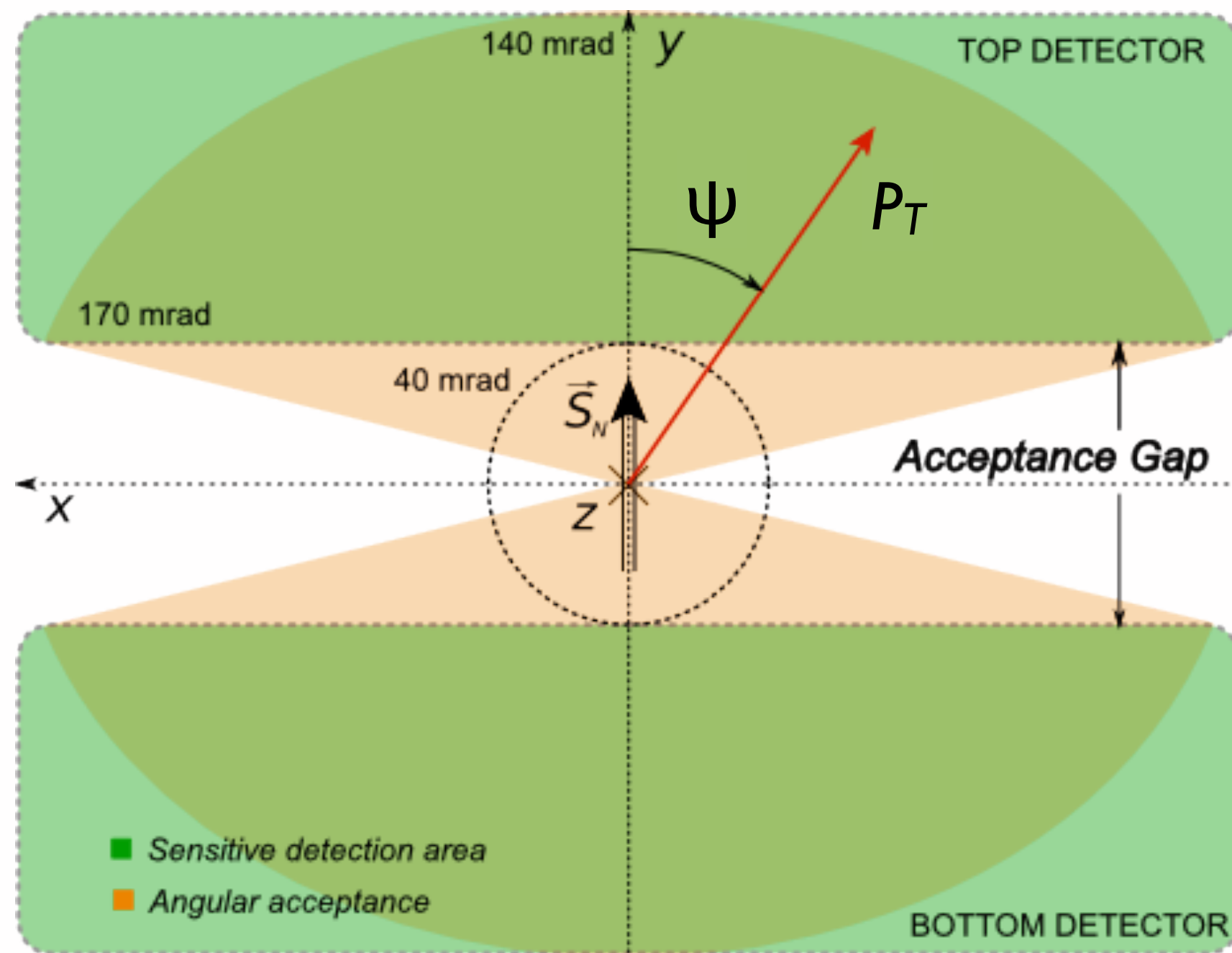
π^-
142 M



K^-
7.3 M

Extraction of the Asymmetry Amplitude

Maximum likelihood fit



In every kinematic bin $\{x_F, P_T\}$

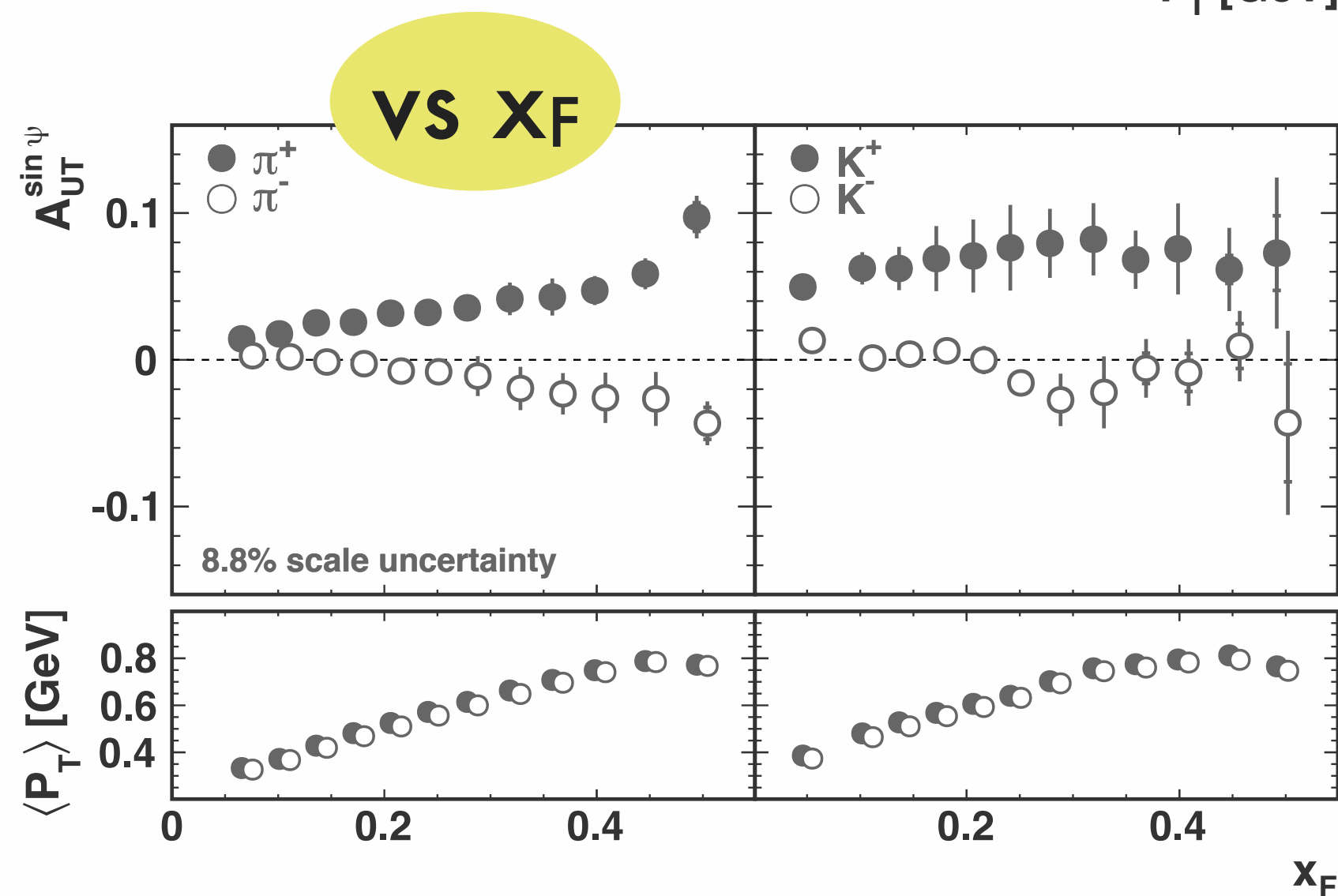
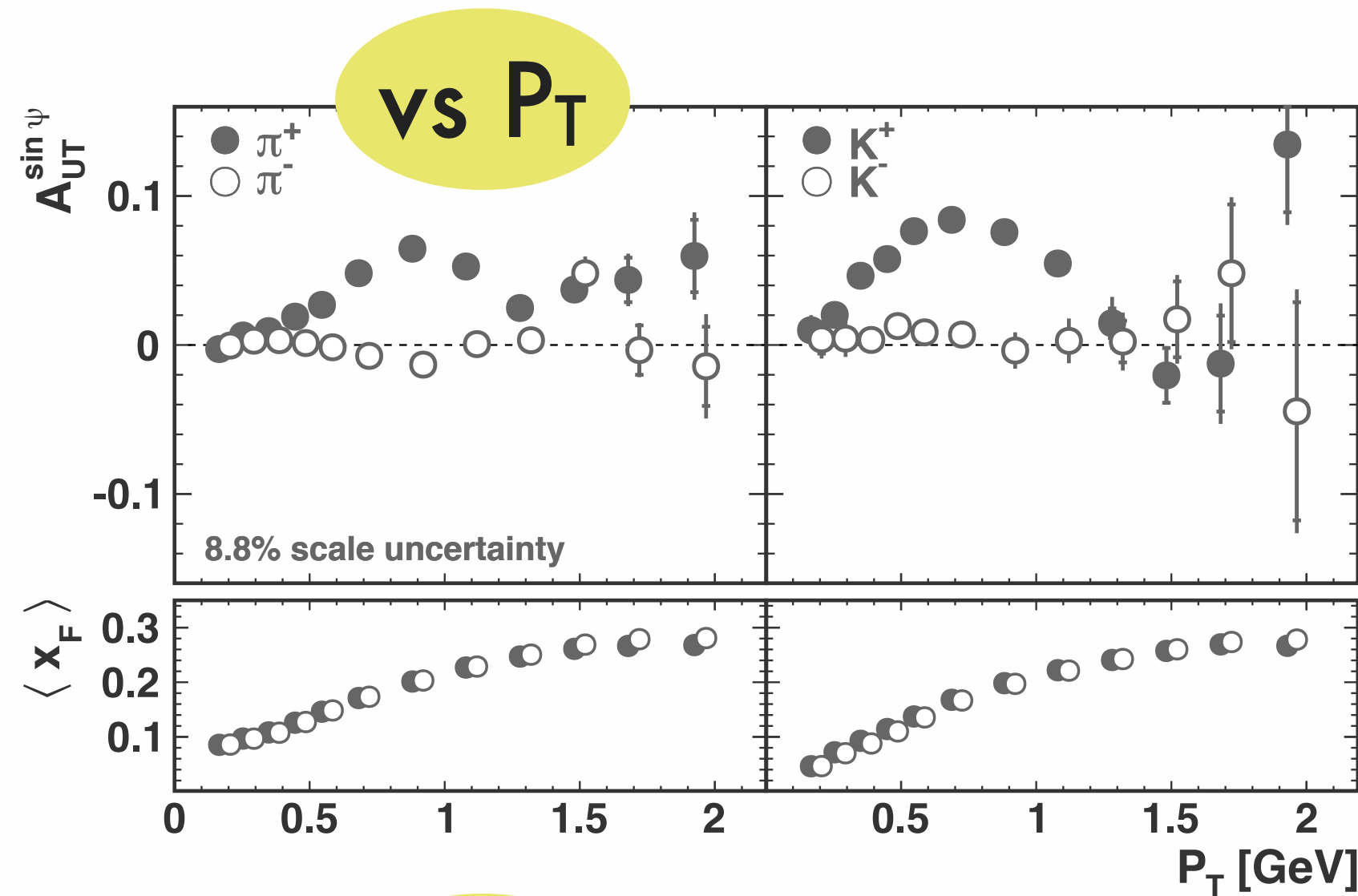
$$d\sigma(\psi) = d\sigma_{UU} \left[1 + S_T A_{UT}^{\sin \psi} \sin \psi \right]$$

The amplitudes are extracted from data by a max. likelihood fit to the differential cross section.

$$\mathcal{L} = \prod_i \left(\frac{1 + S_{T,i} \left[A_0 + A_{UT}^{\sin \psi} \sin \psi_i \right]}{L_i^{\uparrow\downarrow}} \right)^{W_i}$$

Polarization degree $\rightarrow S_{T,i}$
 Asymmetry amplitude $\rightarrow A_{UT}^{\sin \psi}$
 Relative luminosity $\rightarrow L_i^{\uparrow\downarrow}$
 Particle weight $\rightarrow W_i$
 Trigger efficiency $\rightarrow W_i$
 Event weight $\rightarrow W_i$

Detection efficiencies and acceptance effects cancels in the max. likelihood fit.



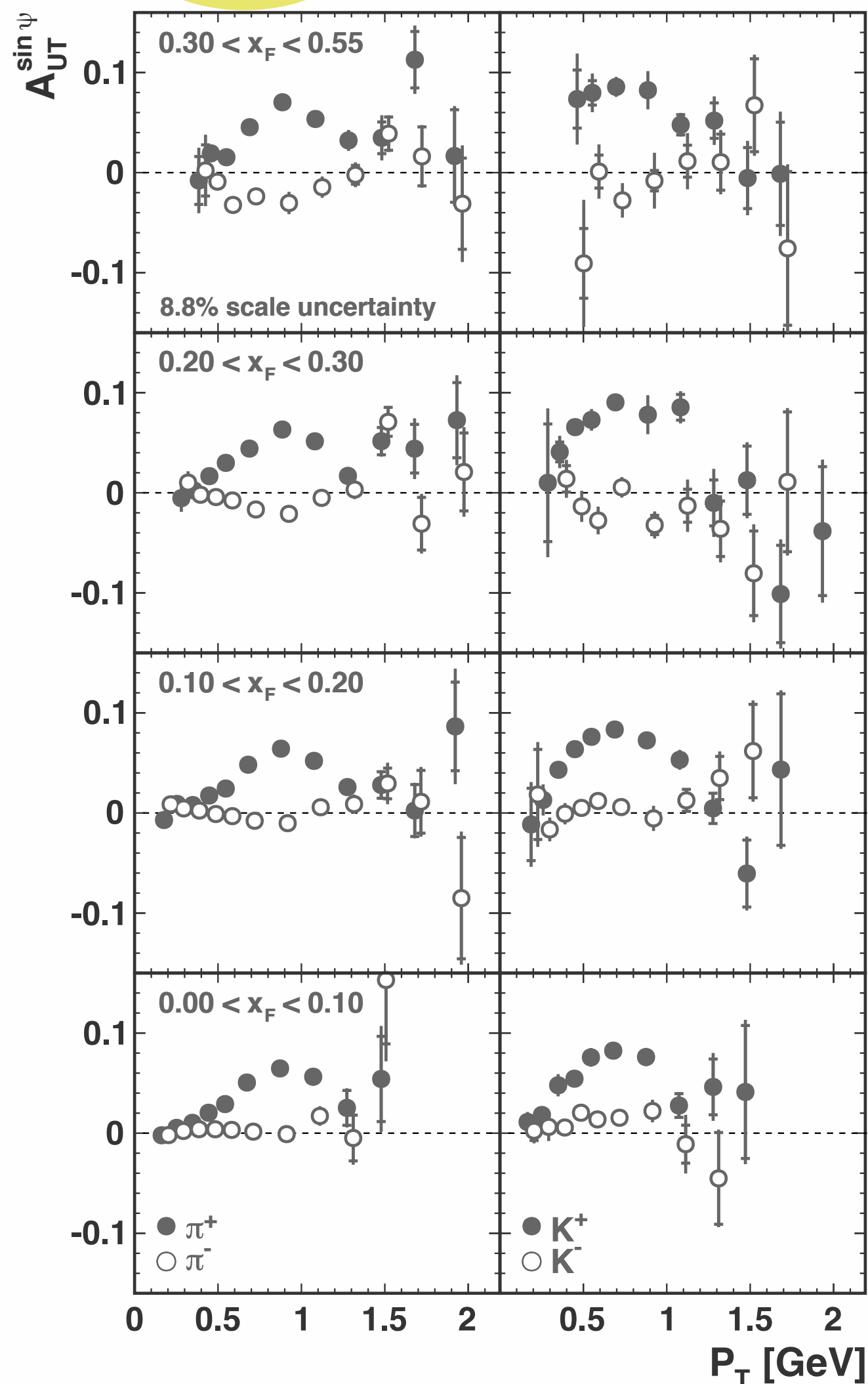
Inclusive SSA (ID)

...as a function of P_T and x_F separately

- ★ The asymmetry amplitudes are **positive** for π^+ and K^+ and compatible with **zero** for π^- and K^- .
- ★ A clear two-fold structure in P_T .
 - >> The amplitudes increase with P_T up to ~ 0.6 (~ 0.9) at $P_T \sim 0.8$ GeV for π^+ (K^+).
 - For this point on, the tendency is inverted.
 - >> For π^+ there is another rise at high P_T , while for K^+ it seems to remain null.
- ★ The amplitudes vs. x_F increase (decrease) nearly linearly for π^+ (π^-).
- ★ For K^+ (K^-) are about constant around 0.7 (0.0).

P_T and x_F are strongly correlated.

vs P_T in x_F bins



Inclusive SSA (2D)

...as a function of P_T and x_F simultaneously

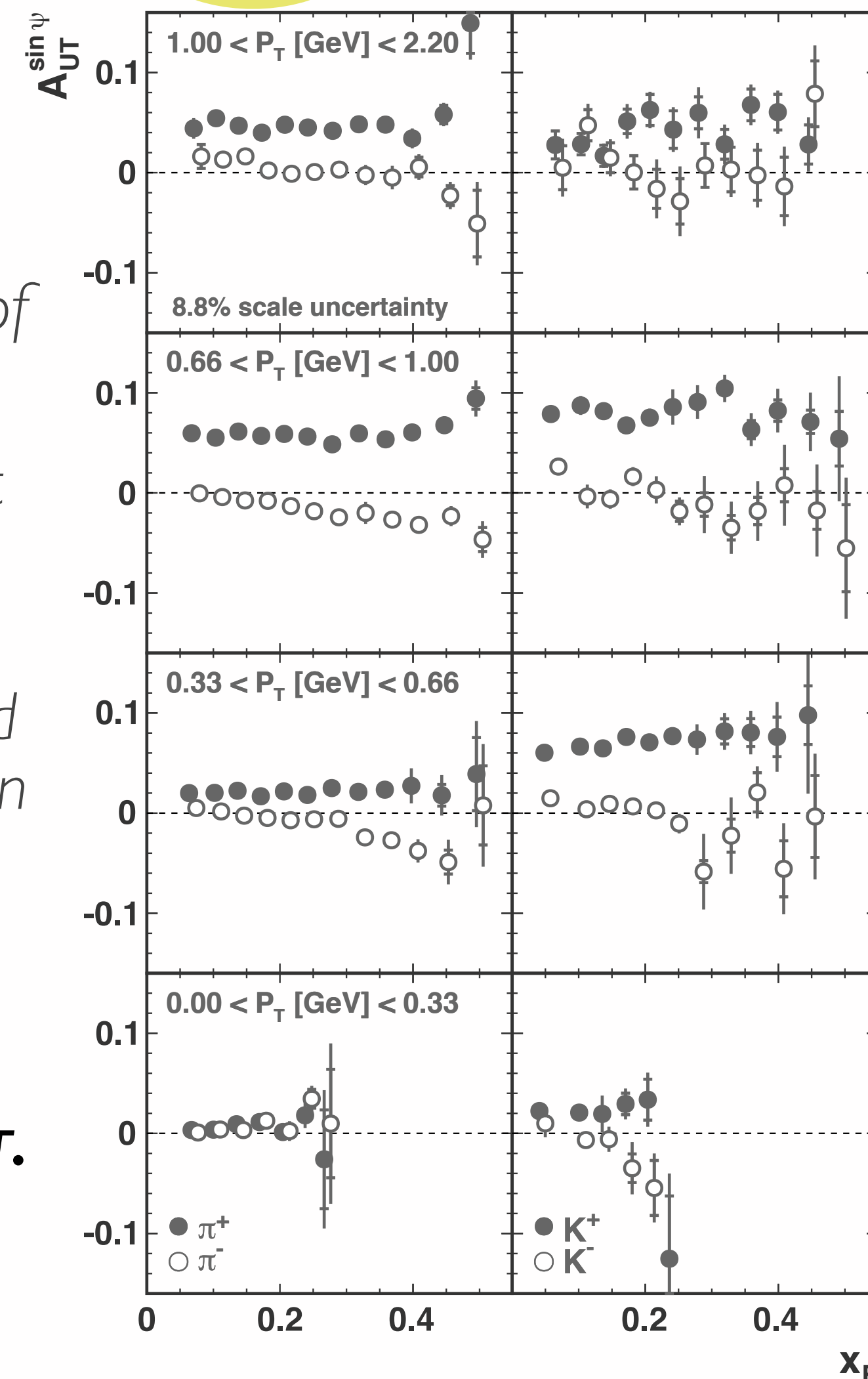
π^+ : Similar dependence on P_T . No dependence on x_F .
The 1D x_F dependence is a reflection of the underlying dependence on P_T .

π^- : Like in the 1D case, there is a slight decreasing linear tendency with x_F .

K^+ and K^- : The dependence on x_F of the kaon amplitudes is less pronounced in 2D, with a slight tendency towards an increase (decrease) with x_F for positive (negative) kaons.

Most of the structure is on P_T .

vs x_F in P_T bins





Systematic uncertainties

- ★ Misalignment of detectors/beam.
- ★ Hadron misidentification.
- ★ Angular and momentum resolution.
- ★ Secondary interactions:
Radiative corrections, M.S., decays, etc.
- ★ Target polarization:
>> 8.8% scale uncertainty.
- ★ Trigger efficiency correction:
>> Difference among two alternative methods.
- ★ Compatibility of data productions:
>> Negligible.

All-in-one Monte Carlo approach:

High-statistics PYTHIA 6.2 +
Spin-dependent model (from fit to data).

$$\mathcal{A}_{UT}^{\sin \psi}(x_F, P_T)$$

Full description of the detector

Extraction of the Asym. in every bin $A_{UT,MC}^{\sin \psi}$

The systematic uncertainty is taken then as the biggest of either:

(i) The deviation

$$|A_{UT,MC}^{\sin \psi} - \mathcal{A}_{UT}^{\sin \psi}(\langle x_F \rangle, \langle P_T \rangle)|$$

(ii) The statistical error of $A_{UT,MC}^{\sin \psi}$



Event Categories

Inclusive sample:

Comprises several subsamples, related to different production channels. These subsamples contribute differently to the measured SSA, and can be differentiated by electron tagging and kinematic constraints on DIS variables:

$$A_{UT}^{\sin \psi} = \sum_i^{\text{samples}} f_i A_{UT,i}^{\sin \psi}$$

Anti-tagged (~98%):

The undetected lepton in most cases had a small scattering angle and remained within the beam pipe. **Photoproduction** regime ($Q^2 \approx 0$).

Tagged or Semi-inclusive:

The scattered lepton is detected.

DIS events:

Non-DIS events

$$(Q^2 > 1 \text{ GeV}^2, W^2 > 10 \text{ GeV}^2, 0.023 < x < 0.4, 0.1 < y < 0.95)$$

mid-z (0.2 < z < 0.7):

Identical to the SIDIS in Sivers and Collins measurements at HERMES.

$$\langle Q^2 \rangle \lesssim \langle P_T^2 \rangle$$

low-z ($z < 0.2$)

$$z \equiv \frac{P \cdot P_h}{P \cdot q} \stackrel{\text{lab}}{=} \frac{E_h}{\nu}$$

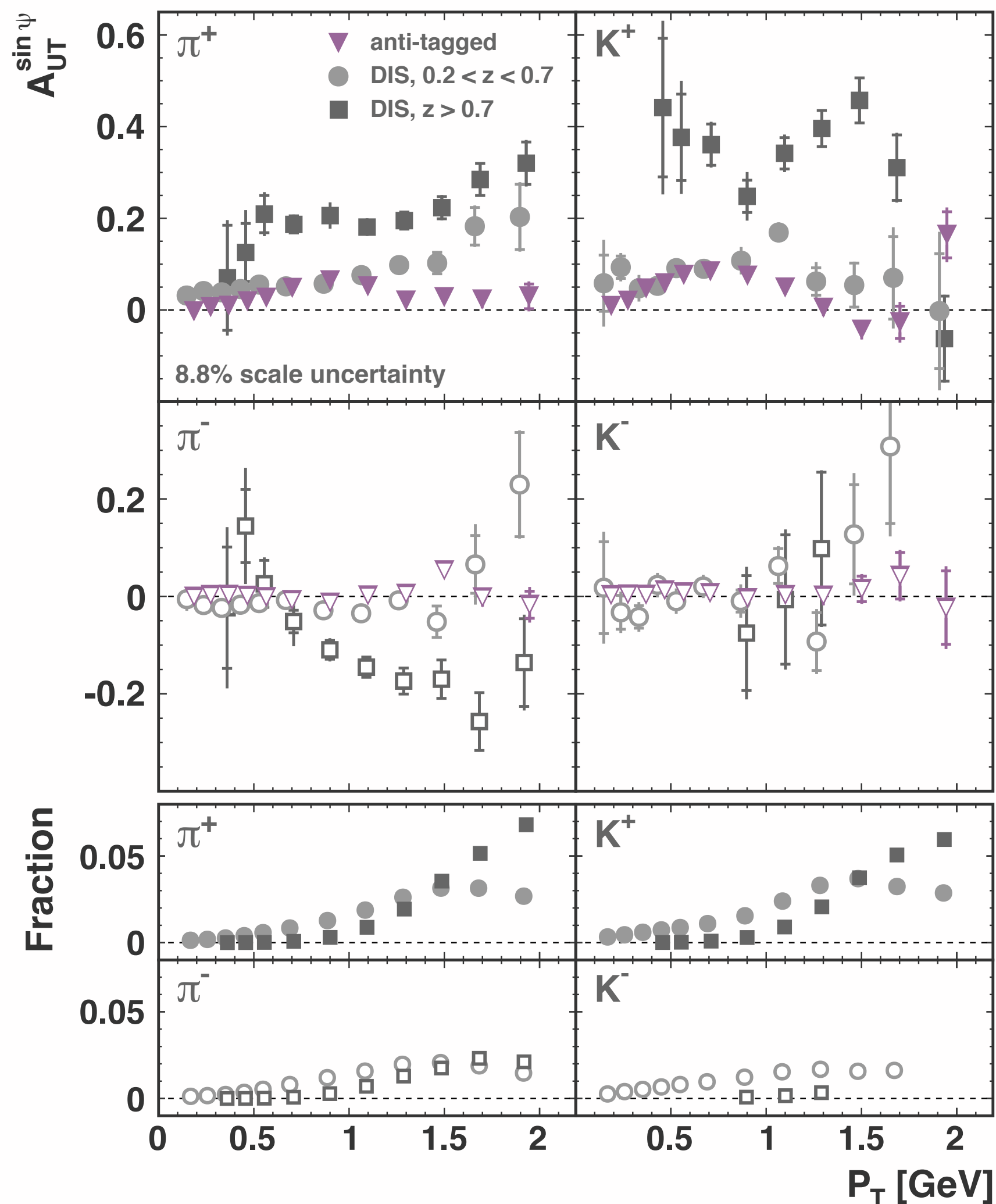
fractional virtual-photon energy carried by the hadron

high-z ($z > 0.7$):

Hadrons in this region might be produced in exclusive reactions.

$$\langle Q^2 \rangle < \langle P_T^2 \rangle$$

.. as a function of P_T



Anti-tagged (*photoproduction*)

- ★ *Inclusive asymmetries dominated by quasi-real photoproduction: they go like $1/P_T$.*
- ★ *At high P_T contribution from DIS becomes sizable.*

DIS, mid-z (*SIDIS*)

- ★ *Q^2 is the largest scale. TMD's and frag. functions holds. Non-vanishing asymmetries at high P_T .*
- ★ *Sivers effect is expected to be the main contribution.*

DIS, high-z (*Exclusive*)

- ★ *Large asymmetries in π^+ , π^- and K^+ .*
- ★ *Exclusive production of ρ mesons in π^+ and π^- channels can be substantial.*
- ★ *Dominance of the struck quark d in π^- production at high z .*

Transverse target Single-Spin Asymmetry (SSA) in inclusive electroproduction of charged pions and kaons.

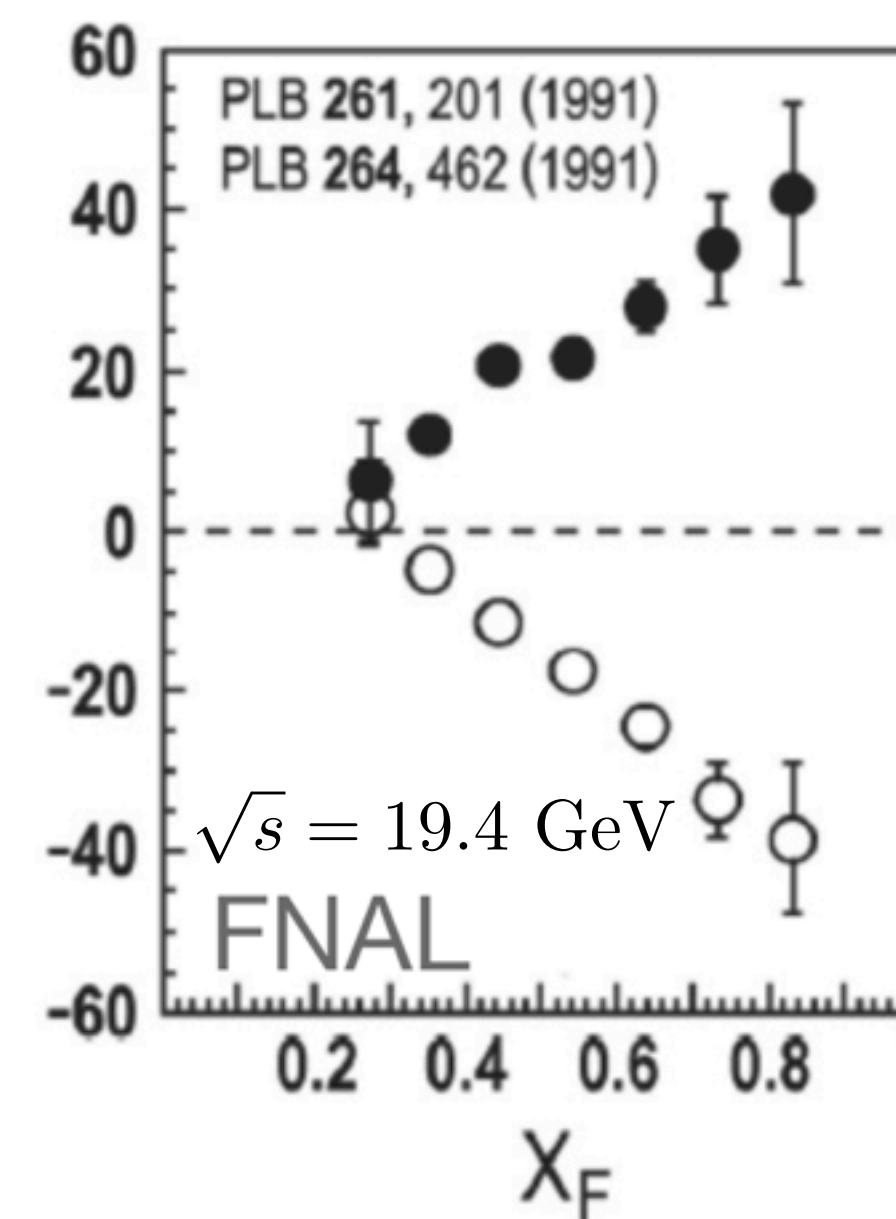
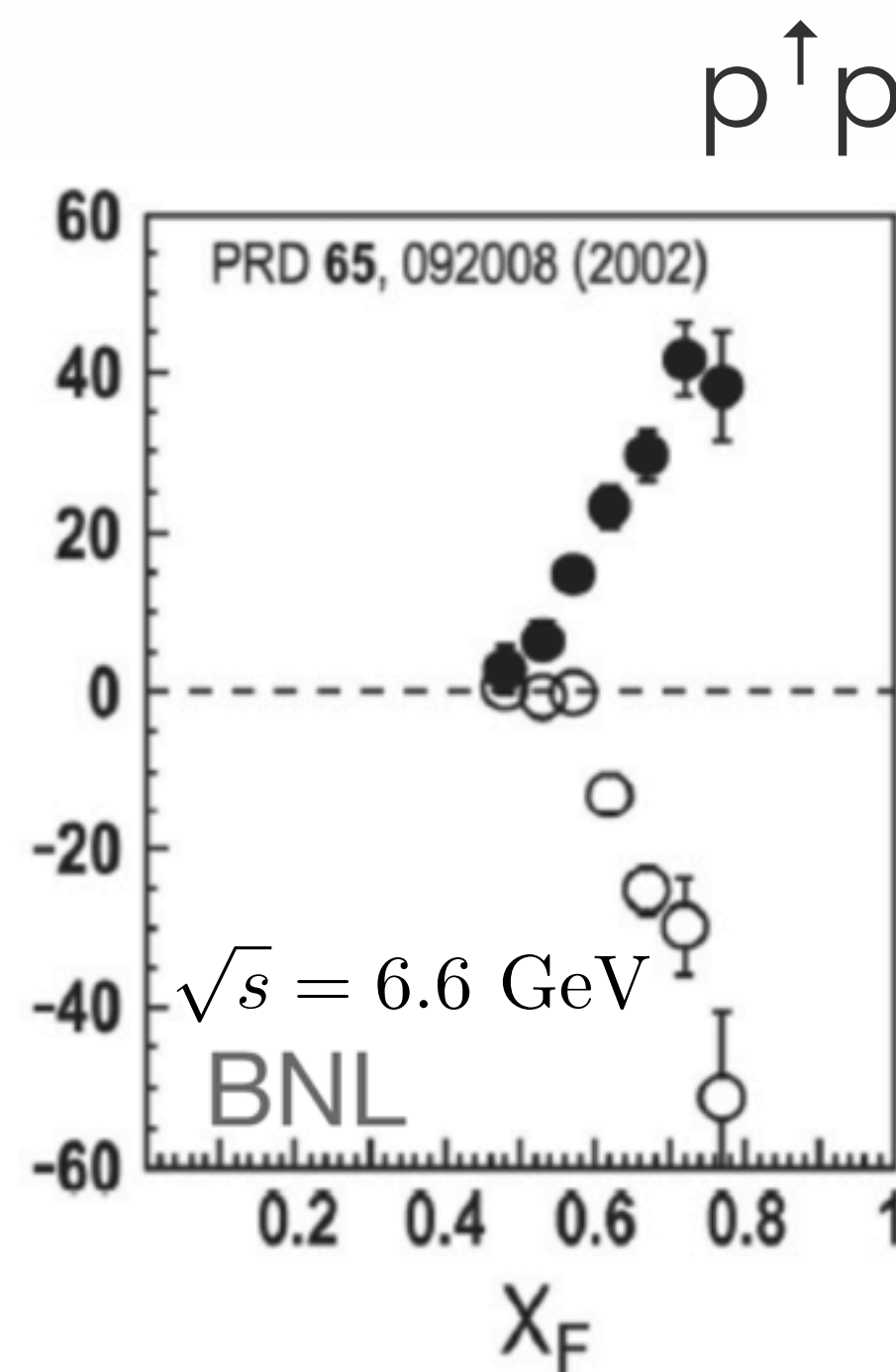
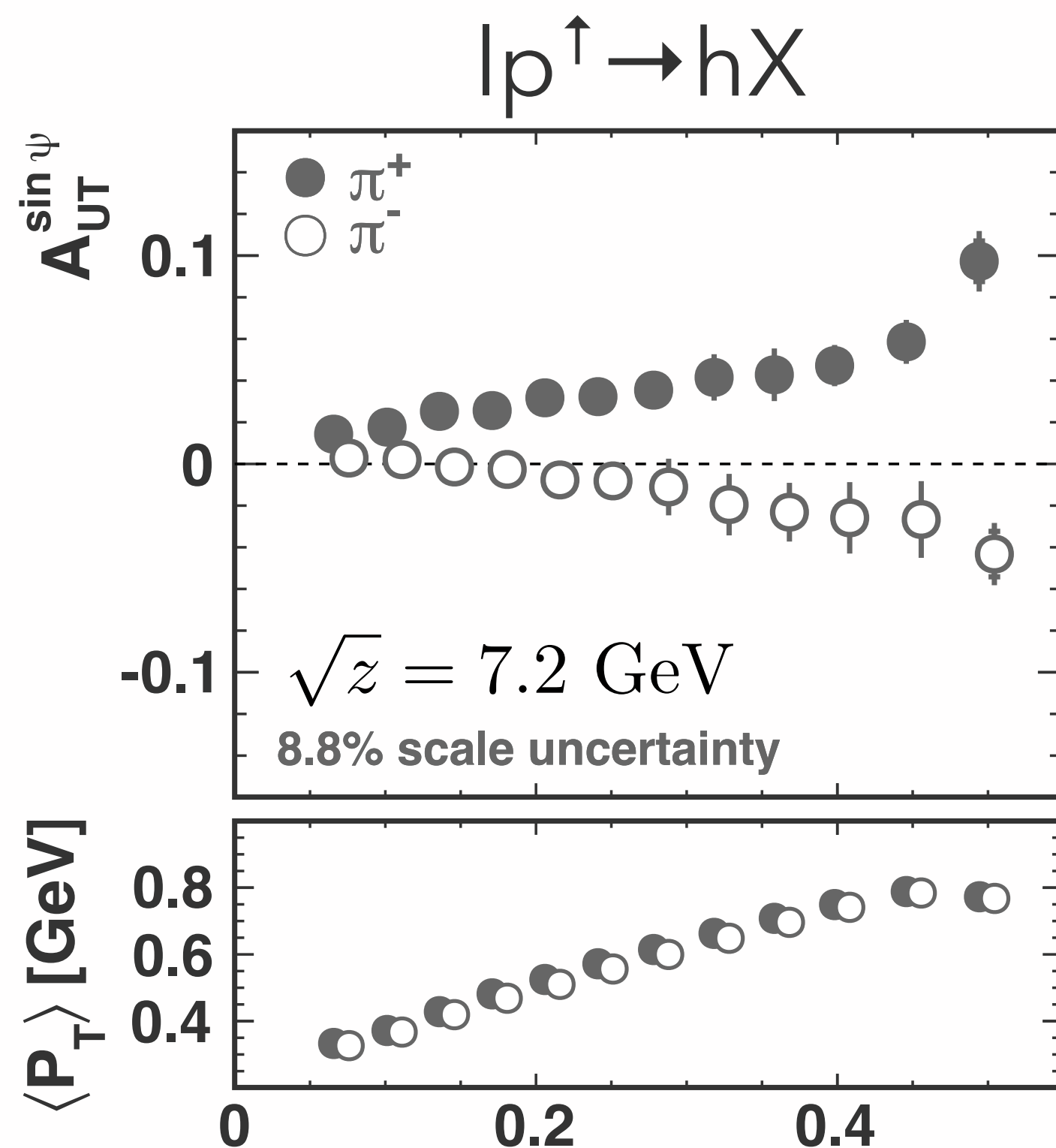
Summary and Outlook

- >> Transverse azimuthal SSA measured in inclusive electroproduction of pions and kaons.
- >> Two-dimensional extraction by binning simultaneously in x_F and P_T :
 - ★ For π^+ , the asymmetry is independent of x_F .
 - ★ For π^- , and less for K^- (K^+), the asymmetry amplitudes decrease (increase) with x_F .
- >> As a function of P_T , the amplitudes are positive for positive mesons.
 - ★ Two-fold structure at low and high P_T .
- >> Inclusive sample dominated by photoproduction :
 - ★ Description in terms of higher-twist effects: Predicts a vanishing A_{UT} with $1/P_T$.
- >> At high- P_T , sizable contribution from SIDIS processes.
 - ★ Description in terms of TMD Sivers distribution function: Non-vanishing A_{UT} at high P_T .
- >> For DIS at high- z , large asymmetry amplitudes have been found.
 - ★ Substantial contribution from exclusive processes.
 - ★ Effect from favored fragmentation of the struck quark dominate.

Backup Slides

Comparison

to A_N in hadron-hadron collisions

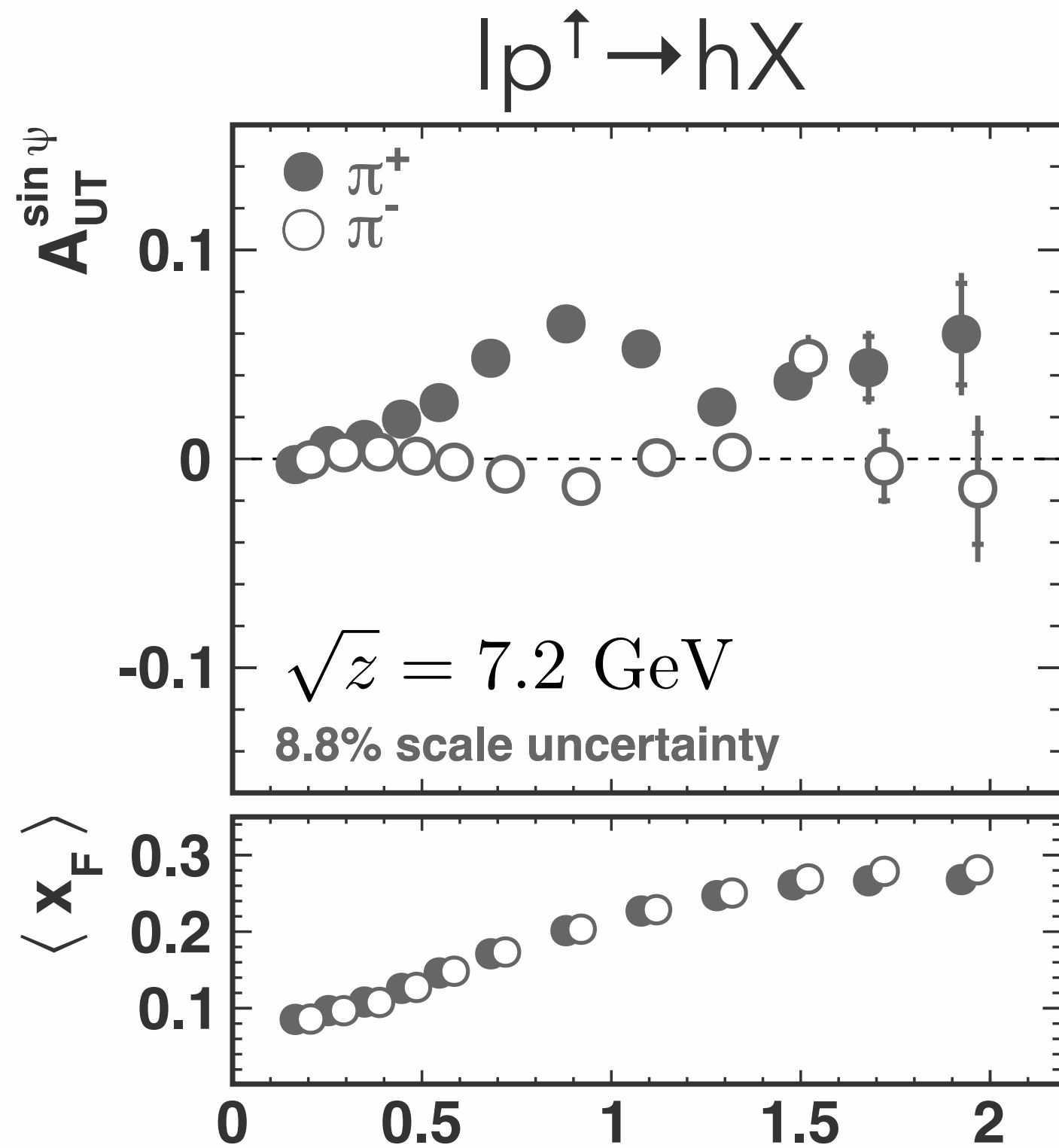


Smaller size, but similar evolution.



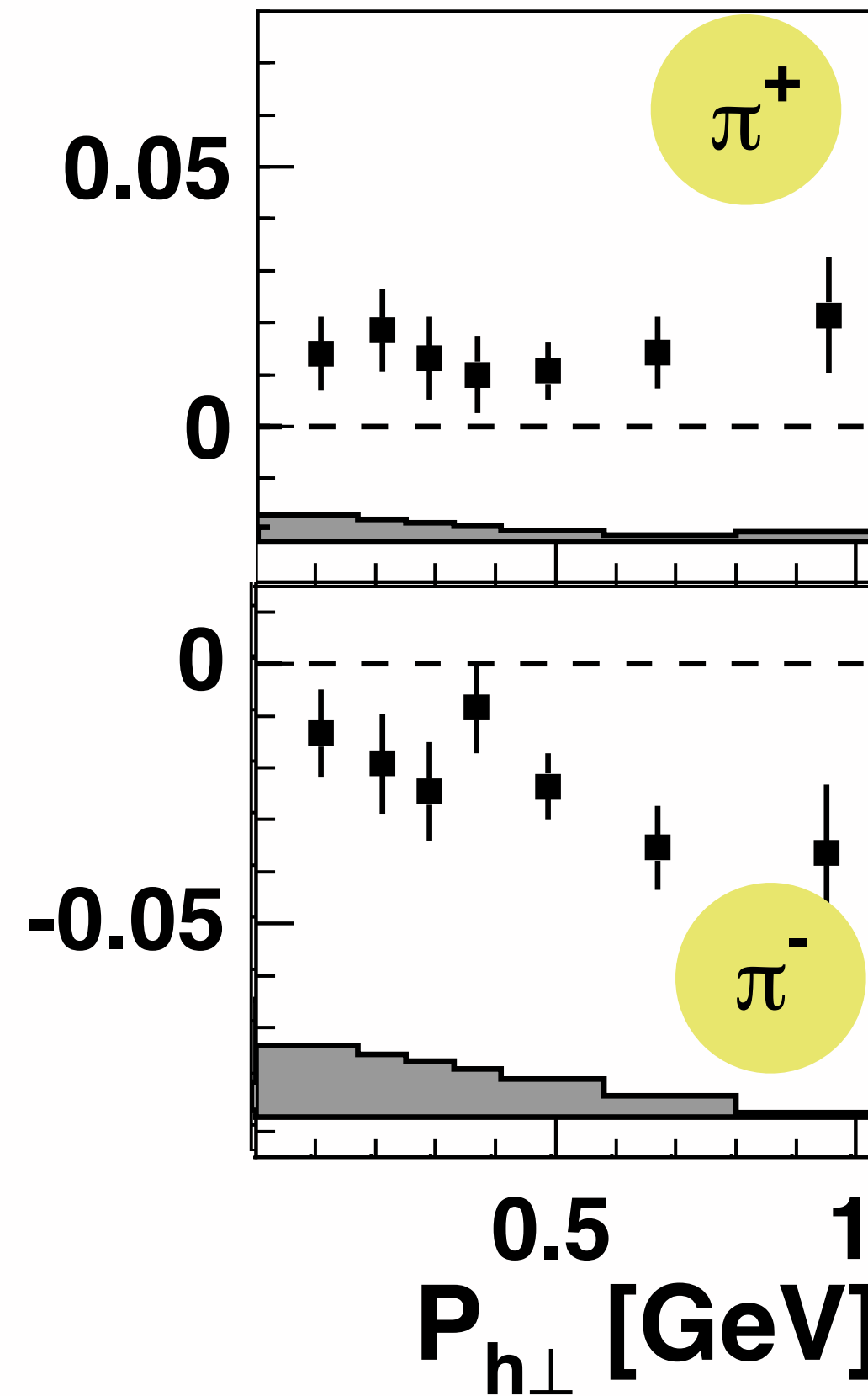
Comparison

to Sivers and Collins amplitudes in SIDIS



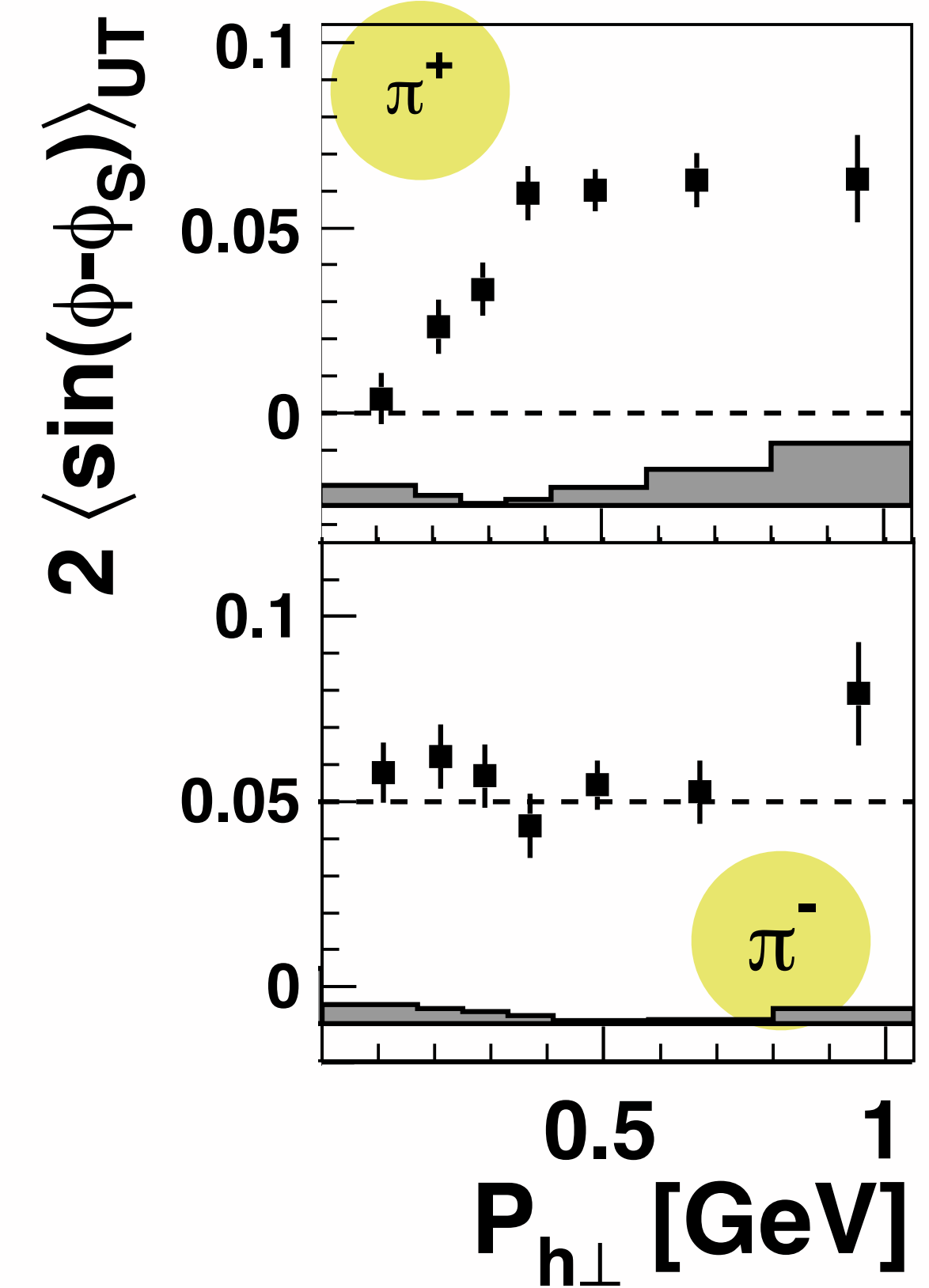
Inclusive SSA similar to Sivers.

COLLINS



$lp^\uparrow \rightarrow l'hX$

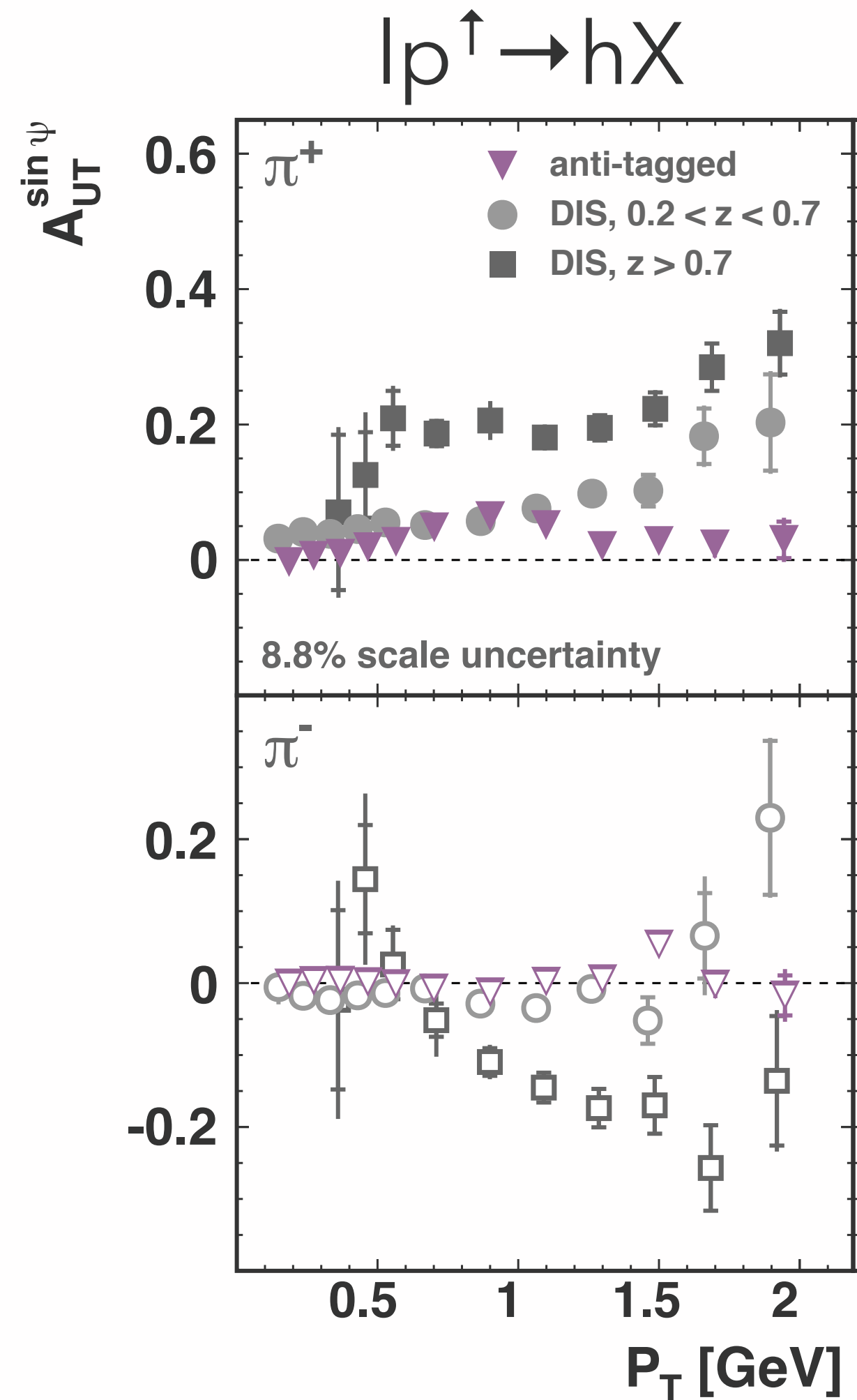
SIVERS



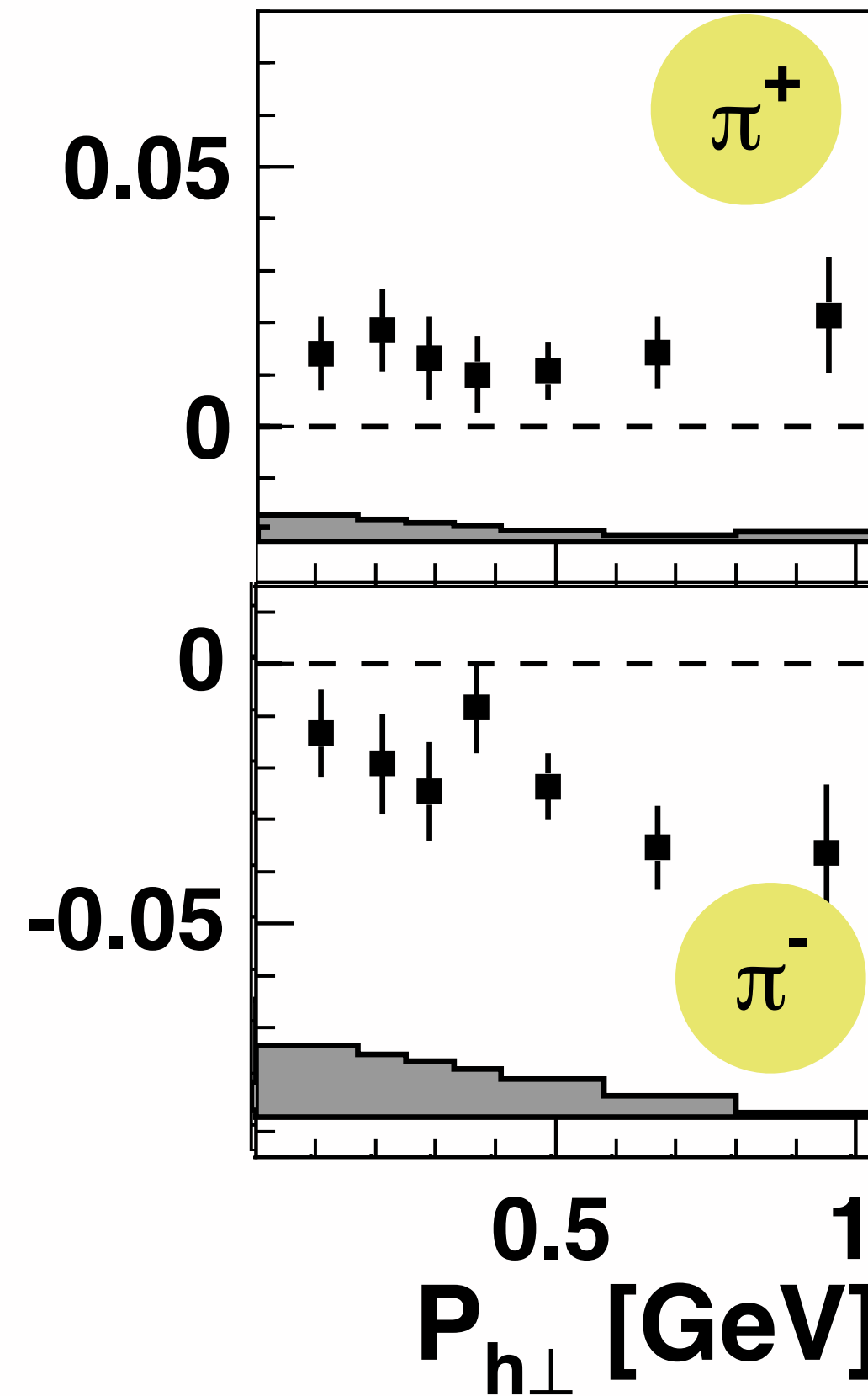


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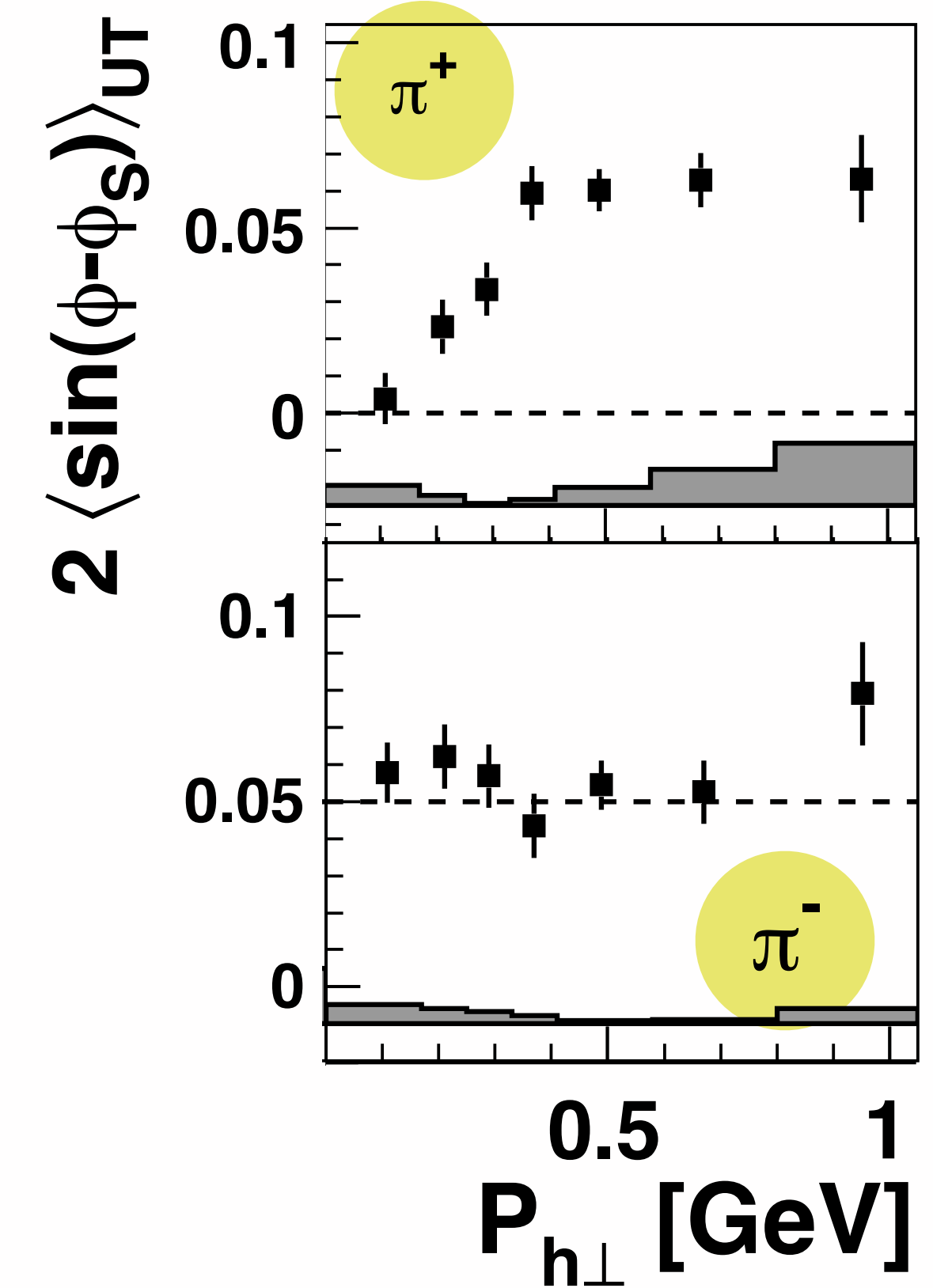


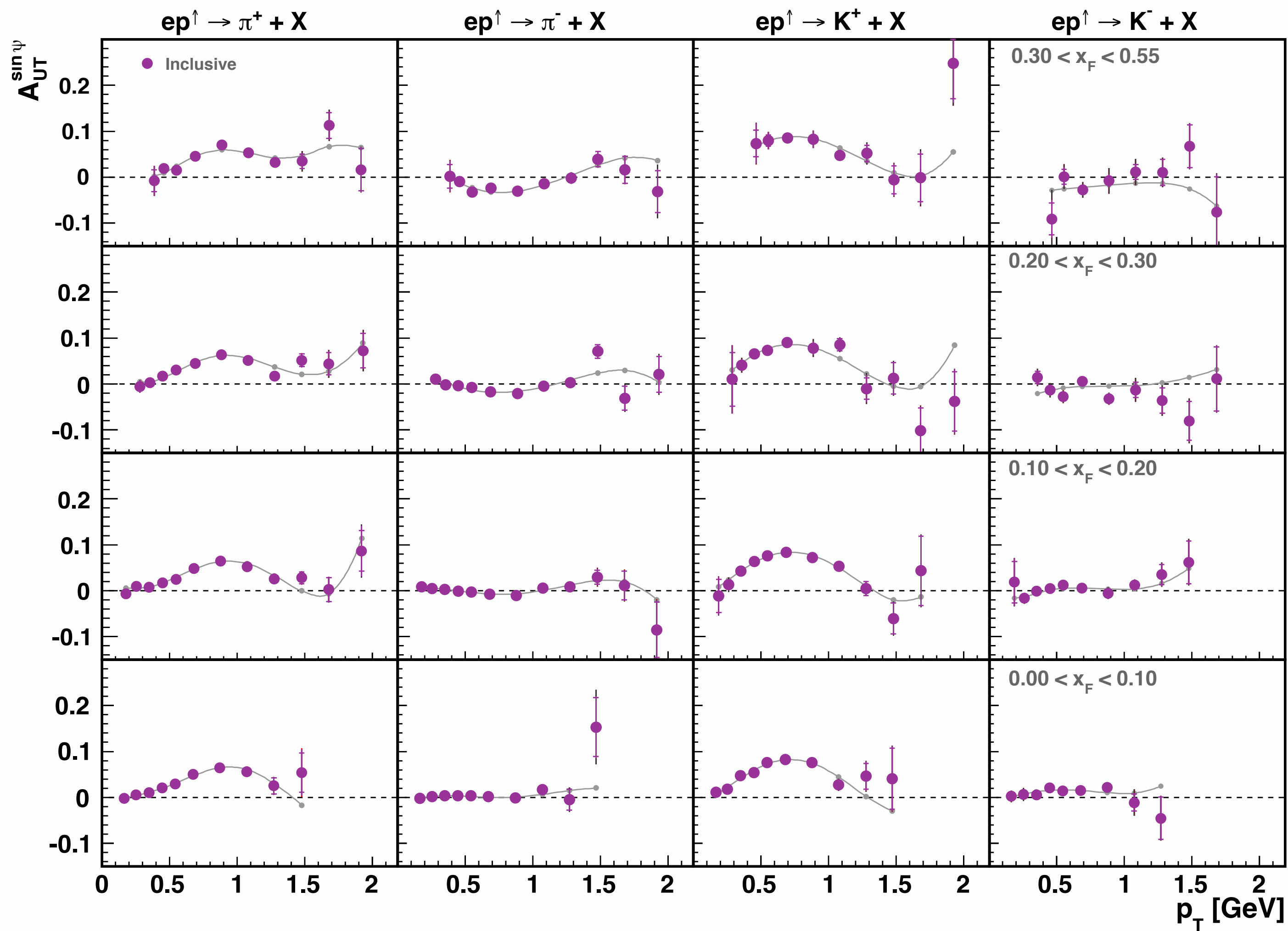
COLLINS



$l p^\uparrow \rightarrow l' h X$

SIVERS





A parametric model

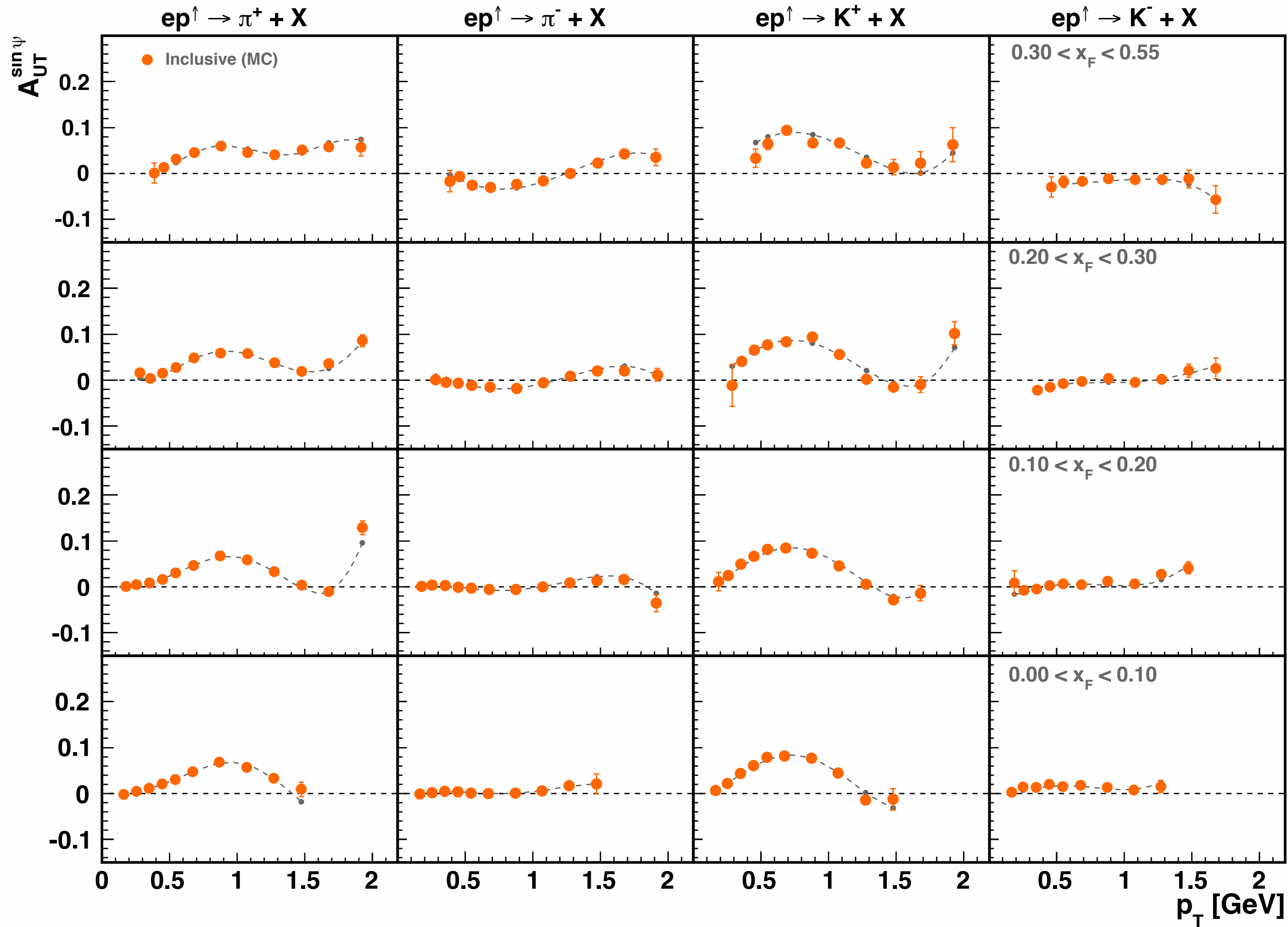
Fit to global data

The functional form is a Taylor expansion in P_T (up to fifth order) and x_F (up to first order) around the average kinematics of the entire experimental data sample.

$$A_{UT}^{\sin \psi} = \sum_{i=0}^5 \hat{P}_T^i c_i + \hat{x}_F \sum_{i=0}^5 \hat{P}_T^i d_i$$

$$\hat{P}_T \equiv P_T - \langle P_T \rangle$$

$$\hat{x}_F \equiv x_F - \langle x_F \rangle$$



A parametric model

Fit to global data

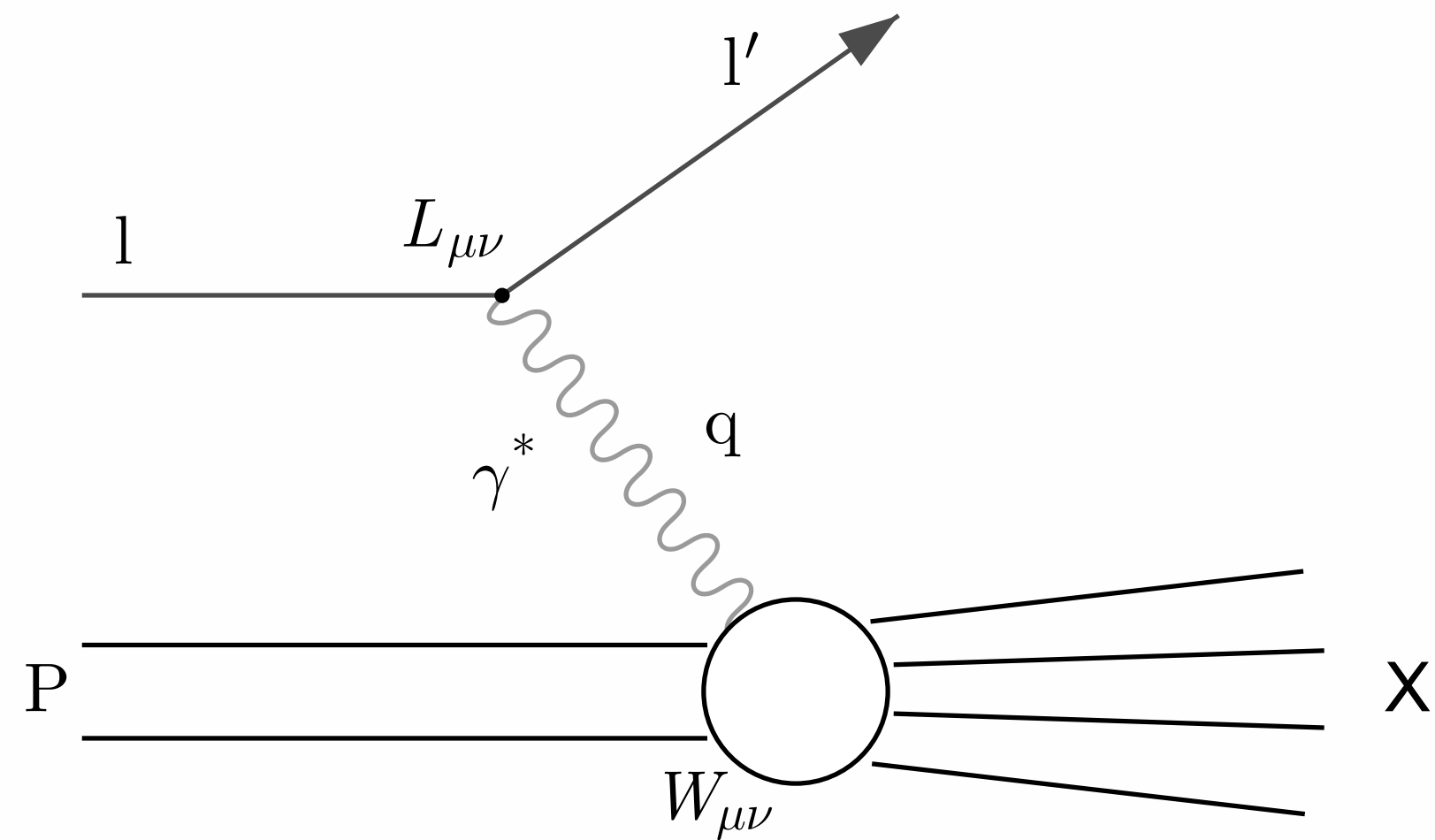
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$$\hat{P}_T \equiv P_T - \langle P_T \rangle$$

$$\hat{x}_F \equiv x_F - \langle x_F \rangle$$

DIS variables



θ	polar scattering angle in the laboratory frame
$\mathbf{q} = \mathbf{l} - \mathbf{l}'$	four-momentum transfer to the target
$s = (\mathbf{P} + \mathbf{l})^2 \stackrel{\text{lab}}{\approx} M^2 + 2ME$	squared centre-of-mass energy
$\nu = \frac{\mathbf{P} \cdot \mathbf{q}}{M} \stackrel{\text{lab}}{=} E - E'$	energy transfer to the target
$y = \frac{\mathbf{P} \cdot \mathbf{q}}{\mathbf{P} \cdot \mathbf{l}} \stackrel{\text{lab}}{=} \frac{\nu}{E}$	fractional energy transfer to the target
$Q^2 = -\mathbf{q}^2 \stackrel{\text{lab}}{\approx} 4EE' \sin^2 \frac{\theta}{2}$	squared invariant mass of the virtual photon
$W^2 = (\mathbf{P} + \mathbf{q})^2$ $\stackrel{\text{lab}}{=} M^2 + 2M\nu - Q^2$	squared mass of the final state
$x = \frac{Q^2}{2\mathbf{P} \cdot \mathbf{q}} \stackrel{\text{lab}}{=} \frac{Q^2}{2M\nu}$	Bjorken scaling variable

$$d\sigma_{UT}^{ep \rightarrow ehX} \sim |\mathbf{S}_\perp| \left[\sin(\phi_h - \phi_S) f_{1T}^\perp \otimes D_1 + \sin(\phi_h + \phi_S) h_1 \otimes H_1^\perp + \dots \right]$$

SIVERS
COLLINS (1993)

SIVERS (1990)

An asymmetric distribution of unpolarized quarks inside a transversely polarized proton leads to an asymmetric distribution of the outgoing hadrons

$$\mathbf{S} \cdot (\hat{\mathbf{P}} \times \hat{\mathbf{k}}_\perp)$$

This happens via a correlation between the spin of the proton, its momentum and the (intrinsic) transverse momentum of the quarks.

COLLINS (1993)

Spin transfer from a transversely polarized quark, fragmenting to an unpolarized hadron should lead to a transverse SSA

$$(\hat{\mathbf{p}}_q \times \hat{\mathbf{p}}_{h\perp}) \cdot \mathbf{s}_q$$

This happens via a correlation between the spin of the fragmenting quark, its momentum and the transverse momentum of the produced hadron.