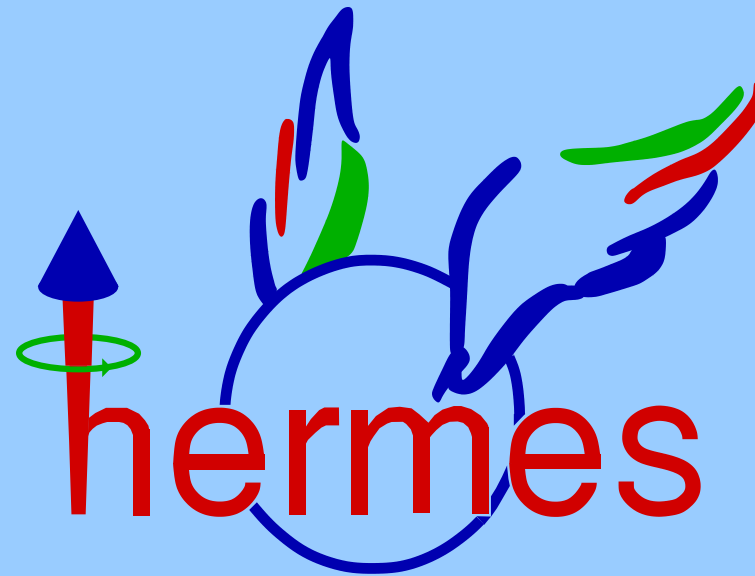


Exclusive Vector Meson Production at HERMES

Armine Rostomyan

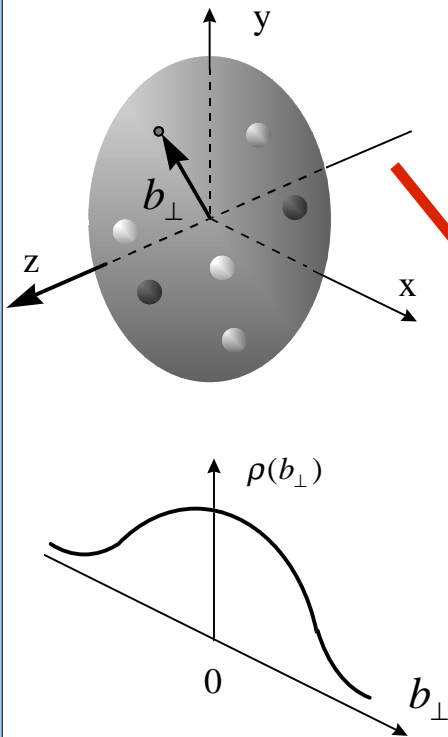


München, DPG 06
22.03.2006

Generalized information

$$l N \rightarrow l' N'$$

Form factor



Proton form factors,
transverse localization
of partons

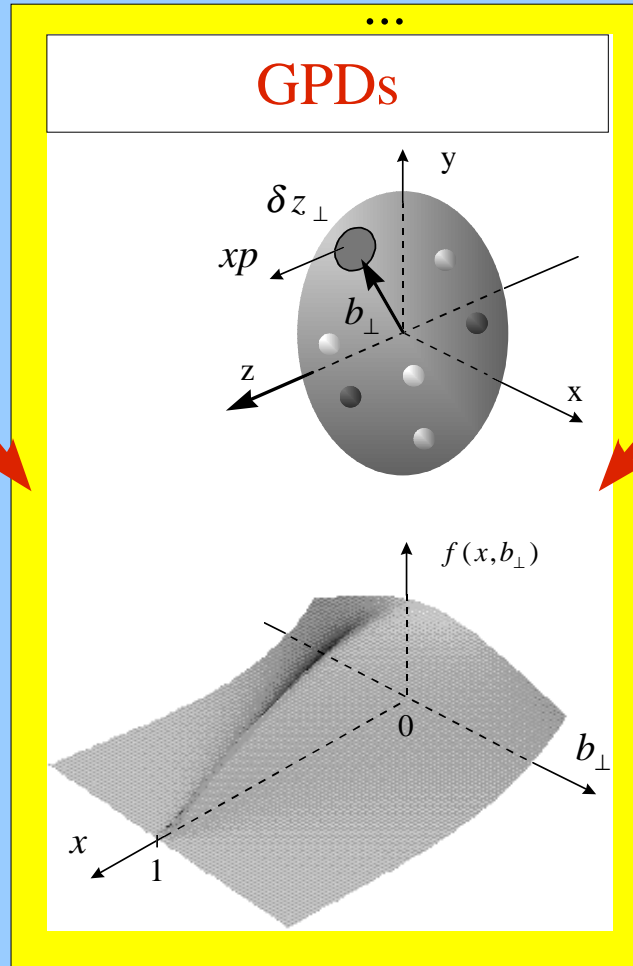
$$\gamma$$

$$l N \rightarrow l' \rho N'$$

$$\pi$$

...

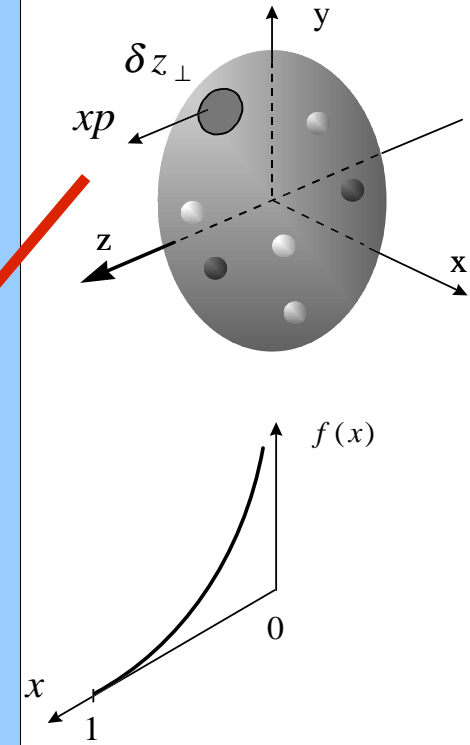
GPDs



Correlated quark momentum and helicity
distribution in transverse space

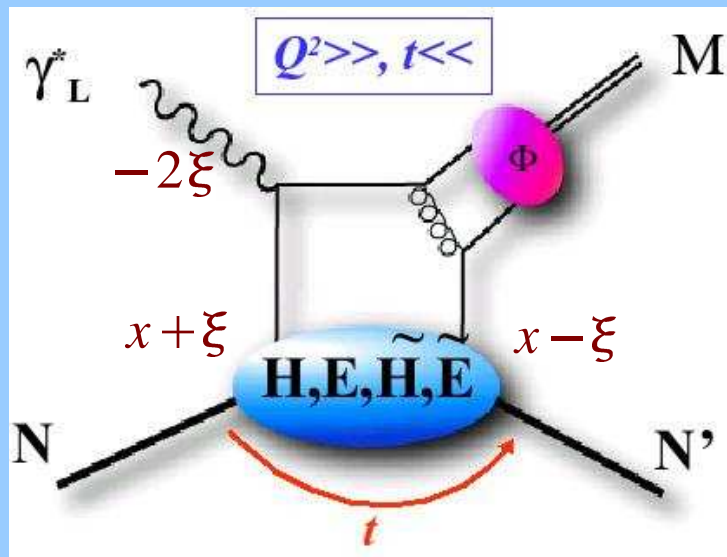
$$l N \rightarrow l' X$$

Parton density

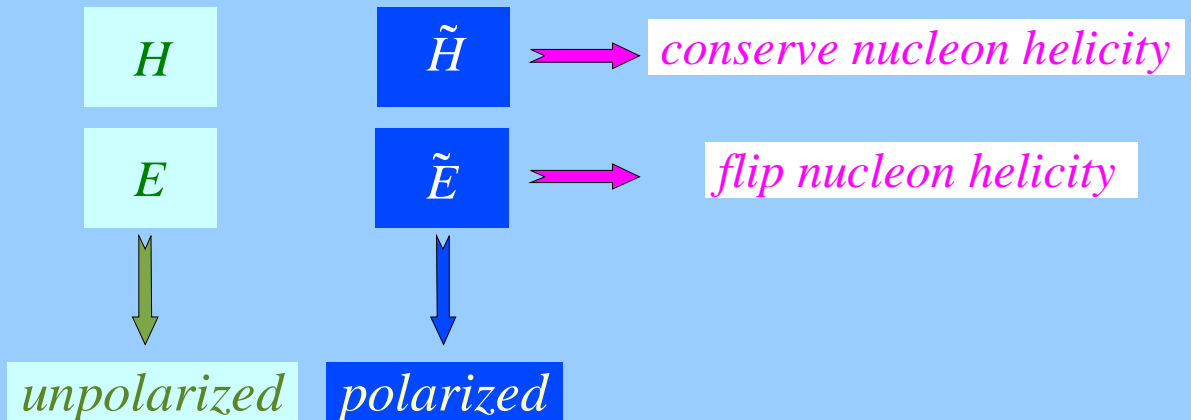


Structure functions,
longitudinal quark
momentum and
helicity distributions

Probabilistic interpretation of GPDs



4 GPDs defined for each quark flavour:



Described by 3 variables: x, ξ, t

$x + \xi$ longitudinal momentum fraction of the quark

-2ξ exchanged longitudinal momentum fraction

t squared momentum transfer

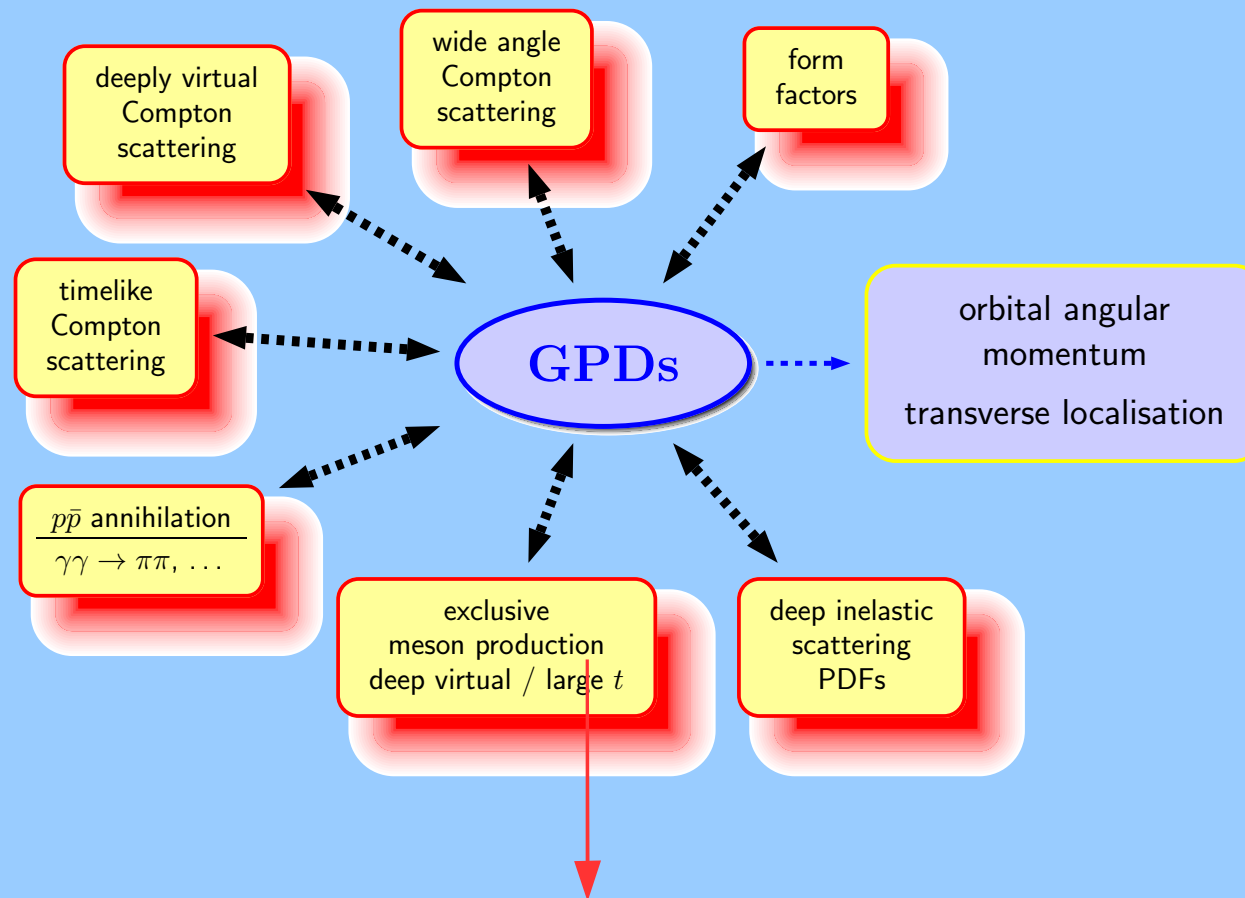
$$x \neq x_B$$

$$\xi = \frac{x_B}{2 - x_B}$$

$$-t = -(N - N')^2$$

GPDs = probability amplitude for N to emit a parton $x + \xi$
and for N' to absorb it $x - \xi$

Access to GPDs

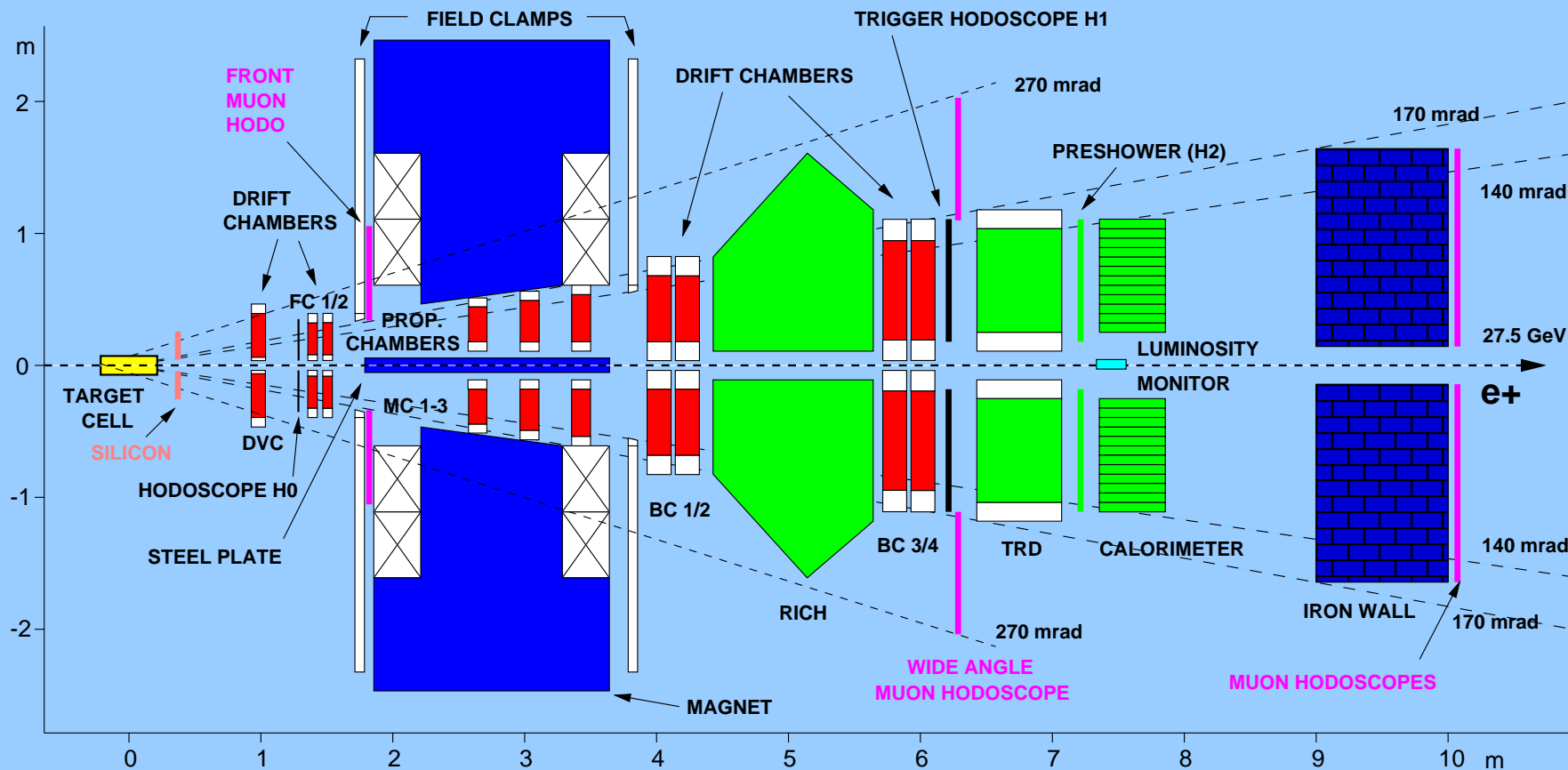


Quantum numbers of final state selects different GPDs

- vector mesons (ρ, ω, ϕ) → unpolarized GPDs: $H E$
- pseudoscalar mesons (π, η) → polarized GPDs: $\tilde{H} \tilde{E}$

Factorization for longitudinal photons only

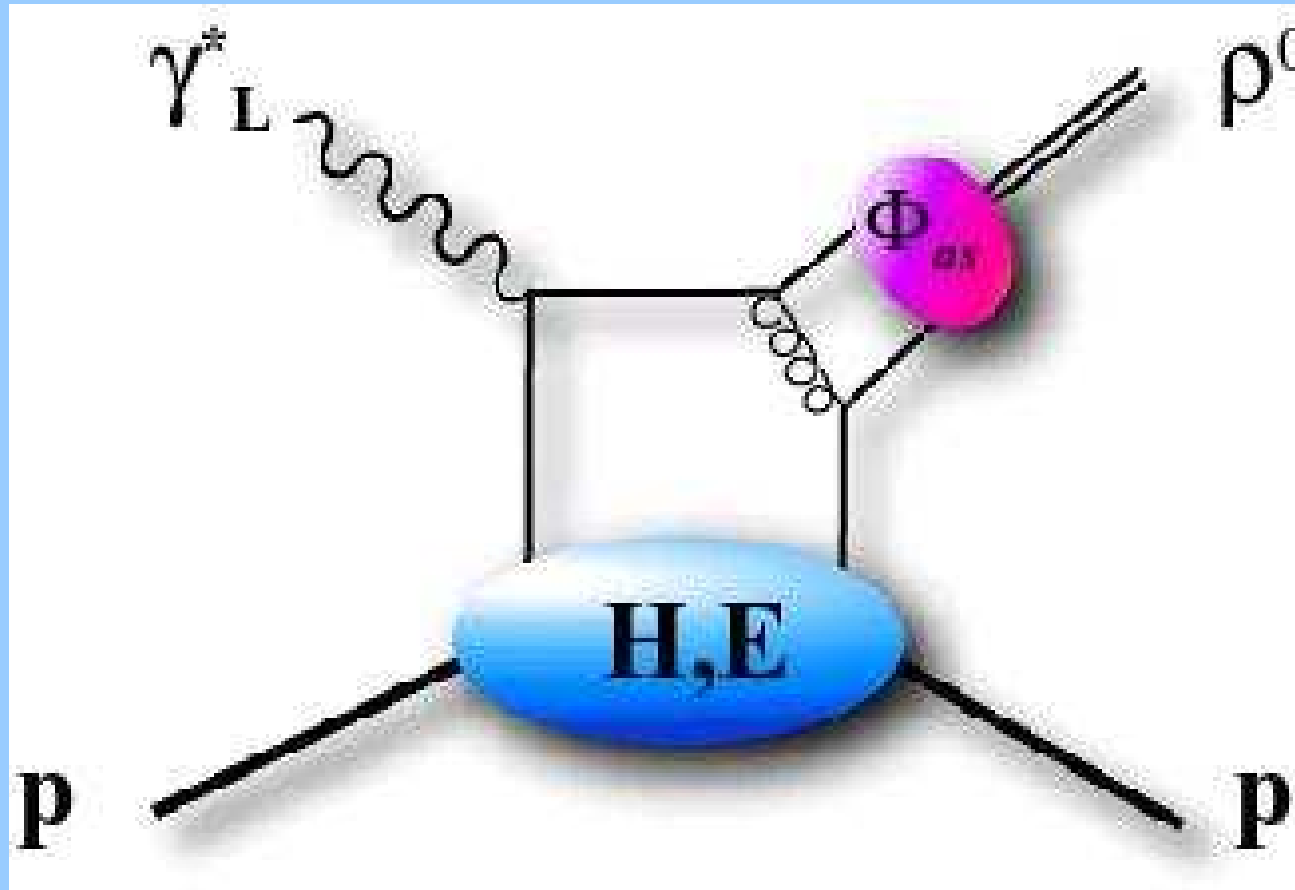
The Spectrometer



- fixed target experiment
- forward spectrometer
- no recoil detection

- **Acceptance:** $40 \text{ mrad} < \theta < 140 \text{ mrad}$
- **Tracking:** $\delta P_e / P_e < 2\%$, $\delta \theta < 0.6 \text{ mrad}$
- **Particle identification:**
- ♦ **TRD, Calo, Preshower:** hadron/lepton separation:
 $\epsilon_e > 99\%$
- ♦ **Rich:** hadron identification (p, π, K)

Cross-section of Exclusive Vector Mesons



Exclusive Vector Meson Selection $ep \rightarrow e'V(p)$

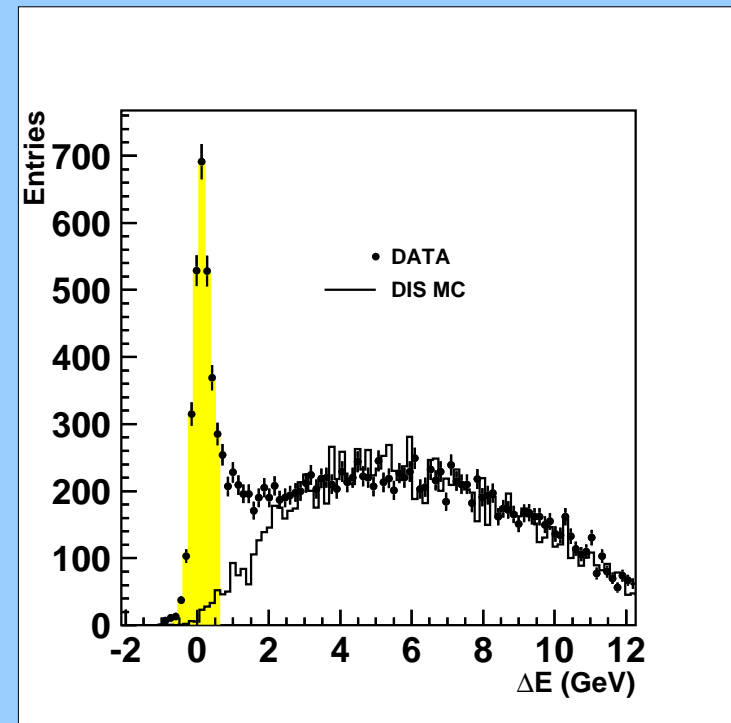
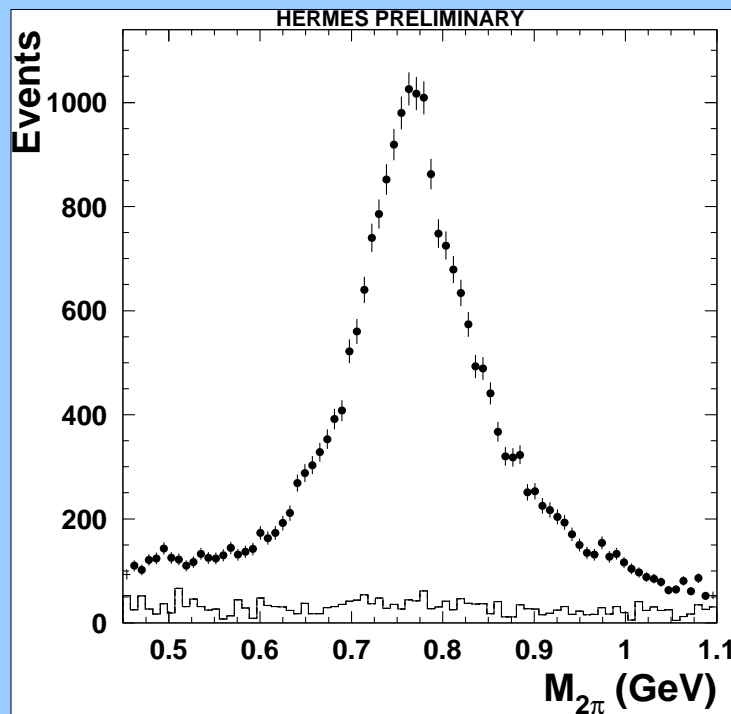
$$\rho^0 \rightarrow \pi^+ \pi^-$$

$$\phi \rightarrow K^+ K^-$$

- no recoil detection
- exclusive ρ^0 and ϕ reaction through the energy and momentum transfer:

$$\Delta E = \frac{M_x^2 - M_p^2}{2 M_p}$$

$$t' = t - t_0$$



σ_L/σ_T Separation

- GPD calculations only for longitudinal component of the cross-section (σ_L):

$$\sigma_L = \frac{R}{1 + \varepsilon R} \sigma_{\gamma^+ p \rightarrow V p}$$

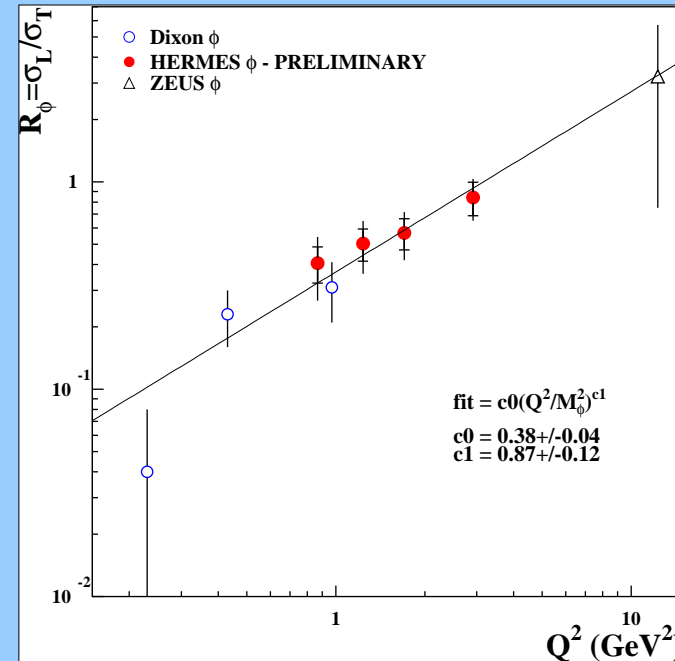
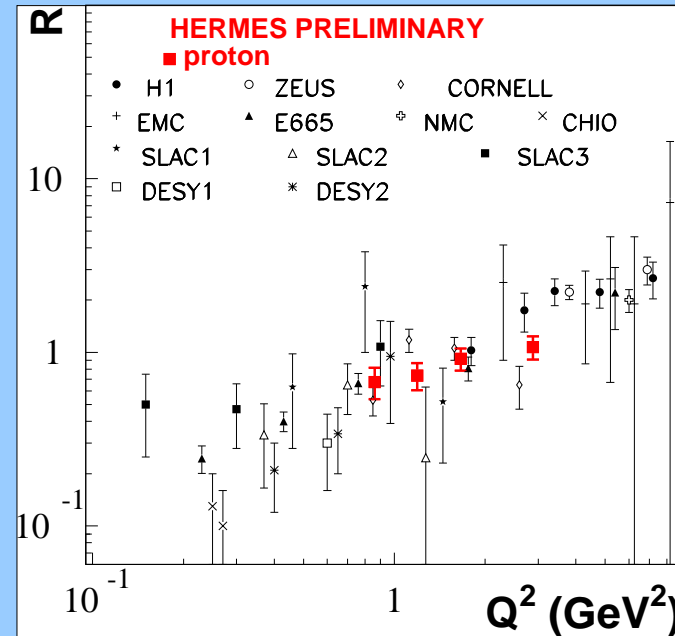
$$R = \frac{\sigma_L}{\sigma_T}$$

where ε - polarization of γ^*

- Assuming SCHC

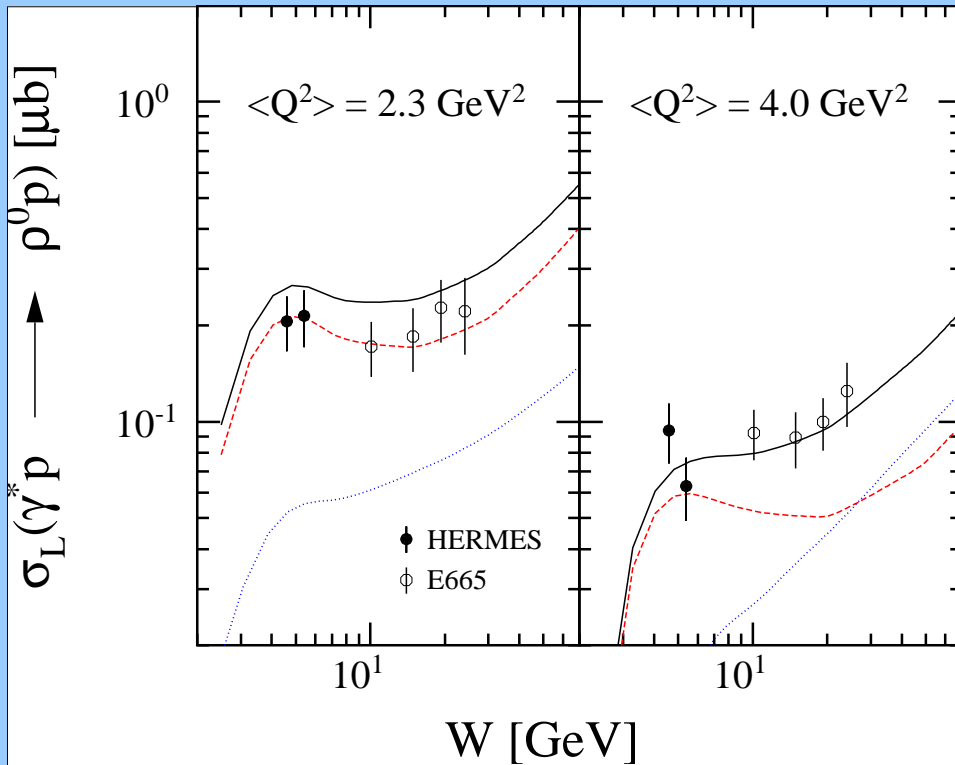
$$R = \frac{1}{\varepsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}}$$

$$r_{00}^{04} \rightarrow W(\cos \theta)$$



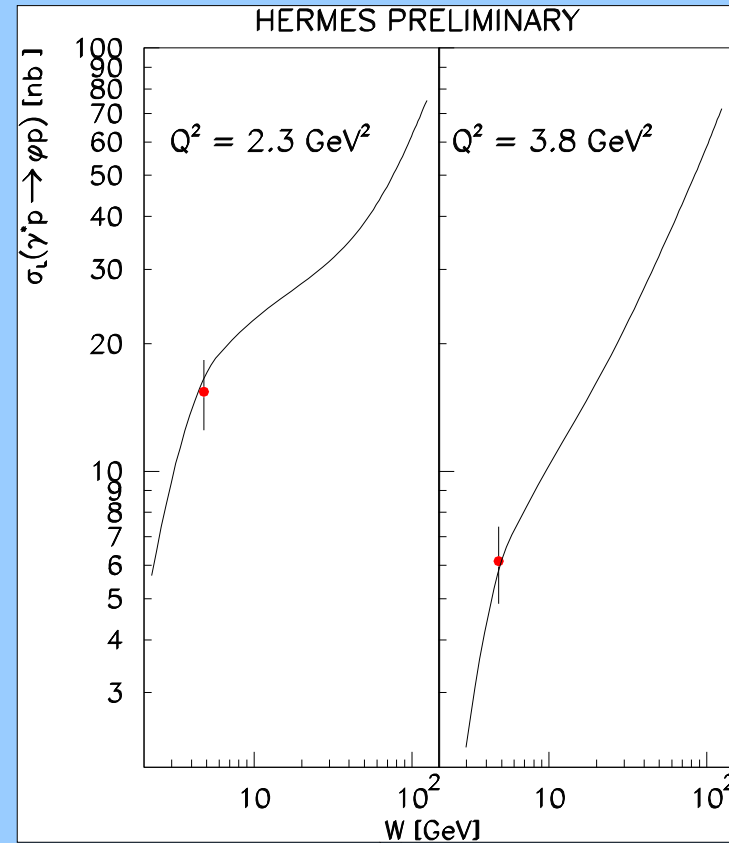
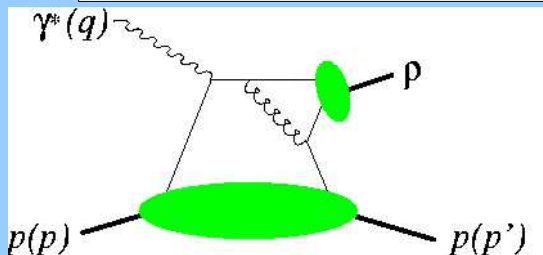
$\sigma_L^{\gamma^* p \rightarrow \rho^0 p}$ and $\sigma_L^{\gamma^* p \rightarrow \phi p}$

Vanderhaeghen, Guichon, Guidal (1999)

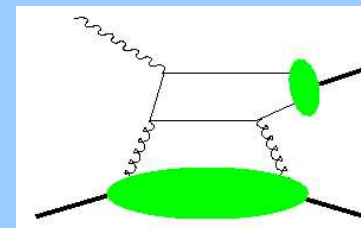


- quarks and gluons: at the same order of α_s
- gluon GPDs can be accessed

dominated by quark exchange

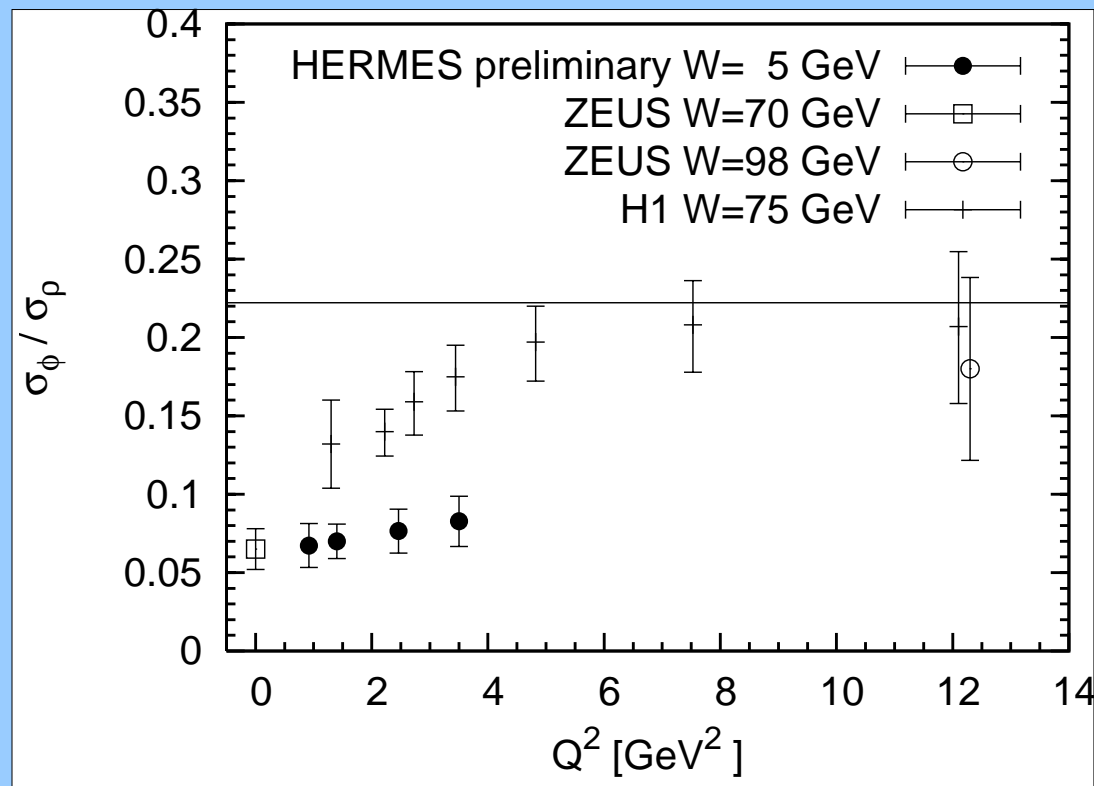


gluon exchange



$$\sigma \gamma^* p \rightarrow \phi p / \sigma \gamma^* p \rightarrow \rho^0 p$$

Diehl, Vinnikov (2005)



$$\sigma_{\phi} / \sigma_{\rho} \approx \frac{2}{9} \frac{|\tau_g|^2}{|\tau_q|^2 + 2|\tau_q||\tau_g|\cos\alpha + |\tau_g|^2}$$

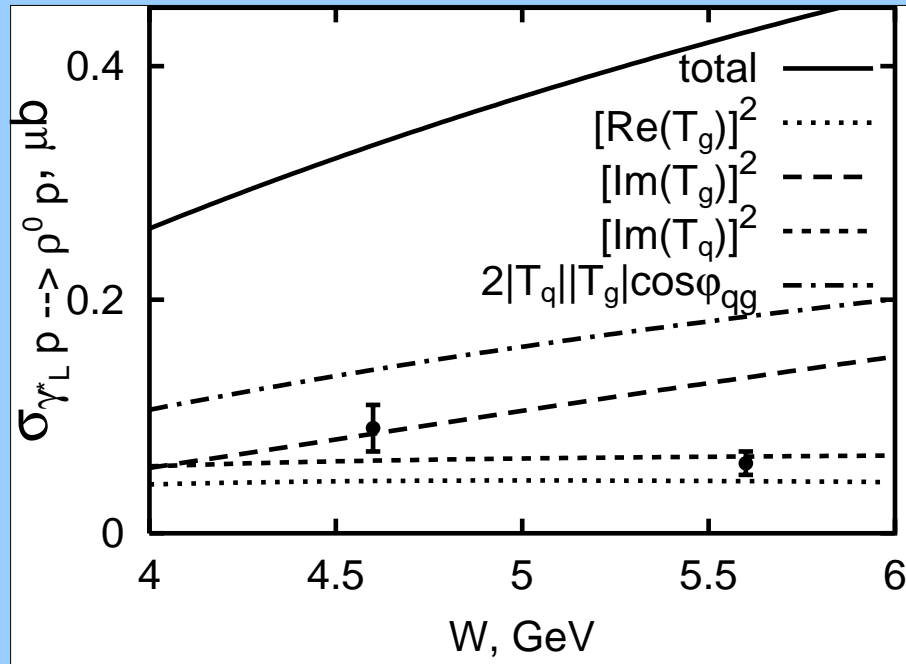
$$0.38 < |\tau_g / \tau_q| < 1.5$$

gluon contribution and quark-gluon interference can not be neglected

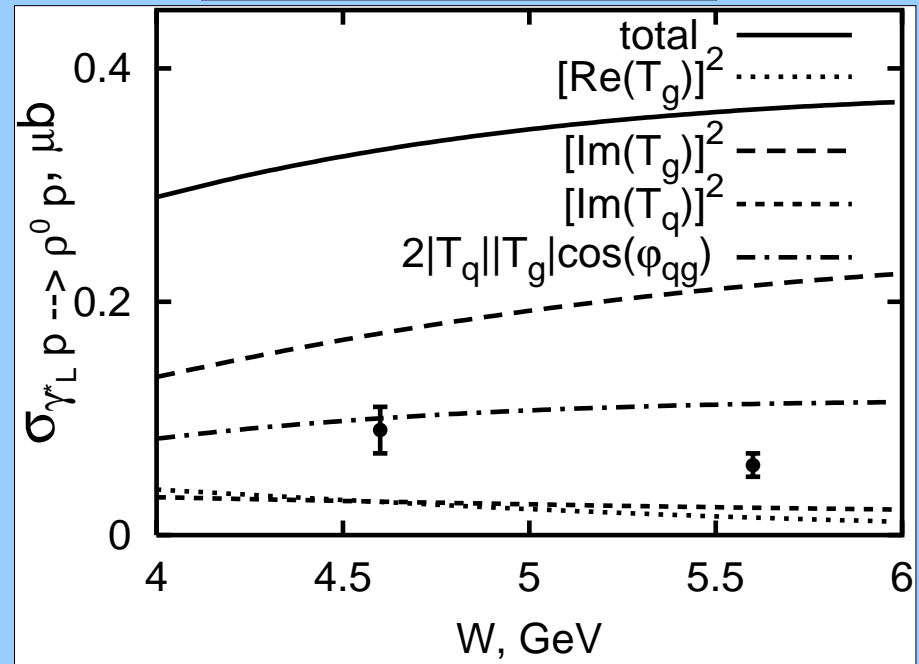
$$\sigma_{L}^{\gamma^* p \rightarrow \rho^0 p}$$

parametrizations for: H^q, H^g, E^q , no hint for E^g

factorizes GPD model



Regge GPD model



Ellinghaus, Nowak, Vinnikov, Ye(2005)

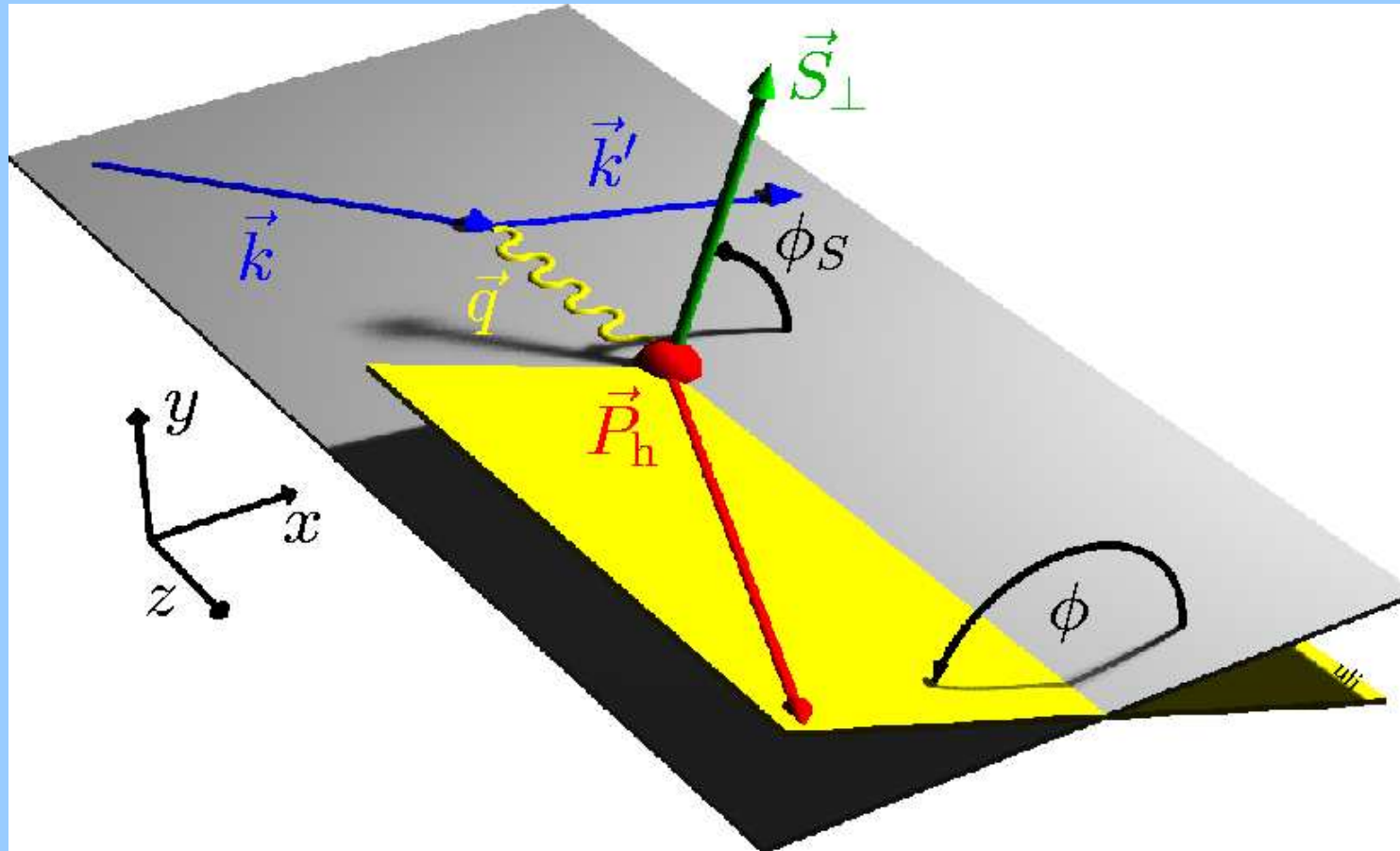
Theory

- the calculation overshoots the experimental data
- k_{\perp} is not taken into account
- quark and gluon amplitudes have to be scaled down in a similar proportion, a factor of 5 suppression of the cross section

Experiment

- provide new results with higher statistics

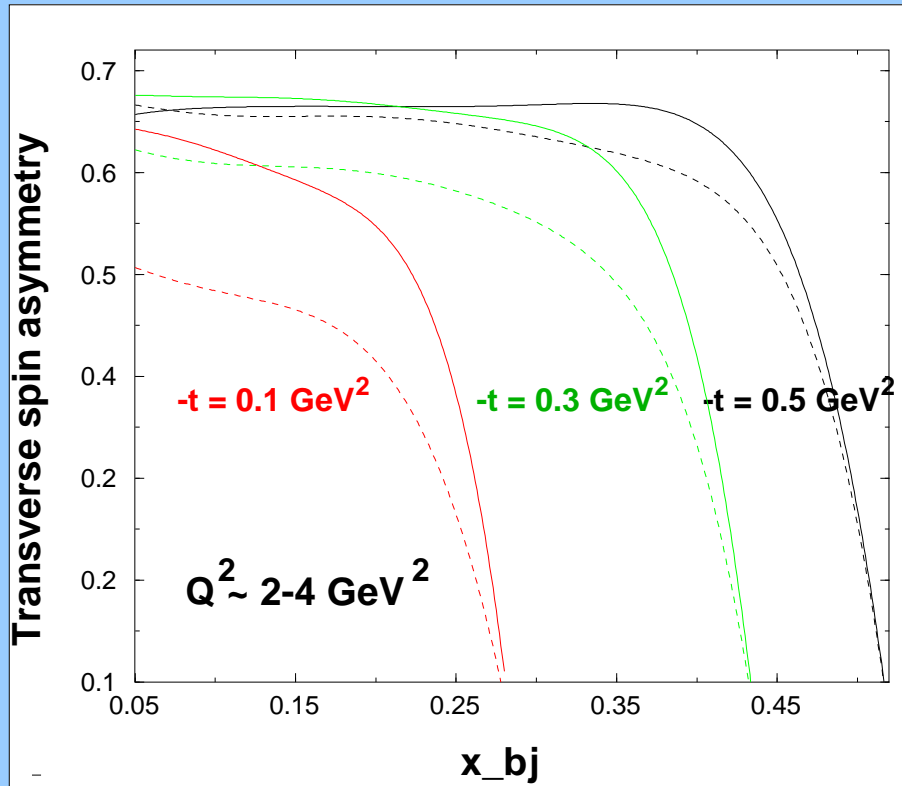
Transverse Target Spin Asymmetries (TTSA)



TTSA of Exclusive π^+ and ρ^0

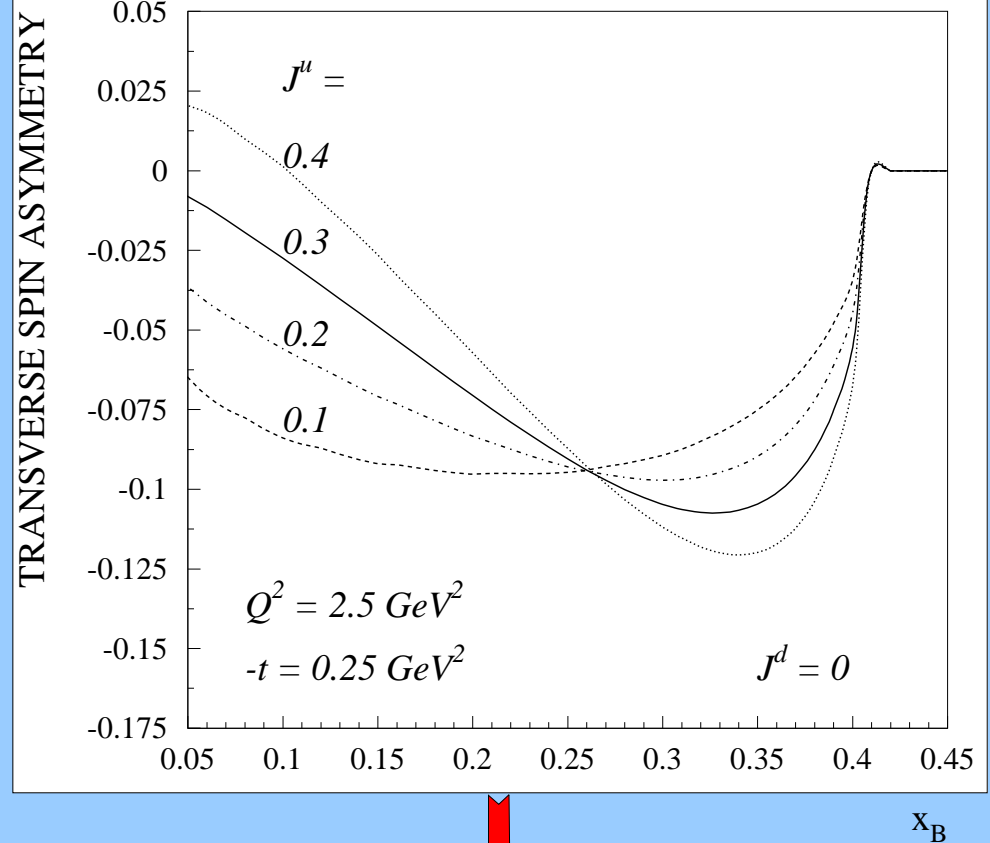
Frankfurt, Polyakov, Strikman, Vanderhaeghen (2000)

$$e p \rightarrow e' \pi^+ n$$



Goetze, Polyakov, Vanderhaeghen (2001)

$$\gamma_L^* + p \rightarrow \rho_L^0 + p$$



$$|S_T| \sin(\phi - \phi_s) \tilde{E} \tilde{H}$$

- linear dependence on GPDs
- higher order corrections cancel

$$|S_T| \sin(\phi - \phi_s) E H$$

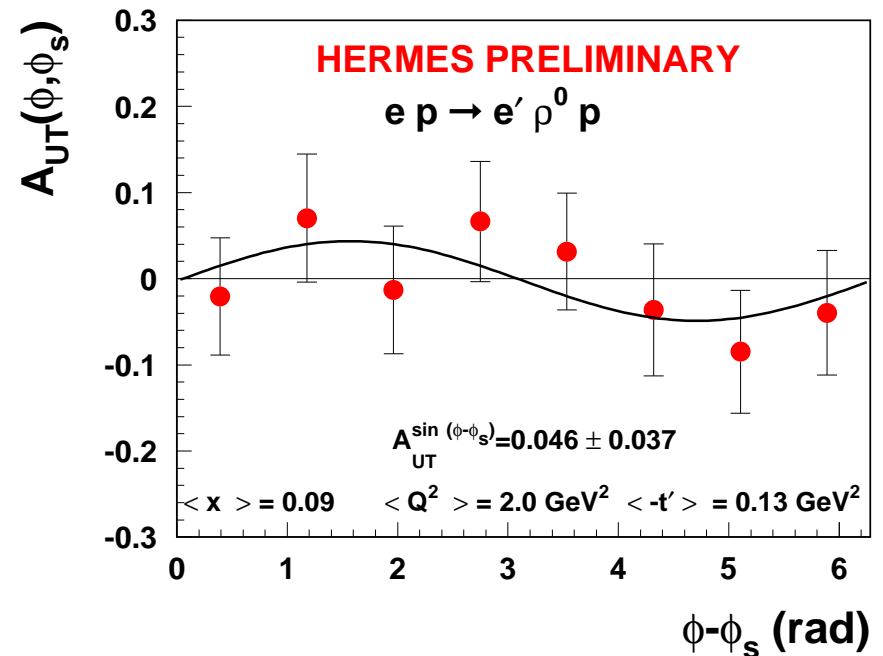
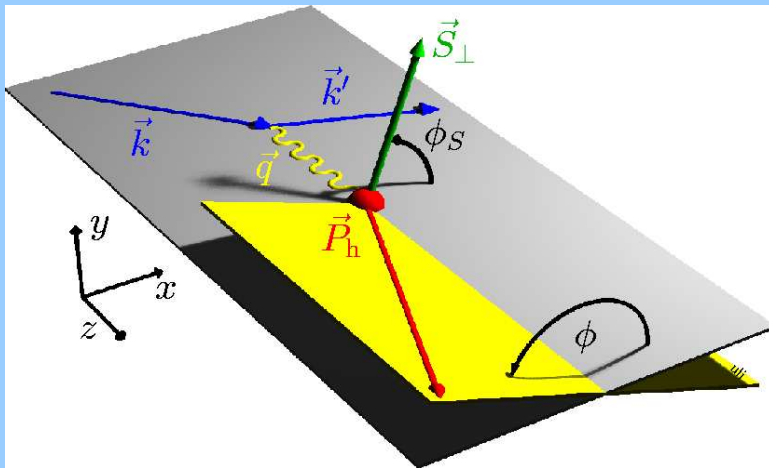
- E is kinematically not suppressed
- TTSA promising observable which allow an access to E
- E related to J^q through Ji sum rule

TTSA of Exclusive ρ^0

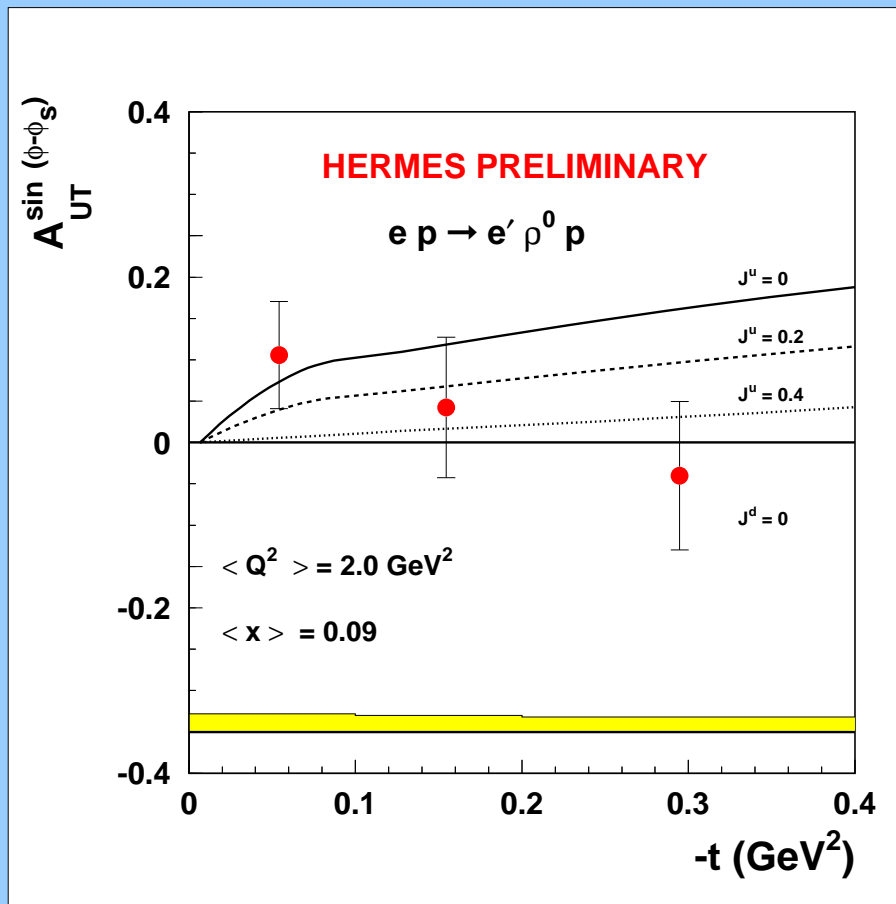
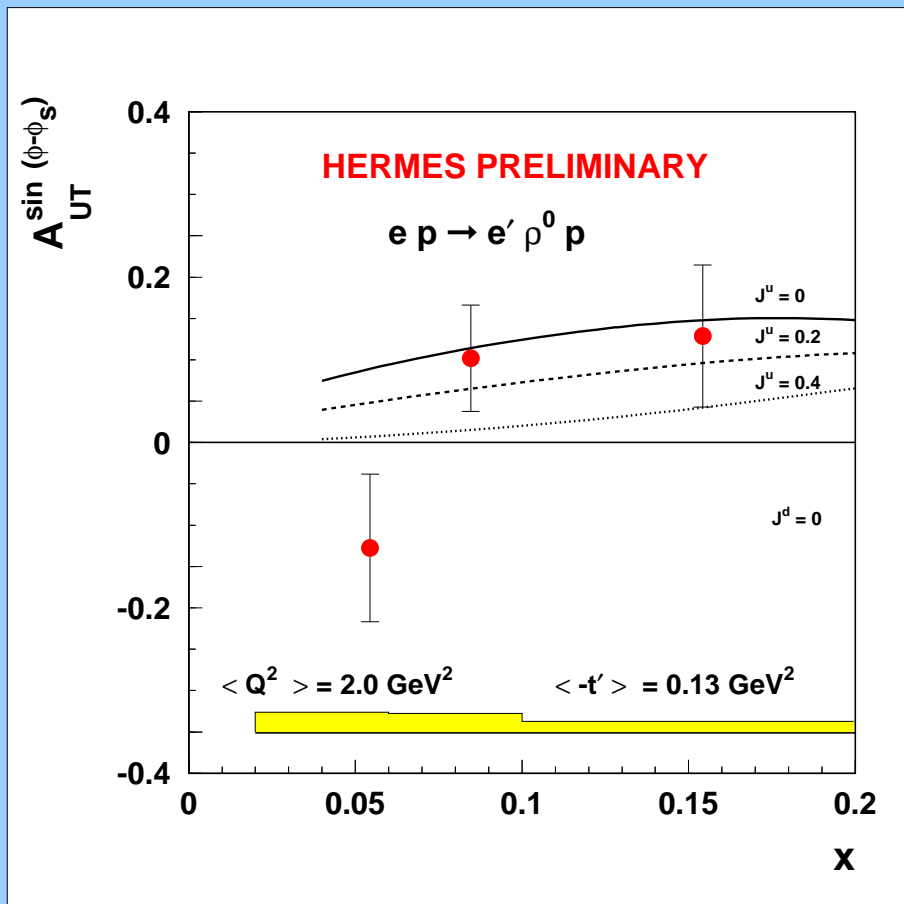
$$A_{UT} = -\frac{\pi}{2} A_{theor}$$

$$A_{UT}(\phi - \phi_s) = \frac{1}{|P|} \frac{N^\uparrow(\phi - \phi_s) - N^\downarrow(\phi - \phi_s)}{N^\uparrow(\phi - \phi_s) + N^\downarrow(\phi - \phi_s)}$$

$$A_{UT}(\phi - \phi_s) = A_{UT}^{\sin(\phi - \phi_s)} \cdot \sin(\phi - \phi_s) + const$$



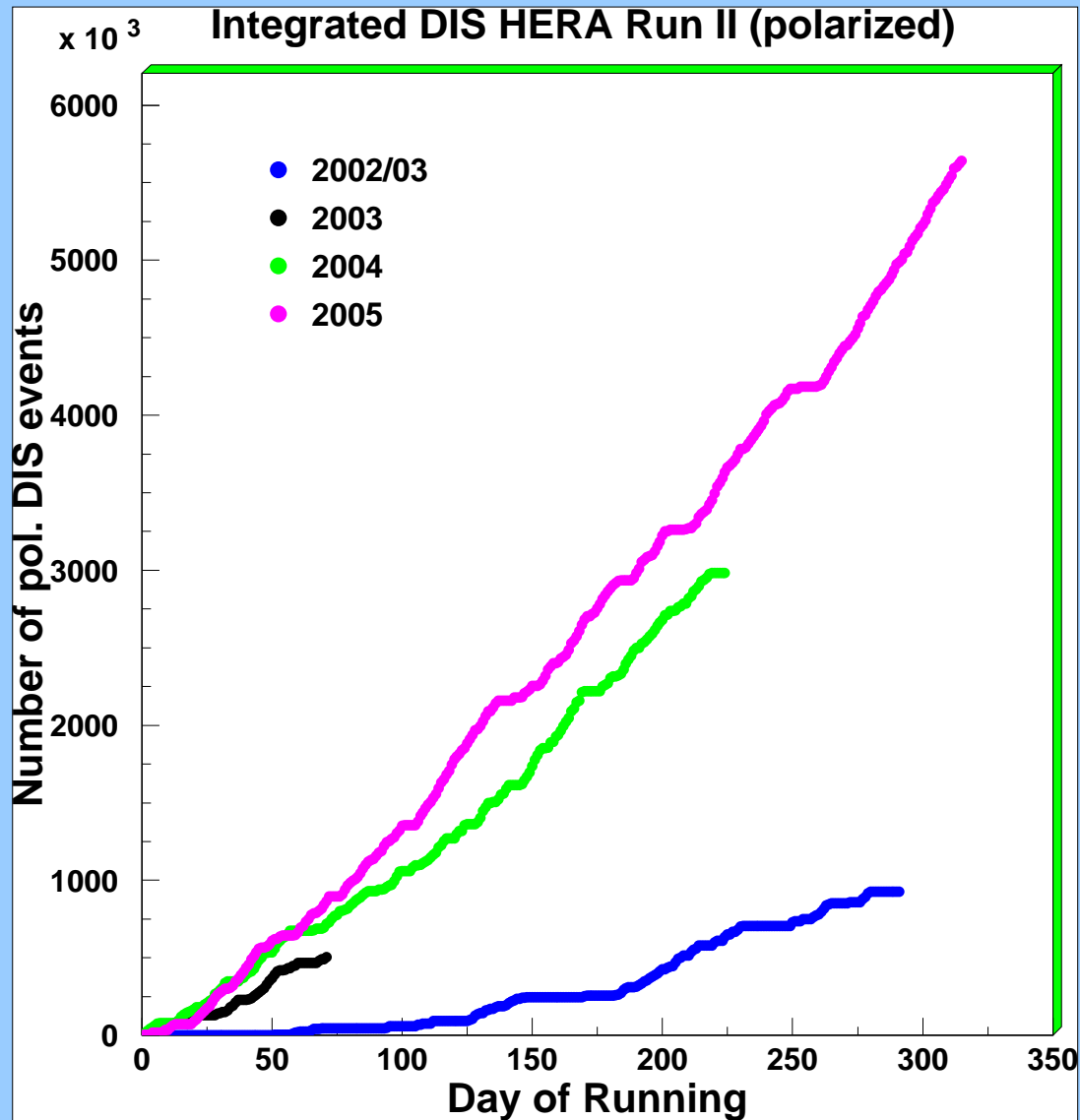
Kinematic Dependences



- L/T separation has not yet been done
- transverse component is suppressed at high Q^2
- within the statistical errors in agreement with theoretical calculations
- the statistics is not yet enough to make a statement about J^u
- more data is available

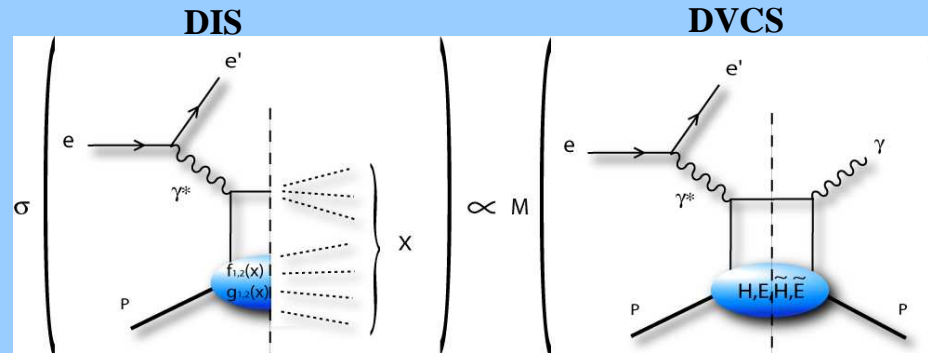
*indication of
 small E_g*

More Results are Coming...



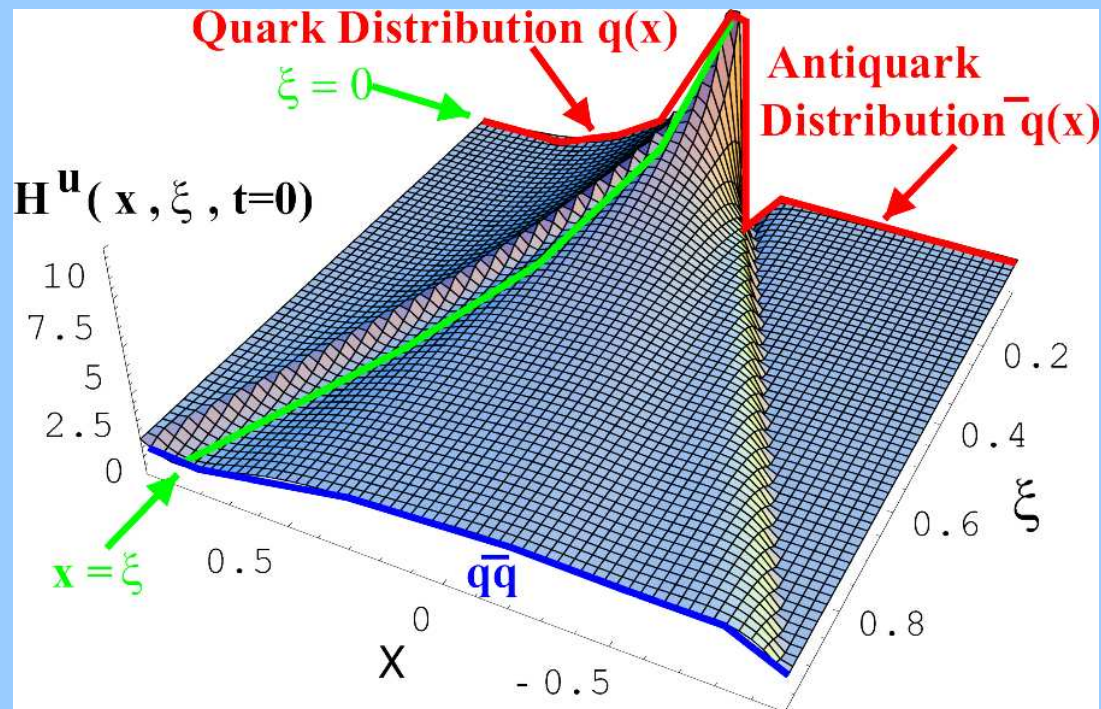
BACKUP: GPD and DIS

Forward limit ($t \rightarrow 0, \xi \rightarrow 0$)



$$H^q(x, \xi=0, t=0) = \mathbf{q}(x)$$

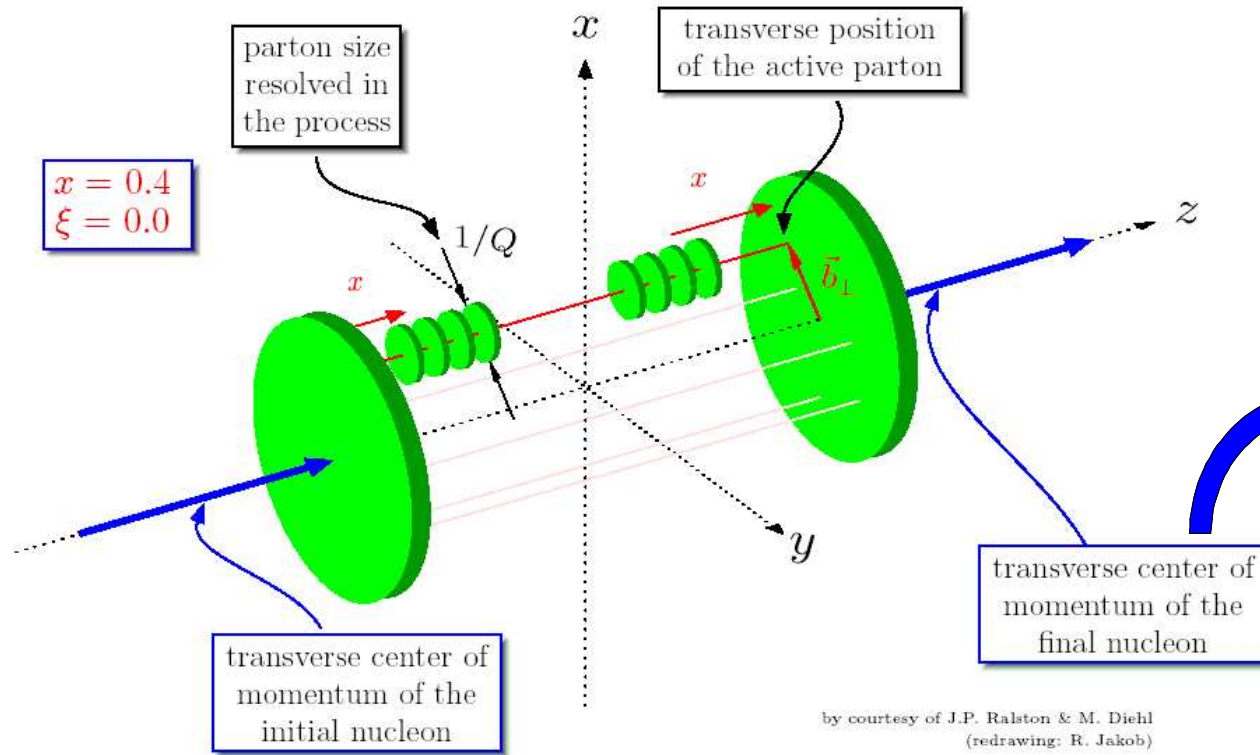
$$\tilde{H}^q(x, \xi=0, t=0) = \Delta \mathbf{q}(x)$$



BACKUP: Geometrical interpretation of GPDs

$$q(x, \vec{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} H(x, \xi=0, -\Delta_\perp^2) e^{i\vec{b}_\perp \vec{\Delta}_\perp}$$

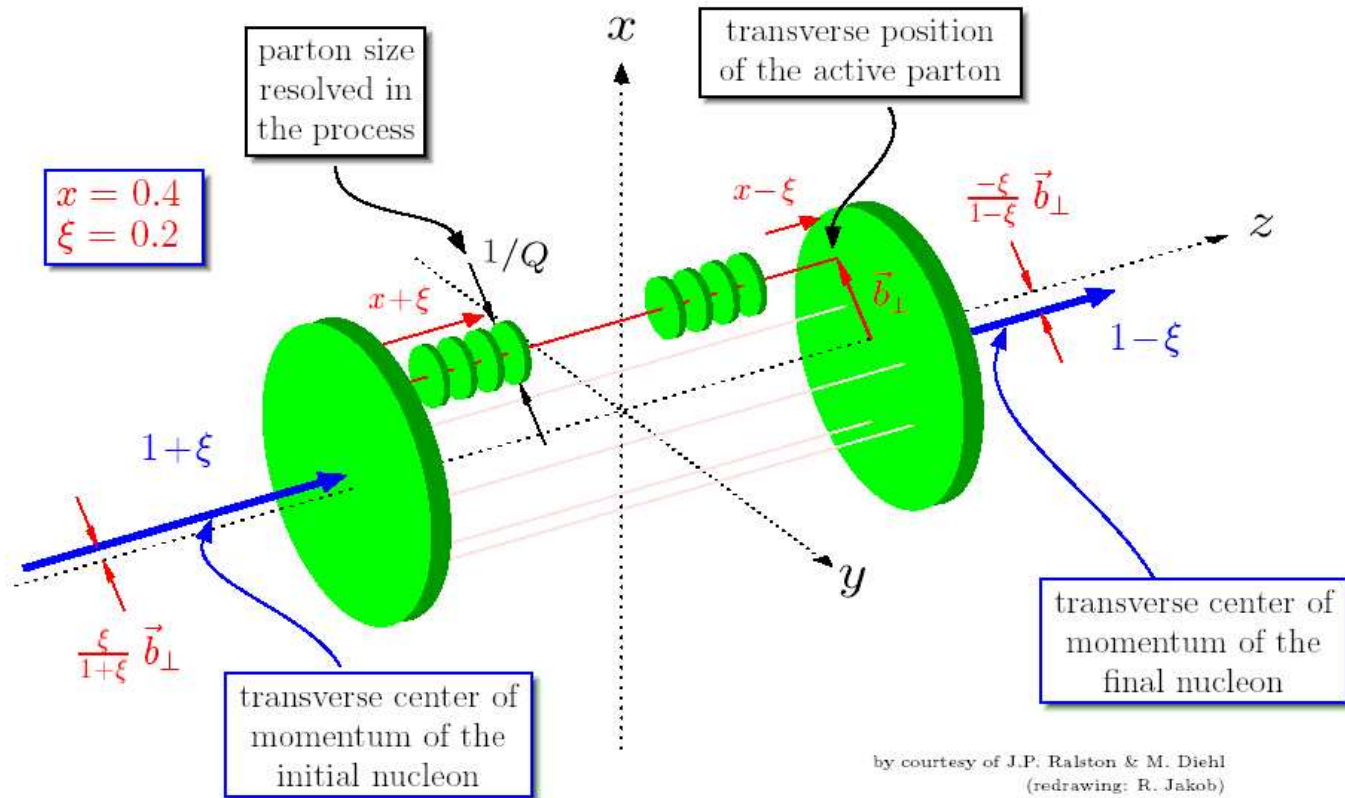
$$\int q(x, b_\perp) d^2 b_\perp = q(x)$$



$$\vec{b}_\perp = \sum_{i=1}^N x_i \vec{b}_{\perp i}$$

BACKUP: Geometrical interpretation of GPDs

the general case: $\xi \neq 0$, by M. Diehl



BACKUP: Limiting cases and sum rules

Forward limit ($t \rightarrow 0, \xi \rightarrow 0$)

$$\begin{array}{ll}
 H^q(x, 0, 0) = q(x) & \tilde{H}^q(x, 0, 0) = \Delta q(x) \quad \text{for } x > 0 \\
 H^q(x, 0, 0) = -\bar{q}(-x) & \tilde{H}^q(x, 0, 0) = \Delta \bar{q}(-x) \quad \text{for } x < 0 \\
 H^g(x, 0, 0) = xg(x) & \tilde{H}^g(x, 0, 0) = x\Delta g(x) \quad \text{for } x > 0
 \end{array}$$

$E^q, \tilde{E}^q, E^g, \tilde{E}^g$



- no corresponding relations
- are visible only in exclusive processes

Sum rules

$$\begin{array}{ll}
 \int_{-1}^{+1} H^q(x, \xi, t) dx = F_1^q(t) & \int_{-1}^{+1} E^q(x, \xi, t) dx = F_2^q(t) \\
 \int_{-1}^{+1} \tilde{H}^q(x, \xi, t) dx = g_A^q(t) & \int_{-1}^{+1} \tilde{E}^q(x, \xi, t) dx = h_A^q(t)
 \end{array}$$

Ji sum rule

$$\int_{-1}^{+1} (H(x, \xi, t=0) + E(x, \xi, t=0) x) dx = \frac{1}{2} \Delta \Sigma + L_q$$

30% (DIS)