



### Spin Density Matrix Elements in exclusive production of $\omega$ mesons at Hermes

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- HERMES experiment and data processing
- SDMEs, helicity amplitudes and angular distribution
- Results.
  - SDMEs at average kinematics
  - Unnatural-Parity Exchange for  $\omega$  meson
  - $\checkmark$  Longitudinal to Transverse cross section ratio for  $\omega$  meson
- Summary



#### Hermes Detector was Two Identical Halves of Forward Spectrometer



- Beam  $e^{\pm}$ , P = 27.56 GeV/c longitudinal polarization  $\sim$  55 %.
- Target longitudinally, transversely polarized H or D or unpolarized gas target.
- Acceptance:  $|\Theta_x| < 170$  mrad,  $40 < |\Theta_y| < 140$  mrad.
- **P** Resolution  $\delta P/P \leq 1\%$ ,  $\delta \Theta \leq 0.6$  mrad.
- PID: RICH, TRD, Preshower, Calorimeter.



## **Exclusive** $\omega$ -meson production at HERMES



 $e(k) + N(p) \rightarrow e'(k') + N(p') + \omega \rightarrow (\pi^+ \pi^- \pi^0 (\rightarrow 2\gamma))$  $e \to e' + \gamma^*$  (QED).  $Q^2 = -q^2 = -(k - k')^2 = 1.0 \div 10. \text{ GeV}^2, \langle Q^2 \rangle = 1.9 \text{ GeV}^2$  $I = \sqrt{(q+p)^2} = 3.0 \div 6.3 \text{ GeV}, \langle W \rangle = 4.8 \text{ GeV}$  $I R_B = \frac{Q^2}{2na} = 0.01 \div 0.35, \langle x_B \rangle = 0.08$ ■  $t' = t - t_{min}, 0 \le -t' \le 0.2 \text{ GeV}^2, \langle -t' \rangle = 0.08 \text{ GeV}^2$  $t = (p - p')^2 = (q - v)^2$ Number of  $\omega$  events Hydrogen: -2260, Deuterium: -1332  $D \Delta E = \frac{M_X^2 - M_p^2}{2M_p} \text{ with } M_X^2 = (p + q - p_{\pi^+} - p_{\pi^-} - p_{\pi^0})^2$ and  $M_X$  being missing mass, p, q,  $p_{\pi^+}$  ,  $p_{\pi^-}$   $p_{\pi^0}$  are 4-momenta of proton,  $\gamma^*$  and pions. Exclusive process  $\Delta E = 0$ 

SIDIS background( $\approx 20\%$ ) is subtracted using MC PYTHIA



#### Angular distribution in reaction

 $e + p \rightarrow e' + p' + \omega \rightarrow (\pi^+ \pi^- \pi^0 (\rightarrow 2\gamma))$ 



 $\omega$ -production-plane

$$\begin{split} \omega &\Rightarrow \pi^{+}\pi^{-}\pi^{0} \text{ (conservation of } \vec{J} \text{ )} \\ |\omega; 1m > \rightarrow |\pi^{+}\pi^{-}\pi^{0}; 1m > \Rightarrow Y_{1m}(cos(\theta), \phi) \\ (m = \pm 1, 0). \text{ Angular distribution} \\ \mathcal{W}(r^{\alpha}_{\lambda_{V}\lambda'_{V}}, \Phi, \phi_{n}, \cos \Theta_{n}) \text{ depends linearly on} \\ \text{Spin Density Matrix Elements(SDMEs)} - r^{\alpha}_{\lambda_{V}\lambda'_{V}} \\ \text{and beam polarization } P_{b}. \end{split}$$

SDMEs are bilinear combination of helicity amplitudes.

- For longitudinally polarized beam and unpolarized target there are 23 SDMEs, (15 unpolarized and 8 polarized)
- The SDMEs are determined from the fit of angular distribution of pions from decay  $\omega \Rightarrow \pi^+\pi^-\pi^0$ , by angular distribution  $\mathcal{W}(r^{\alpha}_{\lambda_V\lambda'_V}, \Phi, \phi_n, \cos \Theta_n)$ , with Maximum Likelihood method



 $(\text{SDMEs}) - r^{\alpha}_{\lambda_V \lambda'_V} \text{ are expressed by helicity amplitudes } F_{\lambda_V \lambda'_N; \lambda_\gamma \lambda_N}(W, Q^2, t')$ 

In CM frame of  $\gamma^* N$  they are given by the von Neumann formula:  $r^{\alpha}_{\lambda_V \lambda'_V} \sim \rho_{\lambda_V \lambda'_V} = \frac{1}{2N} \sum_{\lambda_\gamma \lambda'_\gamma \lambda_N \lambda'_N} F_{\lambda_V \lambda'_N} \rho_{\lambda_\gamma \lambda'_\gamma} F^*_{\lambda'_V \lambda'_N} \rho_{\lambda_\gamma \lambda'_\gamma} + \rho_{\lambda_V \lambda'_V} \rho_{\lambda_\gamma \lambda'_V} + \rho_{\lambda_V \lambda'_V} \rho_{\lambda_\gamma \lambda'_V} \rho_{\lambda_\gamma \lambda'_V} + \rho_{\lambda_V \lambda'_V} \rho_{\lambda_\gamma \lambda'_V} \rho_{\lambda_\gamma \lambda'_V} \rho_{\lambda_\gamma \lambda'_V} + \rho_{\lambda_V \lambda'_V} \rho_{\lambda_\gamma \lambda'_V} \rho_{\lambda_\gamma \lambda'_V} \rho_{\lambda_\gamma \lambda'_V} + \rho_{\lambda_V \lambda'_V} \rho_{\lambda_\gamma \lambda'_V}$ 

- On unpolarized target nucleon-helicity-flip amplitudes are suppresed  $F_{\lambda_V \frac{1}{2} \lambda_\gamma \frac{1}{2}} o F_{\lambda_V \lambda_\gamma}$
- Example:  $r_{1-1}^1 = \frac{1}{2} \sum \{ |T_{11}|^2 + |T_{1-1}|^2 |U_{11}|^2 |U_{1-1}|^2 \} / \mathcal{N}$ T natural-parity exchange (NPE) ( $P = (-1)^J$ )
  U unnatural parity exchange (UPE) ( $P = -(-1)^J$ )
  F=T+U
- Helicity conserving  $T_{00}$ ,  $T_{11}$ ,  $U_{11}$ , helicity non concerving  $T_{01}$ ,  $T_{10}$ ,  $T_{1-1}$ ,  $U_{01}$ ,  $U_{10}$ ,  $U_{1-1}$

The dominance of diagonal transitions is called s-channel helicity conservation (SCHC).



## **SDME of exclusive** $\omega$ production at average kinematics



A,  $\gamma_L^* \to \omega_L$  and  $\gamma_T^* \to \omega_T$ B, Interference:  $\gamma_L^* \to \omega_L \& \gamma_T^* \to \omega_T$ C, Spin Flip:  $\gamma_T^* \to \omega_L$ D, Spin Flip:  $\gamma_L^* \to \omega_T$ E, Spin Flip:  $\gamma_T^* \to \omega_{-T}$ 

- Similar magnitudes for SDMEs on proton and deuteron
- if SCHC holds:

$$r_{1-1}^{1} = -\text{Im } r_{1-1}^{2}$$
  
Re  $r_{10}^{5} = -\text{Im } r_{10}^{6}$   
Im  $r_{10}^{7}$  = Re  $r_{10}^{8}$   
for proton

for proton

 $\begin{aligned} r_{1-1}^{1} + \text{Im } r_{1-1}^{2} =& -0.004 \pm 0.038 \pm 0.017 \\ \text{Re } r_{10}^{5} + \text{Im } r_{10}^{6} =& -0.024 \pm 0.013 \pm 0.003 \\ \text{Im } r_{10}^{7} - \text{Re } r_{10}^{8} =& -0.060 \pm 0.010 \pm 0.044 \\ \text{for deuterium} \end{aligned}$ 

$$\begin{split} r_{1-1}^1 + \mathrm{Im} \; r_{1-1}^2 &= 0.033 \pm 0.049 \pm 0.004 \\ \mathrm{Re} \; r_{10}^5 + \mathrm{Im} \; r_{10}^6 &= 0.001 \pm 0.016 \pm 0.015 \\ \mathrm{Im} \; r_{10}^7 - \mathrm{Re} \; r_{10}^8 &= 0.10 \pm 0.11 \pm 0.17 \end{split}$$



# **SDME in exclusive** $\omega$ production at average kinematics





# **Comparison of SDME in exclusive** $\omega$ and $\rho^0$ production at average kinematics



 $\rho^0$  SDMEs, HERMES, Eur. Phys. J. C62 (09) 659.

$$\begin{split} & \blacktriangle \quad \mathsf{A}, \gamma_L^* \to \omega_L \text{ and } \gamma_T^* \to \omega_T \\ r_{1-1}^1 &= \frac{1}{2} \widetilde{\sum} \{ |T_{11}|^2 + |T_{1-1}|^2 - |U_{11}|^2 - |U_{1-1}|^2 \} / \mathcal{N} , \\ \{ \mathbf{f}_{1-1}^2 \} &= \frac{1}{2} \widetilde{\sum} \{ -|T_{11}|^2 + |T_{1-1}|^2 + |U_{11}|^2 - |U_{1-1}|^2 \} / \mathcal{N} \end{split}$$

$$\begin{split} |U_{11}|^2 + |U_{1-1}|^2 &> |T_{11}|^2 + |T_{1-1}|^2 \text{ for } \omega \text{ meson} \\ |T_{1-1}|^2 + |U_{11}|^2 &> |T_{11}|^2 + |U_{1-1}|^2 \text{ for } \omega \text{ meson} \end{split}$$

Assuming  $|T_{1-1}|^2 \approx |\mathsf{U}_{1-1}|^2$  we get  $|U_{11}|^2 > |T_{11}|^2$  for  $\omega$  meson

#### Test of unnatural-parity exchange for $\omega$



#### meson

### Signal of UPE $u_{1} = 1 - r_{00}^{04} + 2r_{1-1}^{04} - 2r_{11}^{1} - 2r_{1-1}^{1}$ $u_{1} = \sum_{\lambda_{N}\lambda_{N}'} \frac{2\epsilon |U_{10}|^{2} + |U_{11} + U_{-11}|^{2}}{N} \quad \text{u1 > 0 means contribution of UPE}$ where $N = N_{T} + \epsilon N_{L}$ , $N_{T} = \sum_{\lambda_{N}\lambda_{N}'} (|T_{11}|^{2} + |T_{01}|^{2} + |T_{-11}|^{2} + |U_{11}|^{2} + |U_{01}|^{2} + |U_{-11}|^{2})$ $N_{L} = \sum_{\lambda_{N}\lambda_{N}'} (|T_{00}|^{2} + |T_{10}|^{2} + |T_{-10}|^{2} + |U_{10}|^{2} + |U_{-10}|^{2}).$



### Longitudinal to Transverse cross section hermes ratio

$$\begin{split} R_{L/T} &= \frac{\sigma_L}{\sigma_T} = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}}, \qquad r_{00}^{04} = \widetilde{\sum} \{\epsilon |T_{00}|^2 + |T_{01}|^2 + |U_{01}|^2, \} / \mathcal{N} \qquad \mathcal{N} = \epsilon \sigma_L + \sigma_T \\ \sigma_L &= |T_{00}|^2 + |T_{10}|^2 + |T_{-10}|^2 + |U_{10}|^2 + |U_{-10}|^2 \\ \sigma_T &= |U_{11}|^2 + |U_{01}|^2 + |U_{-11}|^2 + |T_{11}|^2 + |T_{01}|^2 + |T_{-11}|^2 \end{split}$$





at average kinematics

$$\begin{split} R_{L/T}(p) &= \textbf{0.25}{\pm}0.03 \pm 0.07 \\ R_{L/T}(d) &= \textbf{0.024}{\pm}0.04 \pm \textbf{0.0.08} \end{split}$$





- The SDMEs were extracted for electroproduction of  $\omega$  vector meson on proton and deuteron at HERMES.
- They are presented divided into five classes according to the helicity transition.
- The hypothesis SCHC in  $\omega$  meson production seems to be slightly violates.
- The UPE contribution seems to be very large (dominant) for  $\omega$  meson production.
- Longitudinal to Transverse cross section ratio for  $\omega$  meson is smaller than for  $\rho^0$ .