

Recent results from the hermes experiment

Charlotte Van Hulse, on behalf of the HERMES collaboration
University of the Basque Country - UPV/EHU

eman ta zabal zazu



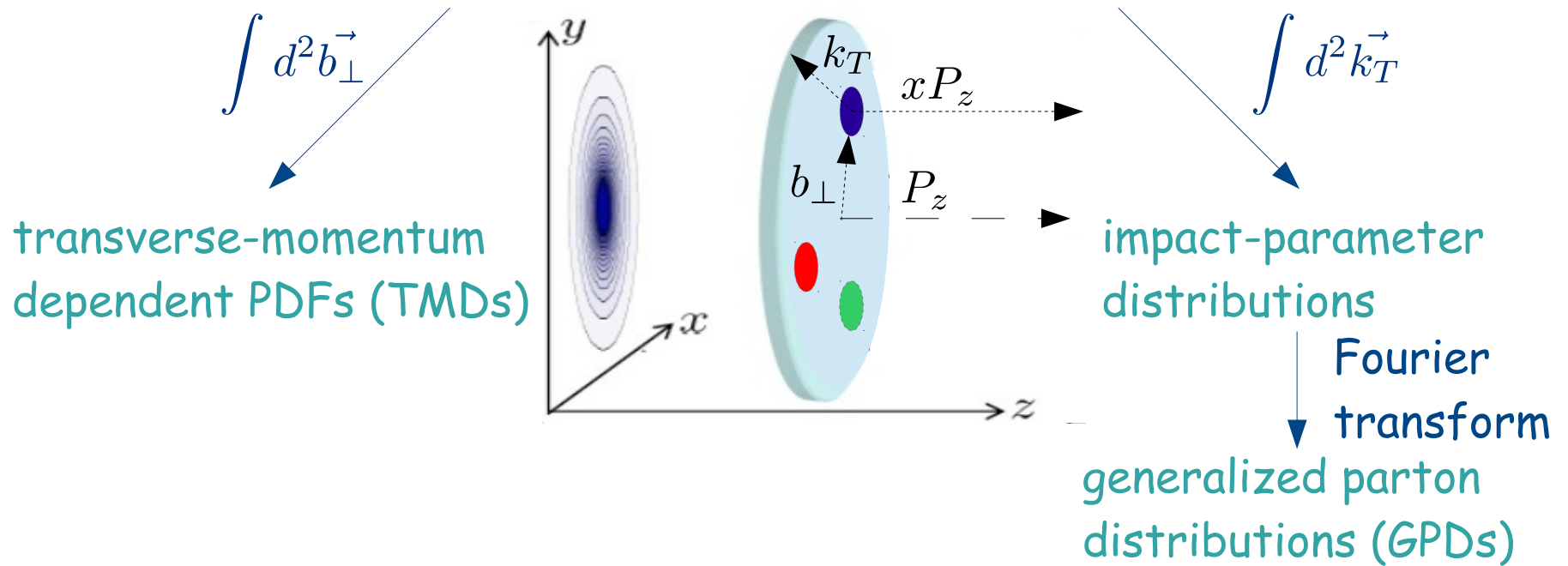
Universidad del País Vasco Euskal Herriko Unibertsitatea

ikerbasque
Basque Foundation for Science

PacSpin 2013, Ji'nan

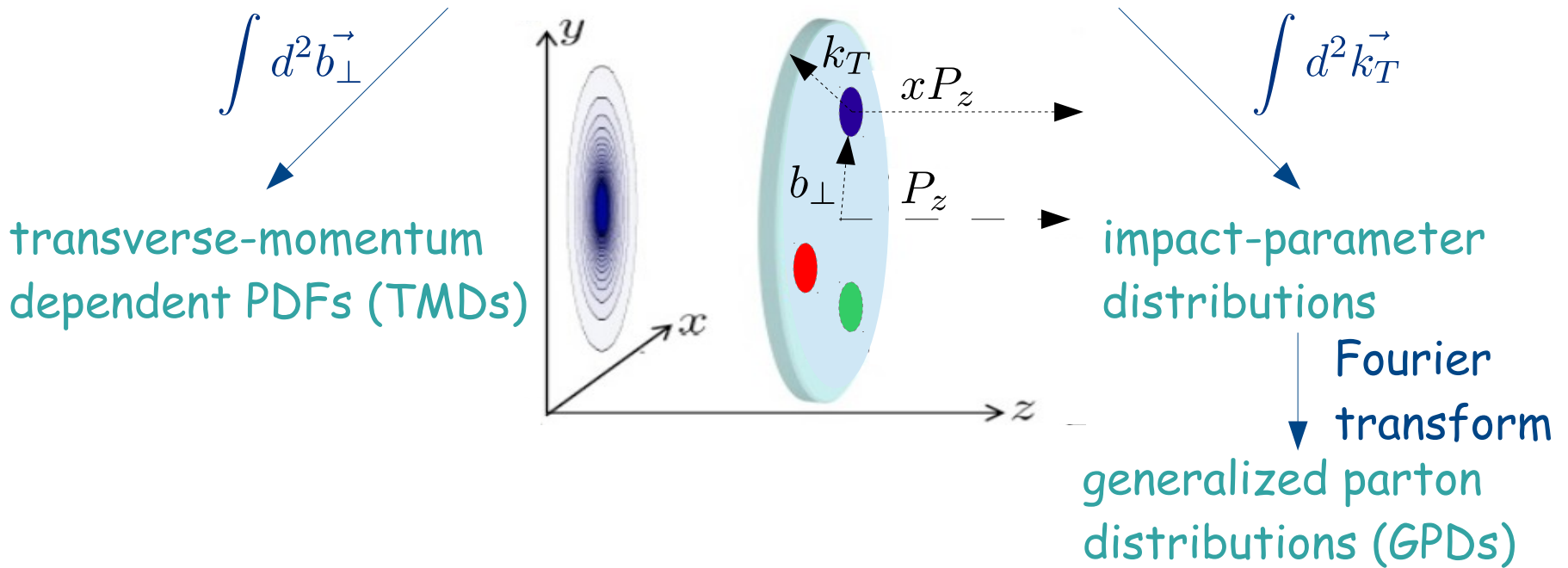
The nucleon in multiple dimensions

Wigner distributions $W(x, \vec{k}_T, \vec{b}_\perp)$

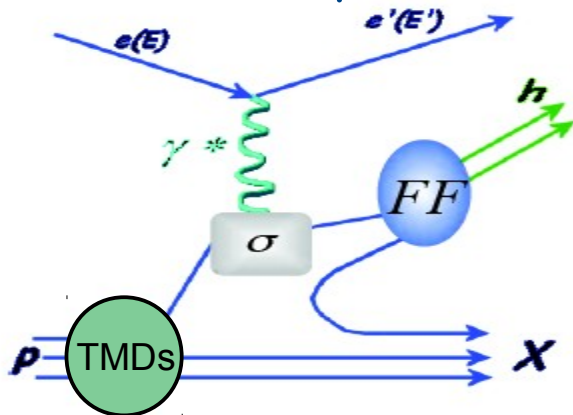


The nucleon in multiple dimensions

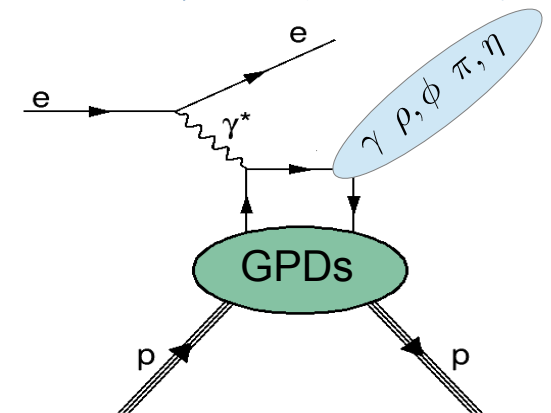
Wigner distributions $W(x, \vec{k}_T, b_\perp)$



semi-inclusive deep-inelastic scattering (DIS)



hard exclusive reactions



Semi-inclusive DIS

Semi-inclusive DIS cross section

$$\frac{d\sigma}{dx dy dz d\phi_h dP_{h\perp}^2 d\phi_S} = \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x} \right)$$

$$\left\{ F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi_h) F_{UU}^{\cos(\phi_h)} + \epsilon \cos(2\phi_h) F_{UU}^{\cos(2\phi_h)} \right.$$

beam polarization

$$+ \lambda_e \sqrt{2\epsilon(1-\epsilon)} \sin(\phi_h) F_{LU}^{\sin(\phi_h)}$$

longitudinal target polarization

$$+ S_L \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi_h) F_{UL}^{\sin(\phi_h)} + \epsilon \sin(2\phi_h) F_{UL}^{\sin(2\phi_h)} \right]$$

$$+ S_L \lambda_e \left[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_h) F_{LL}^{\cos(\phi_h)} \right]$$

transverse target polarization

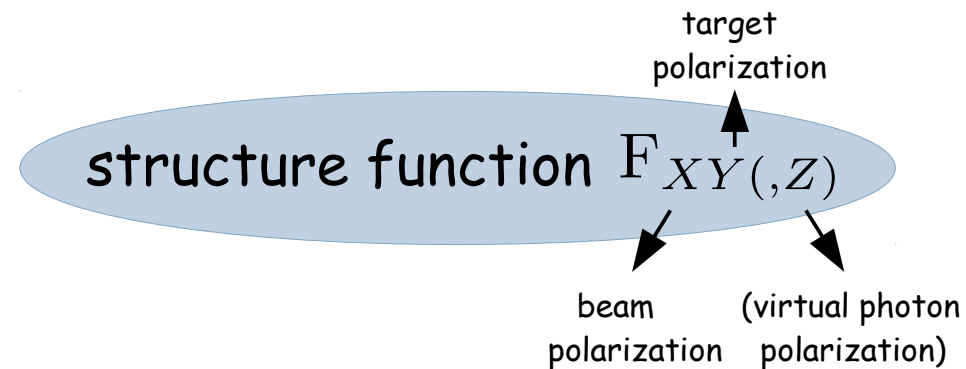
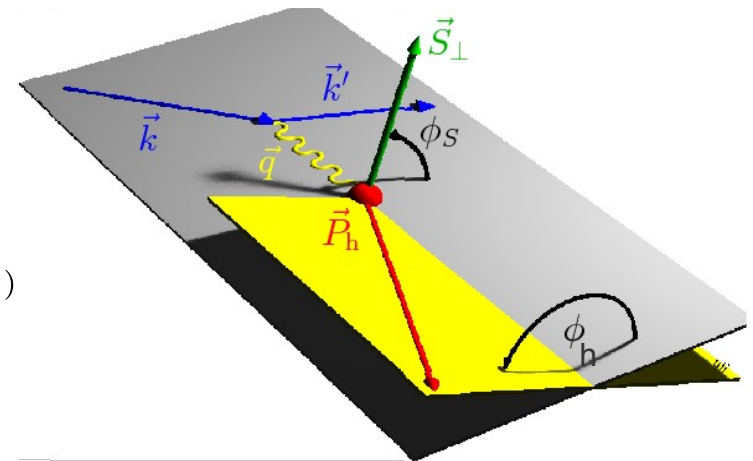
$$+ S_T \left[\sin(\phi_h - \phi_S) \left(F_{UT,T}^{\sin(\phi_h - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right) \right.$$

$$+ \epsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \epsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)}$$

$$\left. + \sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} + \sqrt{2\epsilon(1+\epsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right]$$

$$+ S_T \lambda_e \left[\sqrt{1-\epsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \right.$$

$$\left. + \sqrt{2\epsilon(1-\epsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right\}$$



Semi-inclusive DIS cross section

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$$\left\{ F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi_h) F_{UU}^{\cos(\phi_h)} + \epsilon \cos(2\phi_h) F_{UU}^{\cos(2\phi_h)} \right.$$

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transverse target polarization

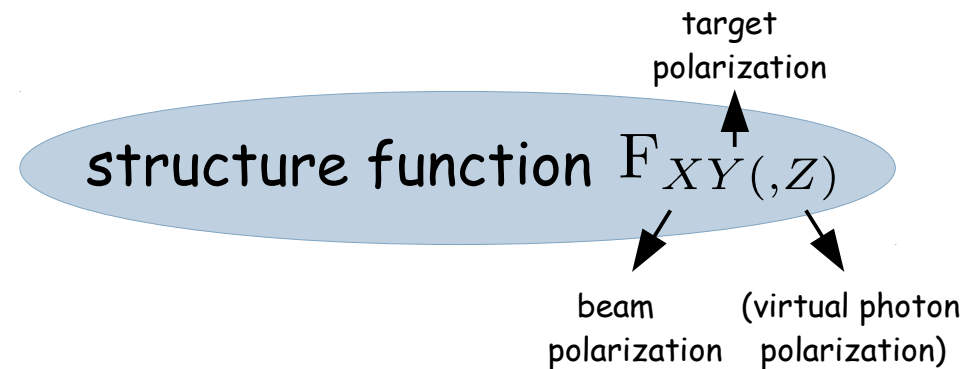
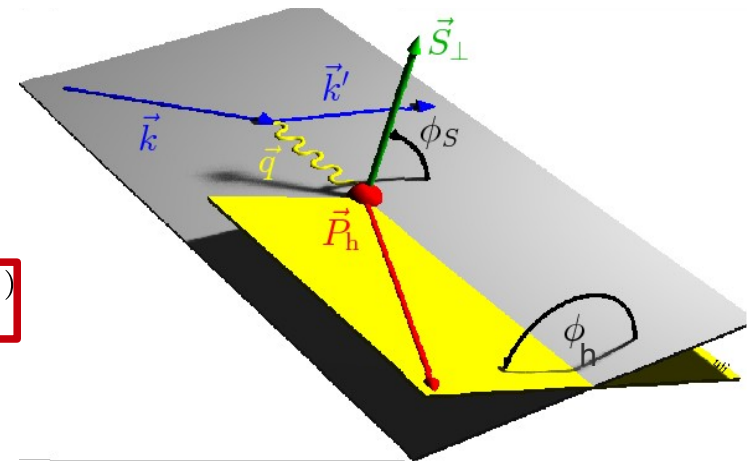
$$+ S_T \left[\sin(\phi_h - \phi_S) F_{UT,T}^{\sin(\phi_h - \phi_S)} + \epsilon F_{UT,L}^{\sin(\phi_h - \phi_S)} \right]$$

$$+ \epsilon \sin(\phi_h + \phi_S) F_{UT}^{\sin(\phi_h + \phi_S)} + \epsilon \sin(3\phi_h - \phi_S) F_{UT}^{\sin(3\phi_h - \phi_S)}$$

$$+ \left[\sqrt{2\epsilon(1+\epsilon)} \sin(\phi_S) F_{UT}^{\sin(\phi_S)} + \sqrt{2\epsilon(1+\epsilon)} \sin(2\phi_h - \phi_S) F_{UT}^{\sin(2\phi_h - \phi_S)} \right]$$

$$+ S_T \lambda_e \left[\sqrt{1-\epsilon^2} \cos(\phi_h - \phi_S) F_{LT}^{\cos(\phi_h - \phi_S)} + \sqrt{2\epsilon(1-\epsilon)} \cos(\phi_S) F_{LT}^{\cos(\phi_S)} \right]$$

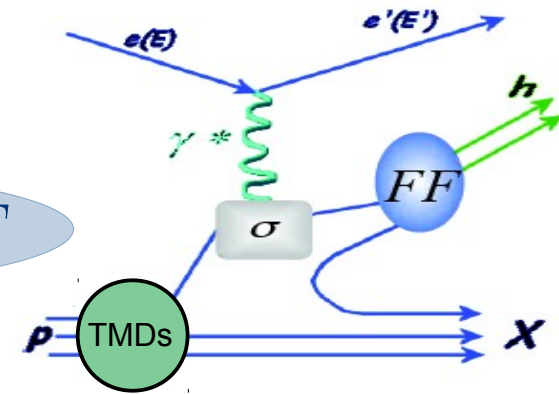
$$+ \left. \left[\sqrt{2\epsilon(1-\epsilon)} \cos(2\phi_h - \phi_S) F_{LT}^{\cos(2\phi_h - \phi_S)} \right] \right\}$$



leading twist

Semi-inclusive DIS cross section

structure function $F_{XY} \propto TMD \otimes FF$



transverse momentum distributions (TMDs)

		quark		
		U	L	T
nucleon	U	f_1		h_1^\perp
	L		g_1	h_{1L}^\perp
	T	f_{1T}^\perp	g_{1T}^\perp	h_1 h_{1T}^\perp

fragmentation functions (FFs)

		quark		
		U	L	T
h	U	D_1		H_1^\perp

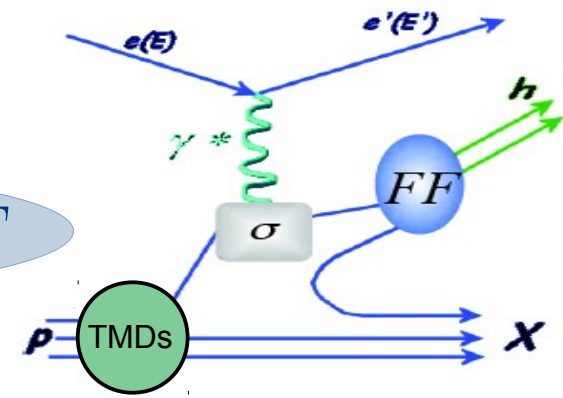
nucleon with transverse/longitudinal spin

quark with transverse/longitudinal spin

quark transverse momentum

Semi-inclusive DIS cross section

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transverse momentum distributions (TMDs)

fragmentation functions (FFs)

		quark			quark			
		U	L	T	spin-independent FF	Collins FF	T	
nucleon	U	f_1		h_1^\perp	h U D_1		H_1^\perp	
	L		g_1	h_{1L}^\perp				
	T	f_{1T}^\perp	g_{1T}^\perp	h_1 h_{1T}^\perp				

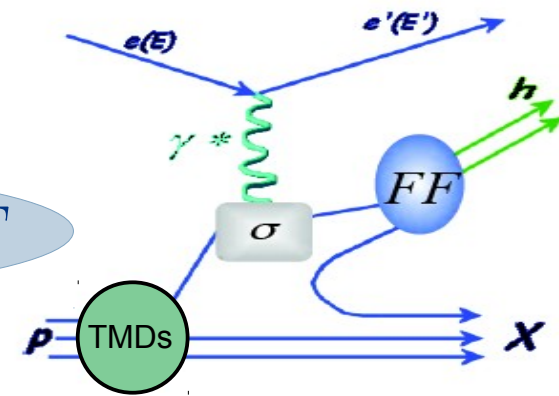
nucleon with transverse/longitudinal spin

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quark transverse momentum

Semi-inclusive DIS cross section

structure function $F_{XY} \propto TMD \otimes FF$



transverse momentum distributions (TMDs)

fragmentation functions (FFs)

		quark			quark		
		U	D	T	U	D	T
nucleon	U	f_1	Boer-Mulders DF h_1^\perp		spin-independent FF D_1	Collins FF H_1^\perp	
	L		g_1	h_{1L}^\perp			
	T	Sivers DF f_{1T}^\perp	g_{1T}^\perp	h_1			

nucleon with transverse/longitudinal spin

quark with transverse/longitudinal spin

quark transverse momentum

Hadron multiplicities

Hadron multiplicities

$$\frac{d\sigma}{dx dy dz d\phi_h dP_{h\perp}^2} \quad \int d\phi_h$$

$$M^h(x_B, Q^2, z, P_{h\perp}) = \frac{1}{d^2 N^{DIS}(x_B, Q^2)} \frac{d^4 N^h(x_B, Q^2, z, P_{h\perp})}{dz dP_{h\perp}}$$

$$\propto \frac{F_{UU,T} + \epsilon F_{UU,L}}{F_T + \epsilon F_L}$$

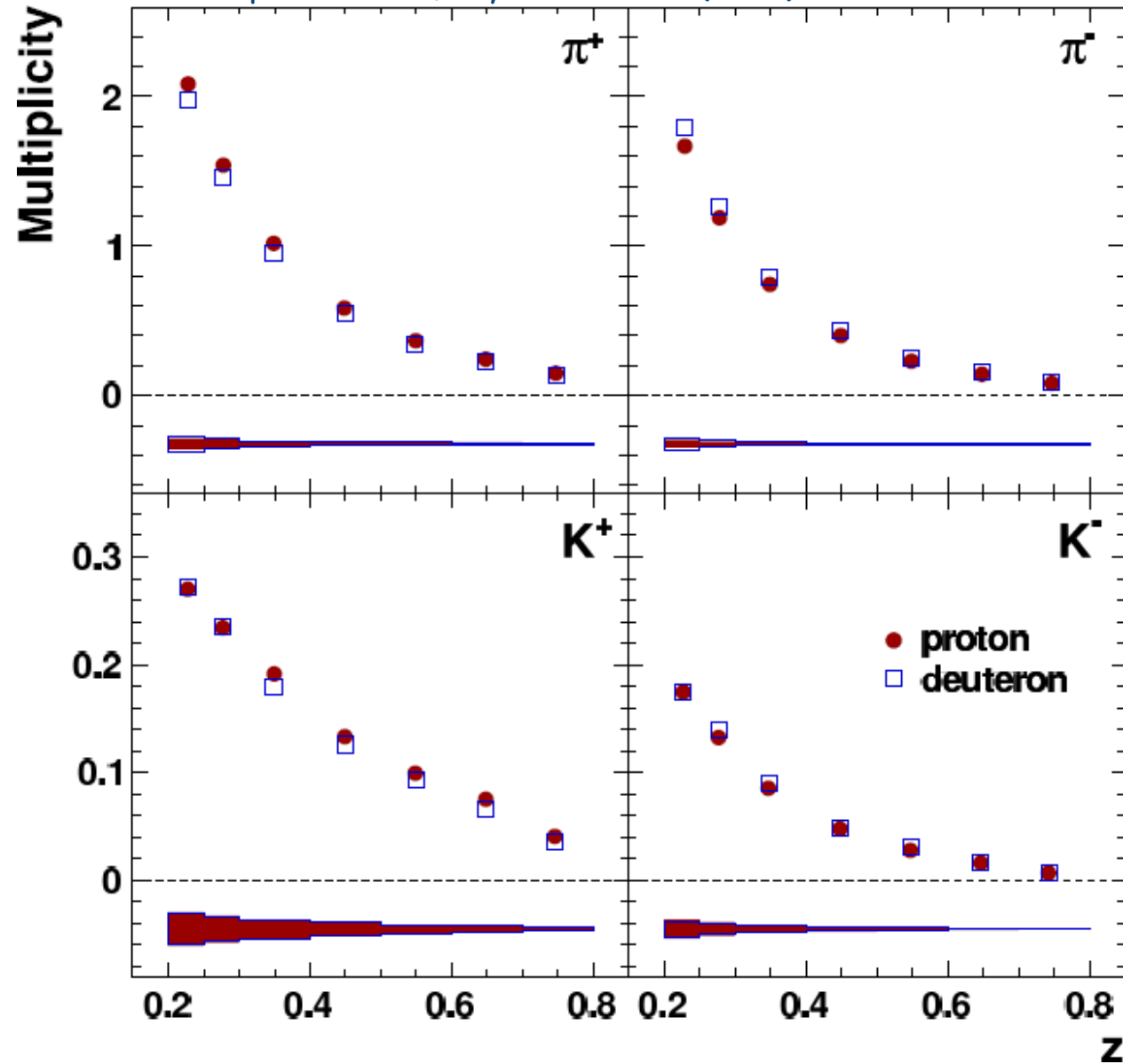
$$\propto \frac{\sum_q e_q^2 f_1^q(x_B, k_T^2, Q^2) \otimes D_1^q(z, p_T^2, Q^2)}{\sum_q e_q^2 f_1^q(x_B, Q^2)}$$

k_T : transverse momentum of struck quark

p_T : transverse momentum of fragmenting quark

Results projected in z

A. Airapetian et al., Phys. Rev. D **87** (2013) 074029

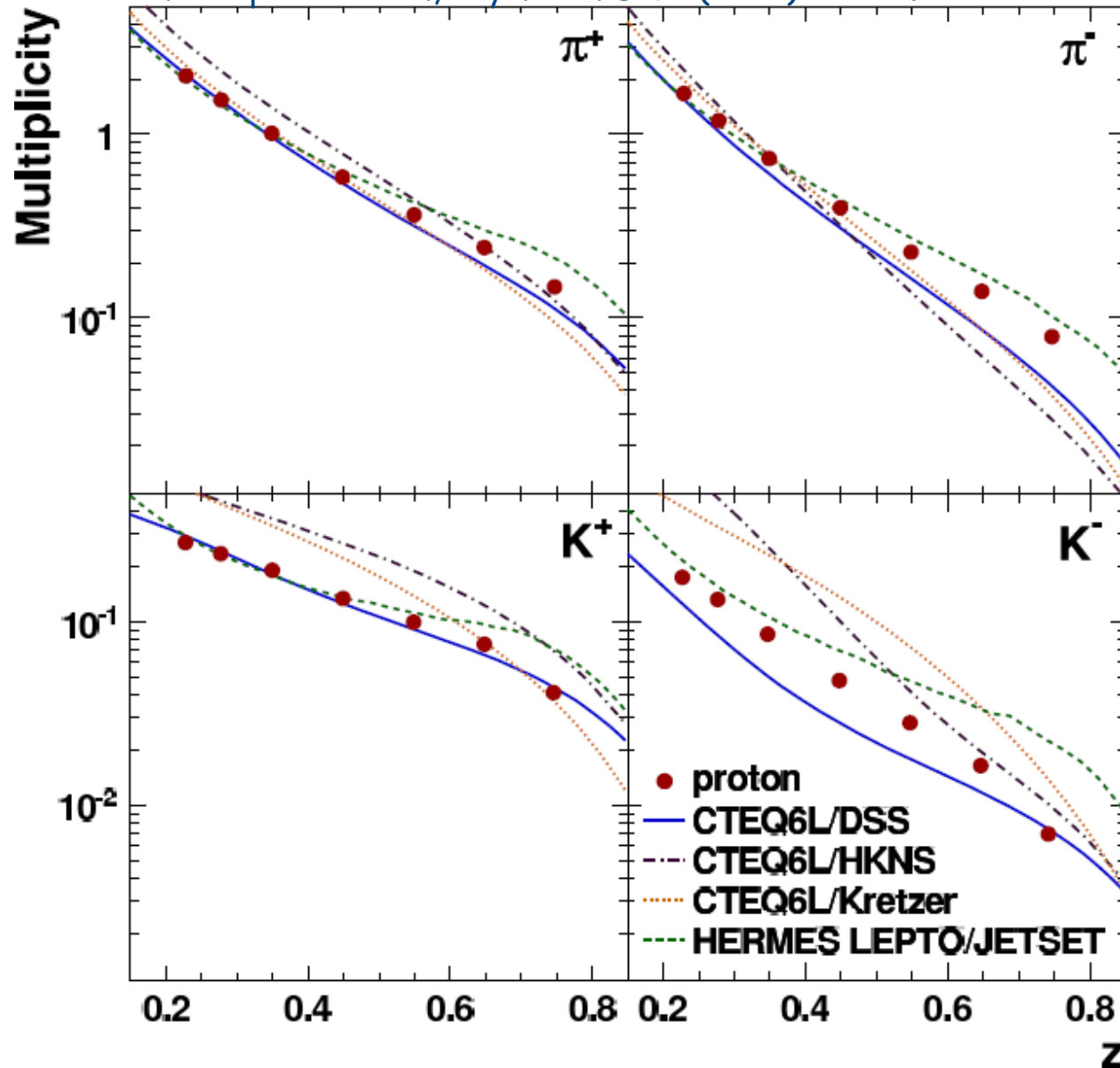


multiplicities reflect

- nucleon valence-quark content (u-dominance)
- favored \leftrightarrow unfavored fragmentation

Comparison to models

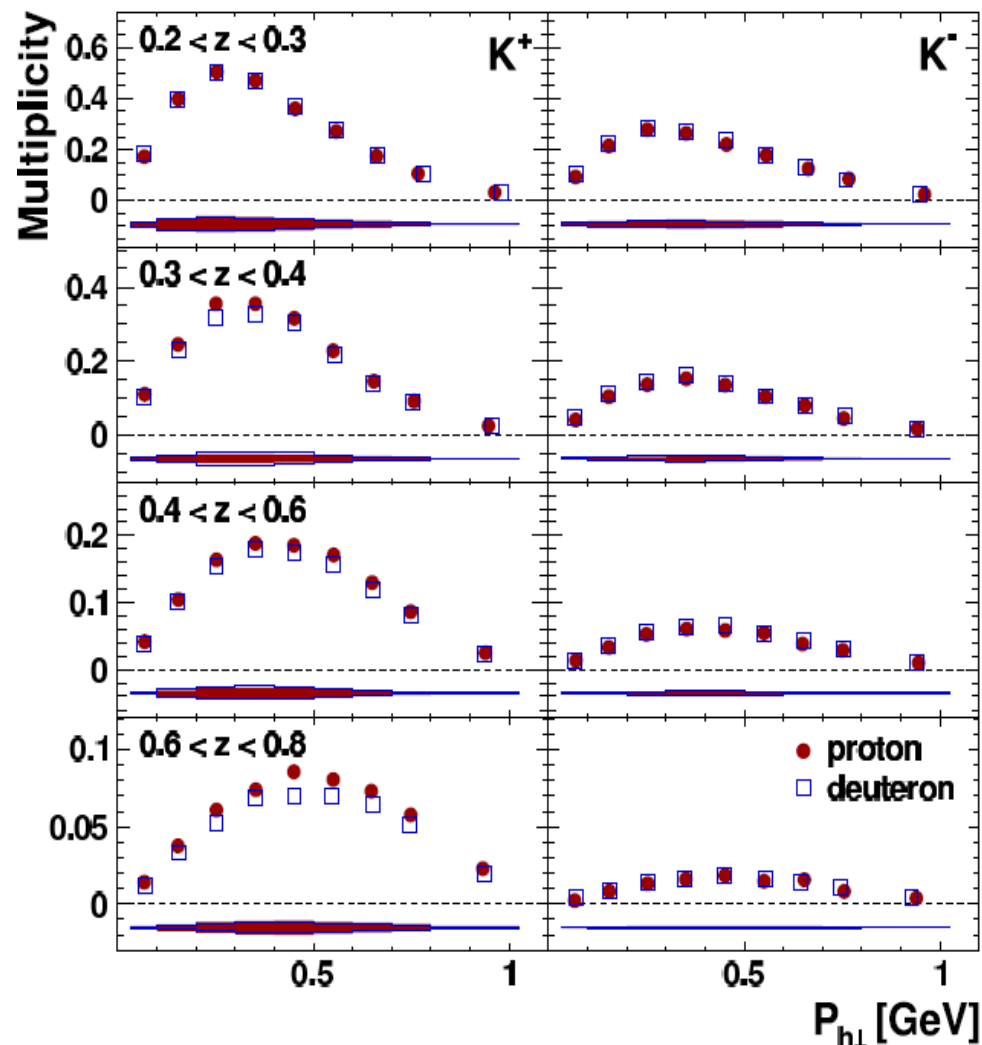
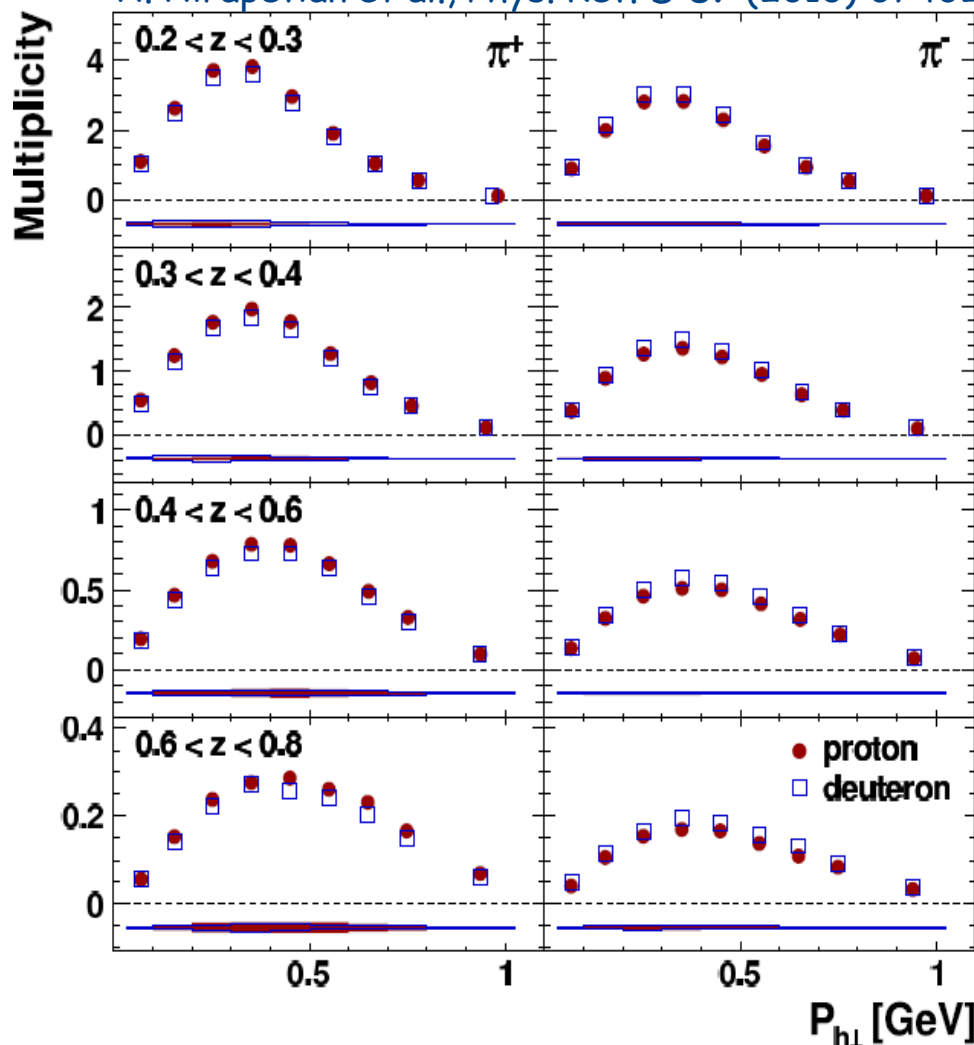
A. Airapetian et al., Phys. Rev. D **87** (2013) 074029



- LO in α_S
- CTEQ6L PDFs
JHEP **0207** (2002) 012
- DSS FFs
Phys. Rev. D **75** (2007) 114010
- Kretzer FFs
Phys. Rev. D **62** (2000) 054001
- reasonable agreement between DSS and positive mesons
- poor agreement for negative mesons

Results projected in z and $P_{h\perp}$

A. Airapetian et al., Phys. Rev. D **87** (2013) 074029



- $P_{h\perp}$: - transverse intrinsic struck-quark momentum
- transverse momentum from fragmentation process
- K^- : broader distribution

Re-evaluation of the strange quark distribution

K^\pm multiplicities from unpolarized deuterium

$$\frac{d^2 N^K(x)}{d^2 N^{DIS}(x)} = \frac{Q(x) \int_{0.2}^{0.8} \mathcal{D}_Q^K(z) dz + S(x) \int_{0.2}^{0.8} \mathcal{D}_S^K(z) dz}{5Q(x) + 2S(x)}$$

$$Q(x) \equiv u(x) + \bar{u}(x) + d(x) + \bar{d}(x)$$

$$S(x) \equiv s(x) + \bar{s}(x)$$

$$\mathcal{D}_Q^K(z) \equiv 4D_u^K(z) + D_d^K(z)$$

$$\mathcal{D}_S^K(z) \equiv 2D_s^K(z)$$

Re-evaluation of the strange quark distribution

K^\pm multiplicities from unpolarized deuterium

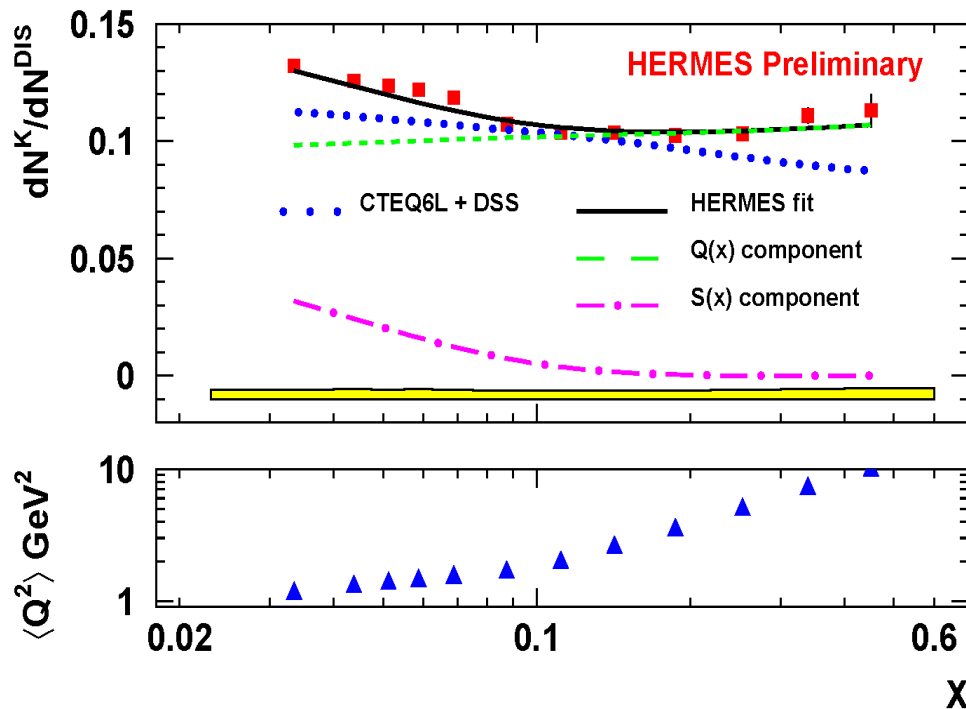
$$\frac{d^2 N^K(x)}{d^2 N^{DIS}(x)} = \frac{Q(x) \int_{0.2}^{0.8} \mathcal{D}_Q^K(z) dz + S(x) \int_{0.2}^{0.8} \mathcal{D}_S^K(z) dz}{5Q(x) + 2S(x)}$$

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$$\mathcal{D}_Q^K(z) \equiv 4D_u^K(z) + D_d^K(z)$$

$$\mathcal{D}_S^K(z) \equiv 2D_s^K(z)$$



$S(x) \approx 0$ for $x > 0.1$

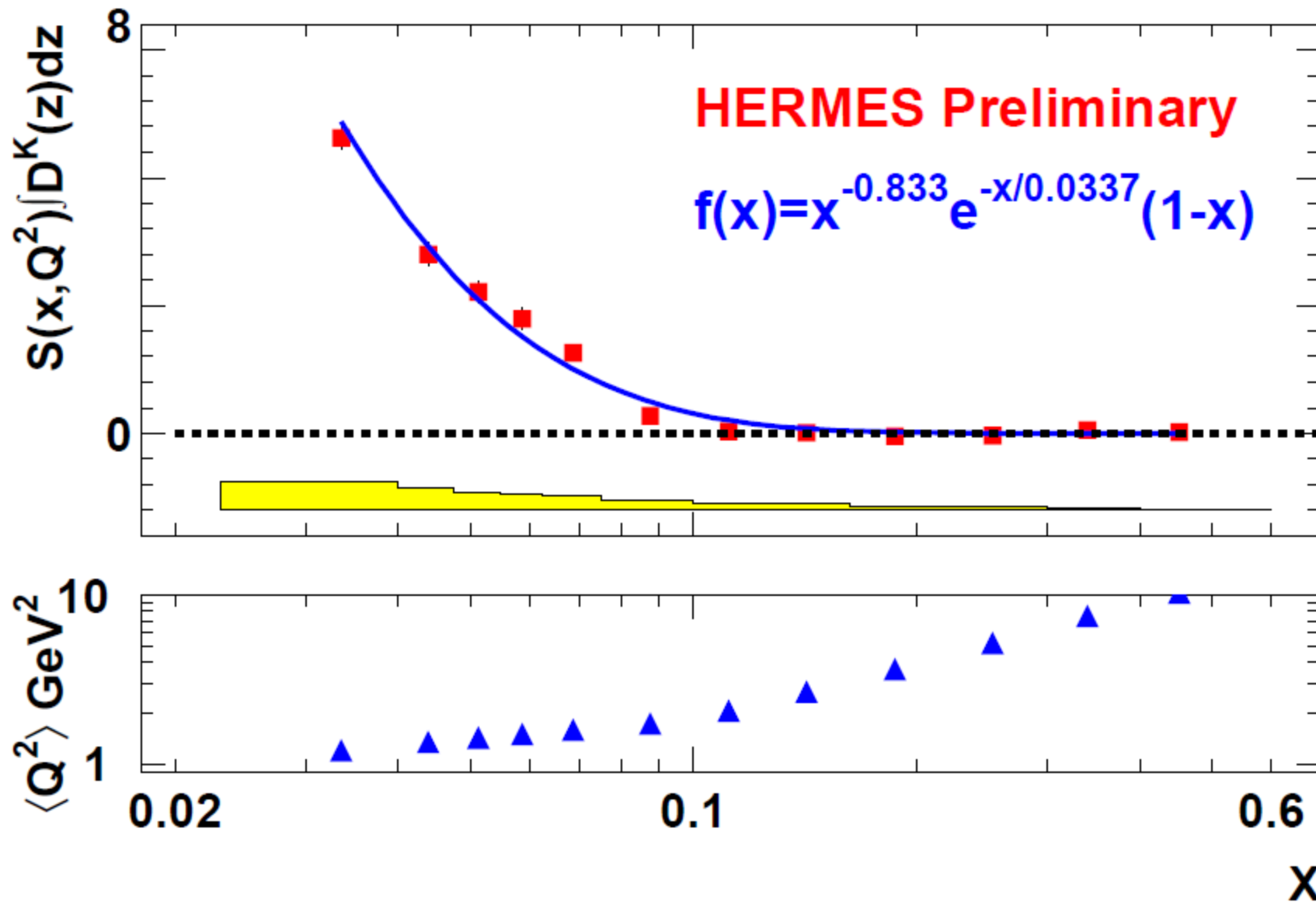
fit

$$f(x) = (0.102 \pm 0.002) + (0.013 \pm 0.010) x$$

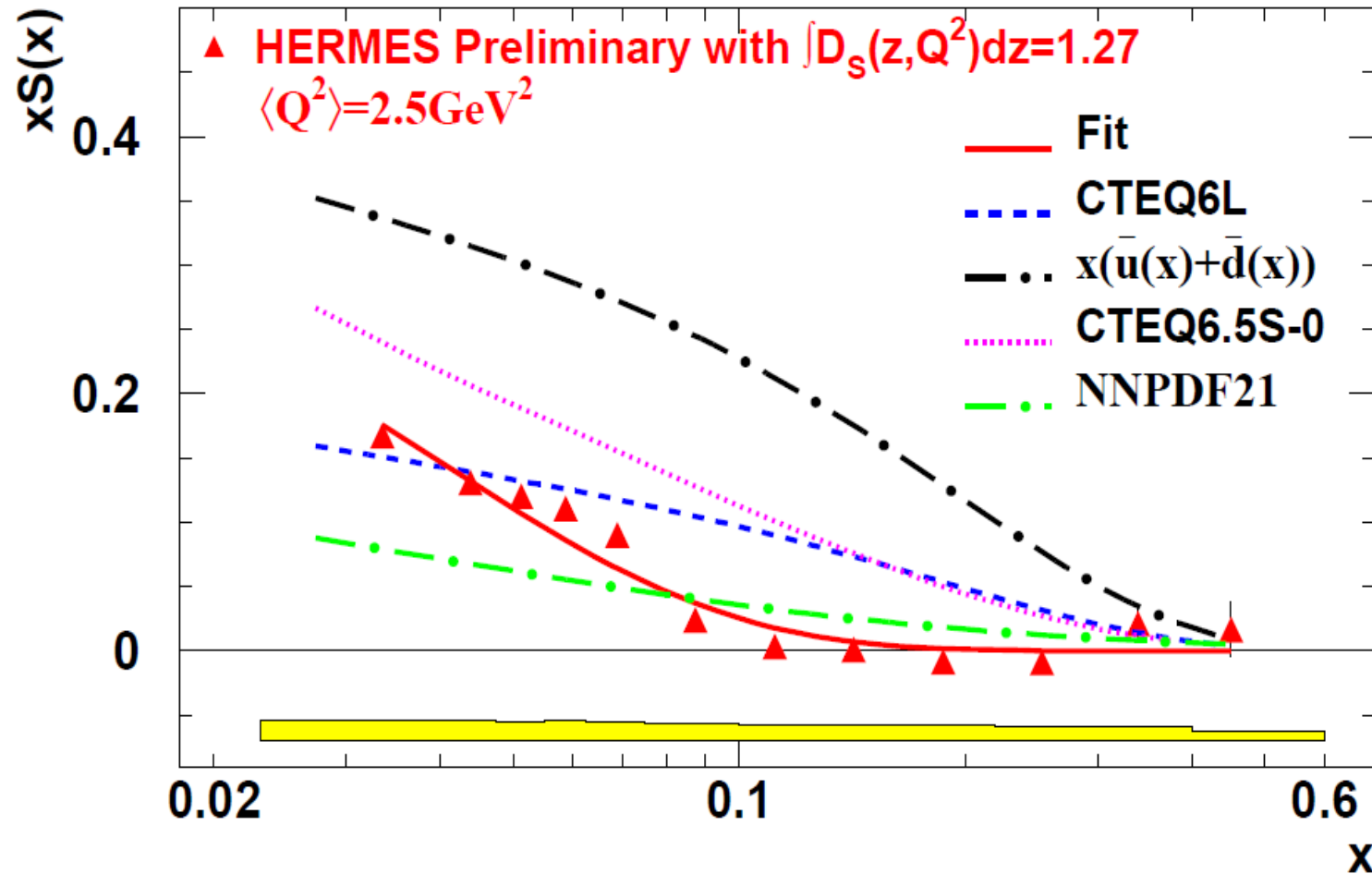
$$\int_{0.2}^{0.8} \mathcal{D}_Q^K(z) dz = 0.398 \pm 0.010$$

Re-evaluation of the strange quark distribution

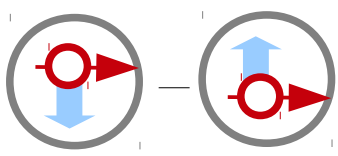
$Q(x)$ from CTEQ6L



Re-evaluation of the strange quark distribution



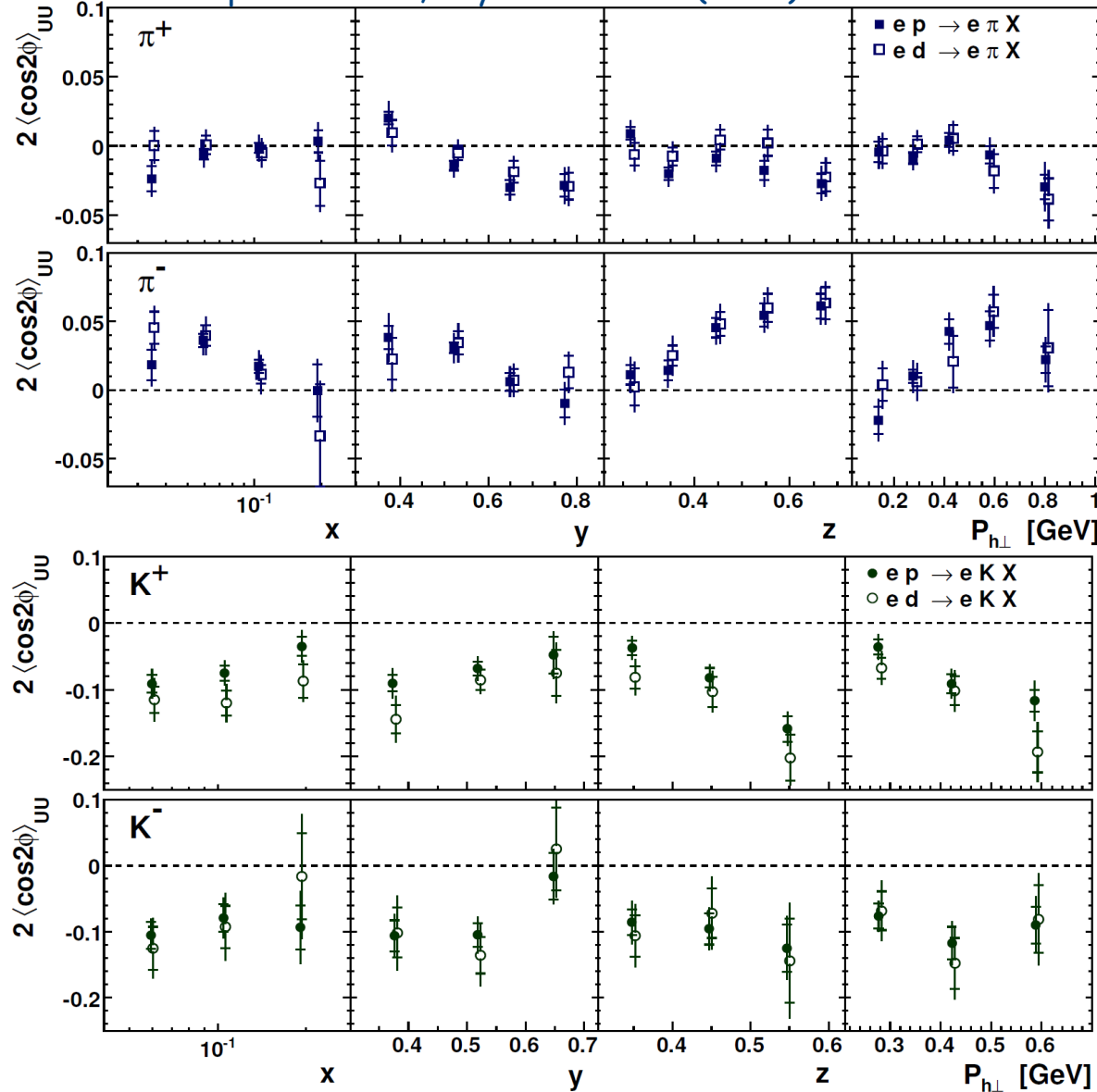
- $xS(x)$ for certain value of $\int \mathcal{D}_S^K(z) dz$
- independent of value, shape of $xS(x)$ incompatible with predictions



Boer-Mulders function h_1^\perp

$$F_{UU}^{\cos(2\phi_h)} \propto h_1^\perp \otimes H_1^\perp$$

A. Airapetian et al., Phys. Rev. D **87** (2013) 012010

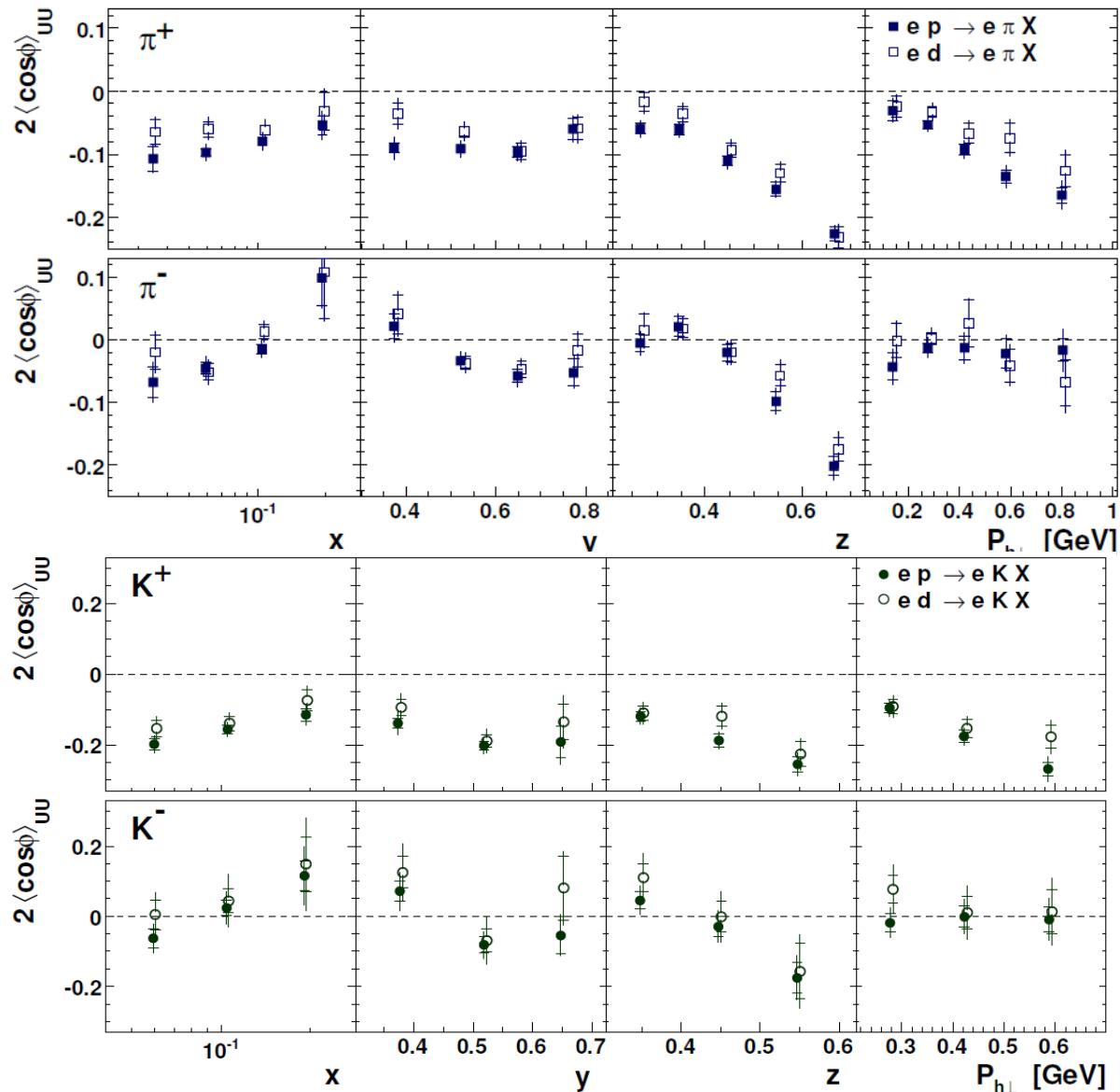


- significant amplitudes \rightarrow orbital angular momentum
- H-D comparison: $h_1^{\perp,u} \approx h_1^{\perp,d}$
- $\pi^- > 0 \iff \pi^+ \leq 0$, consistent with $H_1^{\perp,fav} \approx -H_1^{\perp,unfav}$
- K^+ & K^- : large negative amplitudes
- $K^- \approx K^+$
- K^\pm very different from π^\pm

Twist-3: Cahn, Boer-Mulders, ...

$$F_{UU}^{\cos\phi_h} \propto \frac{M}{Q} [f_1 \times D_1, h_1^\perp \times H_1^\perp, \dots]$$

A. Airapetian et al., Phys. Rev. D **87** (2013) 012010



- H-D comparison:
weak flavour dependence
- strong flavour dependence

DIS with transversely polarized target

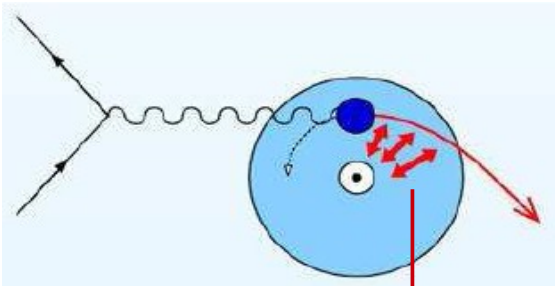
Structure function g_2

- $$\frac{d^3 \sigma_{LT}}{dx dy d\phi'} \propto -h_k \left(\frac{y}{2} g_1(x, Q^2) + g_2(x, Q^2) \right) \cos(\phi')$$

\swarrow
 \searrow
 access via $A_{LT}(x, Q^2, \phi') = h_l \frac{\sigma_{LT}(x, Q^2, \phi')}{\sigma_{UU}(x, Q^2, \phi')}$
 $= -A_T \cos(\phi')$

- $$g_2(x) = g_2^{WW} + \bar{g}_2(x) \quad d_2 = 3 \int_0^1 dx x^2 \bar{g}_2(x)$$

- ## Sivers effect

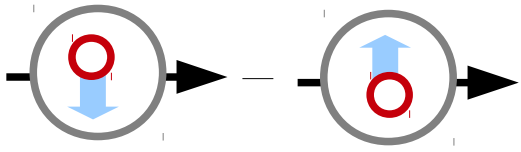


force on struck quark at $t=0$

$$\propto -d_2$$

M. Burkardt, arXiv:0810.3589

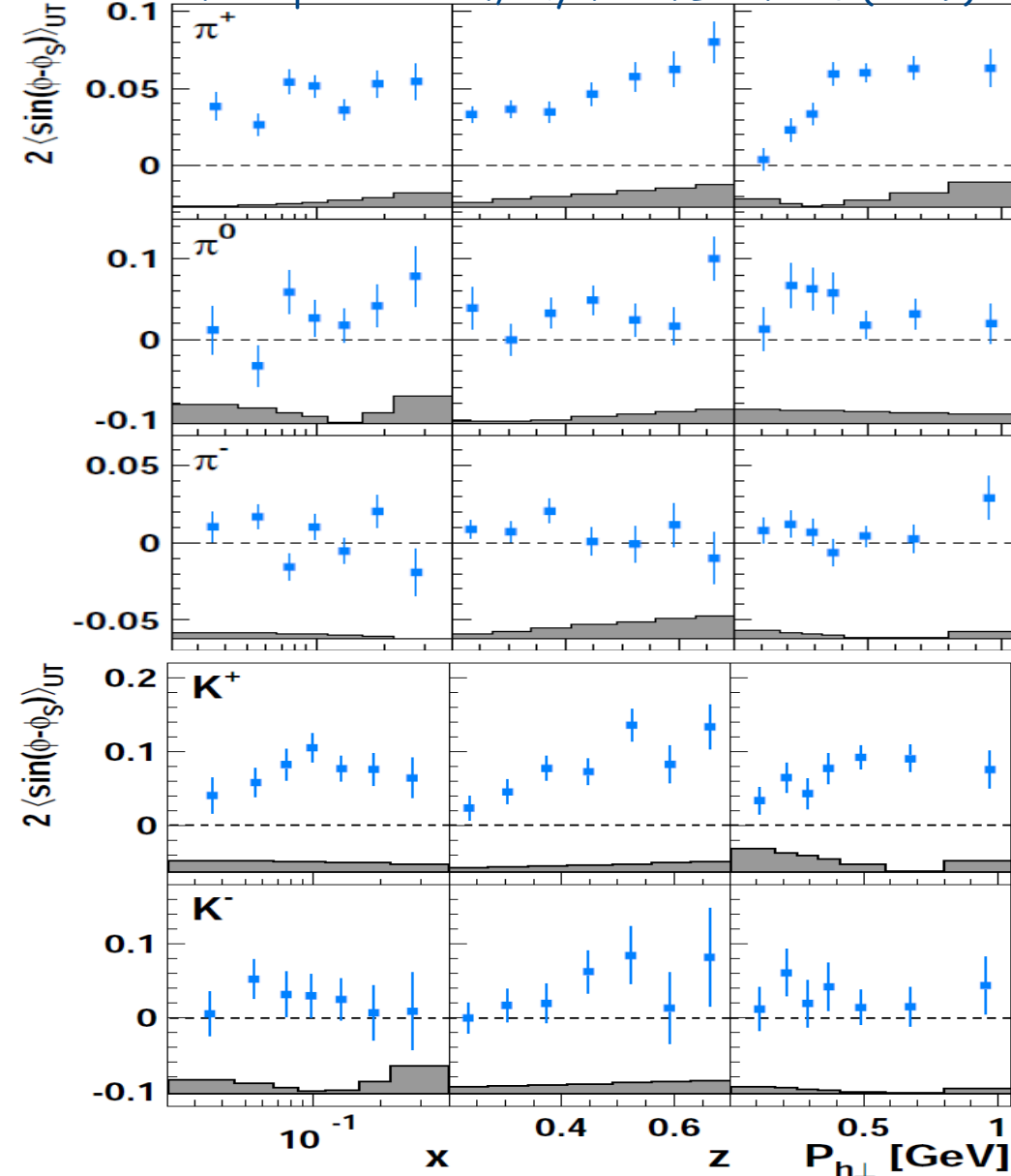
FSI: $t=0 \rightarrow \infty$



Sivers function f_{1T}^\perp

$$F_{UT}^{\sin(\phi_h - \phi_S)} \propto f_{1T}^\perp \otimes D_1$$

A. Airapetian et al., Phys. Rev. Lett. 103 (2009) 152002



- π^+ : significantly positive

→ orbital angular momentum

- u-quark dominance for p^+ amplitude:

$$\approx - \frac{f_{1T}^{\perp,u}(x, k_T^2) \otimes D_1^{u \rightarrow \pi^+}(z, p_T^2)}{f_1^u(x, k_T^2) \otimes D_1^{u \rightarrow \pi^+}(z, p_T^2)}$$

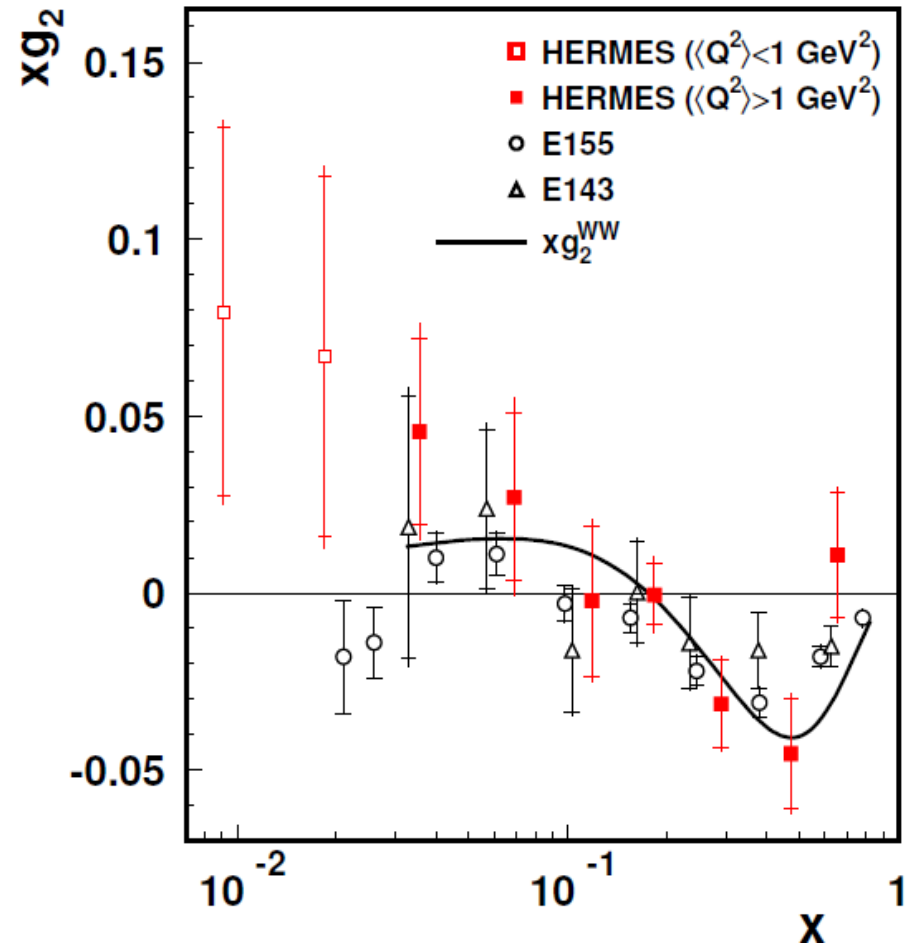
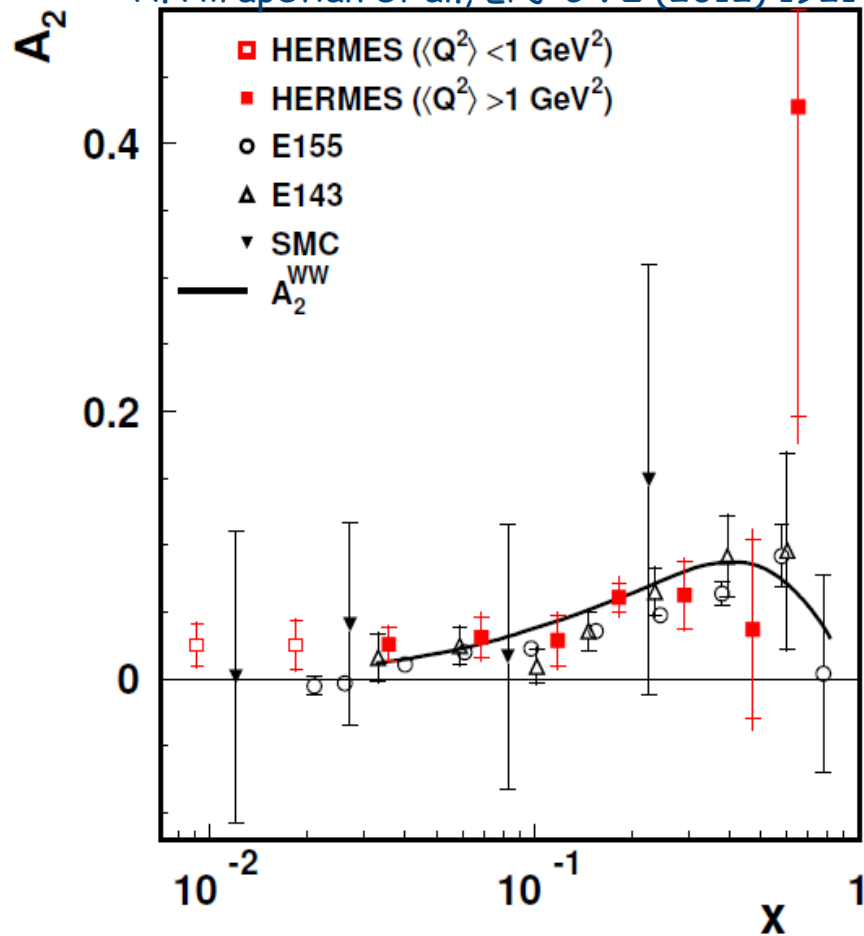
→ $f_{1T}^{\perp,u}(x, k_T^2) < 0$

- π^- : u- and d-quark cancellation

→ $f_{1T}^{\perp,d}(x, k_T^2) > 0$

A_2 and g_2

A. Airapetian et al., EPJ C 72 (2012) 1921



- evaluation of Burkhardt-Cottingham sum rule:

$$\int_{0.023}^{0.9} dx g_2(x) = 0.006 \pm 0.024 \pm 0.017$$

E143+E155:
$$\int_{0.02}^{0.8} dx g_2(x) = -0.042 \pm 0.008$$

- d_2 :

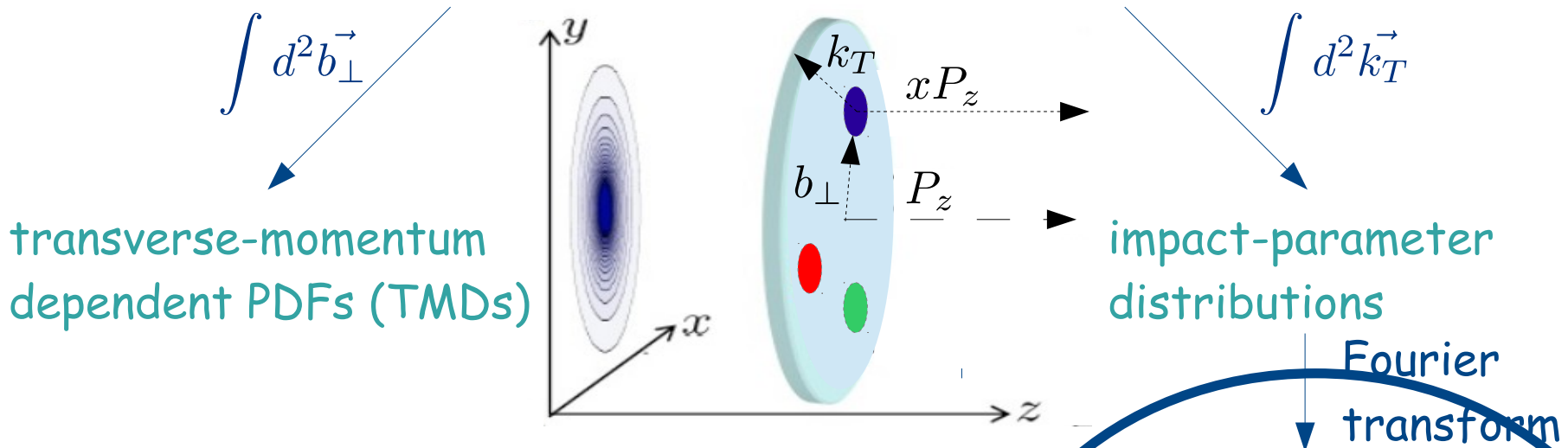
$$d_2 = 0.0148 \pm 0.0096 \pm 0.0048$$

E143+E155:

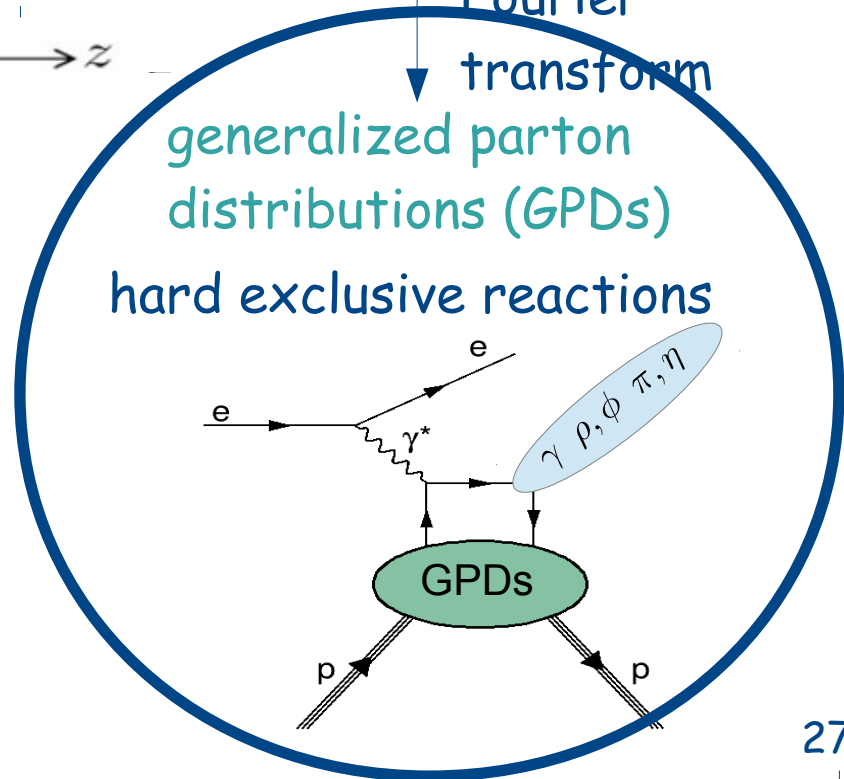
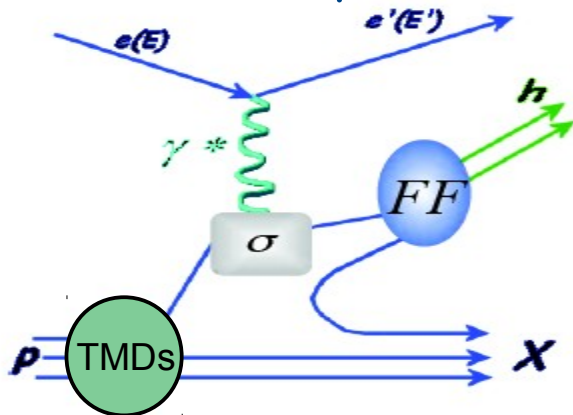
$$d_2 = 0.0032 \pm 0.0017$$

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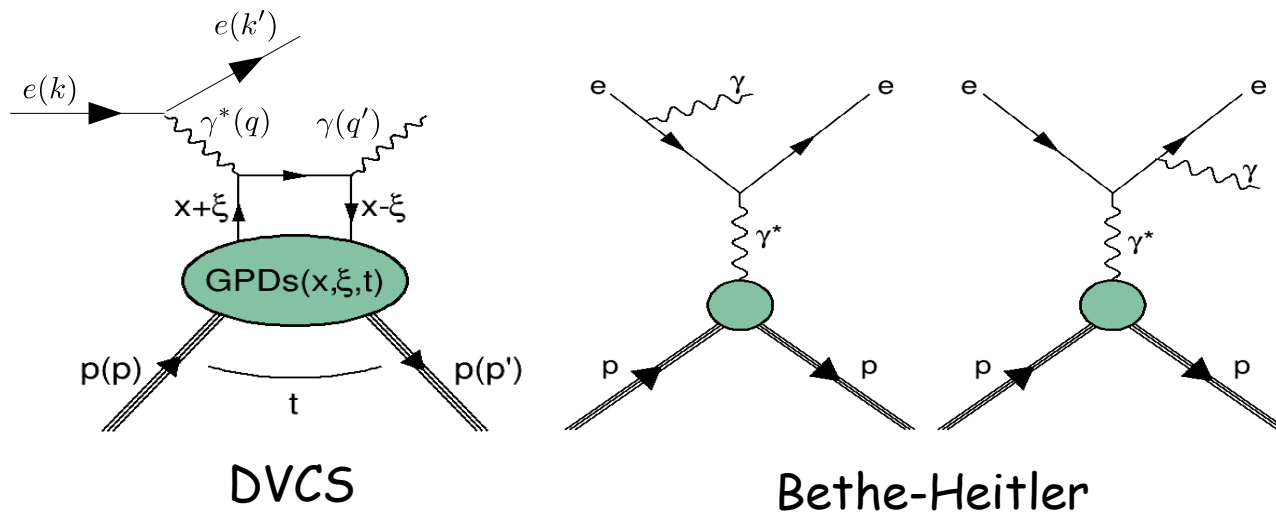


semi-inclusive deep-inelastic scattering (DIS)



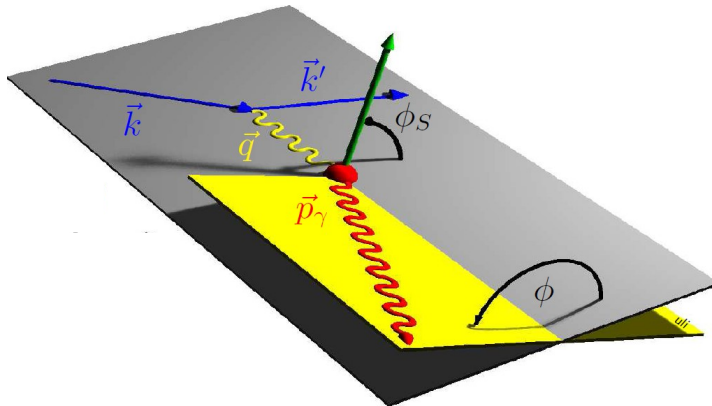
Deeply virtual Compton scattering

Exclusive electroproduction of real photons

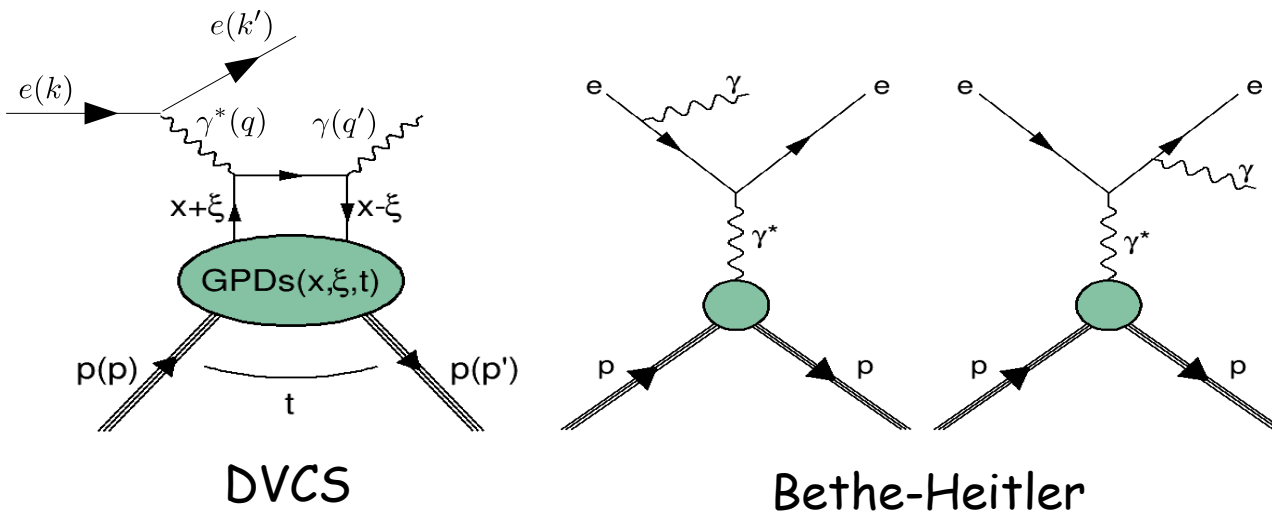


$$d\sigma \propto |\tau_{BH}|^2 + |\tau_{DVCS}|^2 + \tau_{BH} \tau_{DVCS}^* + \tau_{DVCS} \tau_{BH}^*$$

→ access through azimuthal asymmetries

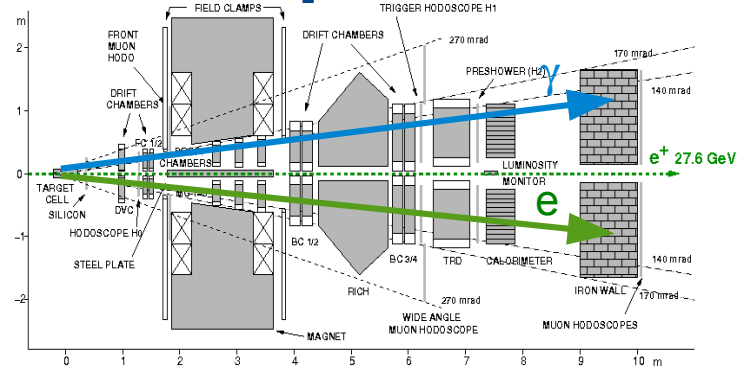
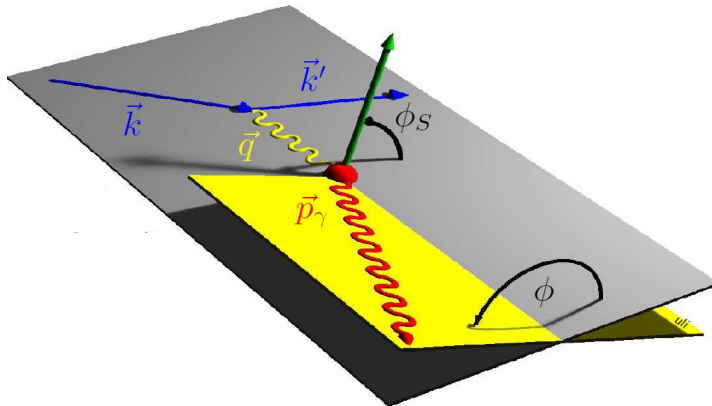


Exclusive electroproduction of real photons



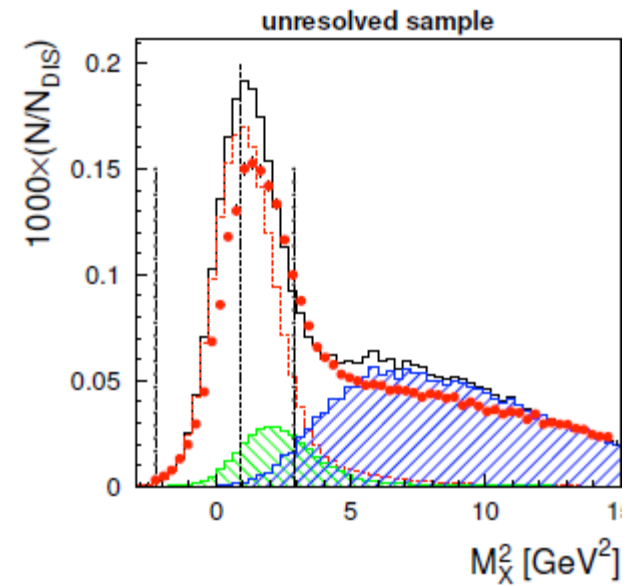
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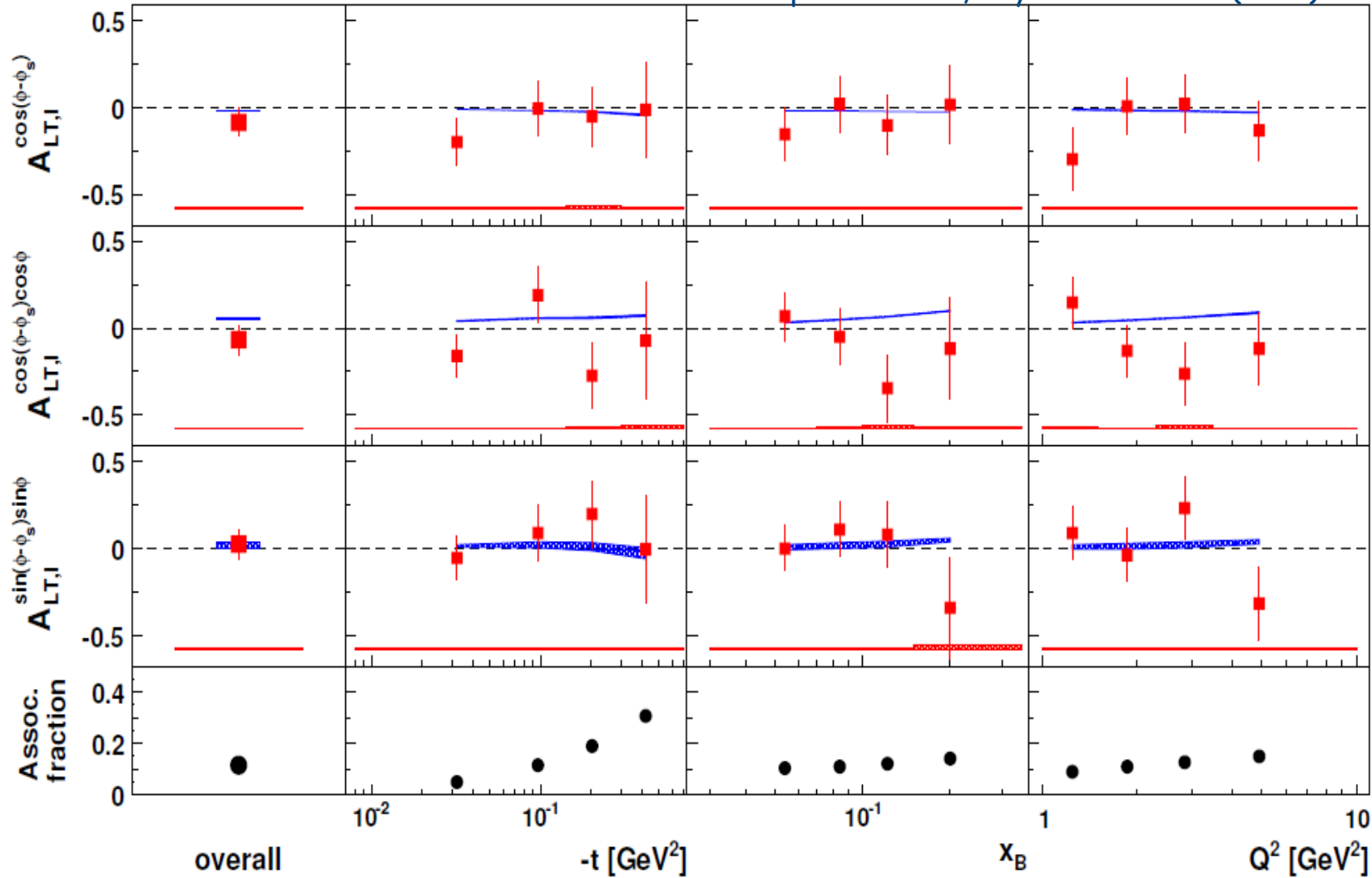
exclusivity

reconstruct proton via missing mass:



Double spin asymmetries on transversely polarized target

A. Airapetian et al., Phys. Lett B 704 (2011) 15



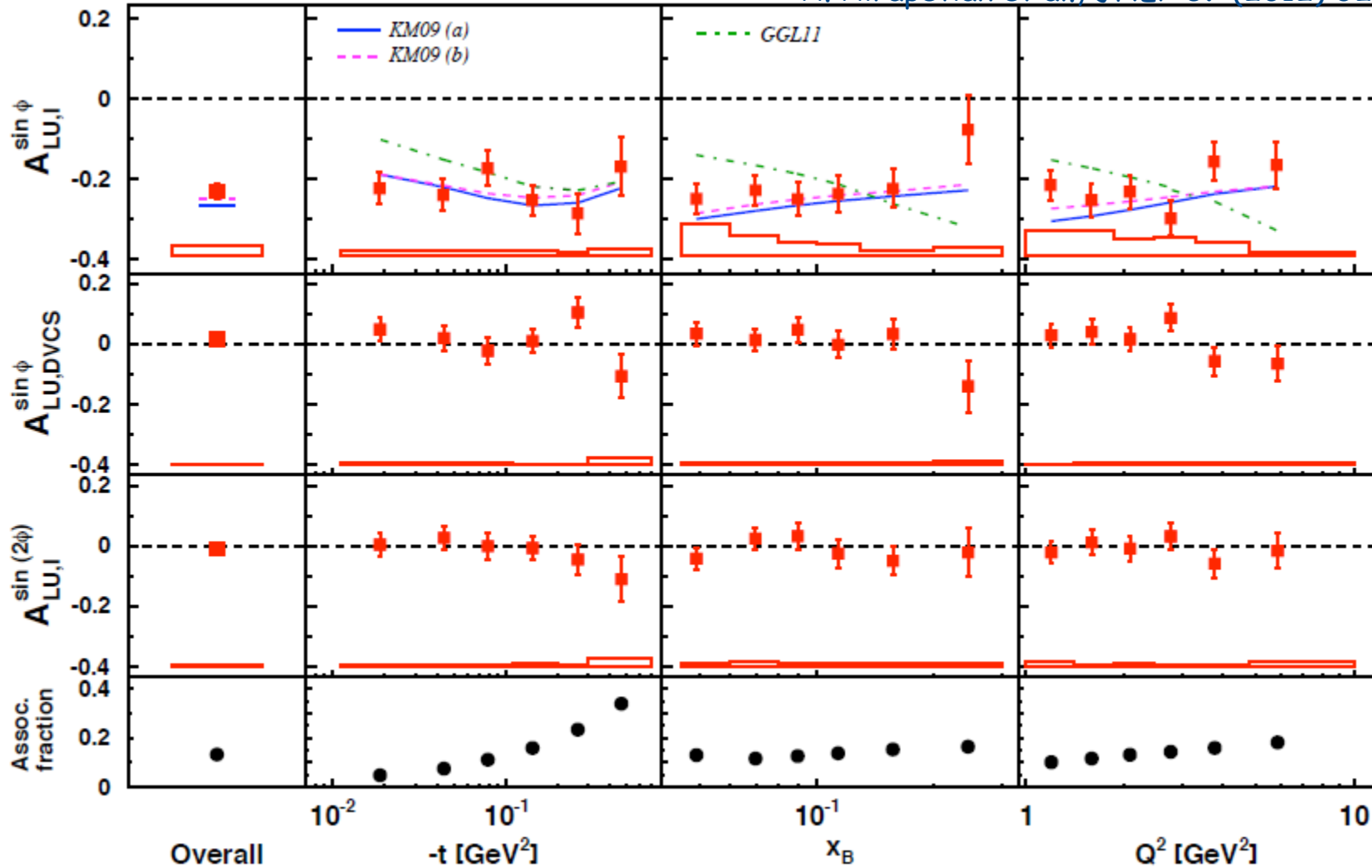
phys. Rev D 60 (1999)
094017 + VGG code:
- double distributions
- Regge t -dependence
- $0.2 < J_u < 0.6$ & $J_d = 0$

- compatible with zero

Charged-separated beam-helicity asymmetry

GPD H

A. Airapetian et al., JHEP 07 (2012) 32



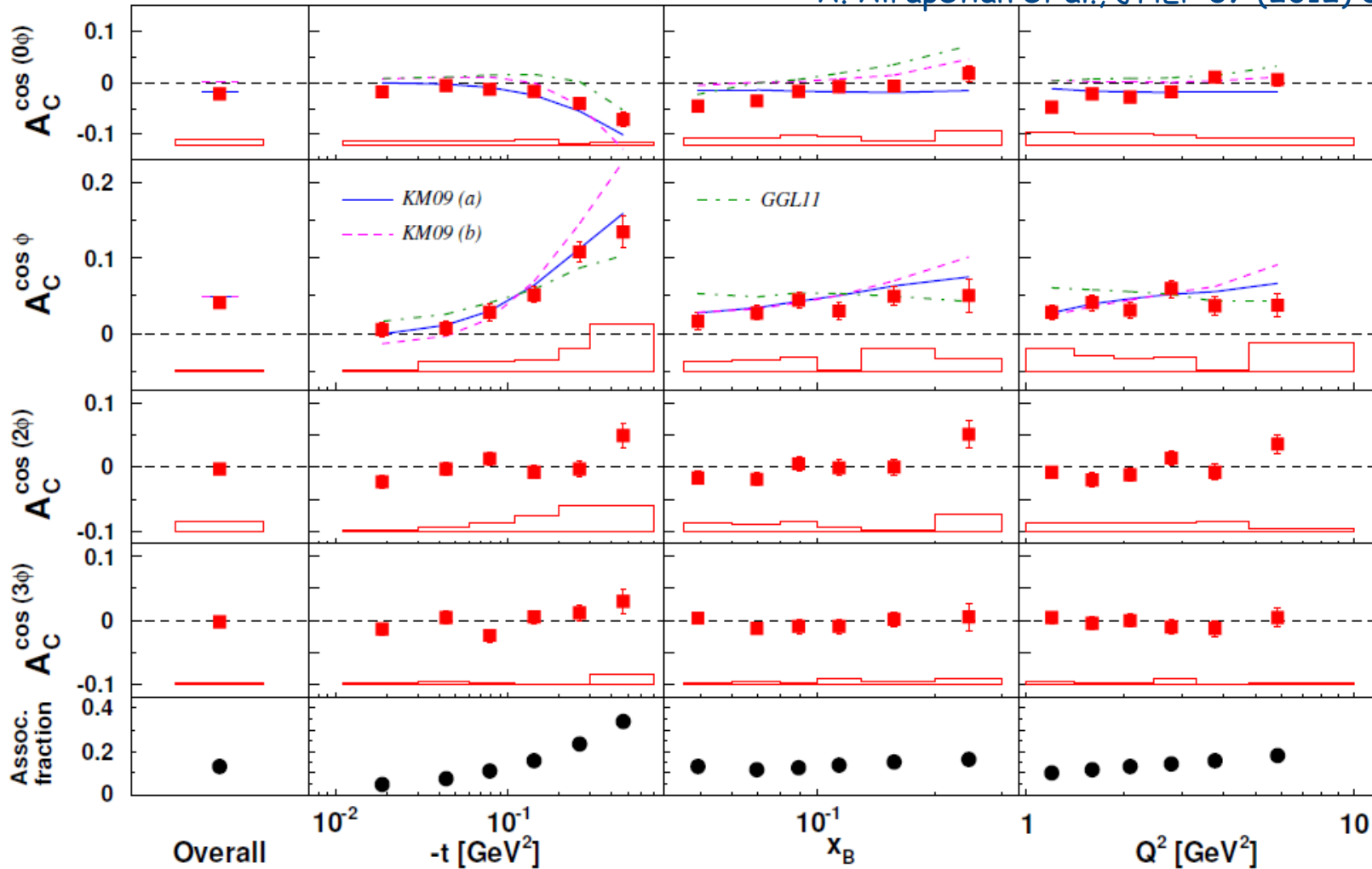
KM09: Nucl. Phys. B
841 (2010) 1:
 fit to HERMES, ZEUS,
 H1 data
 Fit to HERMES, ZEUS,
 H1, Jefferson Lab
 data
 GGL11: Phys. Rev. D
84 (2011) 034007

- data collected from 1996-2007 (74% of data from 2006-2007)
- additional 3.2% scale uncertainty from beam polarization

Beam-charge asymmetry

GPD H

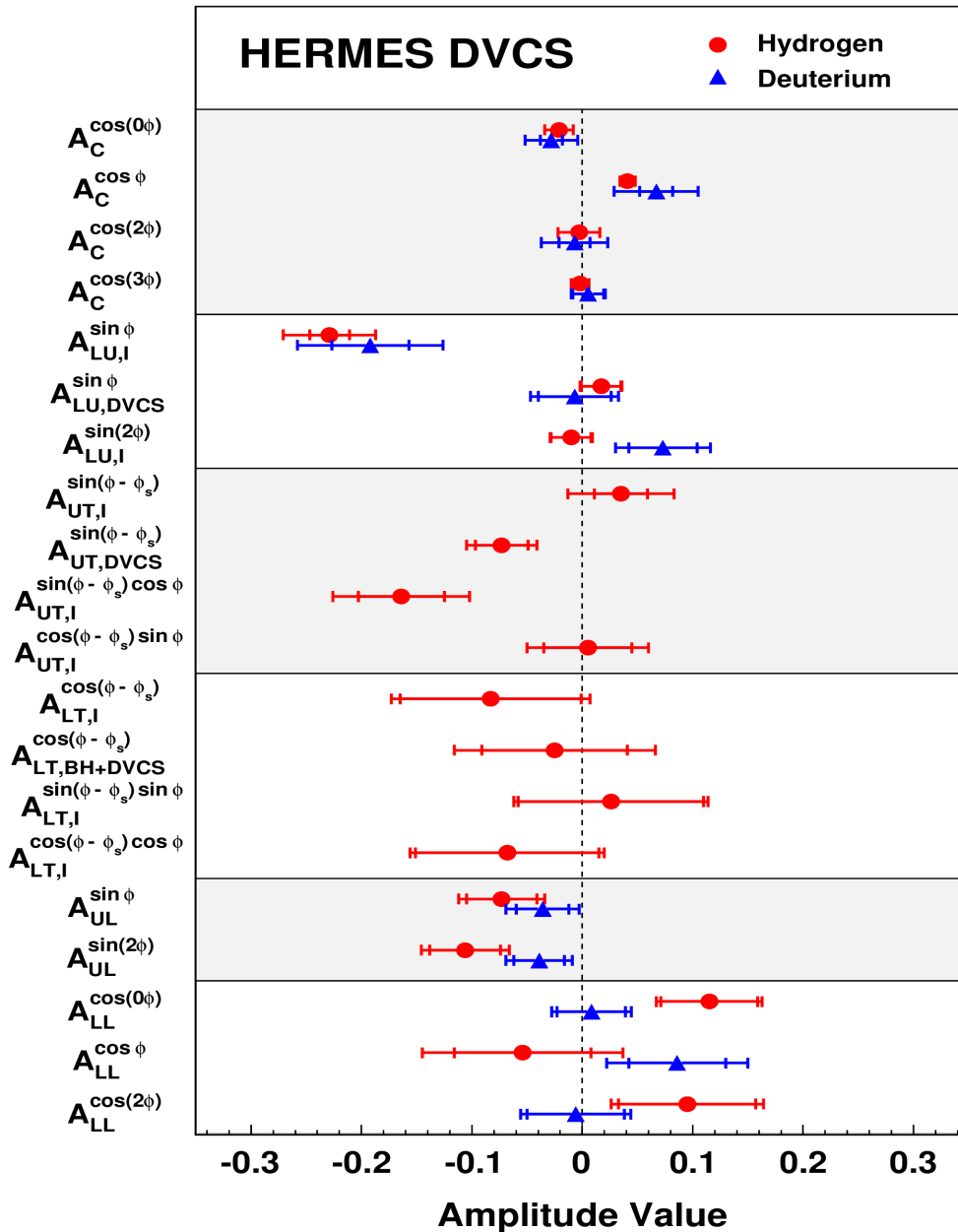
A. Airapetian et al., JHEP 07 (2012) 32



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DVCS at HERMES



beam-charge asymmetry
 JHEP **07** (2012) 32
 Nucl. Phys. B 829 (2010) 1

GPD H

beam-helicity asymmetry
 JHEP **07** (2012) 32
 Nucl. Phys. B 829 (2010) 1

transverse target-spin asymmetry
 JHEP **06** (2008) 066

GPD E

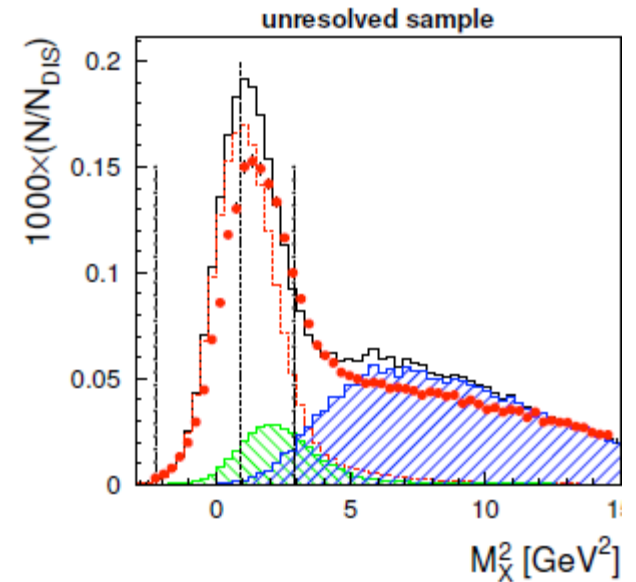
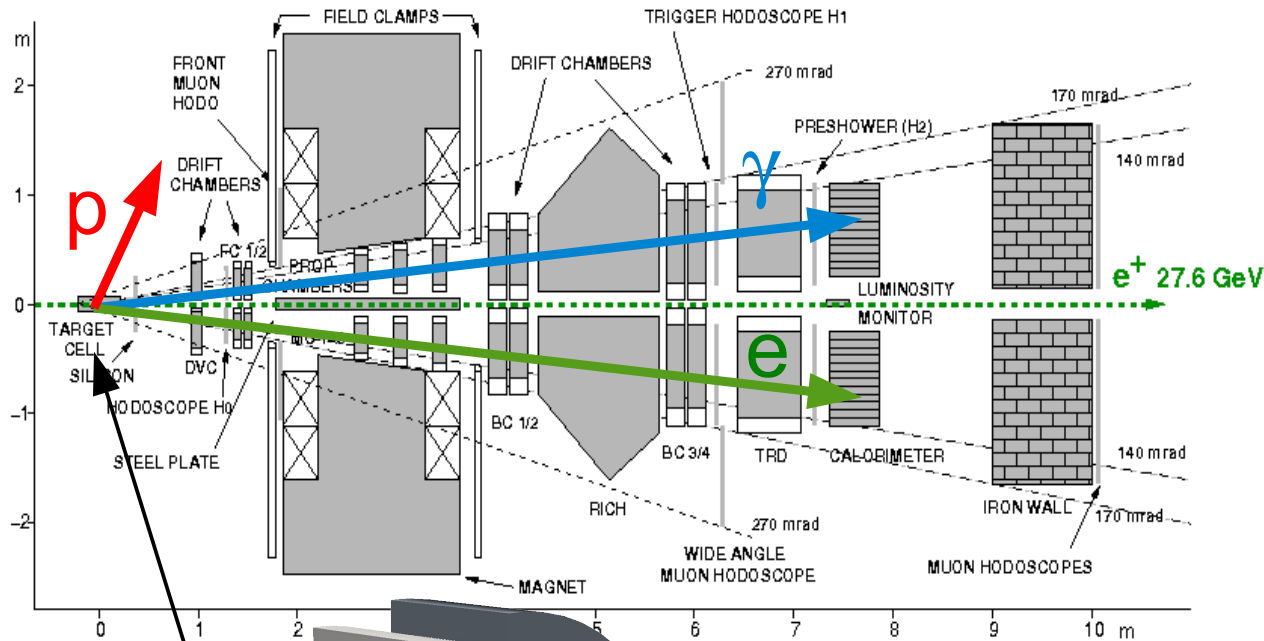
double spin (LT) asymmetry
 Phys. Lett. B **704** (2011) 15

longitudinal target-spin asymmetry
 JHEP **06** (2010) 019
 Nucl. Phys. B 842 (2011) 265

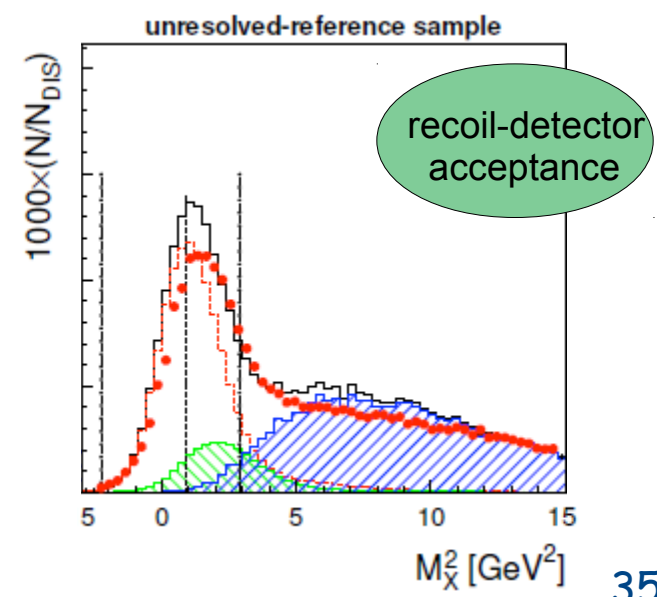
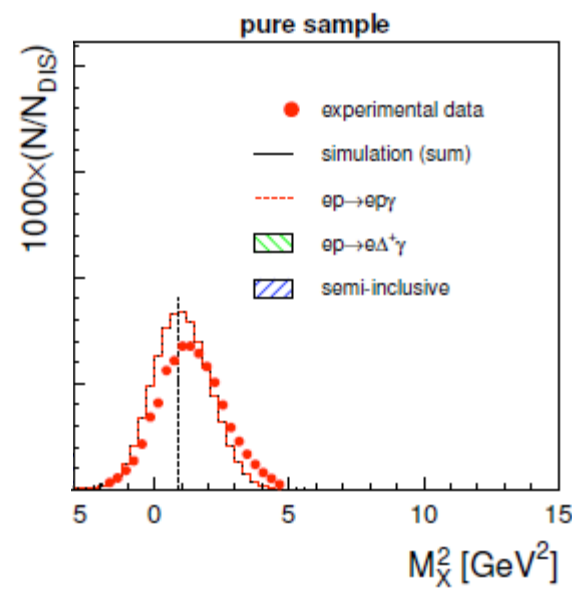
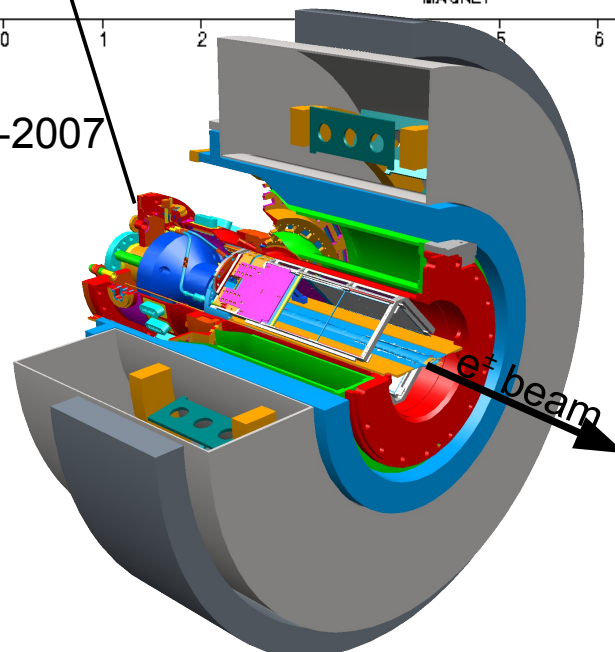
GPD \tilde{H}

double spin (LL) asymmetry
 JHEP **06** (2010) 019
 Nucl. Phys. B 842 (2011) 265

DVCS/BH complete event reconstruction

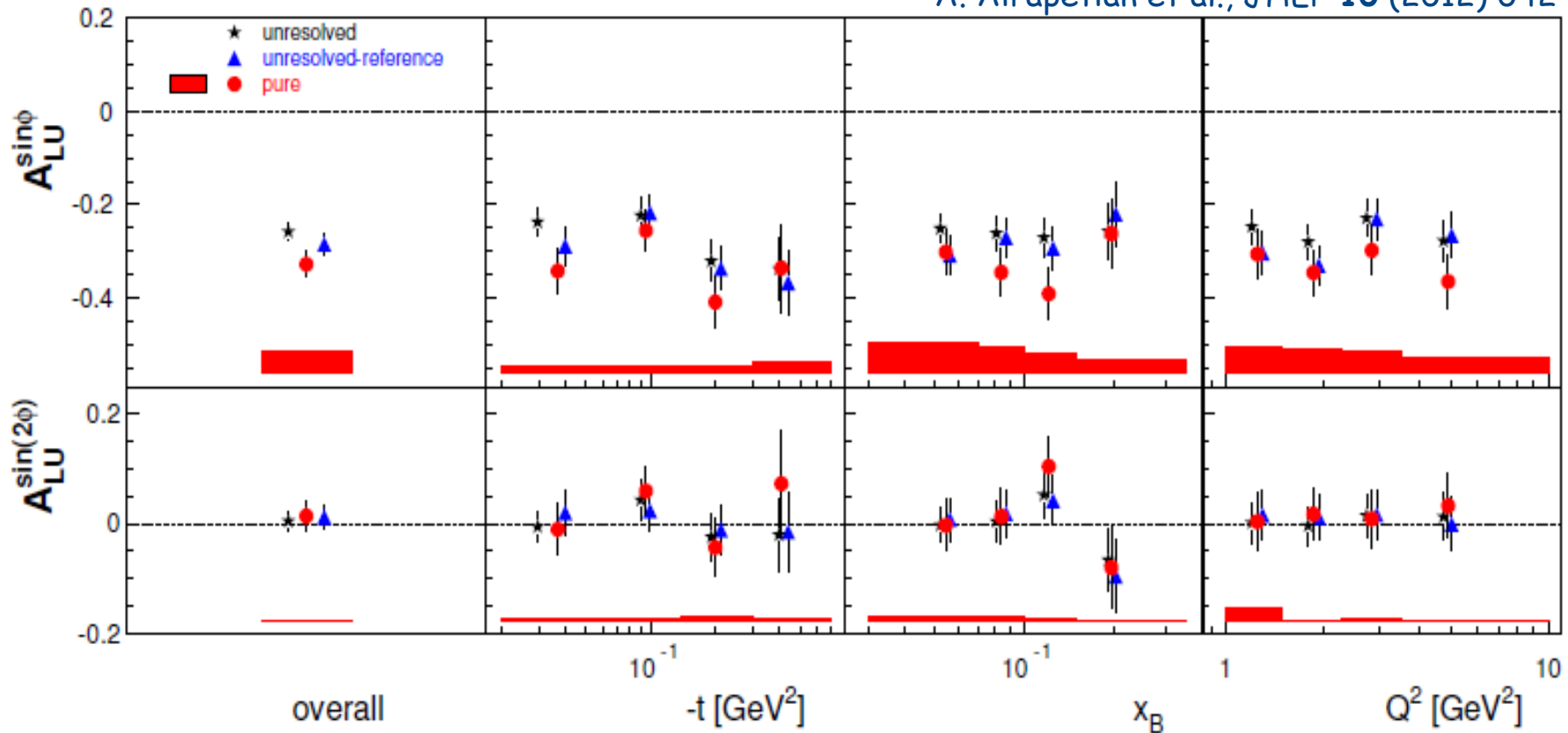


2006-2007



Beam-helicity asymmetry

A. Airapetian et al., JHEP 10 (2012) 042



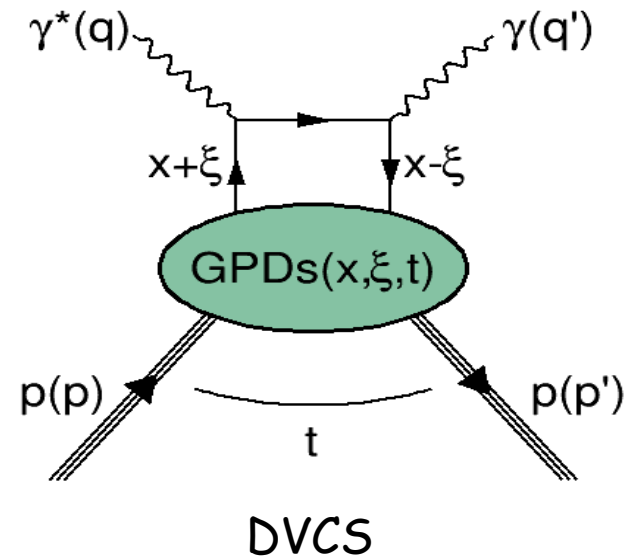
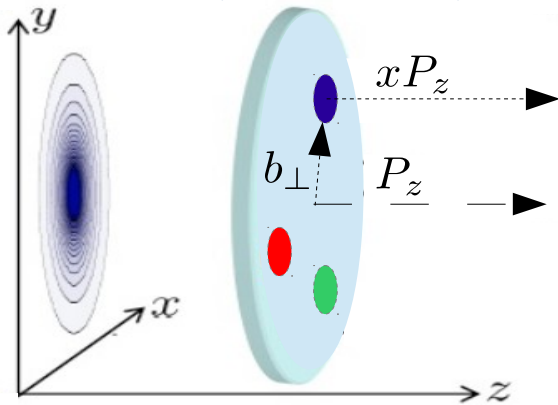
- additional 1.96 % scale uncertainty from beam polarization
- leading asymmetry from pure sample is larger

Beam-helicity asymmetry in

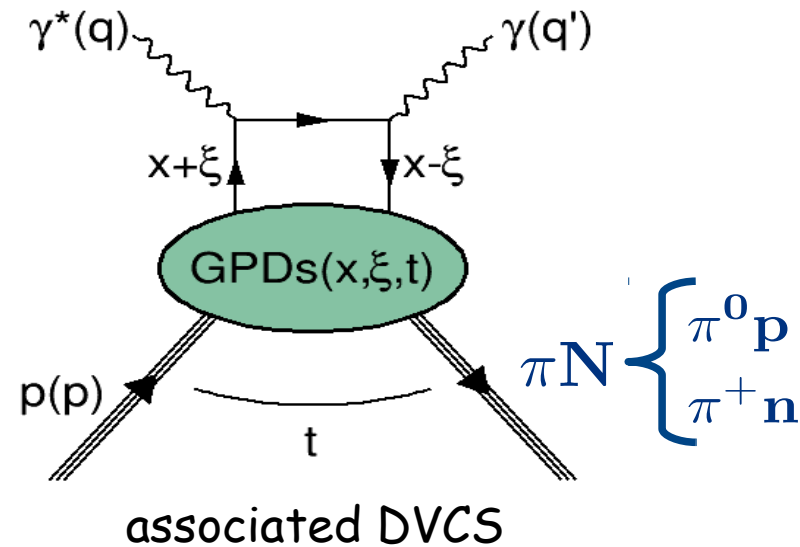
$ep \rightarrow e\gamma\pi N$ in Δ -resonance region

GPDs

quark distribution in longitudinal-momentum and transverse-position space

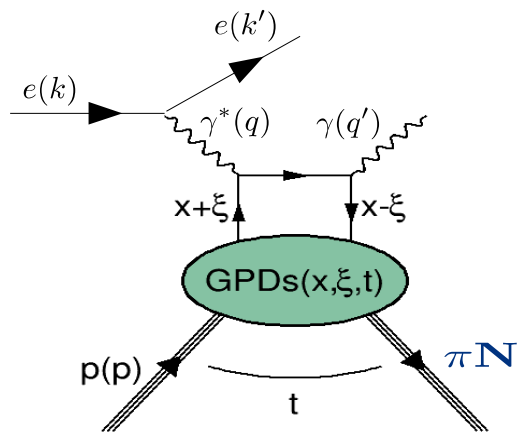


transition GPDs $p \rightarrow \pi N$

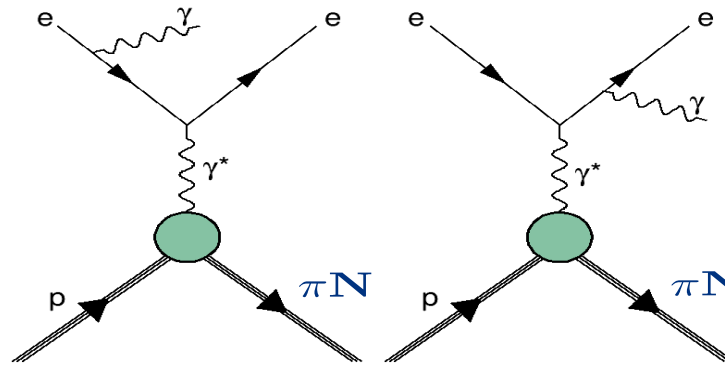


Beam-helicity asymmetry in

$ep \rightarrow e\gamma\pi N$ in Δ -resonance region



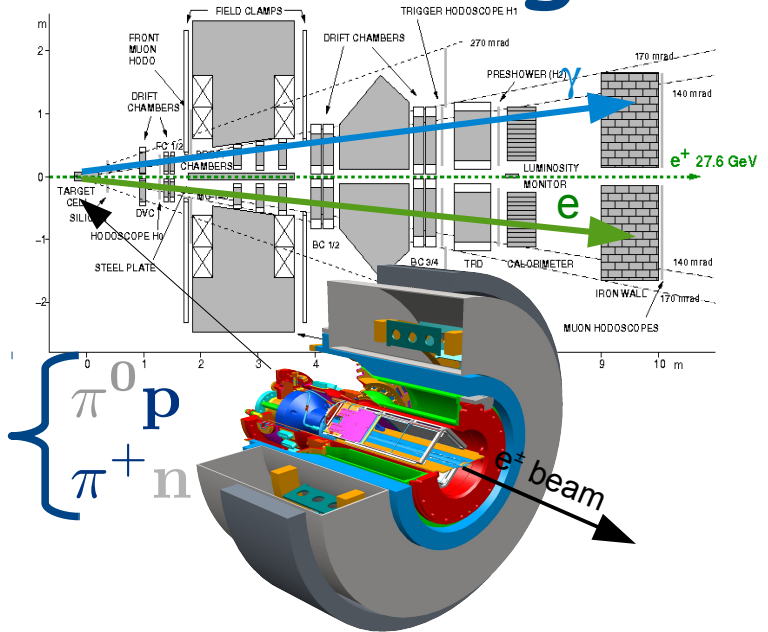
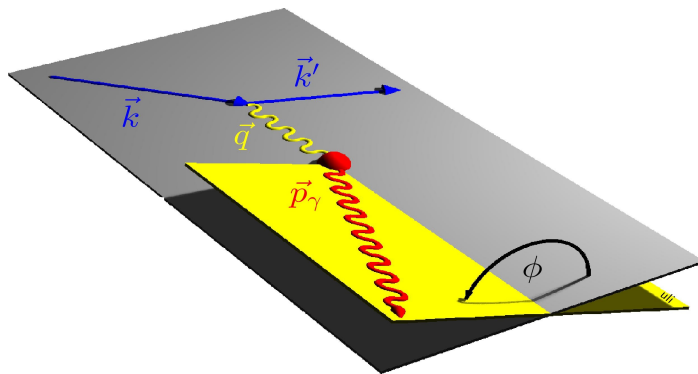
associated DVCS



associated Bethe-Heitler

$$d\sigma \propto |\tau_{BH}|^2 + |\tau_{DVCS}|^2 + \tau_{BH} \tau_{DVCS}^* + \tau_{DVCS} \tau_{BH}^*$$

access through azimuthal asymmetries

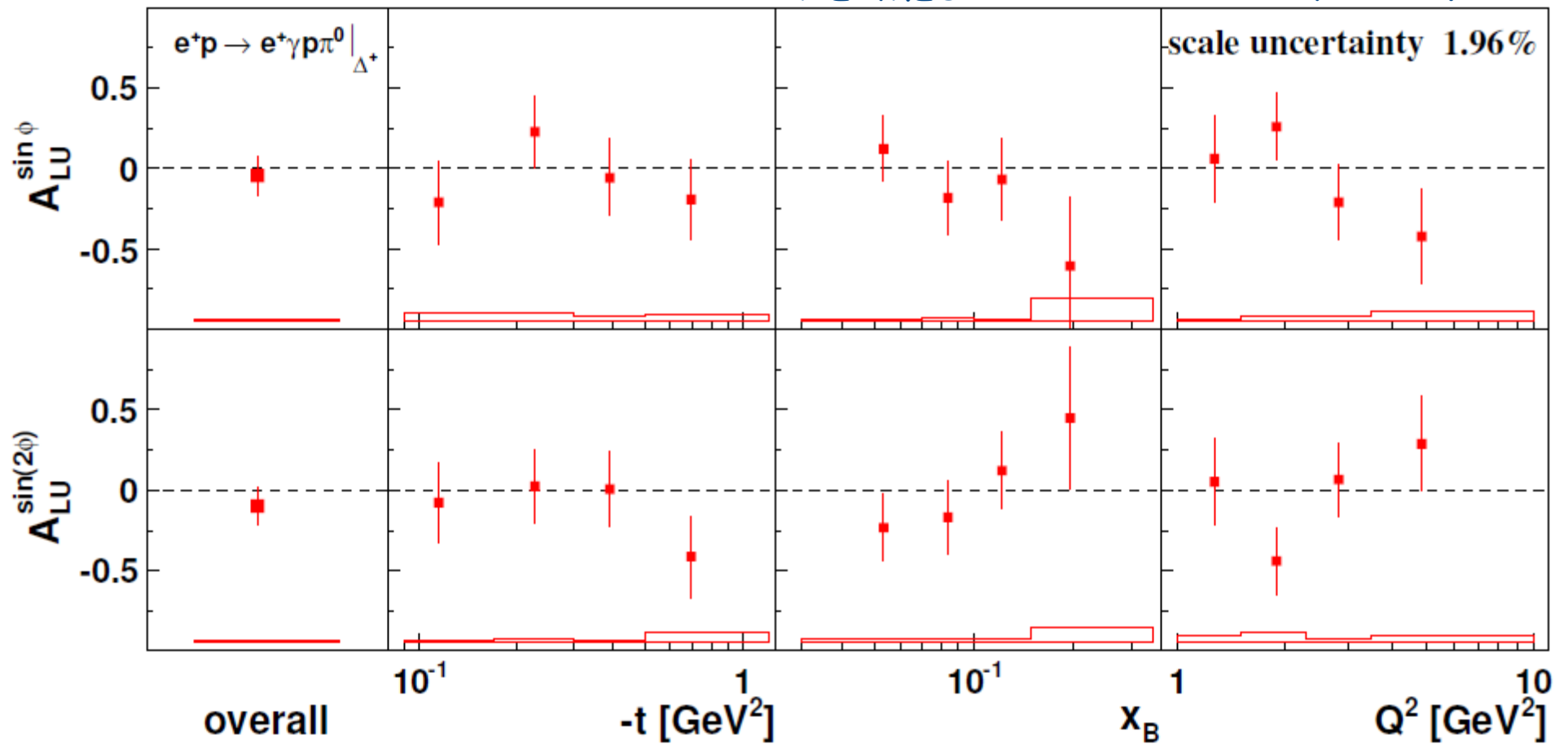


recoil detector

- p and π^+ via PID
- π^0 and n undetected
- kinematic fitting
- select region around Δ -resonance

Beam-helicity asymmetry in $ep \rightarrow e\gamma\pi^0 p$ in Δ -resonance region

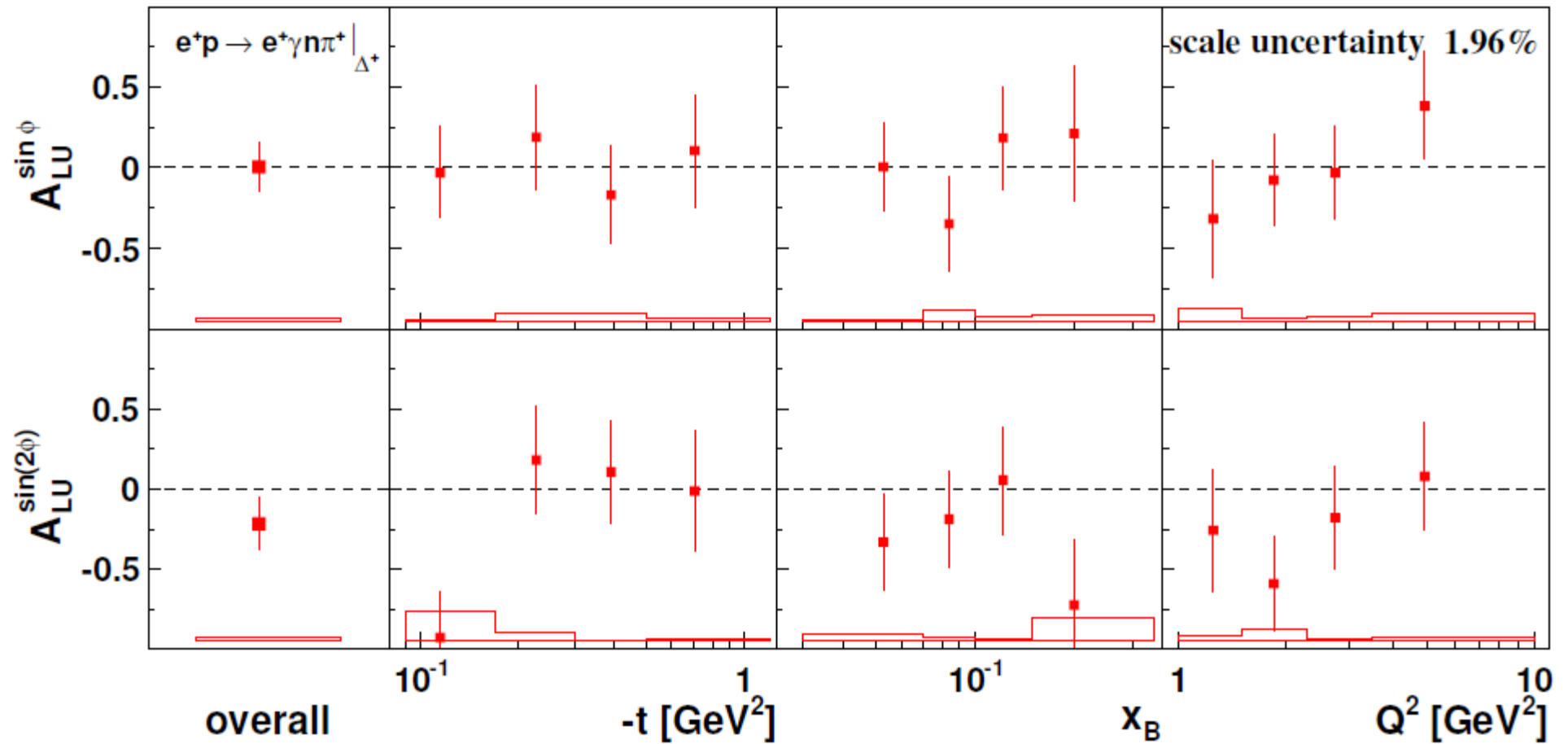
HERMES Collaboration - arXiv:1310.5081



- asymmetry background correction from SIDIS (11%) and $ep \rightarrow e\gamma p$ (4.6%)
- leading asymmetry consistent with zero

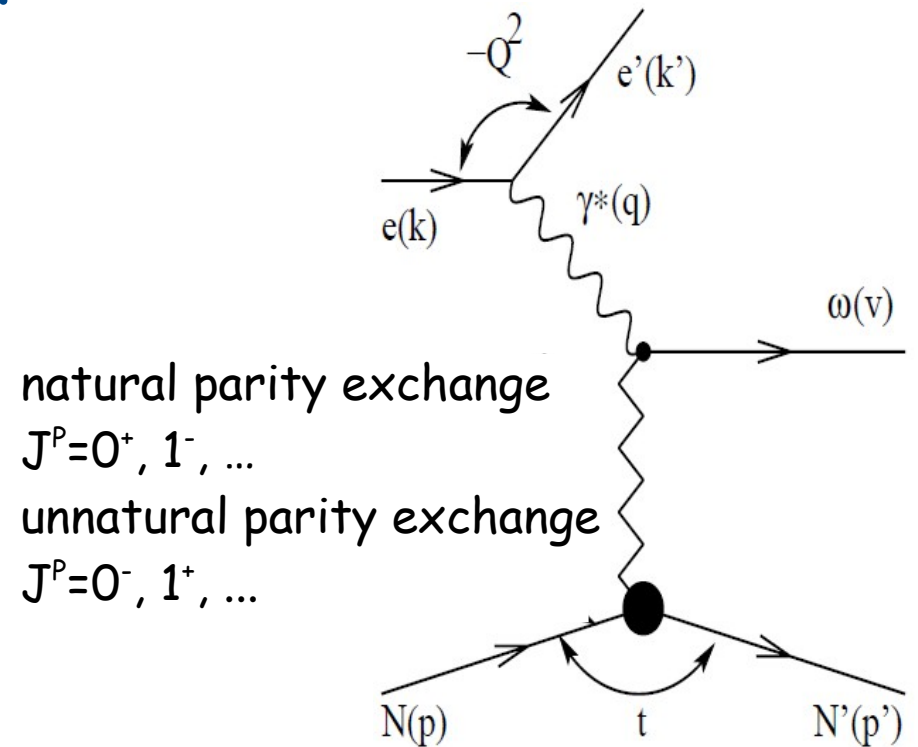
Beam-helicity asymmetry in $ep \rightarrow e\gamma\pi^+n$ in Δ -resonance region

HERMES Collaboration - arXiv:1310.5081



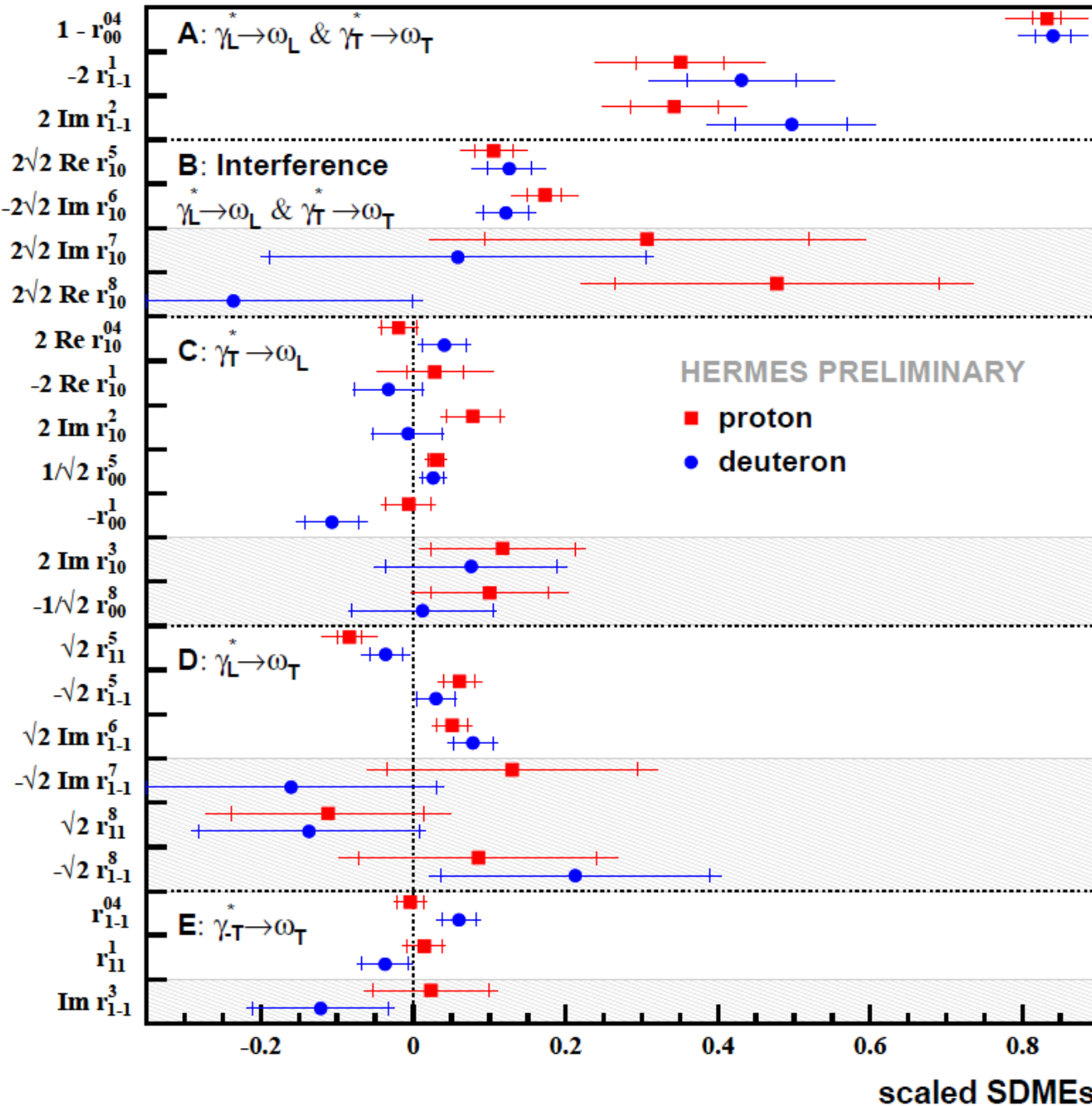
- asymmetry background correction from SIDIS (23%) and $ep \rightarrow e\gamma p$ (0.2%)
- leading asymmetry consistent with zero

Exclusive ω production



Exclusive ω production

- compatible SDMEs for proton and deuteron



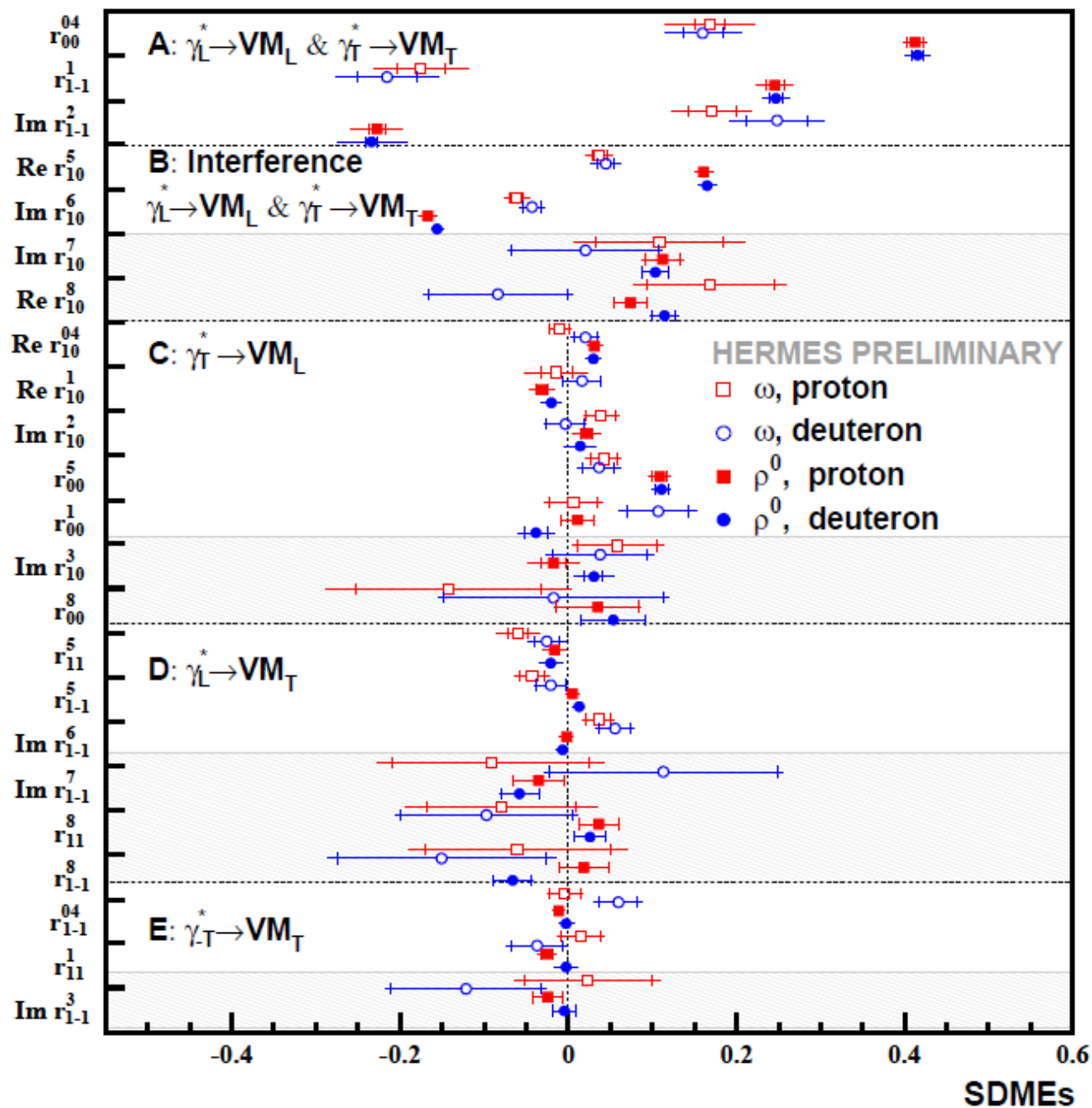
} slight violation of s-channel helicity conservation:

$$r_{11}^5 + r_{1-1}^5 - \Im r_{1-1}^6$$

$$= -0.14 \pm 0.02 \pm 0.04$$

$$= -0.10 \pm 0.03 \pm 0.03$$

Exclusive ω production



• ω - ρ^0 comparison:

• ω : $r_{1-1}^1 < 0$

• ρ^0 : $r_{1-1}^1 > 0$

• ω : $\Im r_{1-1}^2 > 0$

• ρ^0 : $\Im r_{1-1}^2 < 0$

large unnatural parity
exchange for ω production

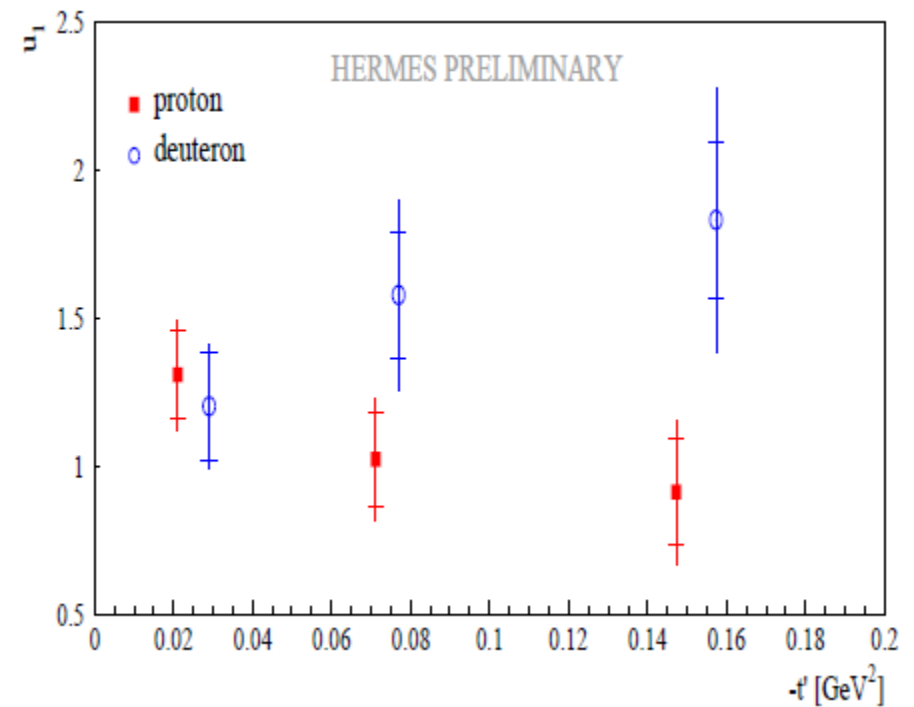
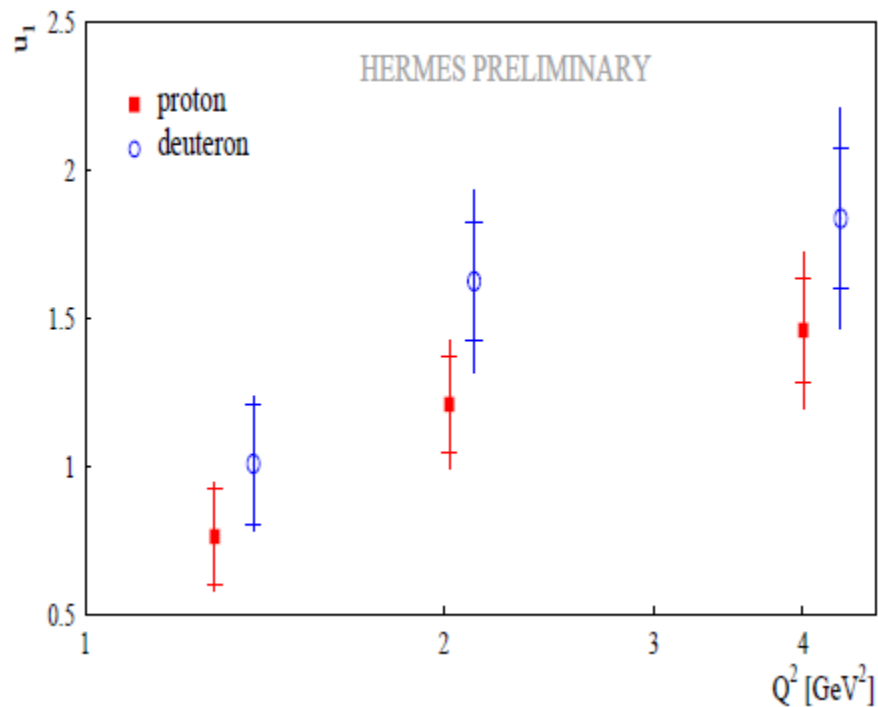
exclusive ρ^0 :

A. Airapetian et al., Eur. Phys. J. C 62
(2009) 659

Exclusive ω production

test of unnatural parity exchange:

$$\begin{aligned} u_1 &= 1 - r_{00}^{04} + 2 r_{1-1}^{04} - 2 r_{11}^1 - 2 r_{1-1}^1 \\ &\propto 2\epsilon |U_{10}|^2 + |U_{11} + U_{-11}|^2 \end{aligned}$$



large unnatural-parity exchange

Summary

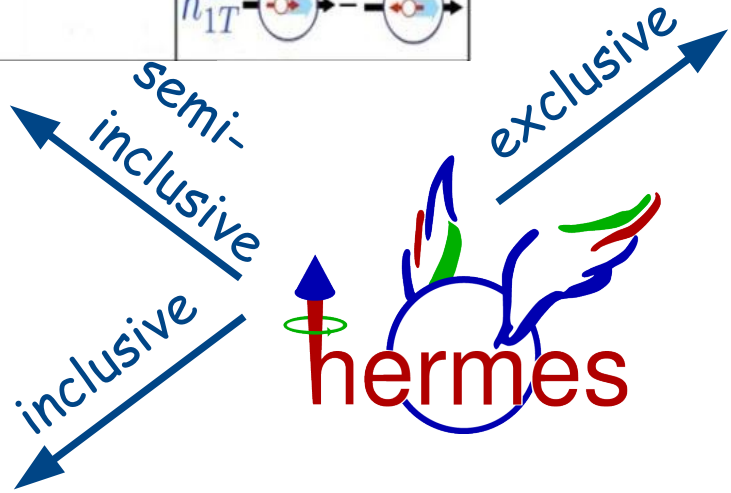
nucleon	quark	L		T		
		U	D	U	D	
U	f_1			h_1^\perp		
L	g_1			h_{1L}^\perp		
T	f_{1T}^\perp			g_{1T}^\perp	h_{1T}^\perp	h_{1T}^\perp

$P_{h\perp}$ dependence:
PRD 87 (2013) 074029

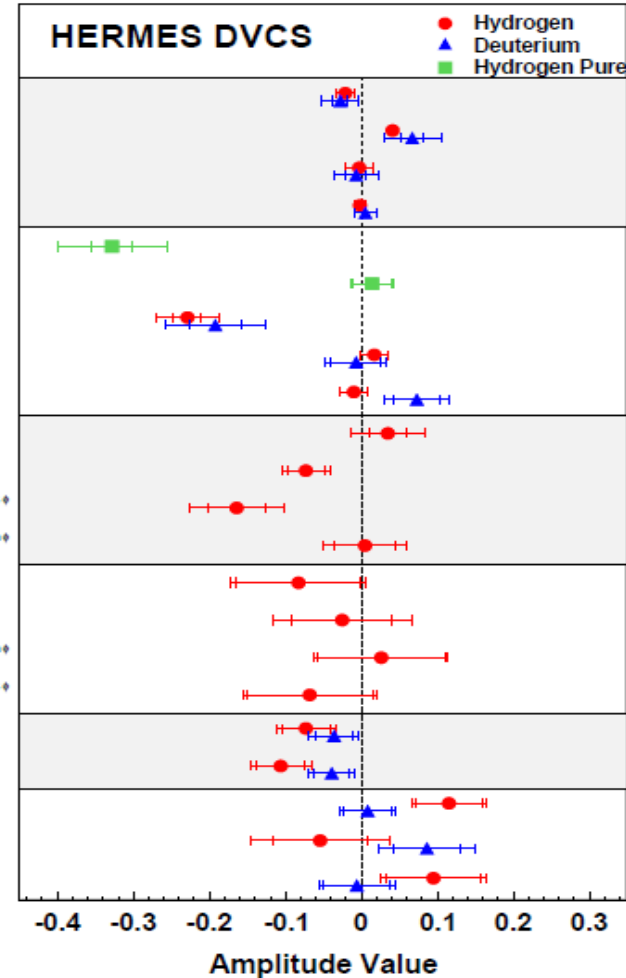
PRD 87 (2013) 012010

PRL 94 (2005) 012002
PRL 103 (2009) 152002

exclusive ω production

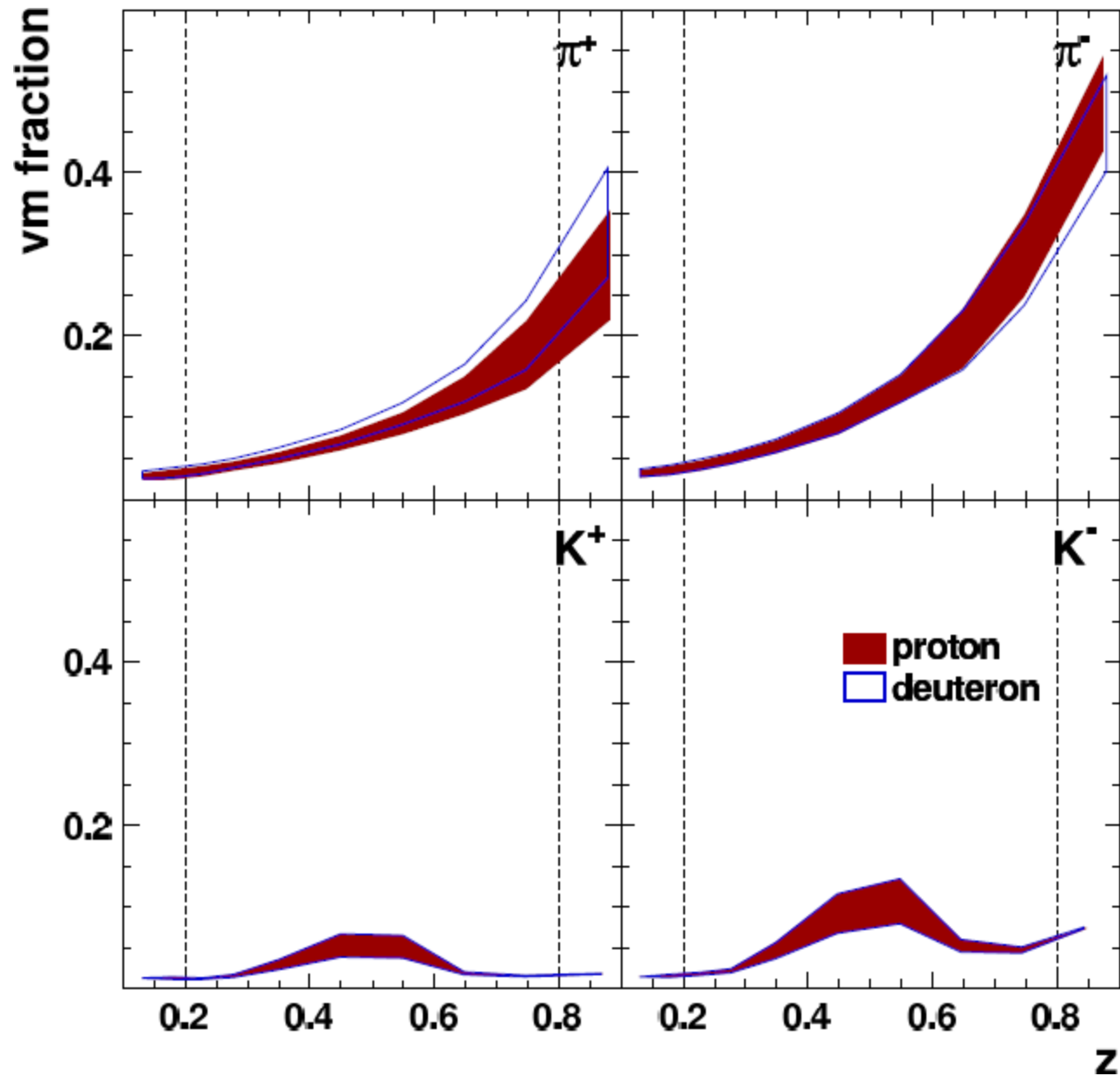


- g_2
- A_{UT} in inclusive electroproduction of pions and protons, see talk T. Shibata-san

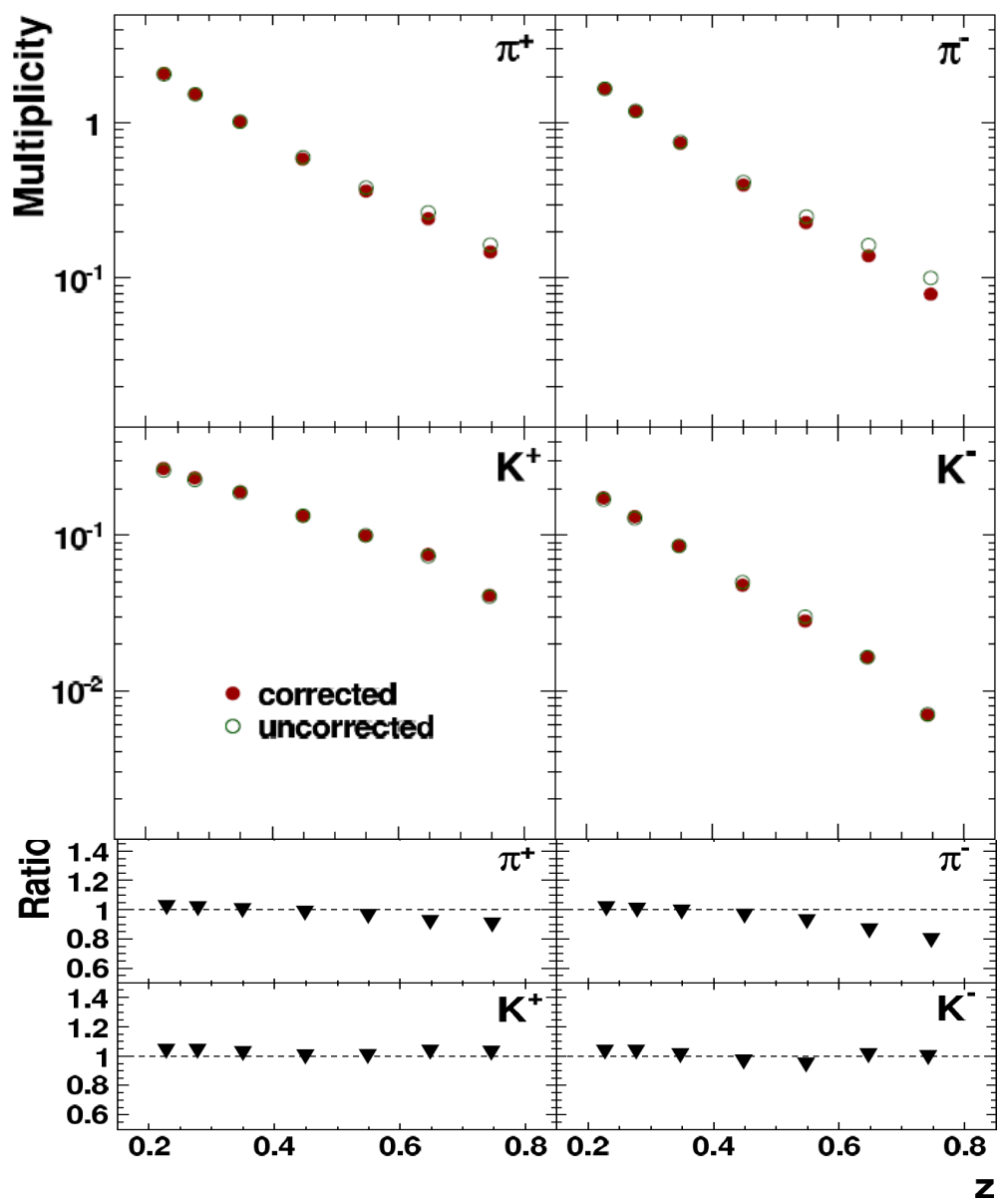


Backup

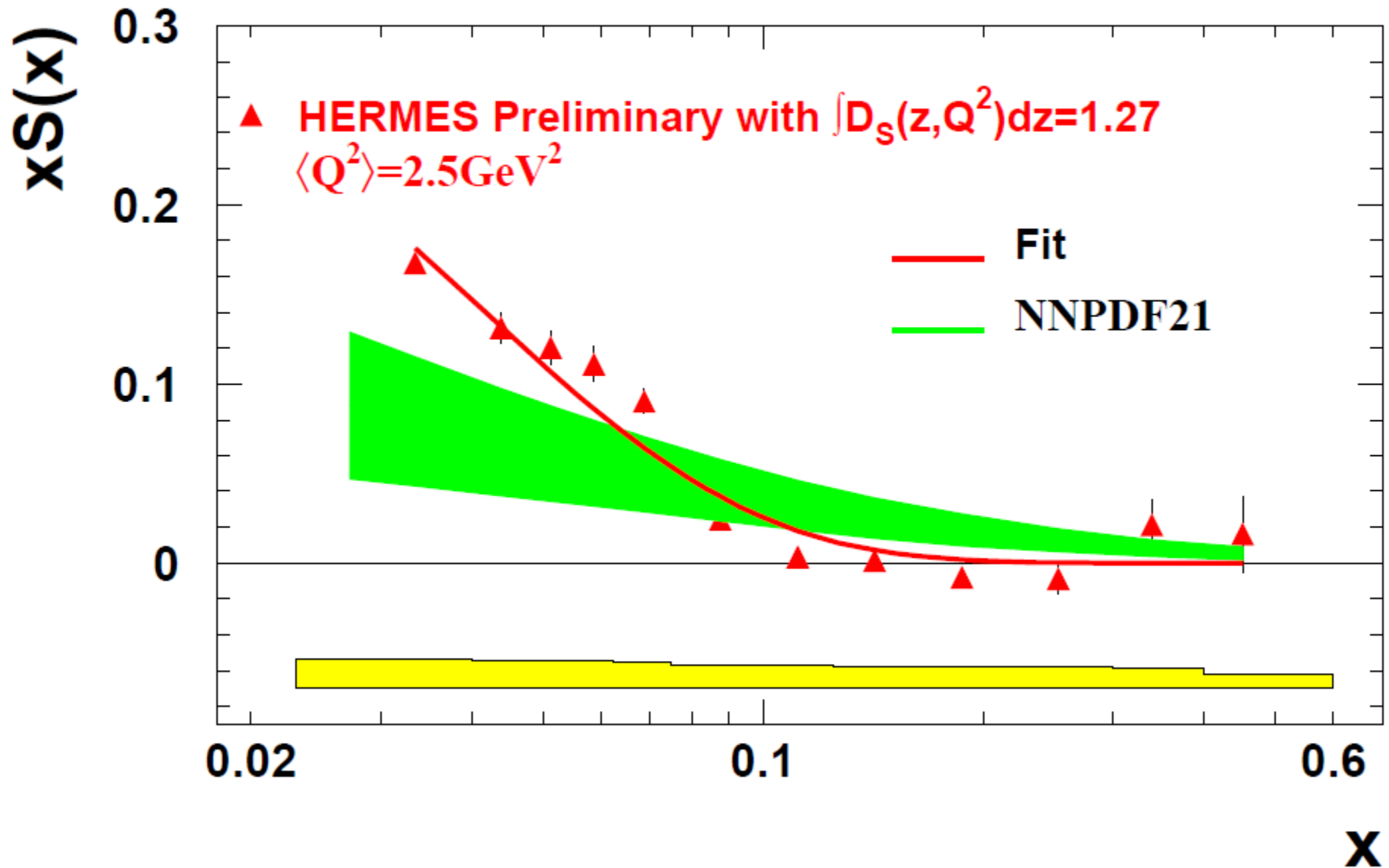
Exclusive vector-meson fraction



Multiplicities corrected in z: without and with exclusive vector-meson correction

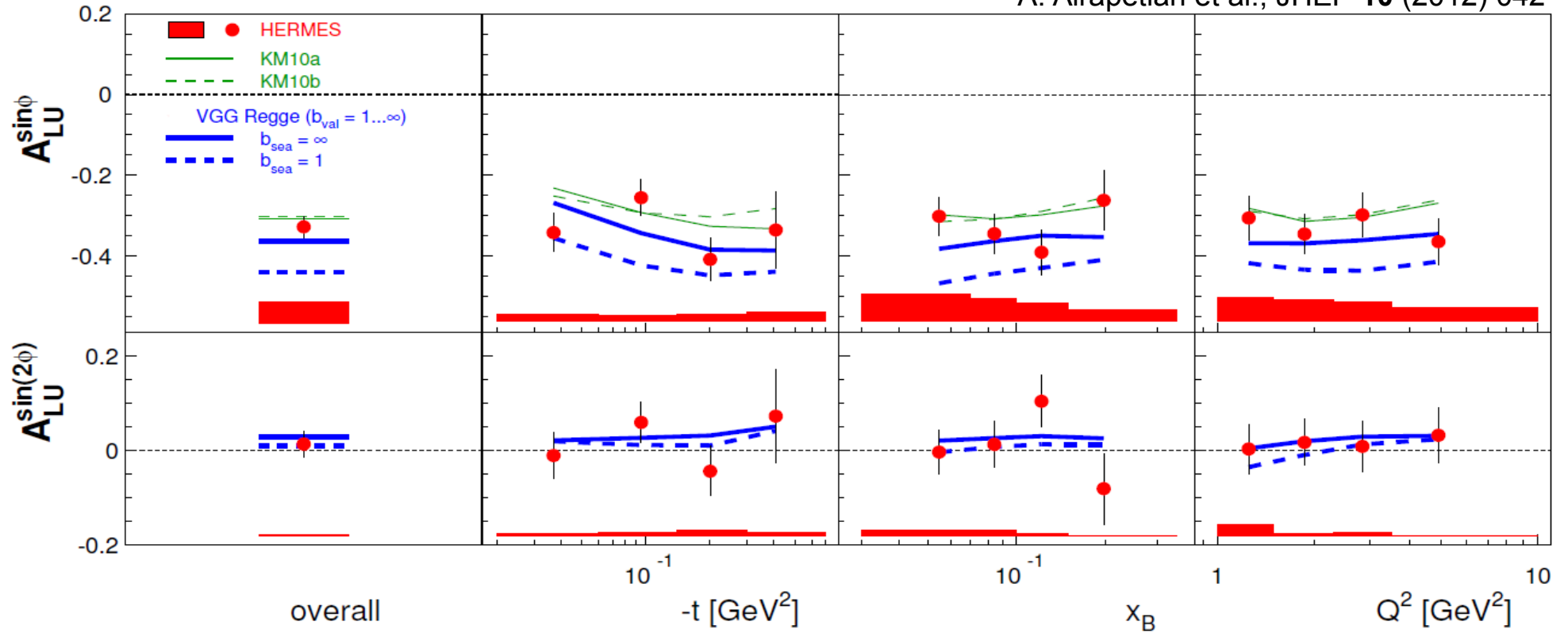


Comparison $S(x)$ with NNPDF21



Beam-helicity asymmetry

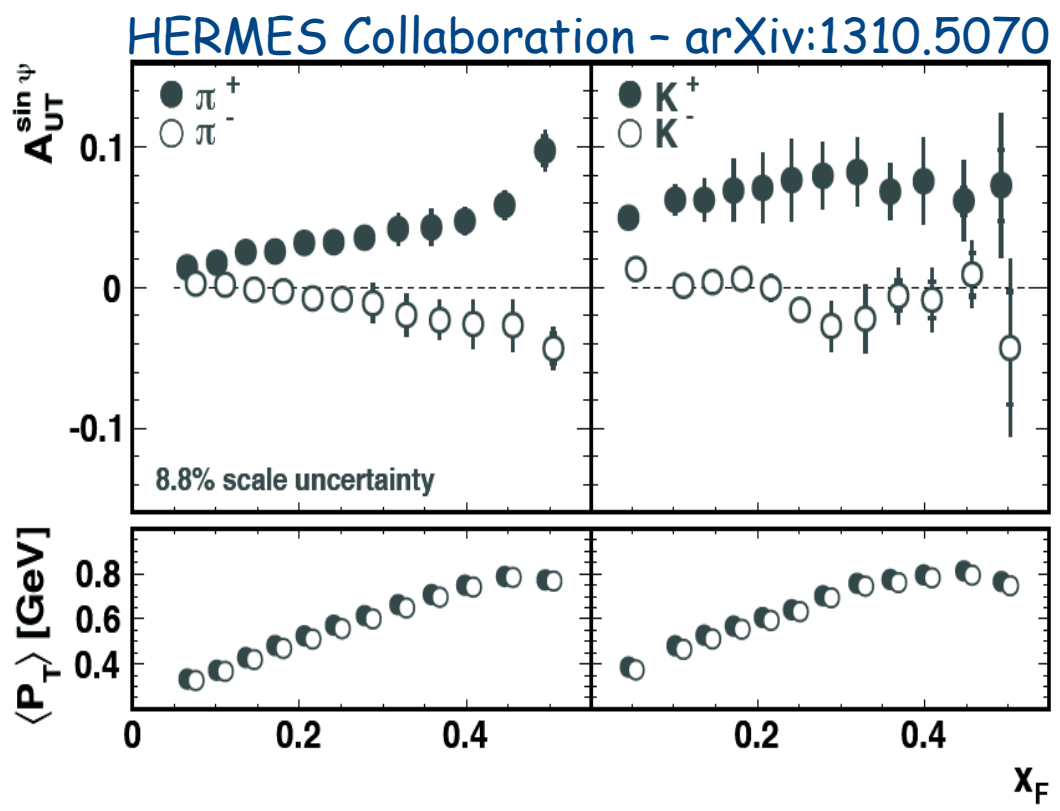
A. Airapetian et al., JHEP **10** (2012) 042



- additional 1.96 % scale uncertainty from beam polarization

Transverse target single-spin asymmetry in inclusive electroproduction of pions and kaons

$x_F = 2P_L/\sqrt{s}$ dependence



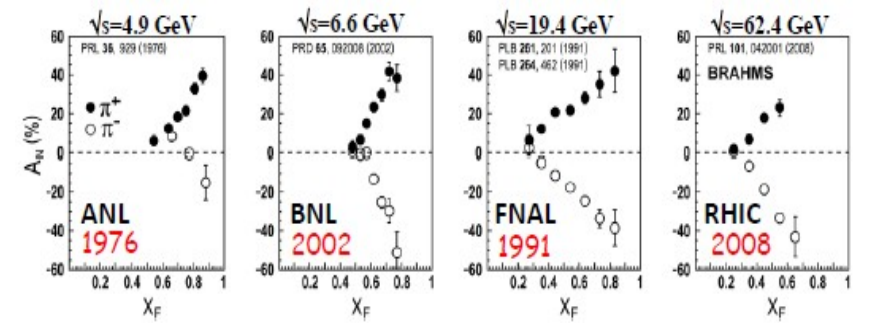
π^+

- positive, increase linearly with x_F

π^-

- negative, decrease linearly with x_F

x_F behavior of pions similar to what observed in hadron-hadron collisions



K^+

- positive, ~constant with x_F

K^-

- compatible with zero, with small variation over x_F

Transverse target single-spin asymmetry in inclusive electroproduction of pions and kaons

disentangle x_F and P_T dependence

π^+

- increase with P_T up to $P_T \approx 0.8$ GeV
- P_T dependence independent of x_F
 $\rightarrow x_F$ increase from P_T dependence

π^-

- small amplitudes, varyingly positive and negative with P_T
- decrease with increasing x_F

K^+

- increase with P_T up to $P_T \approx 0.8$ GeV
- increase with increasing x_F

K^-

- small amplitudes
- decrease with increasing x_F

HERMES Collaboration - arXiv:1310.5070

