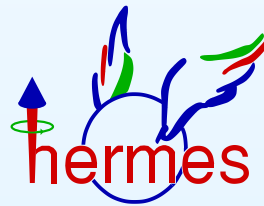


***Measurement of the proton spin structure function g_2^P
and asymmetry A_2^P at the HERMES experiment***

Markus Diefenthaler

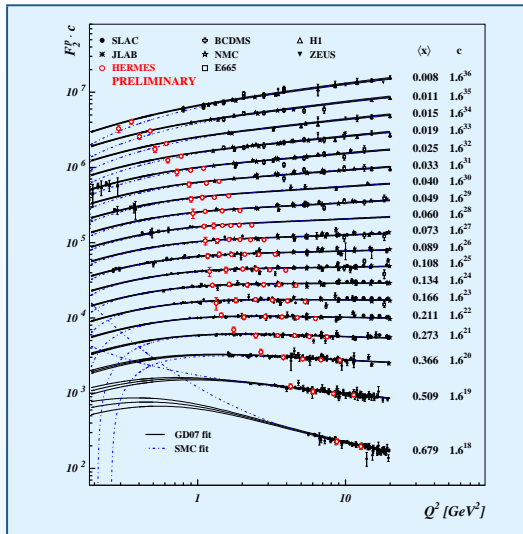


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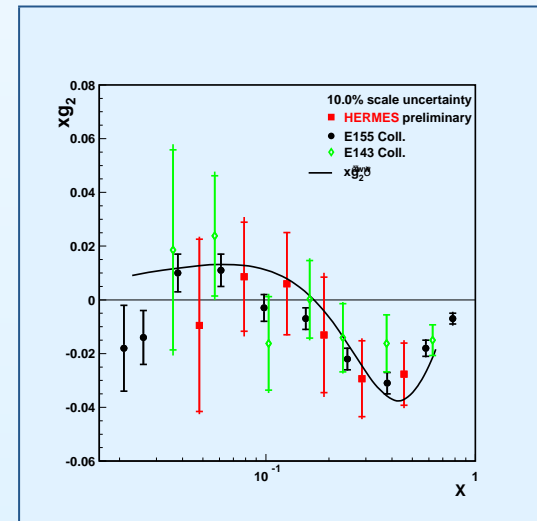
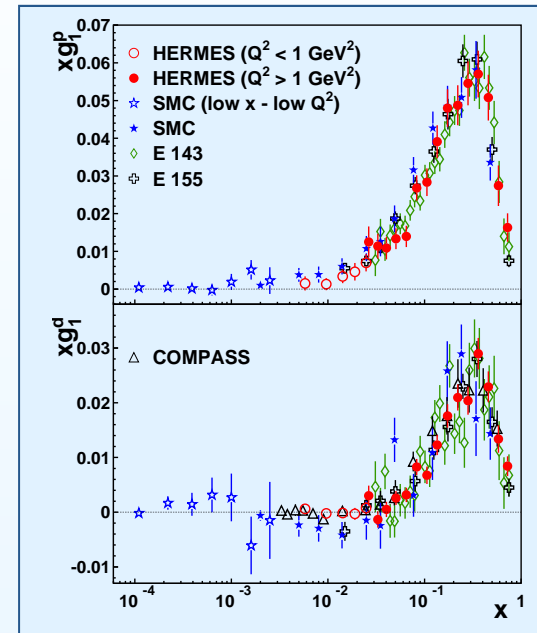
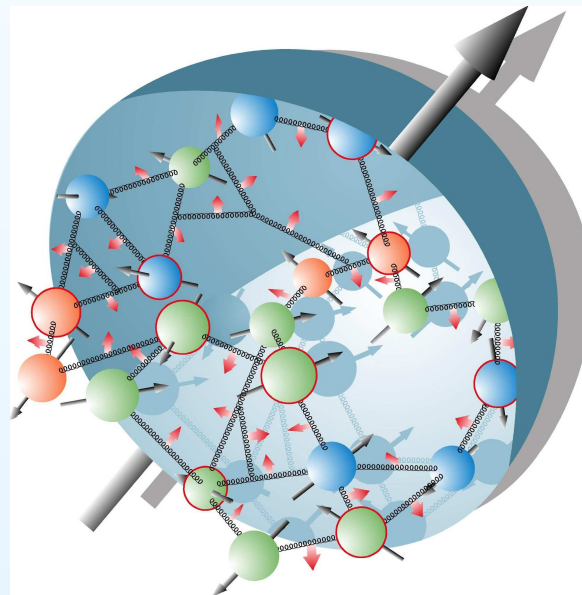


on behalf of the hermes collaboration

The HERA legacy:

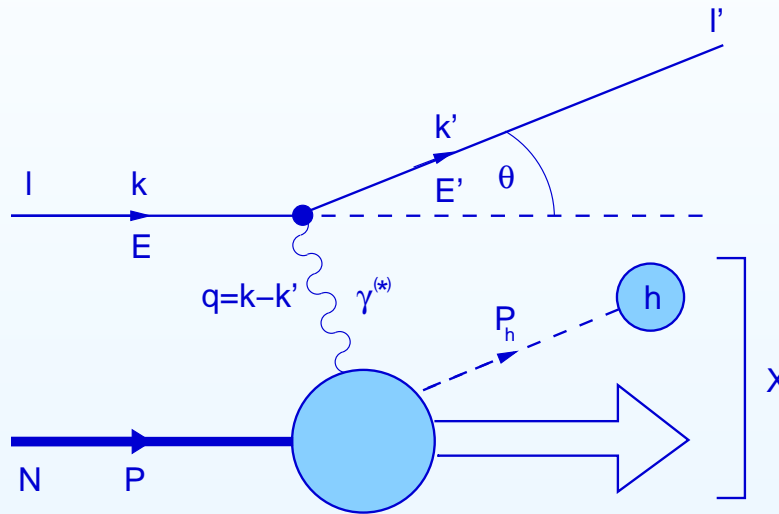


$$F_2(x) = 2xF_1(x)$$



The polarised deep-inelastic scattering process:

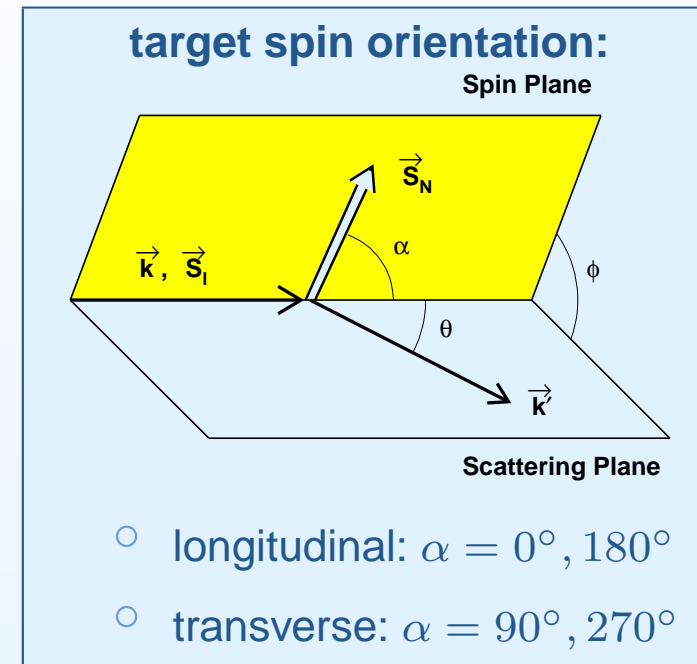
- inclusive measurement ($lN \rightarrow l'X$):



$$Q^2 \equiv -q^2, \quad x = Q^2 / (2P \cdot q)$$

- spin-dependent cross-section contribution:

$$\frac{d^3 (\sigma(\alpha) - \sigma(\alpha + \pi))}{dx dy d\phi_S} = \frac{e^4}{2\pi^2 Q^2} \left\{ \cos \alpha \left[\left(1 - \frac{y}{2} - \frac{y^2 \gamma^2}{4}\right) \mathbf{g}_1(x, Q^2) - \left(\frac{y}{2} \gamma^2\right) \mathbf{g}_2(x, Q^2) \right] - \sin \alpha \cos \phi_S \gamma \sqrt{1 - y - \frac{y^2 \gamma^2}{4}} \left[\frac{y}{2} \mathbf{g}_1(x, Q^2) + \mathbf{g}_2(x, Q^2) \right] \right\}$$



The structure function g_2 :

- **Wandzura-Wilczek decomposition:**

$$g_2(x, Q^2) = g_2^{WW}(x, Q^2) + \bar{g}_2(x, Q^2)$$

- $g_2^{WW}(x, Q^2) = -g_1(x, Q^2) + \int_x^1 \frac{dy}{y} g_2(y, Q^2)$

- pure **twist-three contribution \bar{g}_2**

- **probing quark-gluon correlations**

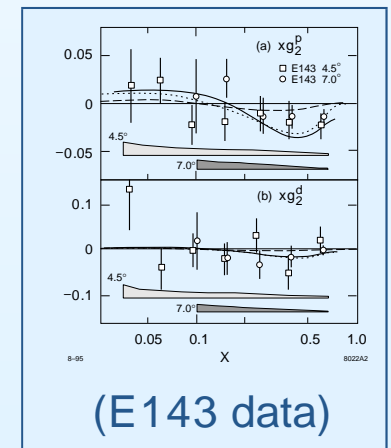
- $\int_0^1 dx x^n \bar{g}_2(x, Q^2) = \frac{n}{4(n+1)} d_n(Q^2)$

- d_2 : Lorentz-force acting on quark

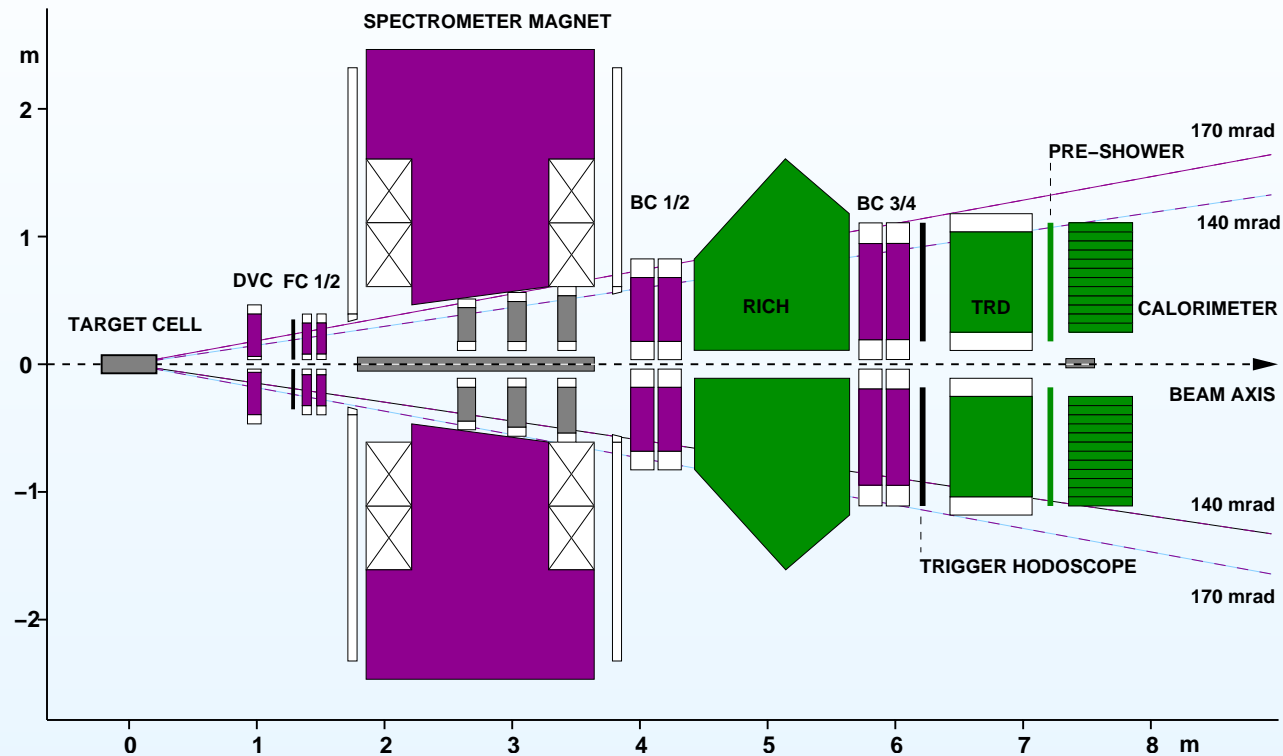
- **Burkhardt-Cottingham sum rule:**

$$\int_0^1 dx g_2(x, Q^2) = 0$$

➔ nodes? (besides $x = 1$ and perhaps $x = 0$)



The HERMES spectrometer:



- polarised **hydrogen target** internal to the HERA storage ring
 - ↳ background-free measurements from highly polarised protons
 - ↳ substantial reduction of time-dependent systematics
- very clean lepton-hadron separation
- 2003–2005: **transversely polarised target**: $\langle P_T \rangle = 0.71 \pm 0.06$

The measurement of double-spin asymmetries:

- **lepton-beam asymmetries:**

$$A_{\parallel} = \frac{\sigma^{\rightarrow\leftarrow} - \sigma^{\rightarrow\Rightarrow}}{\sigma^{\rightarrow\leftarrow} + \sigma^{\rightarrow\Rightarrow}} = D(A_1 + \eta A_2)$$

$$A_{\perp} = \frac{\sigma^{\rightarrow\uparrow} - \sigma^{\rightarrow\downarrow}}{\sigma^{\rightarrow\uparrow} + \sigma^{\rightarrow\downarrow}} = d(A_2 - \zeta A_1)$$

- **virtual-photon asymmetries:**

$$A_1 = \frac{\sigma_{\frac{1}{2}} - \sigma_{\frac{3}{2}}}{\sigma_{\frac{1}{2}} + \sigma_{\frac{3}{2}}} = \frac{g_1 - \gamma^2 g_2}{F_1}$$

$$A_2 = \frac{\sigma_{\text{LT}}}{\sigma_{\text{T}}} = \frac{\gamma(g_1 + g_2)}{F_1}$$

The measurement of A_{\perp} :

- HERA beam facility: $\sigma^{\rightarrow\uparrow}$, $\sigma^{\rightarrow\downarrow}$ and $\sigma^{\leftarrow\uparrow}$, $\sigma^{\leftarrow\downarrow}$

$$\begin{aligned} A_1(\phi_S, x, Q^2) &= \frac{\sigma^{\rightarrow\uparrow}(\phi_S, x, Q^2) - \sigma^{\rightarrow\downarrow}(\phi_S, x, Q^2)}{\sigma^{\rightarrow\uparrow}(\phi_S, x, Q^2) + \sigma^{\rightarrow\downarrow}(\phi_S, x, Q^2)} \\ &= - \frac{\sigma^{\leftarrow\uparrow}(\phi_S, x, Q^2) - \sigma^{\leftarrow\downarrow}(\phi_S, x, Q^2)}{\sigma^{\leftarrow\uparrow}(\phi_S, x, Q^2) + \sigma^{\leftarrow\downarrow}(\phi_S, x, Q^2)} \\ &= \cos\phi_S A_{\perp}(x, Q^2) \end{aligned}$$

- A_{\perp} reconstructed from luminosity-normalised count-rates:

$$A_{\perp}^{(\leftrightarrow)}(\phi_S, x, Q^2) = \frac{1}{|P_B P_T|} \frac{N^{(\leftrightarrow)\uparrow}(\phi_S, x, Q^2) - N^{(\leftrightarrow)\downarrow}(\phi_S, x, Q^2)}{N^{(\leftrightarrow)\uparrow}(\phi_S, x, Q^2) + N^{(\leftrightarrow)\downarrow}(\phi_S, x, Q^2)}$$

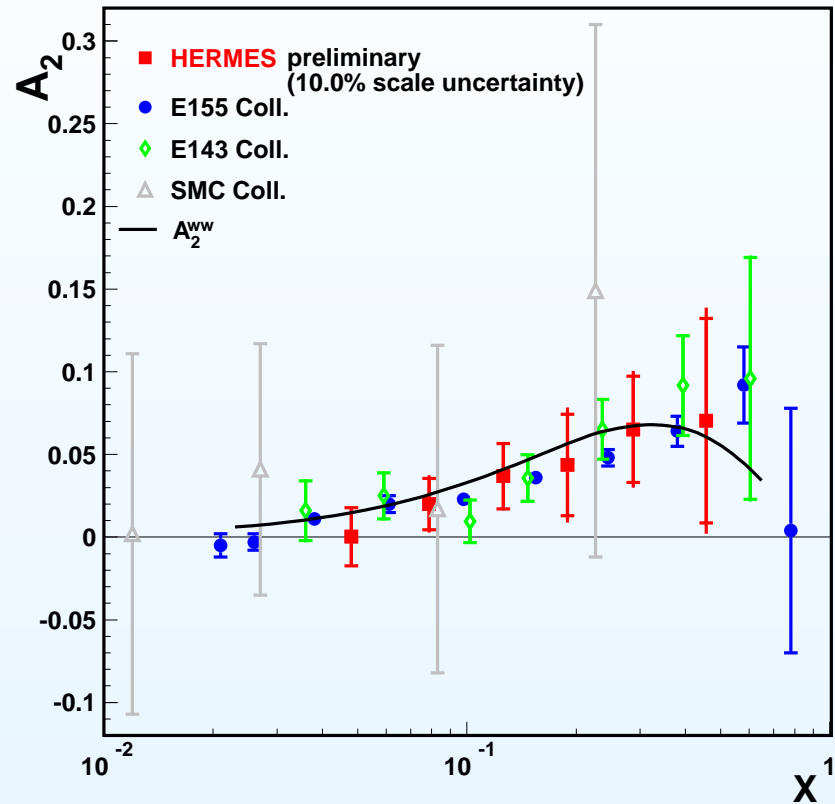
Correction for kinematic smearing effects:

- $A'_\perp(\phi_S, x, Q^2)$ measurement affected by
 - **higher order QED processes**
 - finite detector resolution \rightarrow **kinematic smearing**
- resulting bin migration studied in Monte Carlo (MC)
 - \rightarrow smearing matrix S
 - \rightarrow (model-independent) **unfolding algorithm**

$$A_\perp(j) = -1 + \frac{2}{N_{\text{MC, unpol}}(j)} \sum_{i=1}^{n_{\text{bins}}} (S^{\rightarrow\uparrow} + S^{\rightarrow\downarrow})^{-1}(j, i) \times$$
$$\left(A'_\perp(i) N'_{\text{MC, unpol}}(i) - n_{\text{MC, bg}}(i) + \sum_{k=1}^{n_{\text{bins}}} S^{\rightarrow\uparrow}(i, k) N_{\text{MC, unpol}}(k) \right)$$

\rightarrow statistical correlation of $A_\perp(\phi_S, x, Q^2)$

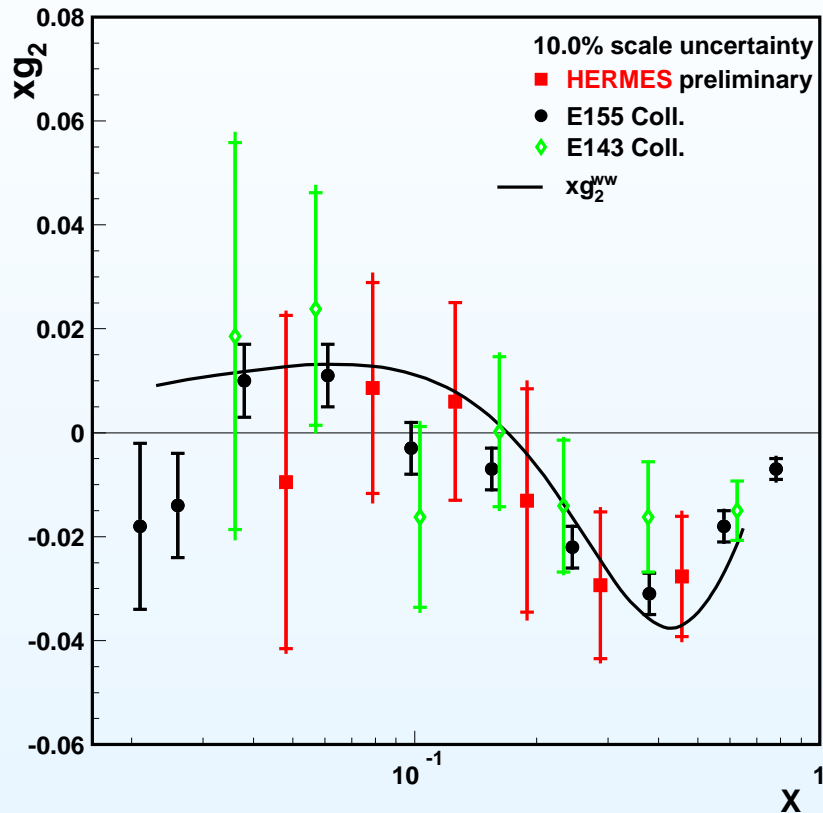
HERMES extraction of $A_2^p(x)$:



$$A_2^p = \frac{1}{1 + \eta\zeta} \left(\frac{A_{\perp}^p}{d} + \frac{\eta\zeta}{\gamma} (1 + \gamma^2) \frac{g_1^p}{F_1^p} \right)$$

R	R1998
F_2^p	GM07
$\frac{g_1^p}{F_1^p}$	E155

HERMES extraction of $g_2^p(x)$:



$$0.0041 < x < 0.9$$

$$0.18 \text{ GeV}^2 < Q^2 < 20 \text{ GeV}^2$$

$$g_2^p = \frac{F_1^p}{\gamma(1 + \eta\zeta)} \left(\frac{A_{\perp}^p}{d} - \left(\gamma - \frac{\eta\zeta}{\gamma} \right) \frac{g_1^p}{F_1^p} \right)$$

R R1998

F_2^p GM07

$\frac{g_1^p}{F_1^p}$ E155

In a nutshell:

- extraction of $A_2^p(x, Q^2)$ in

$$0.0041 < x < 0.9 \text{ and } 0.18 \text{ GeV}^2 < Q^2 < 20 \text{ GeV}^2$$

- extraction of $g_2^p(x, Q^2)$
 - consistent with SMC, E143 and E155
 - ↳ important check of *transverse data*
 - consistent with $g_2^p(x, Q^2) \approx g_2^{p,WW}(x, Q^2)$
 - ↳ no sensitivity to $\bar{g}_2^p(x, Q^2)$?
 - ↳ probing quark-gluon correlation?
- in preparation for final publication
- including a more detailed analysis of $g_2^p(x, Q^2)$