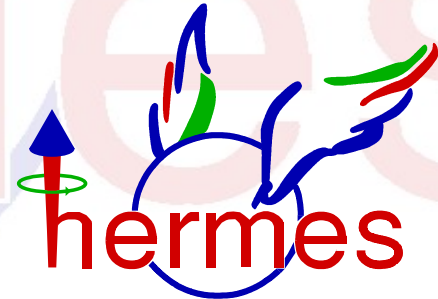


# Bose-Einstein correlations in deep-inelastic scattering from nuclei

**Gevorg Karyan**

**(for the HERMES Collaboration)**



# Hanbury-Brown and Twiss effect

LXXIV. *A New Type of Interferometer for Use in Radio Astronomy*

By R. HANBURY BROWN

Jodrell Bank Experimental Station, Cheshire

and

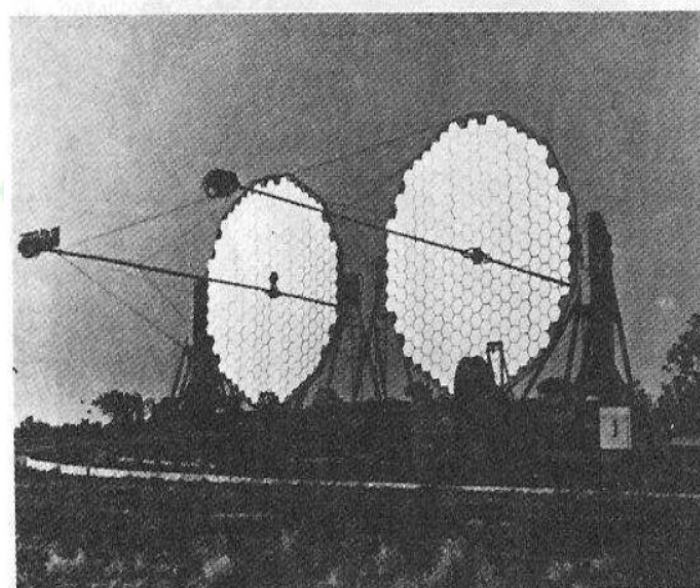
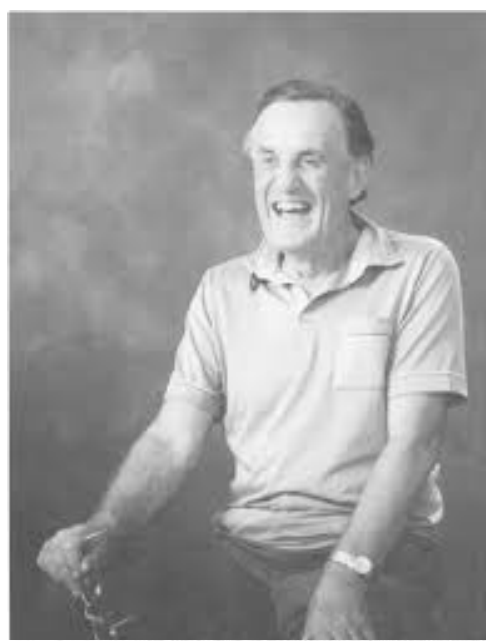
R. Q. TWISS

Services Electronics Research Laboratory, Baldock, Herts.\*

[Received March 20, 1954]

## SUMMARY

A new type of interferometer for measuring the diameter of discrete radio sources is described and its mathematical theory is given. The principle of the instrument is based upon the correlation between the rectified outputs of two independent receivers at each end of a baseline, and it is shown that the cross-correlation coefficient between these outputs is proportional to the square of the amplitude of the Fourier transform of the intensity distribution across the source. The analysis shows that it should be possible to operate the new instrument with extremely long baselines and that it should be almost unaffected by ionospheric irregularities.



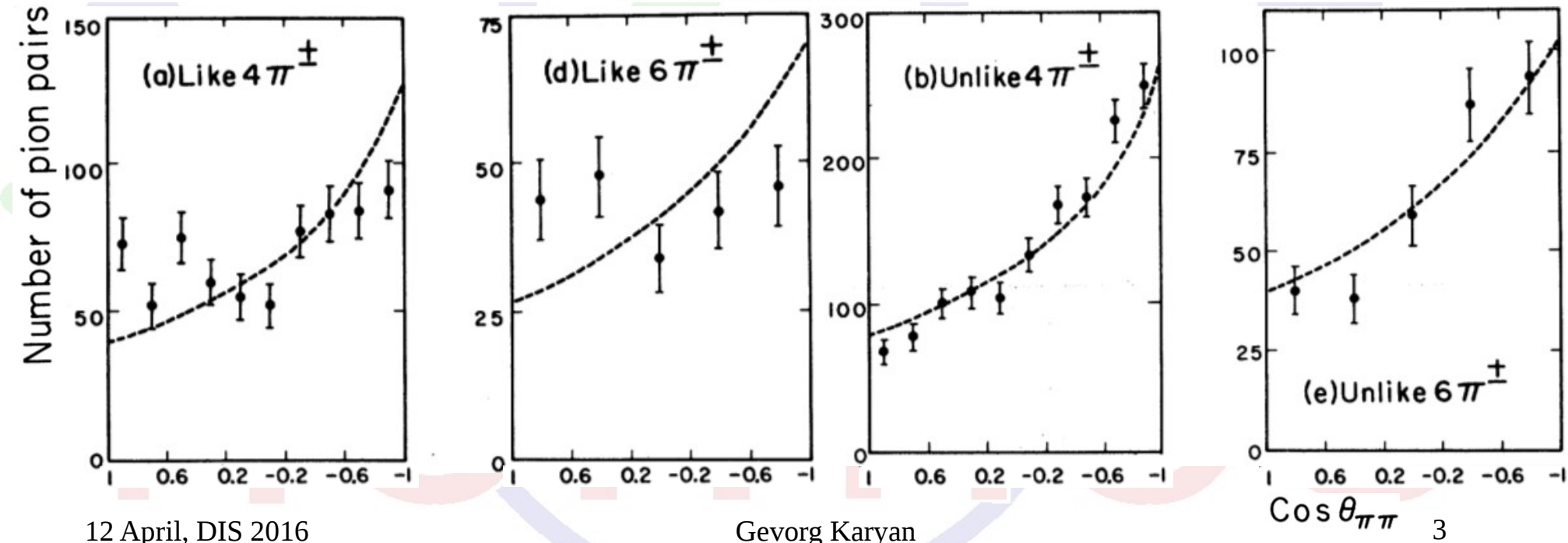
# HBT effect in accelerator physics

## PION-PION CORRELATIONS IN ANTIPROTON ANNIHILATION EVENTS\*

Gerson Goldhaber, William B. Fowler, Sulamith Goldhaber, T. F. Hoang,  
Theodore E. Kalogeropoulos, and Wilson M. Powell

Lawrence Radiation Laboratory and Department of Physics, University of California, Berkeley, California

(Received July 17, 1959)



# Bose-Einstein correlation (GGLP effect)

PHYSICAL REVIEW

VOLUME 120, NUMBER 1

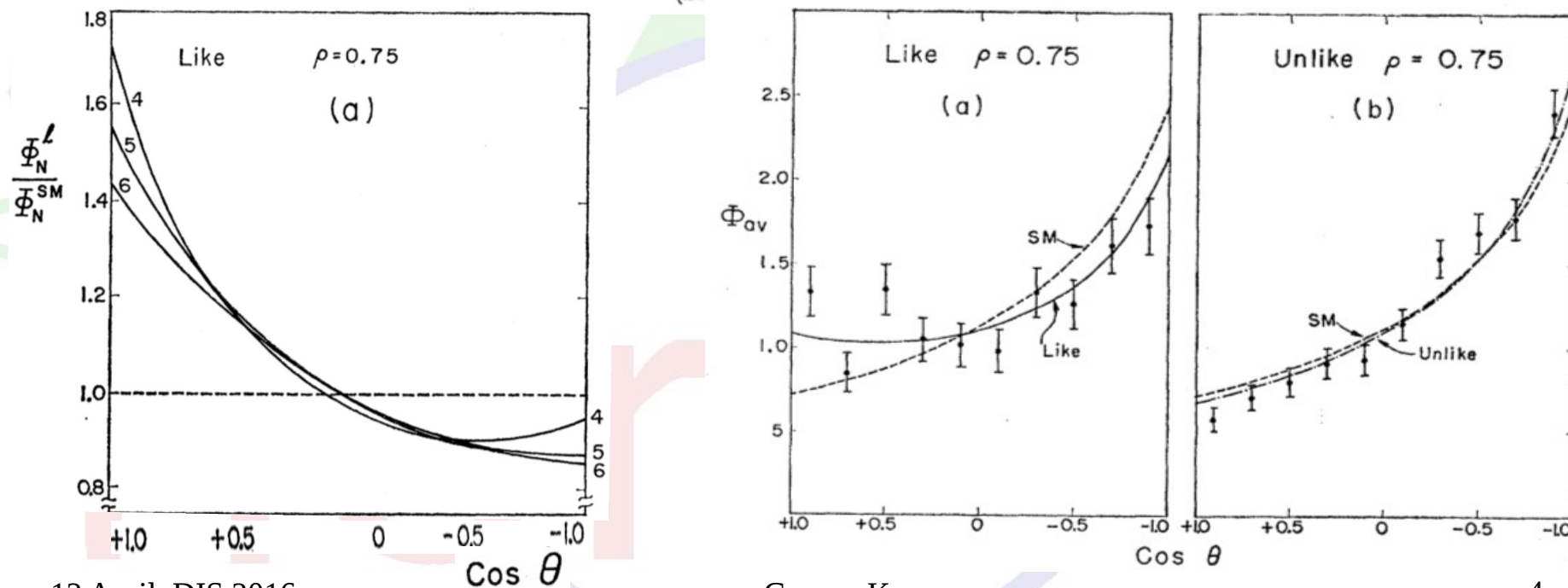
OCTOBER 1, 1960

## Influence of Bose-Einstein Statistics on the Antiproton-Proton Annihilation Process\*

GERSON GOLDHABER, SULAMITH GOLDHABER, WONYONG LEE, AND ABRAHAM PAIS†

*Lawrence Radiation Laboratory and Department of Physics, University of California, Berkeley, California*

(Received May 16, 1960)



# Bose-Einstein correlation (GGLP effect)

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VOLUME 120, NUMBER 1

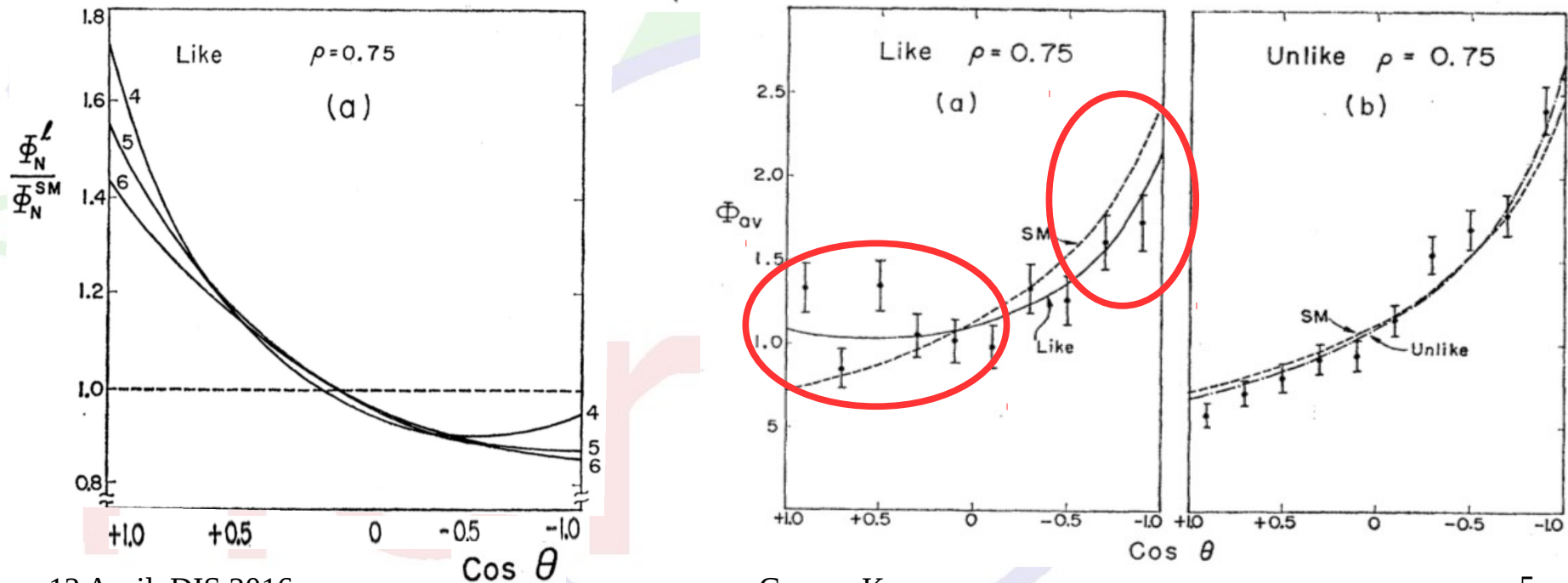
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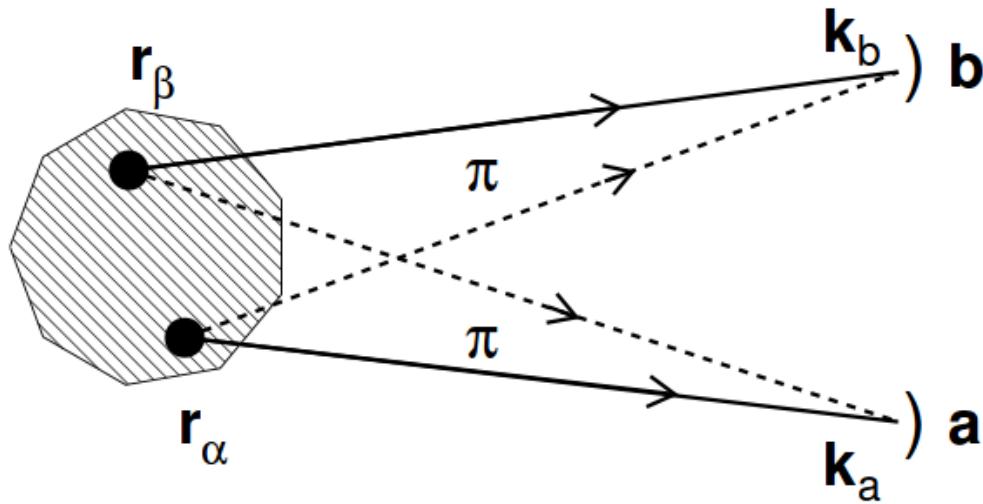


12 April, DIS 2016

Gevorg Karyan

5

# The underlying physics



$$\Psi_{2\pi} = \frac{\Psi_{a\alpha} \Psi_{b\beta} + \Psi_{b\alpha} \Psi_{a\beta}}{\sqrt{2}}$$

$$\Psi_{a\alpha} \approx e^{ik_a r_\alpha}, \Psi_{a\beta} \approx e^{ik_a r_\beta}, \dots$$

$$|\Psi_{2\pi}|^2 = 1 + \cos(\delta \mathbf{k} \delta \mathbf{r}), \quad \delta \mathbf{k} = \mathbf{k}_a - \mathbf{k}_b, \quad \delta \mathbf{r} = \mathbf{r}_\alpha - \mathbf{r}_\beta$$

$$\mathbf{R}(\mathbf{k}_a, \mathbf{k}_b) \approx 1 + \cos(\delta \mathbf{k} \delta \mathbf{r})$$

# Measuring the BEC

$$R(\mathbf{k}_a, \mathbf{k}_b) \approx 1 + \cos(\delta \mathbf{k} \delta \mathbf{r})$$

( two point source )

$$R(\mathbf{p}_1, \mathbf{p}_2) = \frac{D(\mathbf{p}_1, \mathbf{p}_2)}{D(\mathbf{p}_1)D(\mathbf{p}_2)}$$

( continuous space-time distribution )

$D(\mathbf{p}_1, \mathbf{p}_2)$  - two particle probability density

$D(\mathbf{p}_1), D(\mathbf{p}_2)$  - one particle probability density

# Measuring the BEC

$$R(\mathbf{p}_1, \mathbf{p}_2) = \frac{D(\mathbf{p}_1, \mathbf{p}_2)}{D(\mathbf{p}_1)D(\mathbf{p}_2)} \longleftrightarrow \text{experimental observable}$$

Goldhaber parametrization

$$R(\mathbf{p}_1, \mathbf{p}_2) \longrightarrow R(T), \quad T^2 = -(\mathbf{p}_1 - \mathbf{p}_2)^2$$
$$R(T) = 1 + \lambda e^{-T^2 r_G^2}$$

$\lambda$  - chaoticity parameter  
 $r_G$  - source distribution size



# Measuring the BEC

## Goldhaber parametrization



$$R(T) = 1 + \lambda e^{-T^2 r_G^2}$$

$$R(T) = \gamma (1 + \lambda e^{-T^2 r_G^2}) P(T)$$

$$P(T) = 1 + \delta T^2$$

$\gamma$  - normalization parameter

$P_T$  - long range correlation

$r_G, \lambda, \gamma, \delta$  - free parameters

# Measuring the BEC

$$R(\mathbf{p}_1, \mathbf{p}_2) = \frac{D(\mathbf{p}_1, \mathbf{p}_2)}{D_r(\mathbf{p}_1, \mathbf{p}_2)}$$

experimental observable

$$D_r(\mathbf{p}_1, \mathbf{p}_2) \equiv D(\mathbf{p}_1)D(\mathbf{p}_2)$$

$D_r(\mathbf{p}_1, \mathbf{p}_2)$  - two particle probability density reference distribution

experimental two particle distributions without BECs

# Measuring the BEC

$$R(\mathbf{p}_1, \mathbf{p}_2) = \frac{D(\mathbf{p}_1, \mathbf{p}_2)}{D_r(\mathbf{p}_1, \mathbf{p}_2)}$$

experimental observable

$$D_r(\mathbf{p}_1, \mathbf{p}_2) \equiv D(\mathbf{p}_1)D(\mathbf{p}_2)$$

$D_r(\mathbf{p}_1, \mathbf{p}_2)$  - two particle probability density reference distribution

1. Method of Event Mixing (MEM)
2. Method of Unlike-Sign pairs (MUS)

# Measuring the BEC

## 1. Method of Event Mixing (MEM)

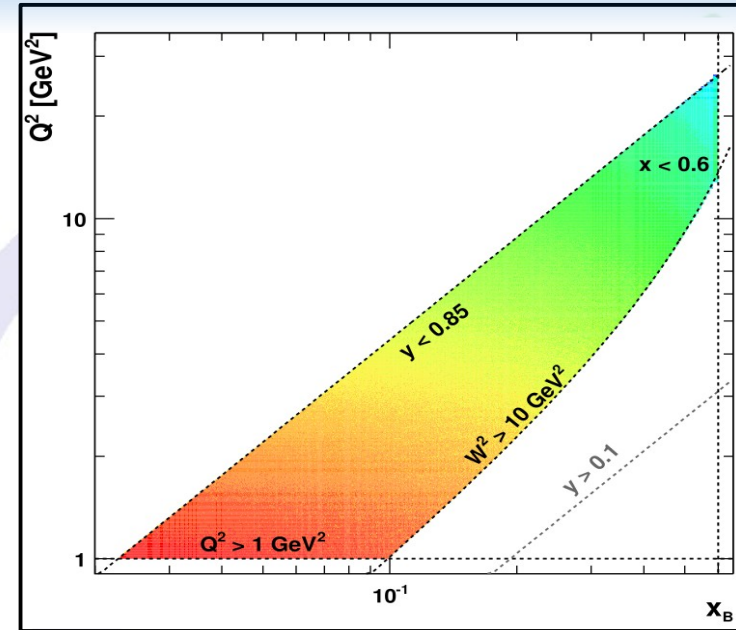
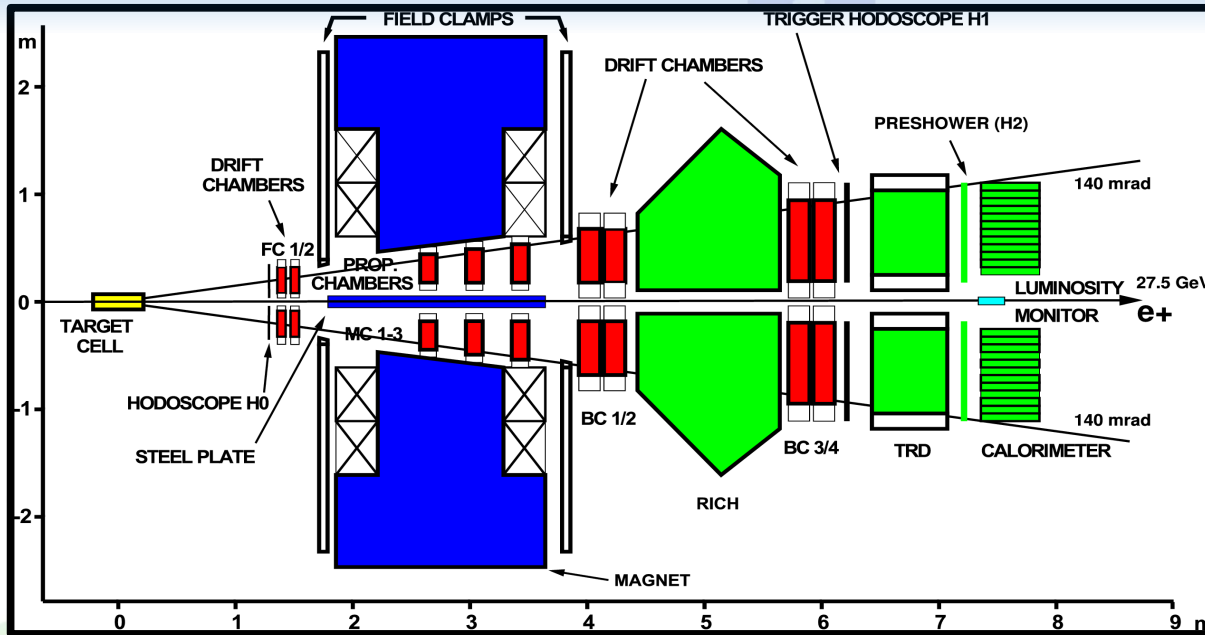
violation of energy and momentum conservation

## 2. Method of Unlike-Sign pairs (MUS)

contribution from resonances excluded in the case of MEM

PYTHIA-based MC tuned to provide an accurate description

# Experiment



**Beam :  $e^-/e^+$  27.6 GeV**

**Target : H, D,  $^3\text{He}$ ,  $^4\text{He}$ , N, Ne, Kr, Xe pure gaseous**

**Good momentum resolution :  $\frac{\delta p}{p} < 2\%$**

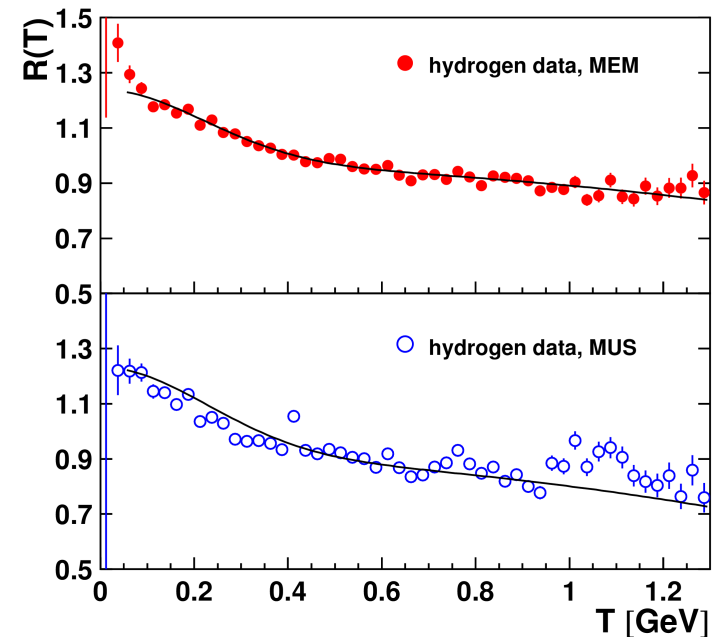
**Excellent particle identification**

# Data selection

- $Q^2 > 1 \text{ GeV}^2$
- $W^2 > 10 \text{ GeV}^2$
- at least two charged hadrons
- $2 \text{ GeV} < p < 15 \text{ GeV}$

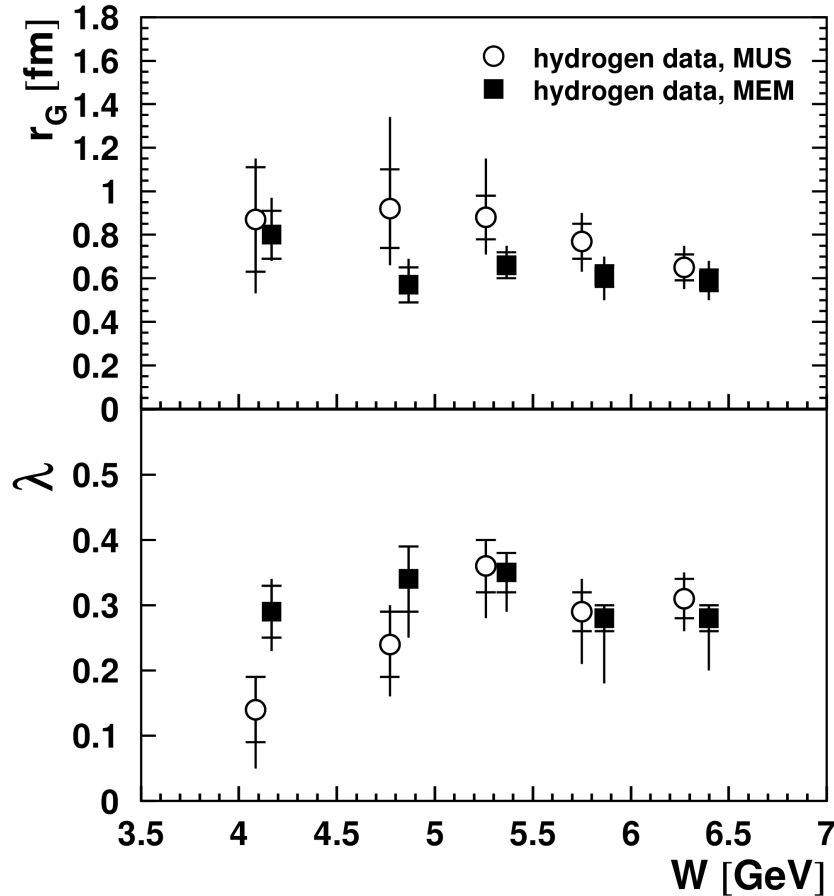
| Nucleus       | $N_{ev}$ | $N^{like}$ | $N^{unlike}$ |
|---------------|----------|------------|--------------|
| $^1\text{H}$  | 1145046  | 478946     | 958185       |
| $^2\text{H}$  | 1297356  | 680143     | 1178797      |
| $^3\text{He}$ | 34391    | 15295      | 29165        |
| $^4\text{He}$ | 79776    | 30539      | 59244        |
| N             | 92968    | 41112      | 78402        |
| Ne            | 175594   | 75898      | 146145       |
| Kr            | 211456   | 91391      | 172946       |
| Xe            | 106274   | 46130      | 87125        |

Black curve – result of fit using the Goldhaber parametrization.



# Results

## A. Airapetian et. al, Eur. Phys. J. C 75 (2015) 361

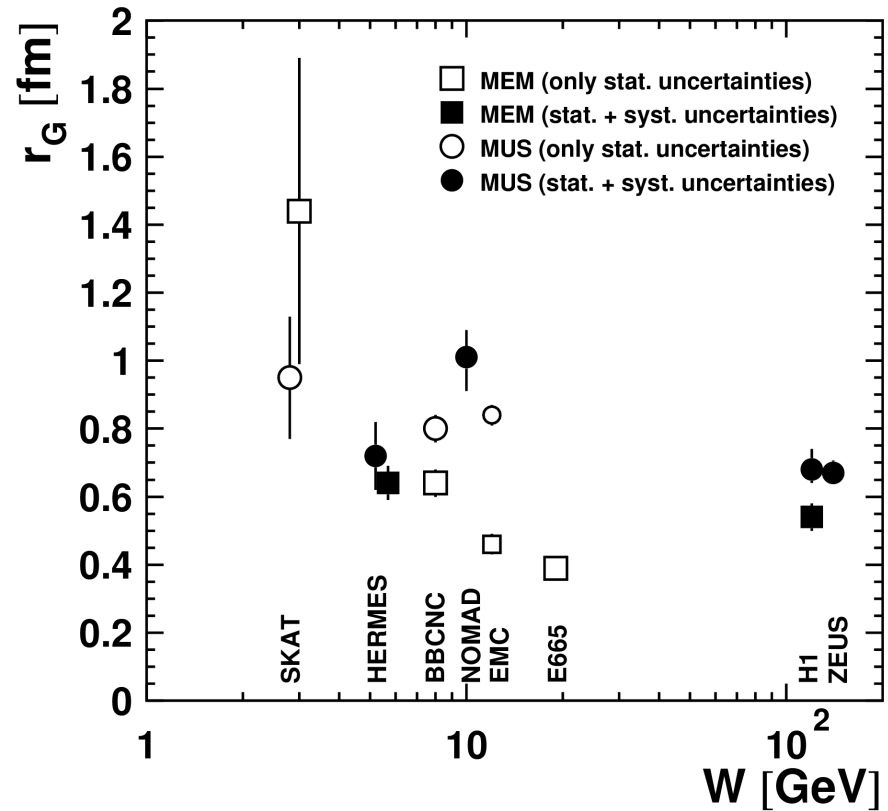
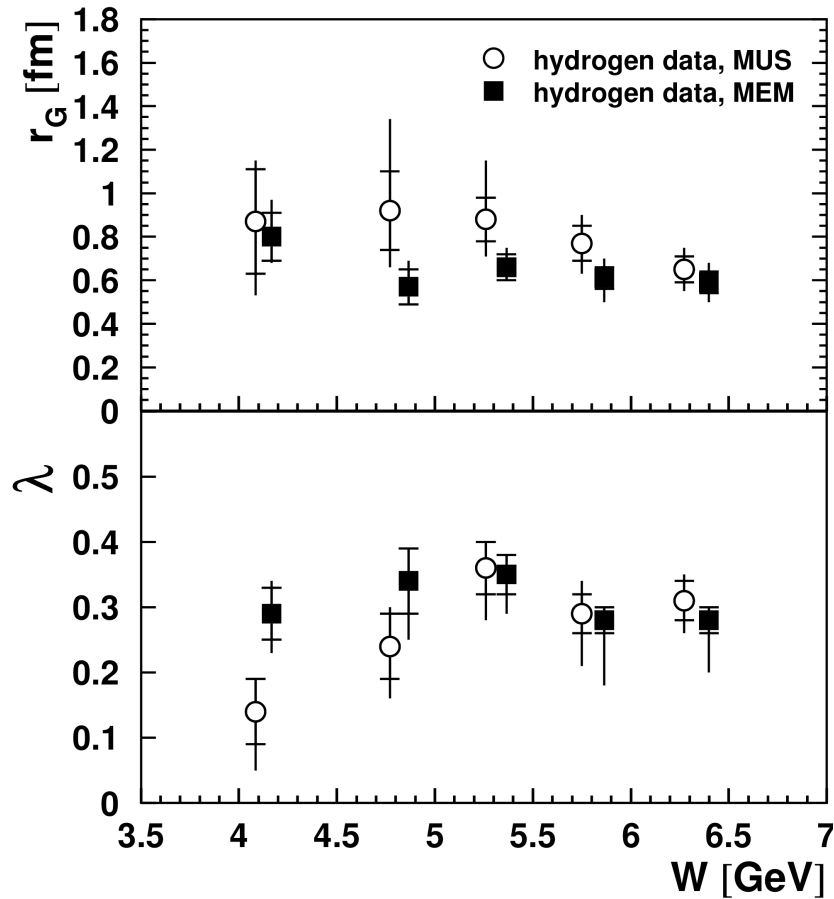


**Table 2.** Results for the Goldhaber parametrization fitted to the HERMES hydrogen data, both for the mixed-event method (*MEM*) and the method of unlike-sign pairs (*MUS*).

| Method     | Goldhaber parameters  |
|------------|---|
| <i>MEM</i> | $r_G = 0.64 \pm 0.03(\text{stat})_{-0.04}^{+0.04}(\text{sys})$ fm<br>$\lambda = 0.28 \pm 0.01(\text{stat})_{-0.05}^{+0.00}(\text{sys})$ |
| <i>MUS</i> | $r_G = 0.72 \pm 0.04(\text{stat})_{-0.09}^{+0.09}(\text{sys})$ fm<br>$\lambda = 0.28 \pm 0.02(\text{stat})_{-0.04}^{+0.02}(\text{sys})$ |

# Results

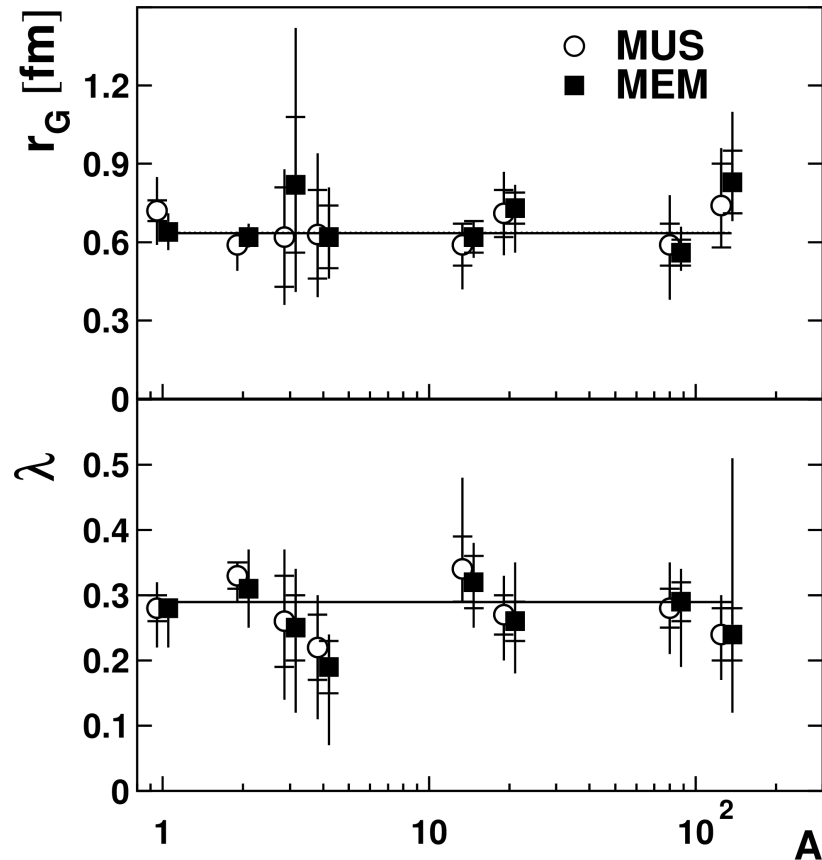
## A. Airapetian et. al, Eur. Phys. J. C 75 (2015) 361





# Results

## A. Airapetian et. al, Eur. Phys. J. C 75 (2015) 361

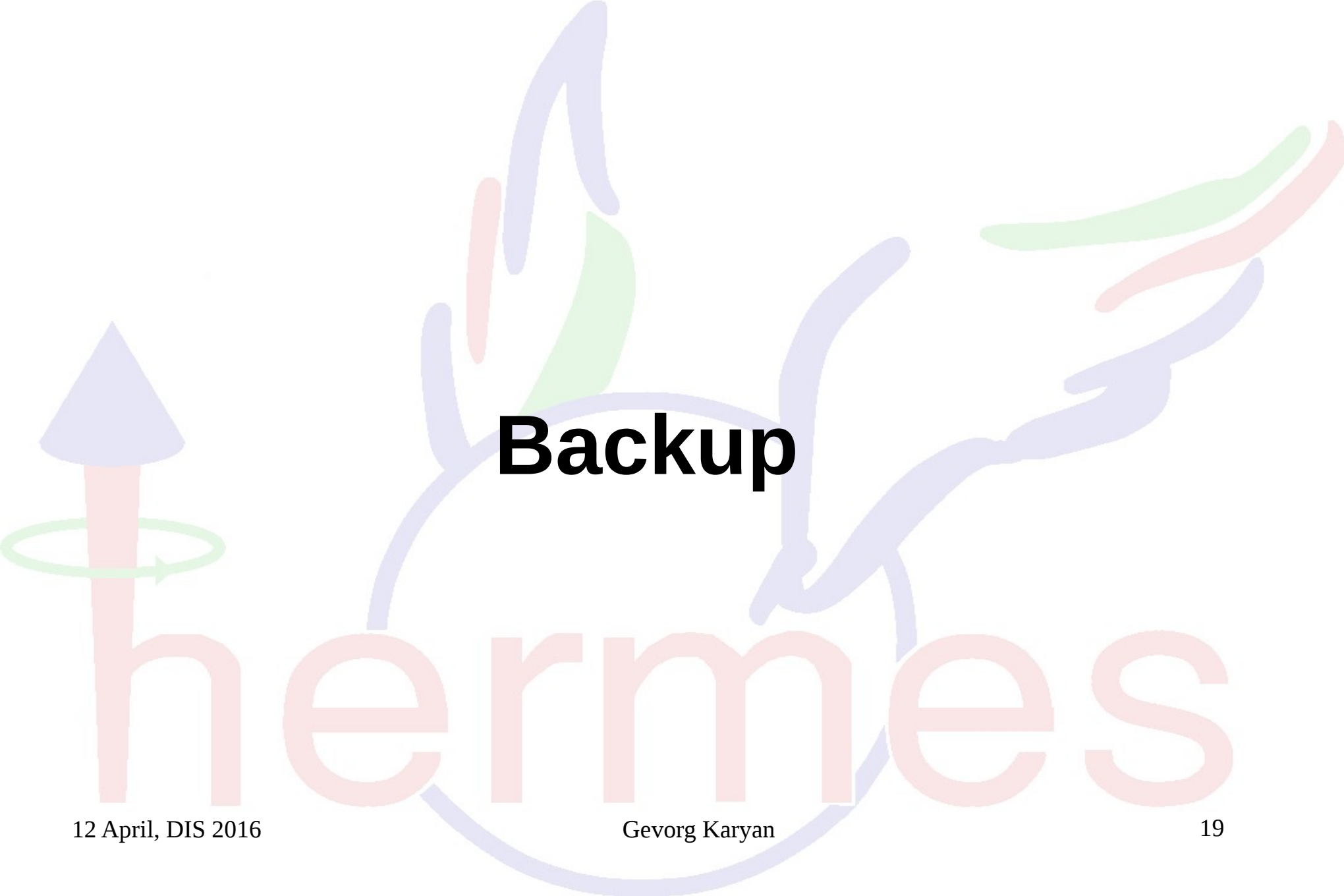


**Table 3.** Fit of a constant to the Goldhaber parameters as a function of the target atomic mass  $A$ . Results are given for both the mixed-event method ( $MEM$ ) and the method of unlike-sign pairs ( $MUS$ ).

| Method | Value   | $\chi^2/\text{NDF}$ |
|--------|---|---------------------|
| $MEM$  | $r_G = 0.634 \pm 0.017$ fm<br>$\lambda = 0.289 \pm 0.006$ | 1.5<br>2.1          |
| $MUS$  | $r_G = 0.636 \pm 0.021$ fm<br>$\lambda = 0.289 \pm 0.011$ | 1.2<br>1.4          |

# Summary

- Bose-Einstein correlation between two like-sign hadrons produced in semi-inclusive deep inelastic electron/positron scattering off nuclear targets ranging from hydrogen to xenon has been measured.
- The results obtained using the two reference sample methods (i.e MUS and MEM ) are in a good agreement.
- Within the total experimental uncertainties, no dependence of the parameters  $r_G$  and  $\lambda$  on the target atomic mass is observed.



# Backup

hermes

# Data/MC comparison

