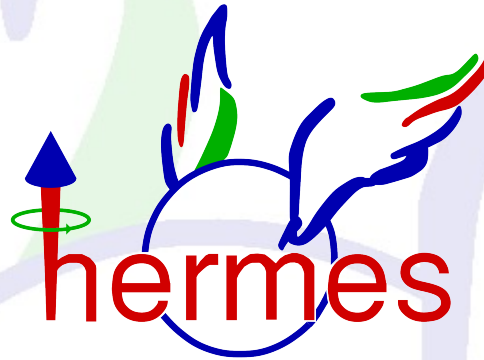


Hadron multiplicities at the HERMES experiment



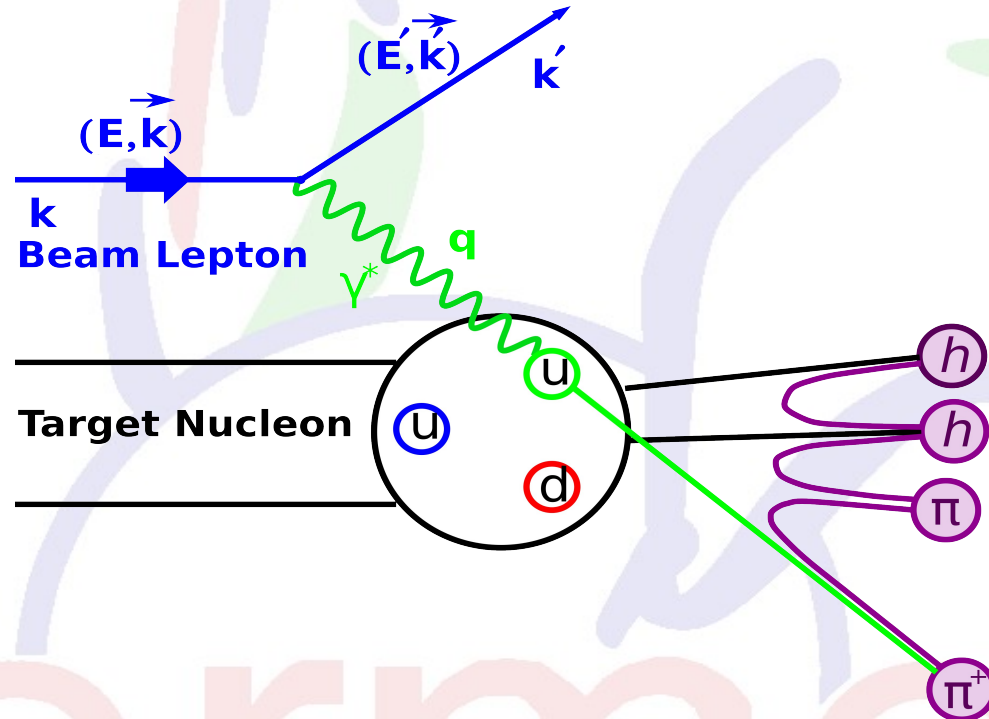
Gevorg Karyan

on behalf of the HERMES Collaboration

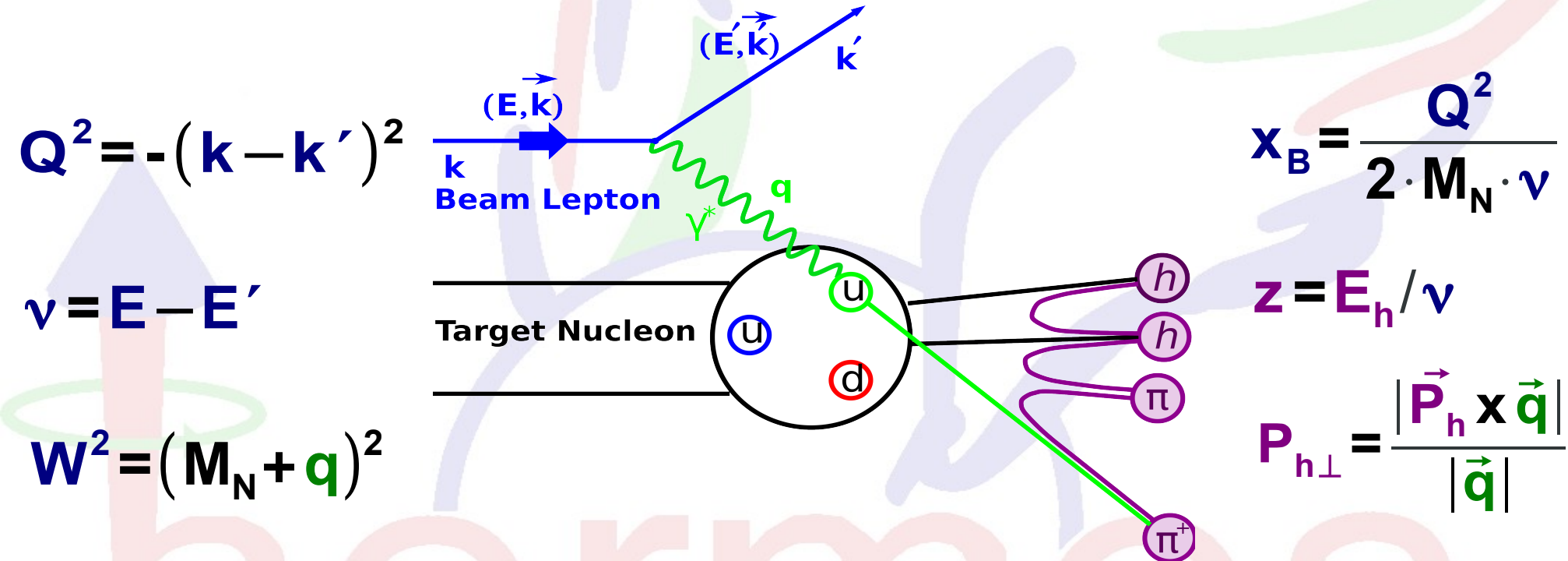
Alikhanyan National Science Laboratory

Yerevan, Armenia

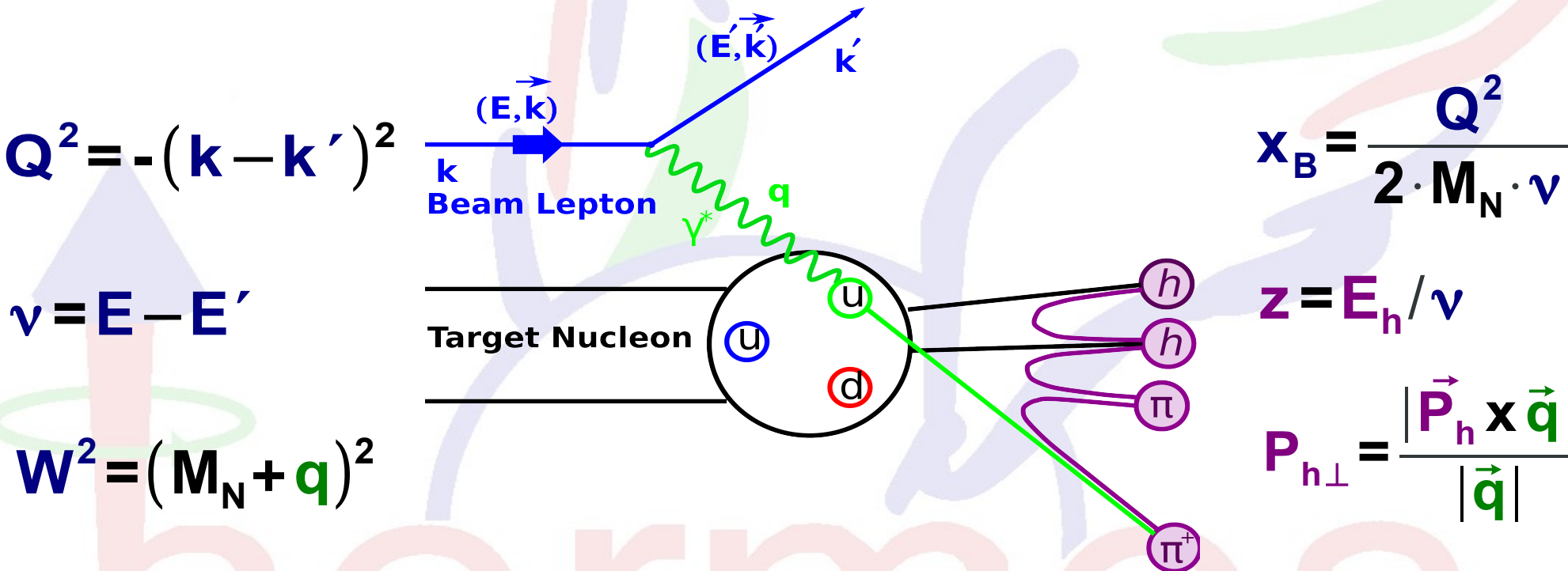
Semi-Inclusive Deep-Inelastic Scattering



Semi-Inclusive Deep-Inelastic Scattering



Semi-Inclusive Deep-Inelastic Scattering



$$Q^2 = -(\vec{k} - \vec{k}')^2$$

$$\nu = E - E'$$

$$W^2 = (M_N + q)^2$$

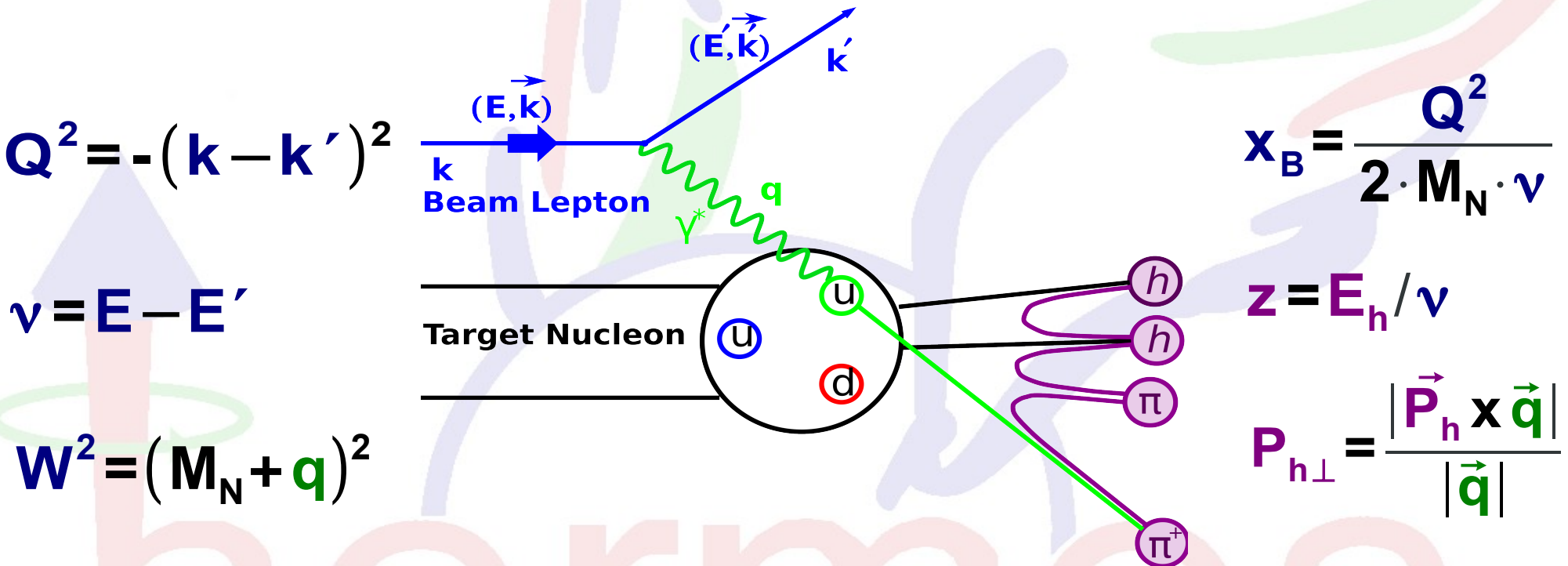
$$x_B = \frac{Q^2}{2 \cdot M_N \cdot \nu}$$

$$z = E_h / \nu$$

$$P_{h\perp} = \frac{|\vec{P}_h \times \vec{q}|}{|\vec{q}|}$$

Leading Order QCD, Factorization

Semi-Inclusive Deep-Inelastic Scattering



$$\sigma_{eN \rightarrow ehX} \sim \sum_f e_f^2 q_f(\mathbf{x}_B, Q^2, \mathbf{p}_T) \sigma^{eq \rightarrow eq} D_f^h(z, Q^2, \mathbf{k}_T)$$

Experimental observable

SIDIS hadron yields

$$M^h(x_B, Q^2, z, P_{h\perp}, \phi) = \frac{N^h(x_B, Q^2, z, P_{h\perp}, \phi)}{N^e(x_B, Q^2)}$$

DIS event yields

Underlying physics

$$M^h \sim \frac{\sum_f e_f^2 q_f(\mathbf{x}_B, Q^2, \mathbf{p}_T) D_f^h(\mathbf{z}, Q^2, \mathbf{k}_T)}{\sum_f e_f^2 q_f(\mathbf{x}_B, Q^2, \mathbf{p}_T)}$$

Underlying physics

$$M^h \sim \frac{\sum_f e_f^2 q_f(x_B, Q^2, p_T) D_f^h(z, Q^2, k_T)}{\sum_f e_f^2 q_f(x_F, Q^2, p_T)}$$

Underlying physics

$$M^h \sim \frac{\sum_f e_f^2 q_f(x_B, Q^2, p_T) D_f^h(z, Q^2, k_T)}{\sum_f e_f^2 q_f(x_F, Q^2, p_T)}$$

The equation above is annotated with blue ovals highlighting the terms $q_f(x_B, Q^2, p_T)$ and $q_f(x_F, Q^2, p_T)$ as PDFs, and a purple oval highlighting $D_f^h(z, Q^2, k_T)$ as FF².

Underlying physics

Collinear Framework

$$M^h \sim \frac{\sum_f e_f^2 q_f(\text{PDF}, Q^2) D_f^h(z, Q^2)}{\sum_f e_f^2 q_f(\text{PDF}, Q^2)}$$

PDFs are well known

FFs are poorly known

Underlying physics

$$M^h \sim \frac{\sum_f e_f^2 q_f(\text{PDF}, Q^2) D_f^h(z, Q^2)}{\sum_f e_f^2 q_f(\text{PDF}, Q^2)}$$

PDFs are well known

FFs are poorly known

Particularly unfavored FFs

Underlying physics

$$M^h \sim \frac{\sum_f e_f^2 q_f(\text{PDF}^2) D_f^h(z, \text{FF}^2)}{\sum_f e_f^2 q_f(\text{PDF}^2)}$$

CTEQ6L, GRV, ...

DSS, Kretzer, ...

Universality

SIDIS(e+N), SLA(e⁺+e⁻), HS(p+p)

Advantage of SIDIS

Charge separated FFs



SIDIS

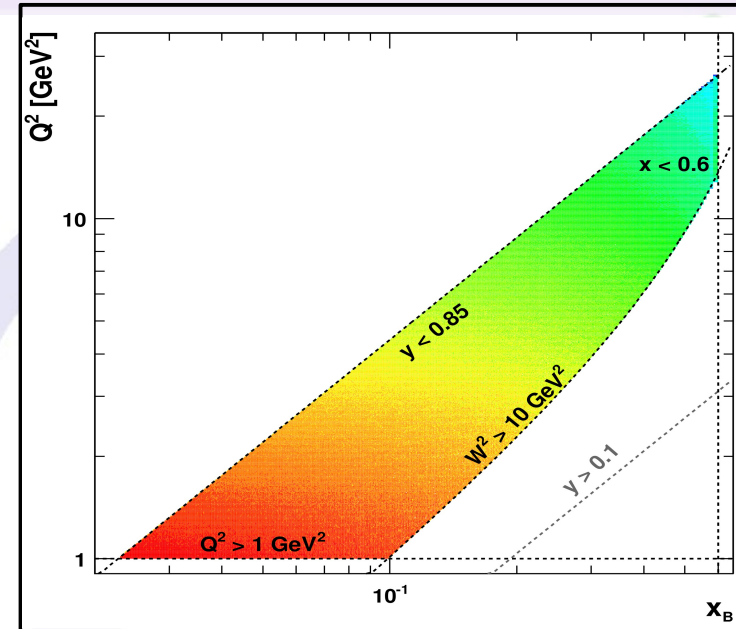
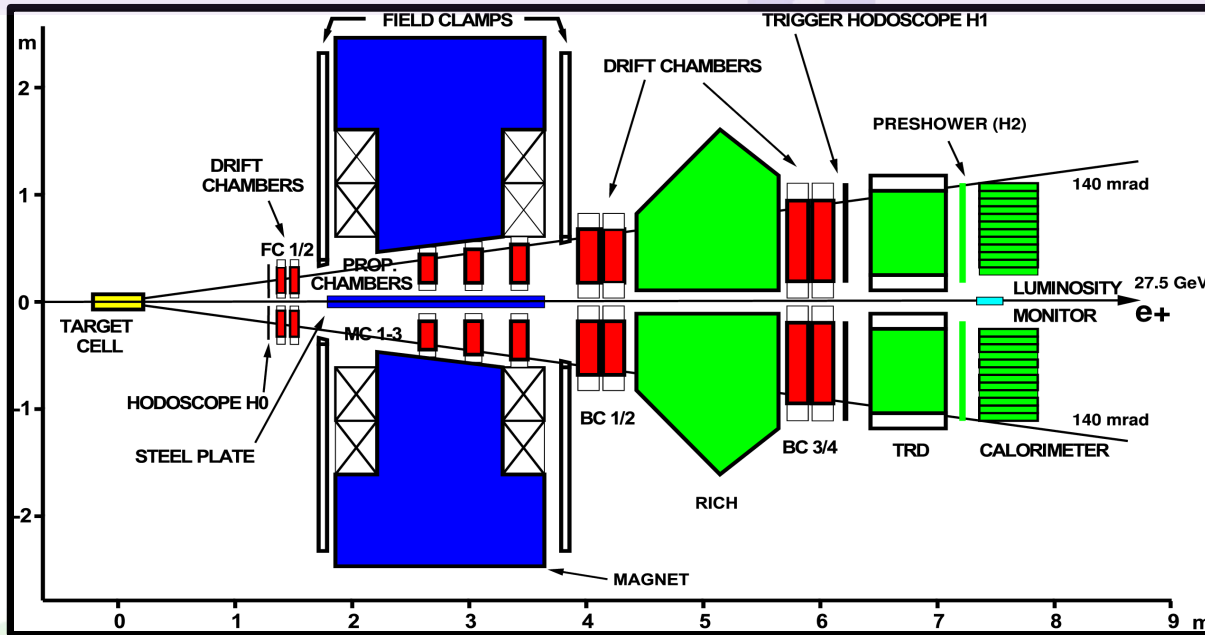
$(D_u^\pi, D_{\bar{u}}^K, \dots)$

(D_u^π, D_s^K, \dots)



Flavour separated FFs

Experiment



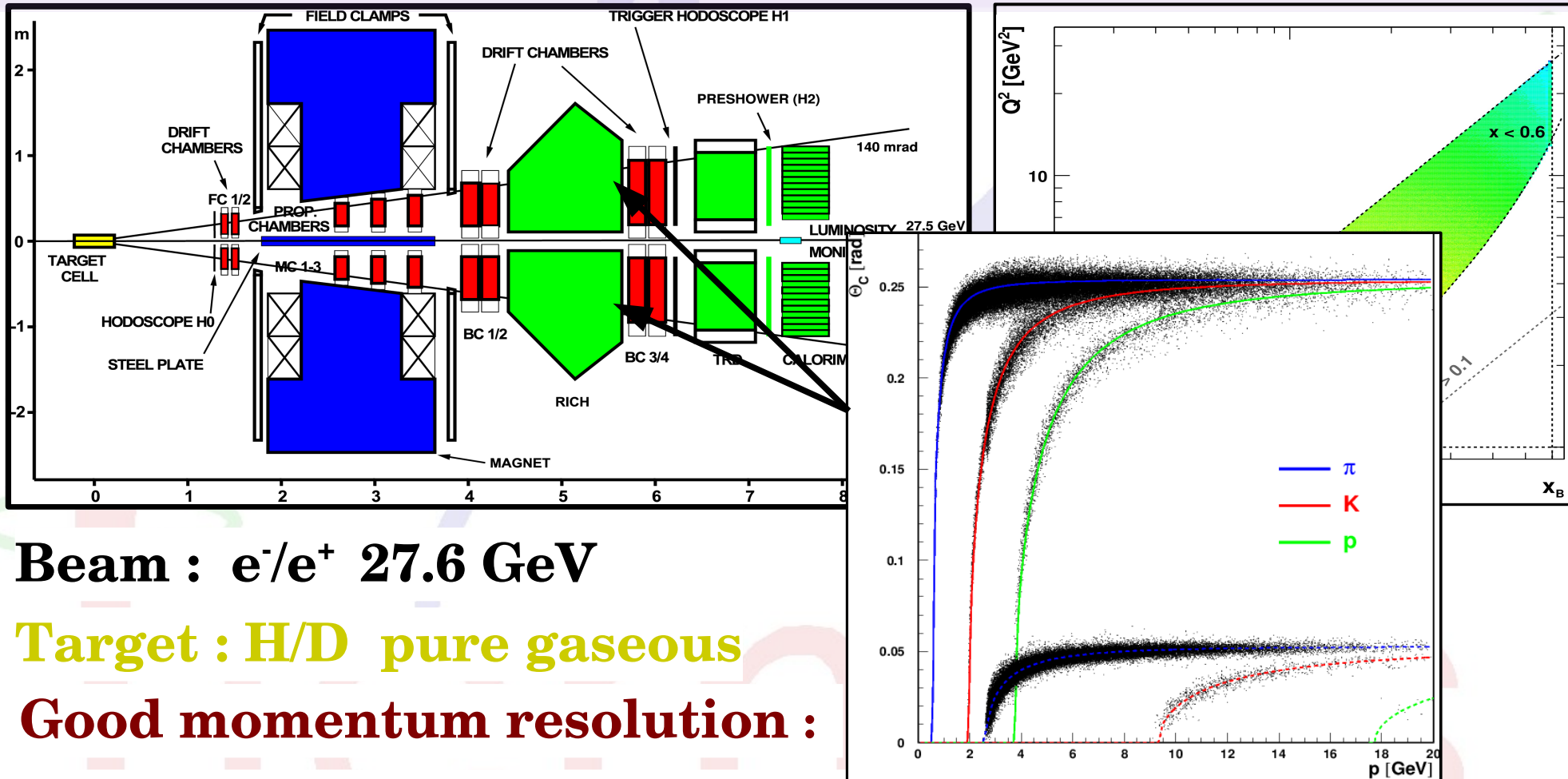
Beam : e^-/e^+ 27.6 GeV

Target : H/D pure gaseous

Good momentum resolution : $\frac{\delta p}{p} < 2 \%$

Excellent particle identification

Experiment



Beam : e^-/e^+ 27.6 GeV

Target : H/D pure gaseous

Good momentum resolution :

Excellent particle identification

Data selection

DIS regime

- $Q^2 > 1 \text{ GeV}^2$
- $W^2 > 10 \text{ GeV}^2$
- $0.1 < \nu/E_{\text{beam}} < 0.85$

SIDIS selection

- $2 \text{ GeV} < p < 15 \text{ GeV}$
- $0.2 < z < 0.8$

Raw Data

Data analysis

Raw Data

**Charge Symmetric Background
(Dalitz decay, $\gamma \rightarrow e^+ e^-$)**

RICH Unfolding

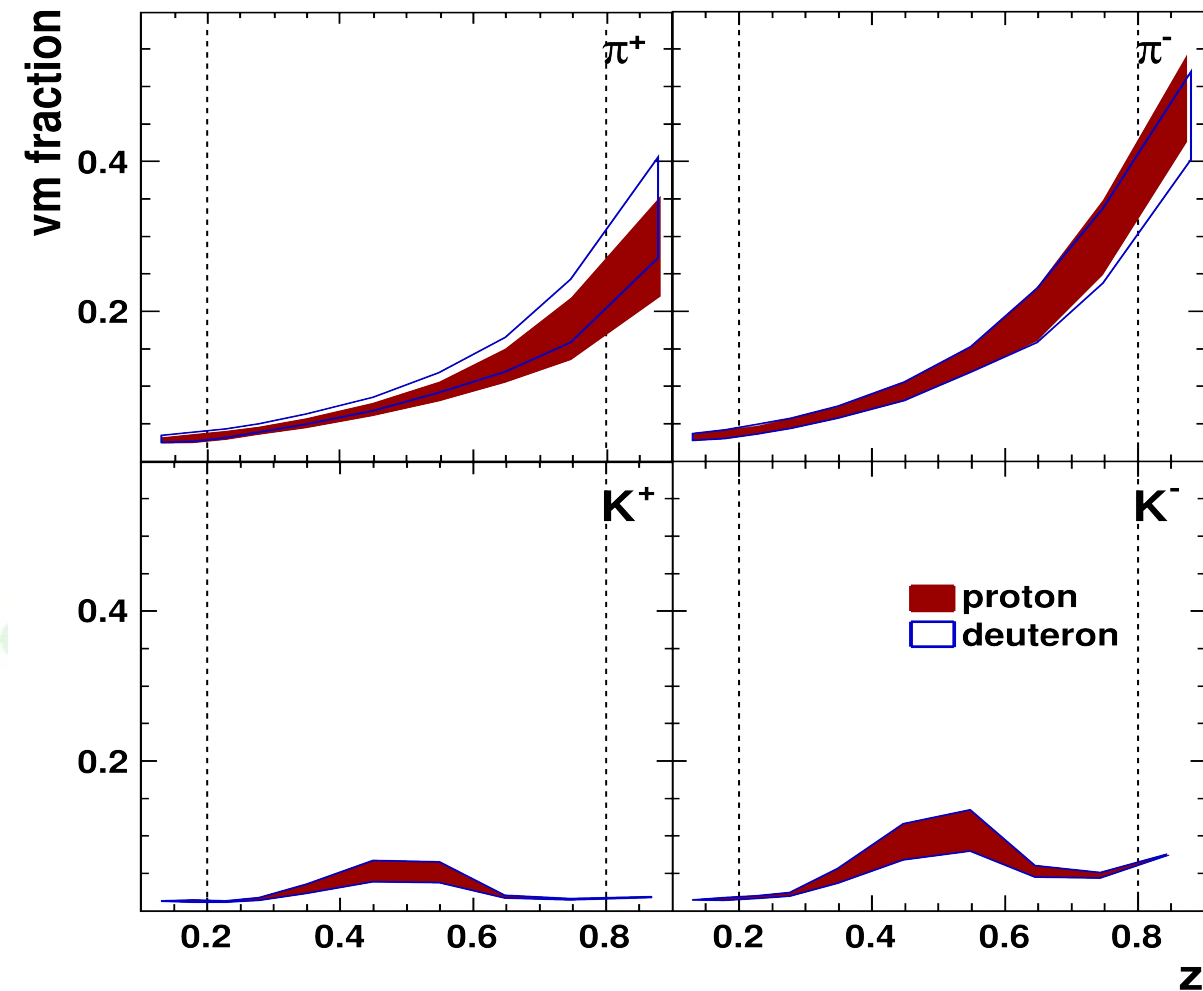
Trigger Efficiencies

**Diffractive Vector Meson
Contribution**

**Detector Smearing & QED Radiative
Effects**

Final Data

Exclusive vector meson contribution



Due to diffractively produced exclusive $\rho^0 \rightarrow \pi^+\pi^-$ and $\phi \rightarrow K^+K^-$

➤ Results with and without VM subtraction

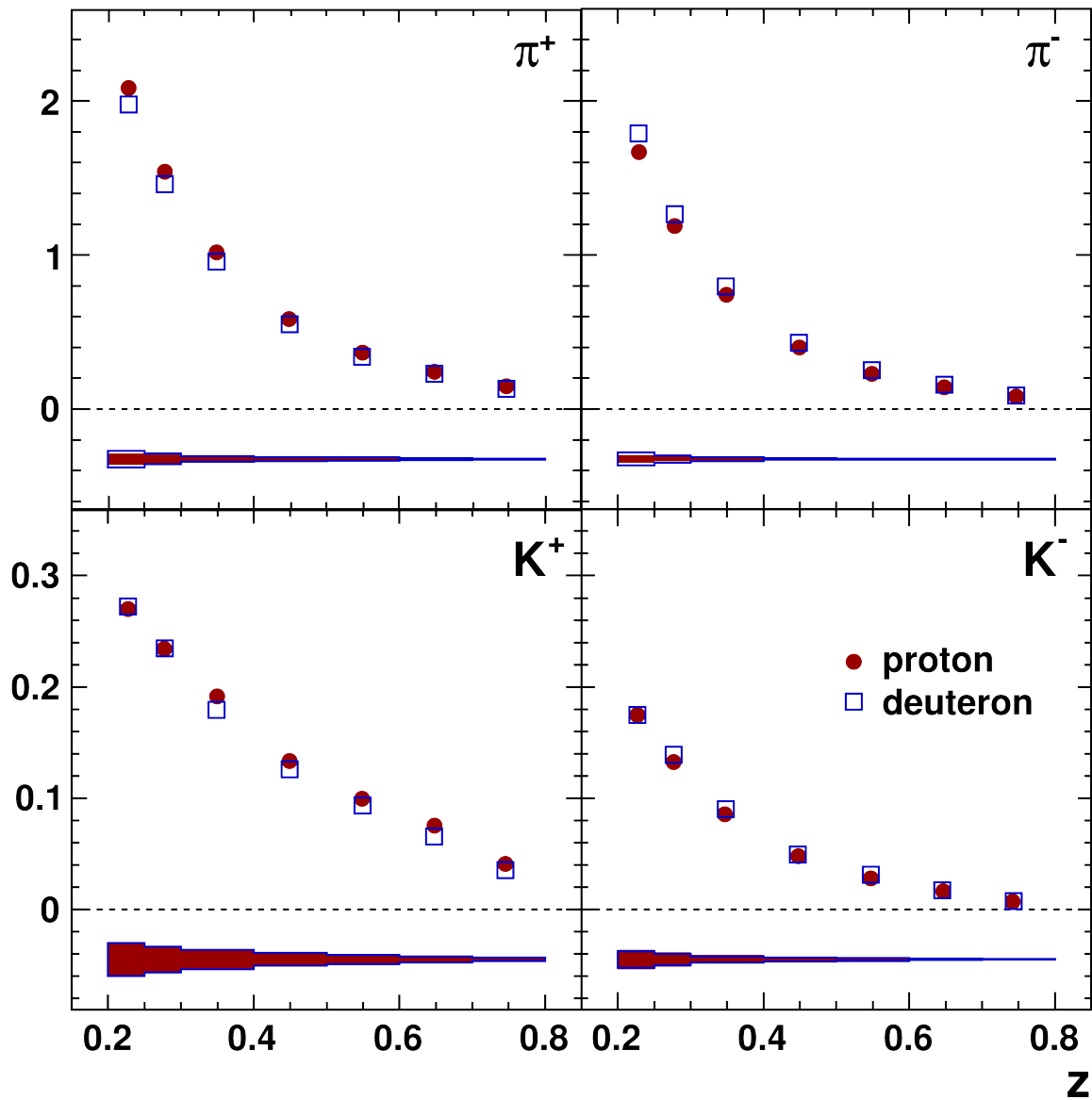
Results

arXiv:1212.5407v1

A. Airapetian et al, Phys. Rev. D87 (2013) 074029

➤ **In this presentation results are with VM subtraction.**

Multiplicity



$$M^h(\mathbf{x}_B, Q^2, \mathbf{z}, \mathbf{P}_{h\perp})$$

$u \rightarrow \pi^+(u\bar{d})$ **favored**

$u \rightarrow \pi^-(d\bar{u})$ **unfavored**

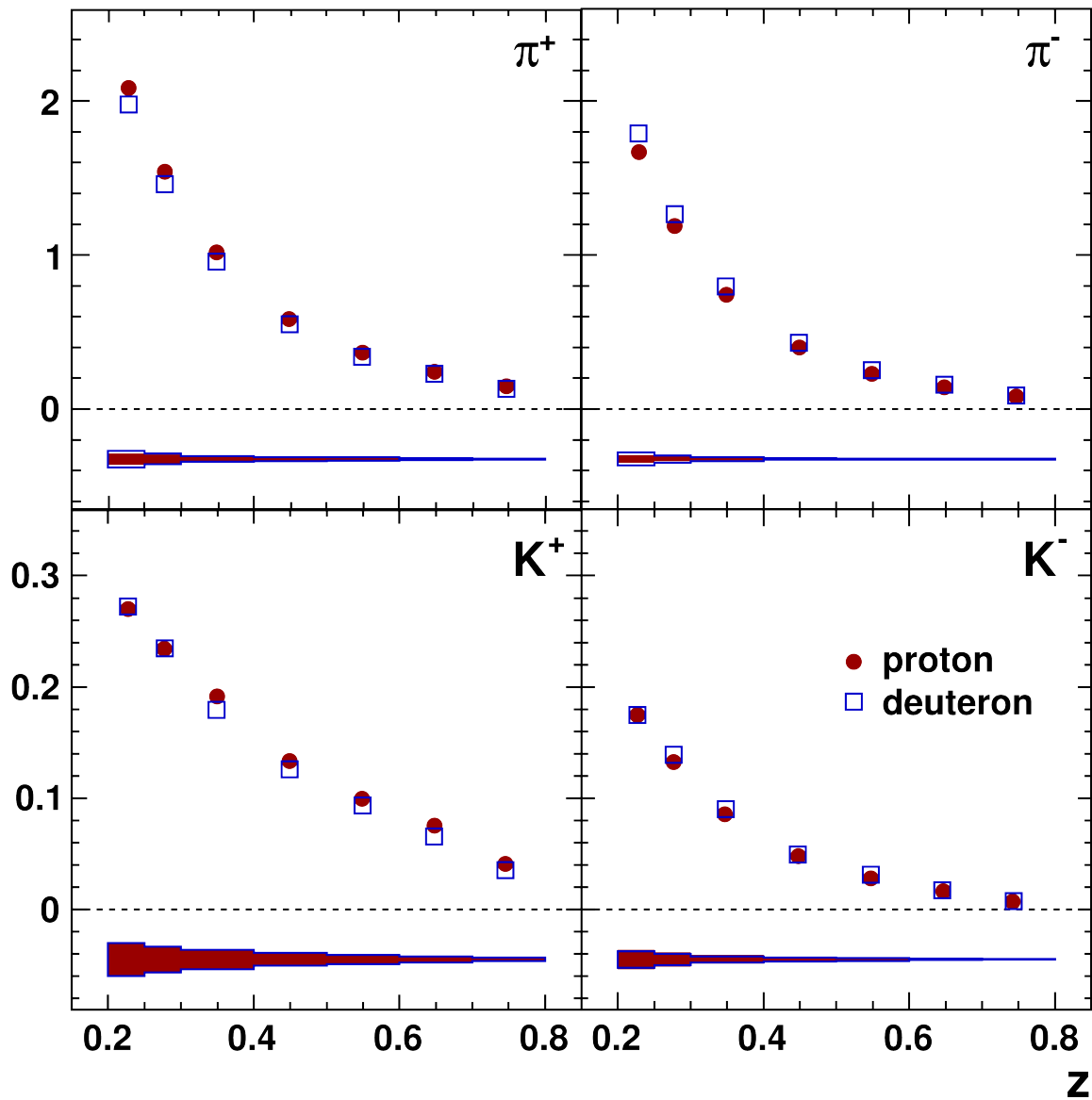
proton \rightarrow u dominance

$$\frac{M_{\text{proton}}^{\pi^+}}{M_{\text{proton}}^{\pi^-}} = 1.2 - 2.6$$

$d \rightarrow \pi^-$ **favored**

$$\frac{M_{\text{deuteron}}^{\pi^-}}{M_{\text{proton}}^{\pi^-}} > 1$$

Multiplicity



$$M^h(x_B, Q^2, z, P_{h\perp})$$

$u \rightarrow K^+(u\bar{s})$ favored

$u \rightarrow K^-(s\bar{u})$ unfavored

proton \rightarrow u dominance

$$\frac{M_{\text{proton}}^{K^+}}{M_{\text{proton}}^{K^-}} = 1.5 - 5.7$$

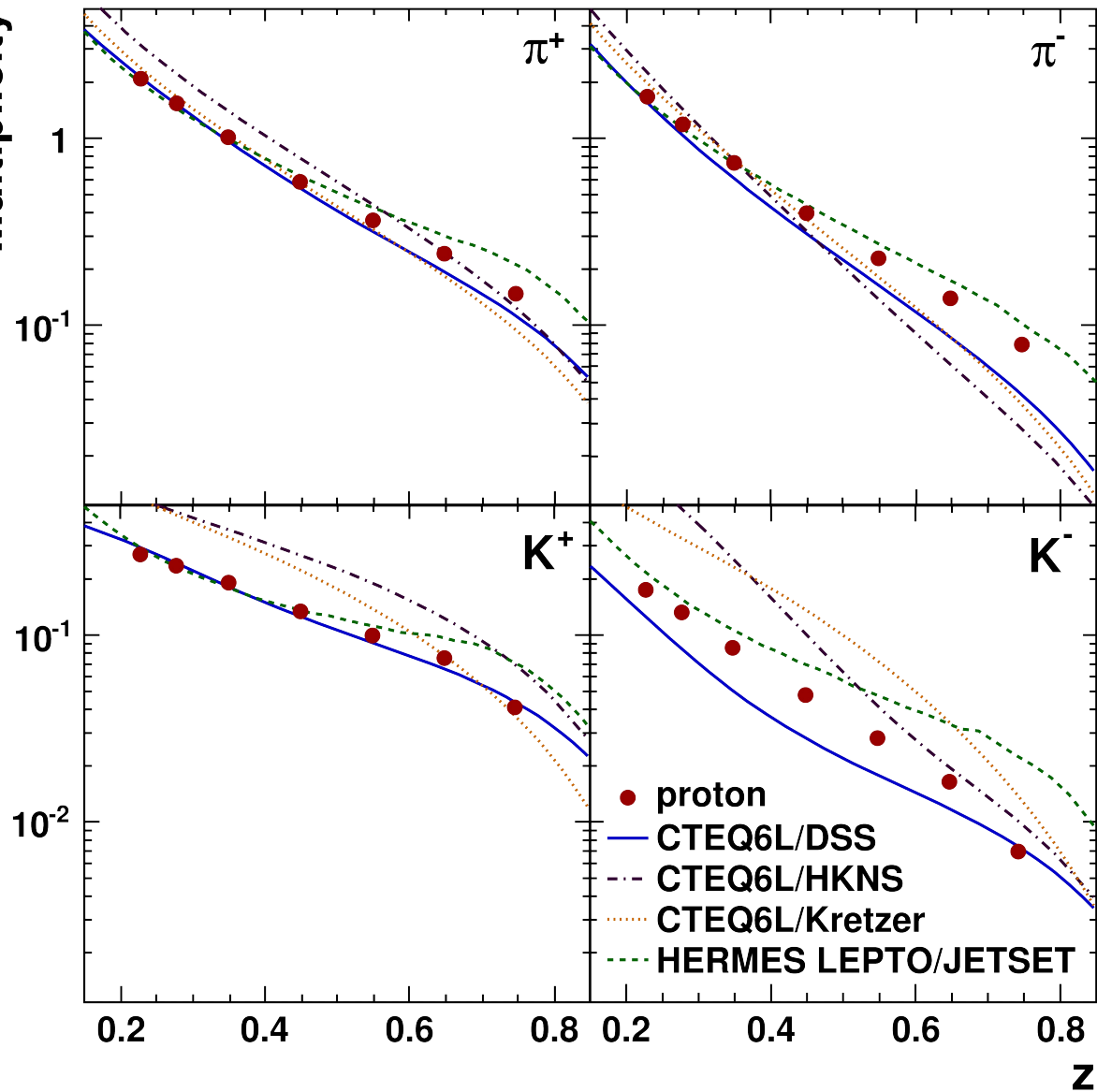
$$M_{\text{proton}}^{K^-}$$

K^- : sea object

$$M_{\text{deuteron}}^{K^-} \simeq M_{\text{proton}}^{K^-}$$

$$M^h(x_B, Q^2, z, P_{h\perp})$$

Multiplicity

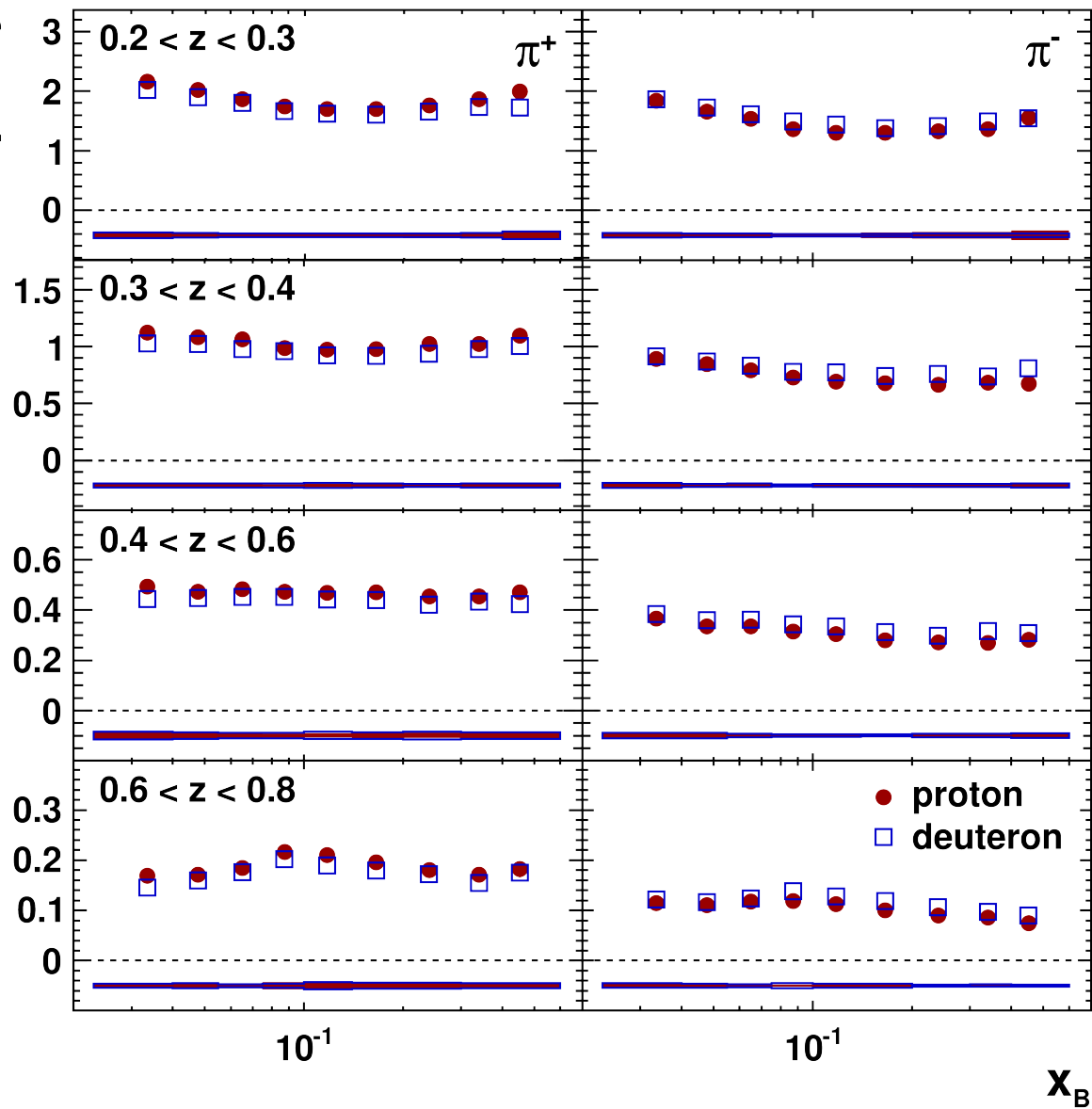


LO calculations

- ◆ Reasonable agreement between **DSS FFs** and **Data** for positively charged pions and kaons.
- ◆ Substantial differences between all FFs and **Data** for negatively charged kaons.

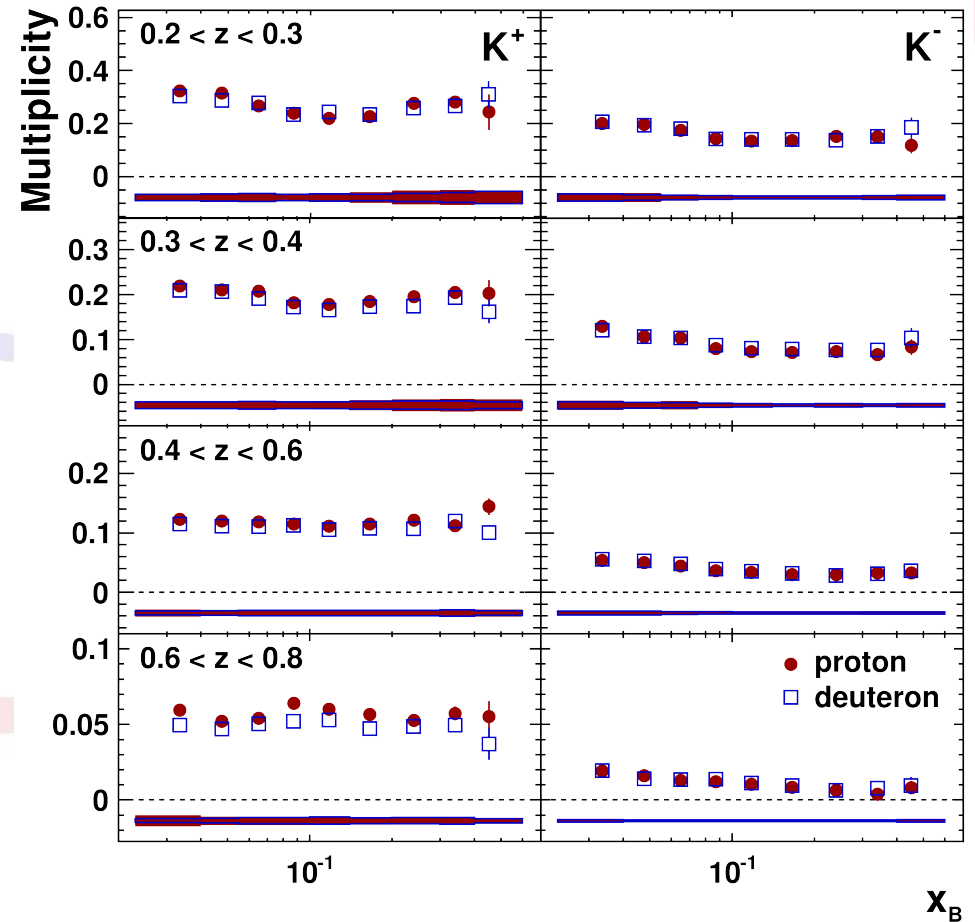
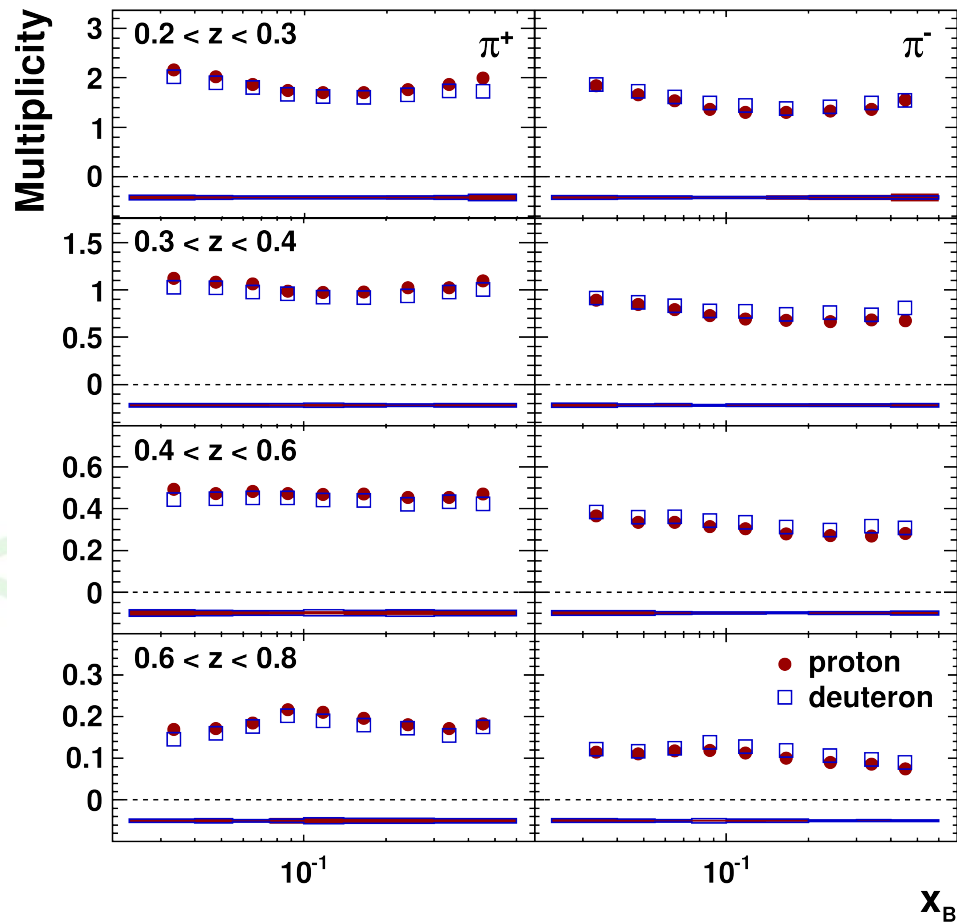
$$M^h(x_B, Q^2, z, P_{h\perp})$$

Multiplicity

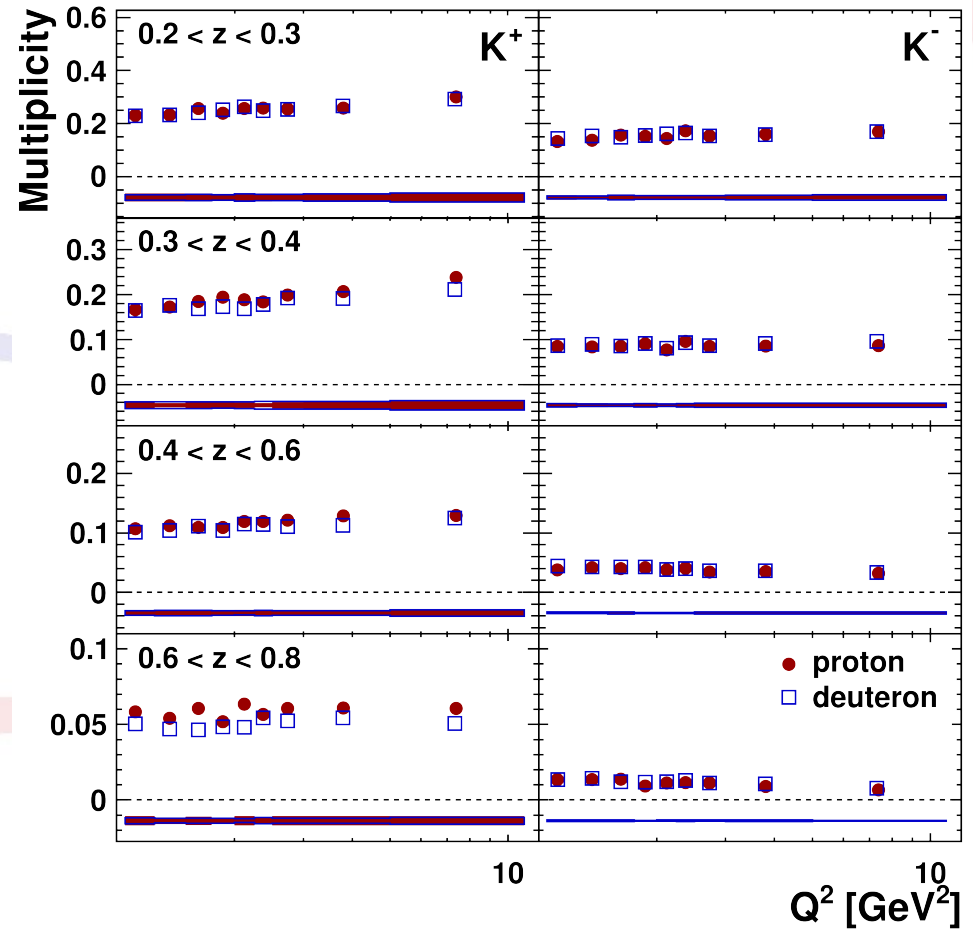
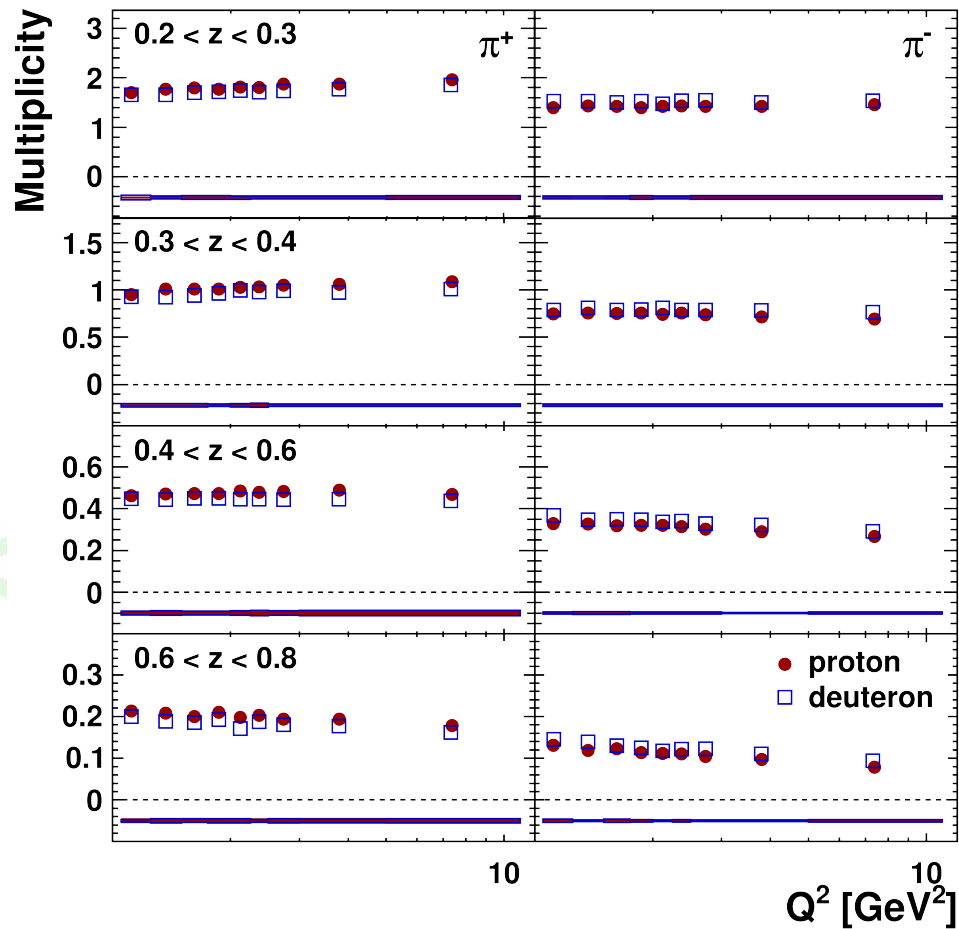


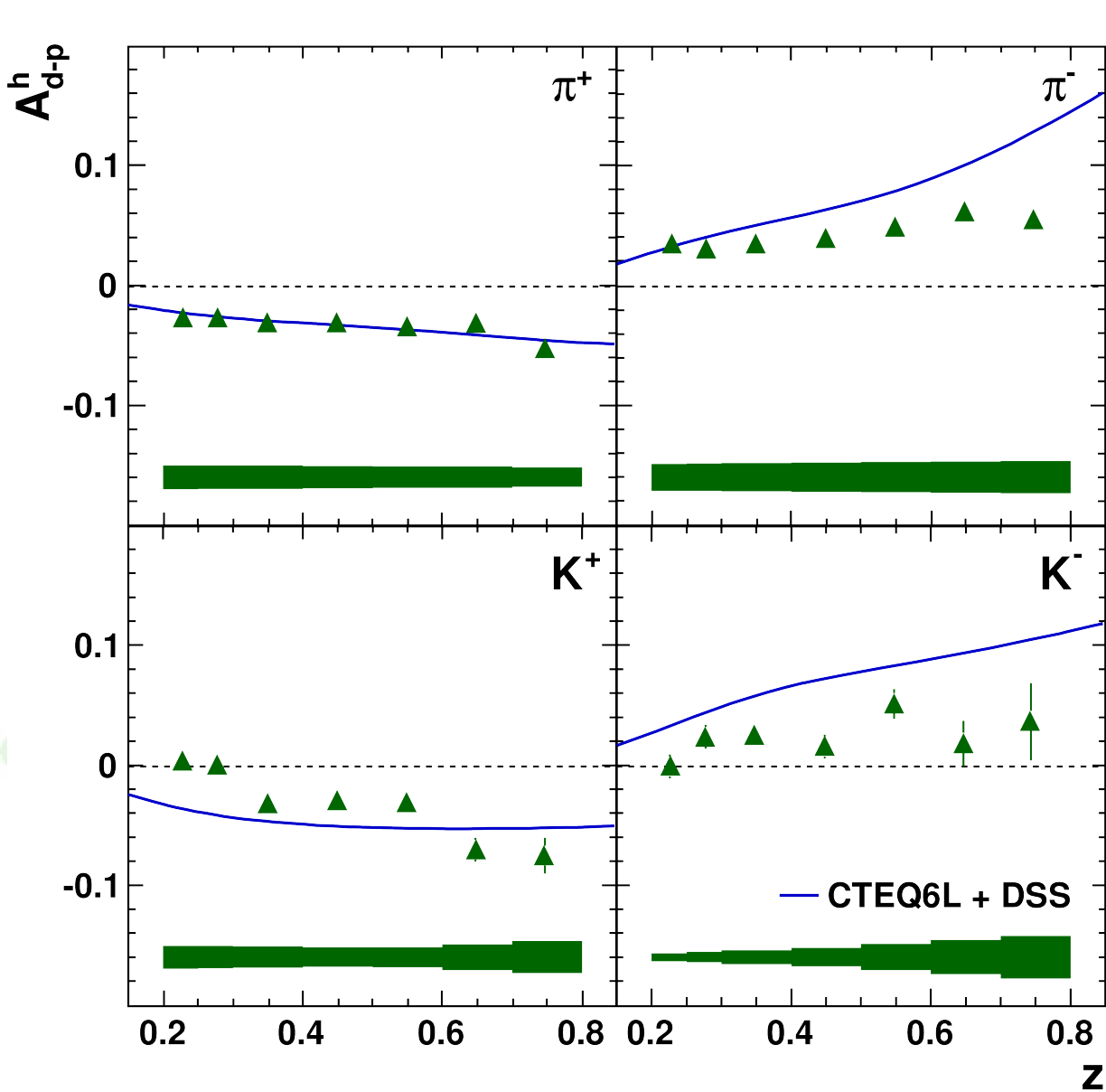
Almost no dependence on x_B

$$M^h(x_B, Q^2, z, P_{h\perp})$$



$$M^h(x_B, Q^2, z, P_{h\perp})$$

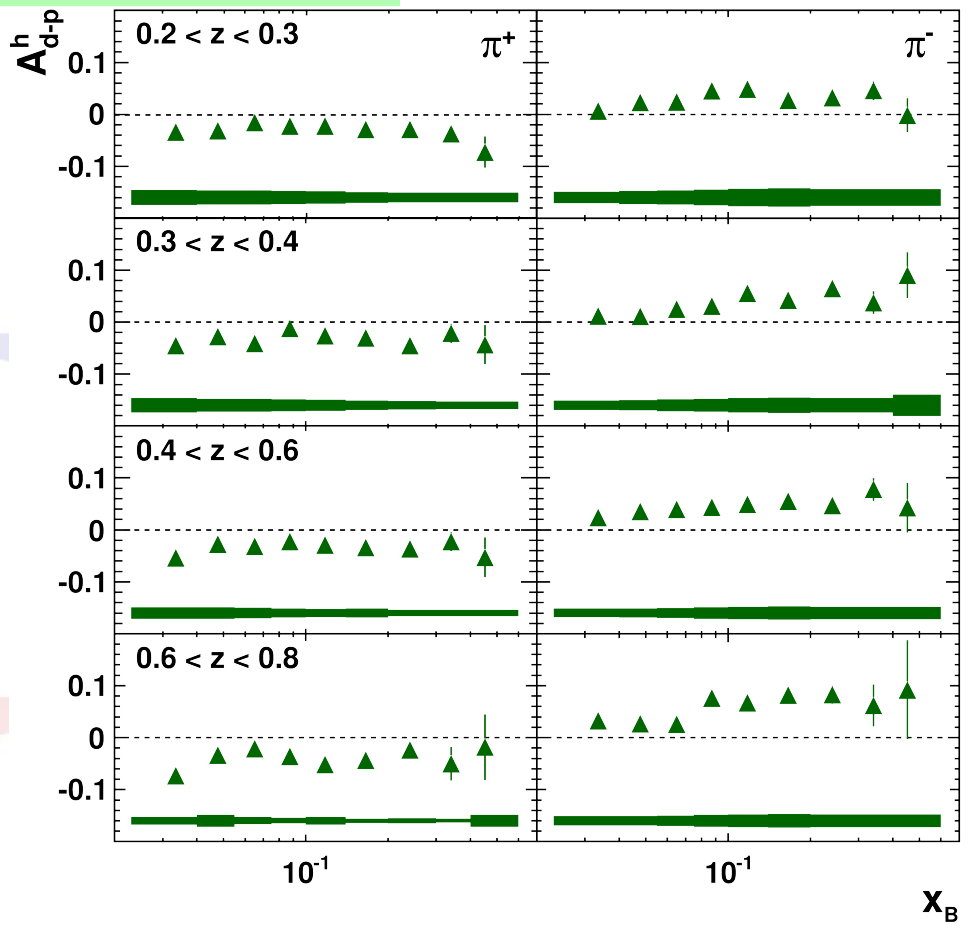
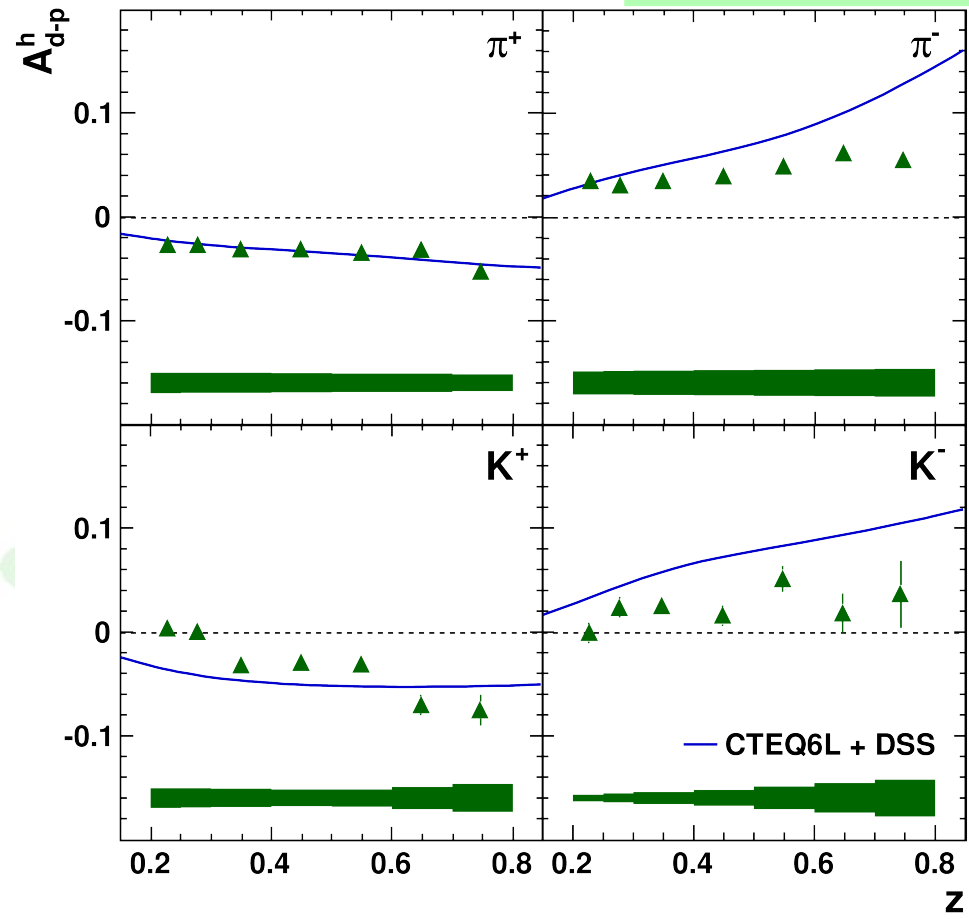




$$A_{d-p}^h = \frac{M_{\text{deuteron}}^h - M_{\text{proton}}^h}{M_{\text{deuteron}}^h + M_{\text{proton}}^h}$$

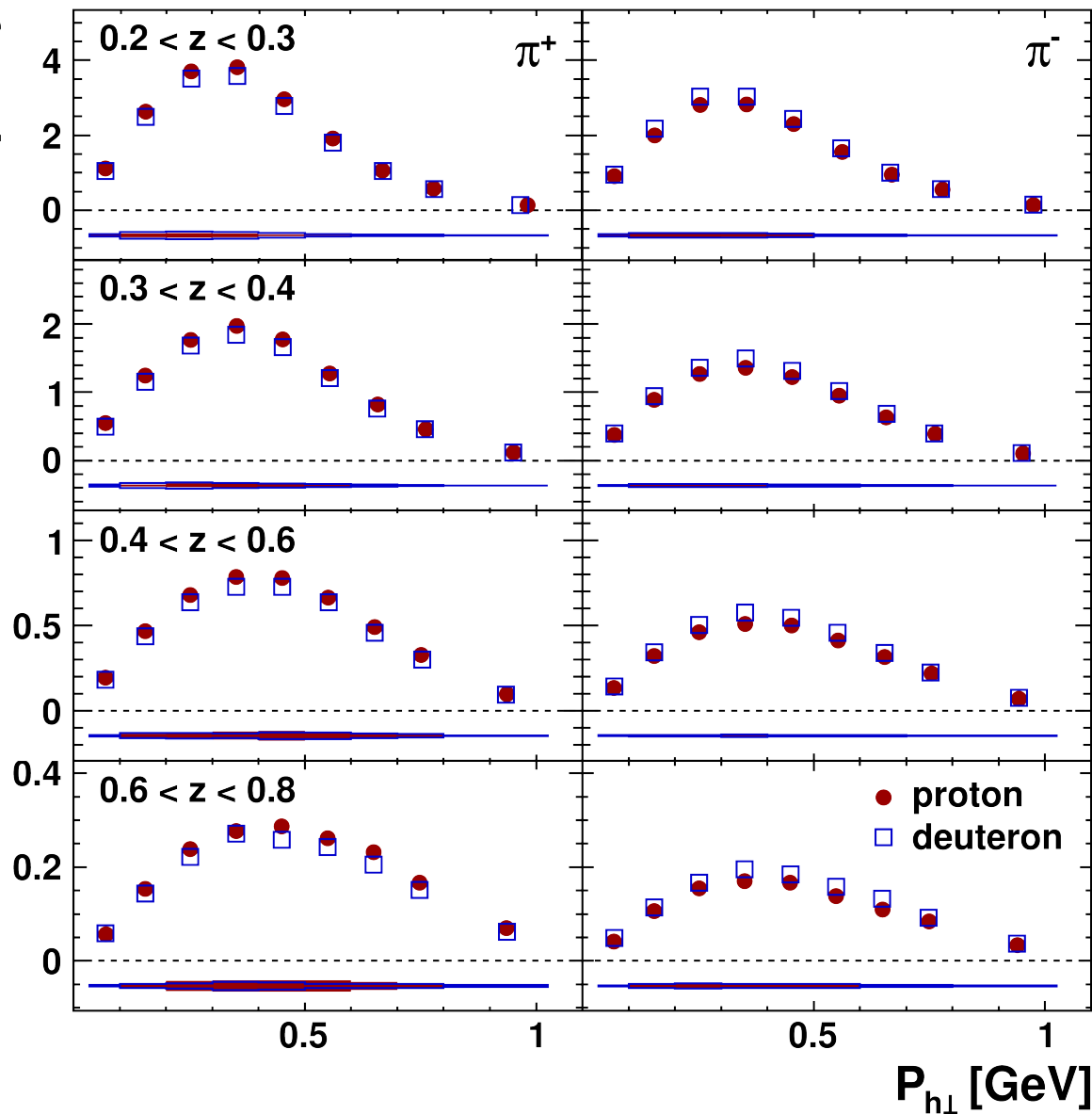
- ◆ Small magnitudes for asymmetries ($M_p^h \approx M_n^h$).
- ◆ A sign of the asymmetries reflects favored/unfavored fragmentation for different hadron species on different nuclei.
- ◆ Good description of data by DSS FF's for positively charged hadrons.

$$A_{d-p}^h = \frac{M_{\text{deuteron}}^h - M_{\text{proton}}^h}{M_{\text{deuteron}}^h + M_{\text{proton}}^h}$$



$$M^h(x_B, Q^2, z, P_{h\perp})$$

Multiplicity

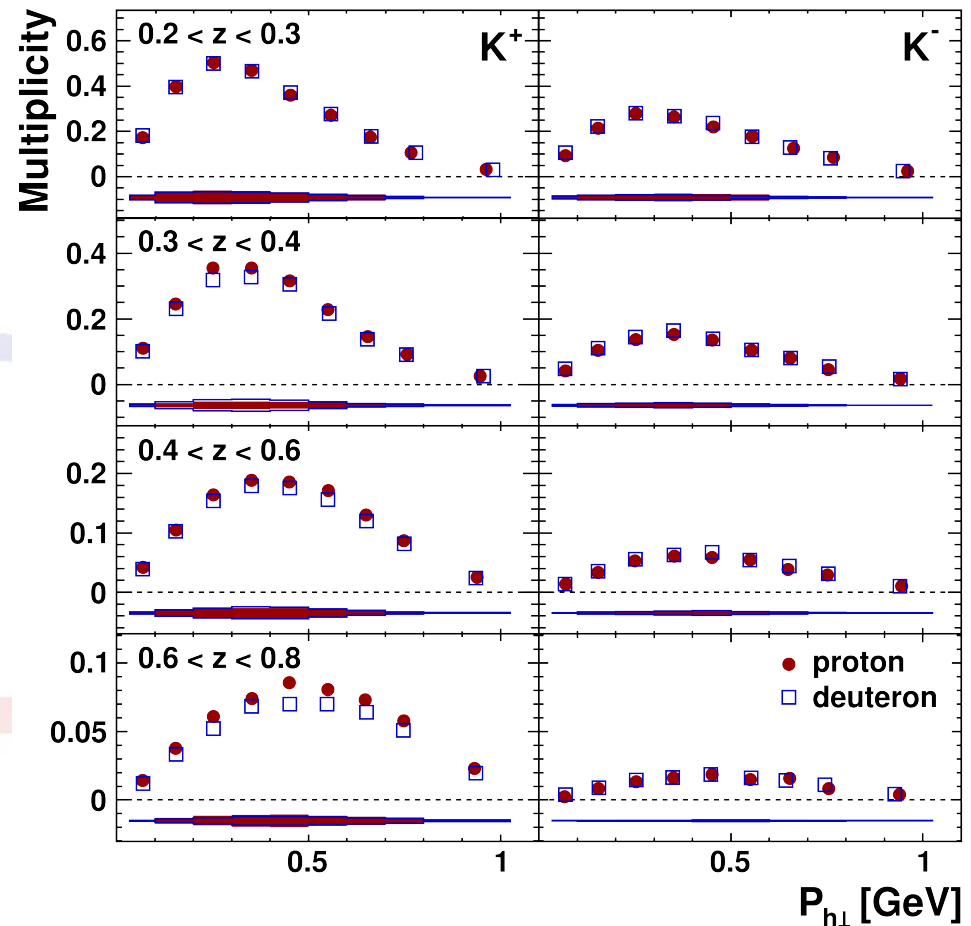
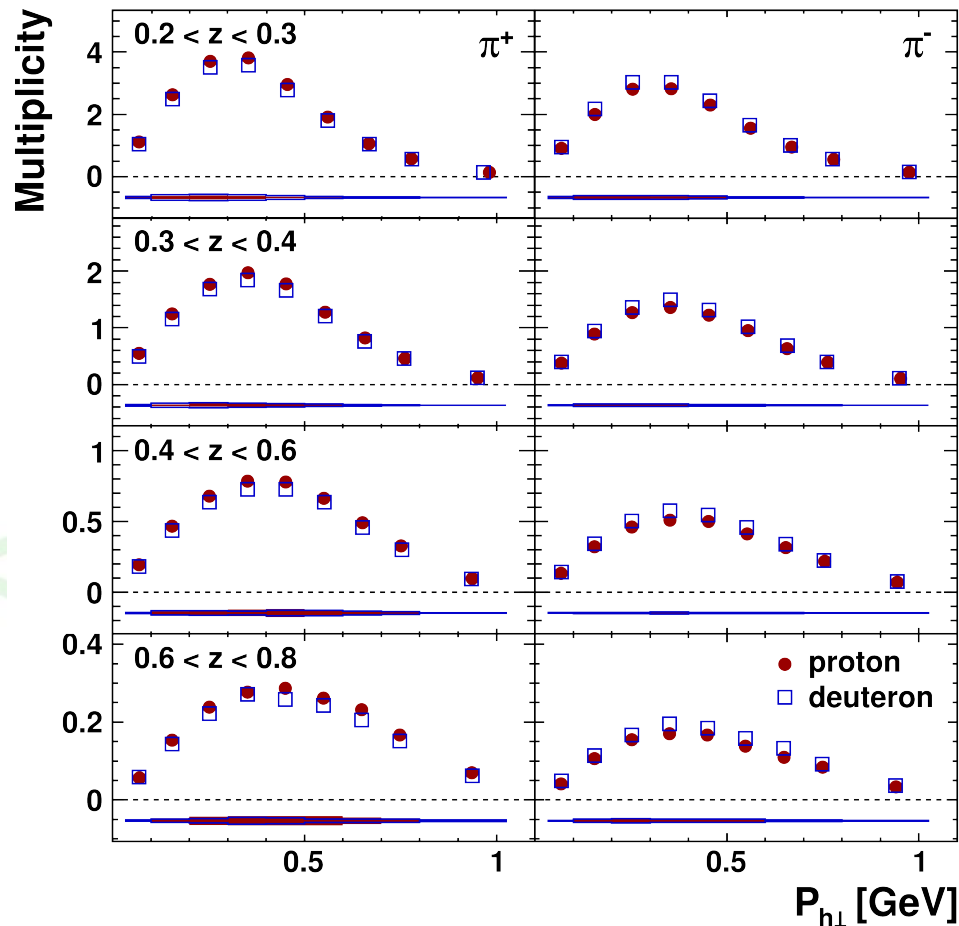


Access to the intrinsic transverse momentum of quark k_T and fragmentation p_T .

Gaussian ansatz :

$$\langle P_{h\perp}^2 \rangle = z^2 \langle k_T^2 \rangle + \langle p_T^2 \rangle$$

$$M^h(\mathbf{x}_B, Q^2, z, P_{h\perp})$$



Summary

- **High statistical data set for positively/negatively charged pion and kaon multiplicities on proton and deuteron.**
- **Fragmentation is favored for the hadrons containing the struck quark as a valence quark.**
- **Data will allow more reliable extraction of unfavored fragmentation functions.**
- **Dependence of multiplicities on hadron transverse momentum will provide constraints on the models of the fragmentation process.**