

INDIANA-ILLINOIS WORKSHOP
ON FRAGMENTATION
FUNCTIONS

BLOOMINGTON, IN,
DECEMBER 12-14, 2013

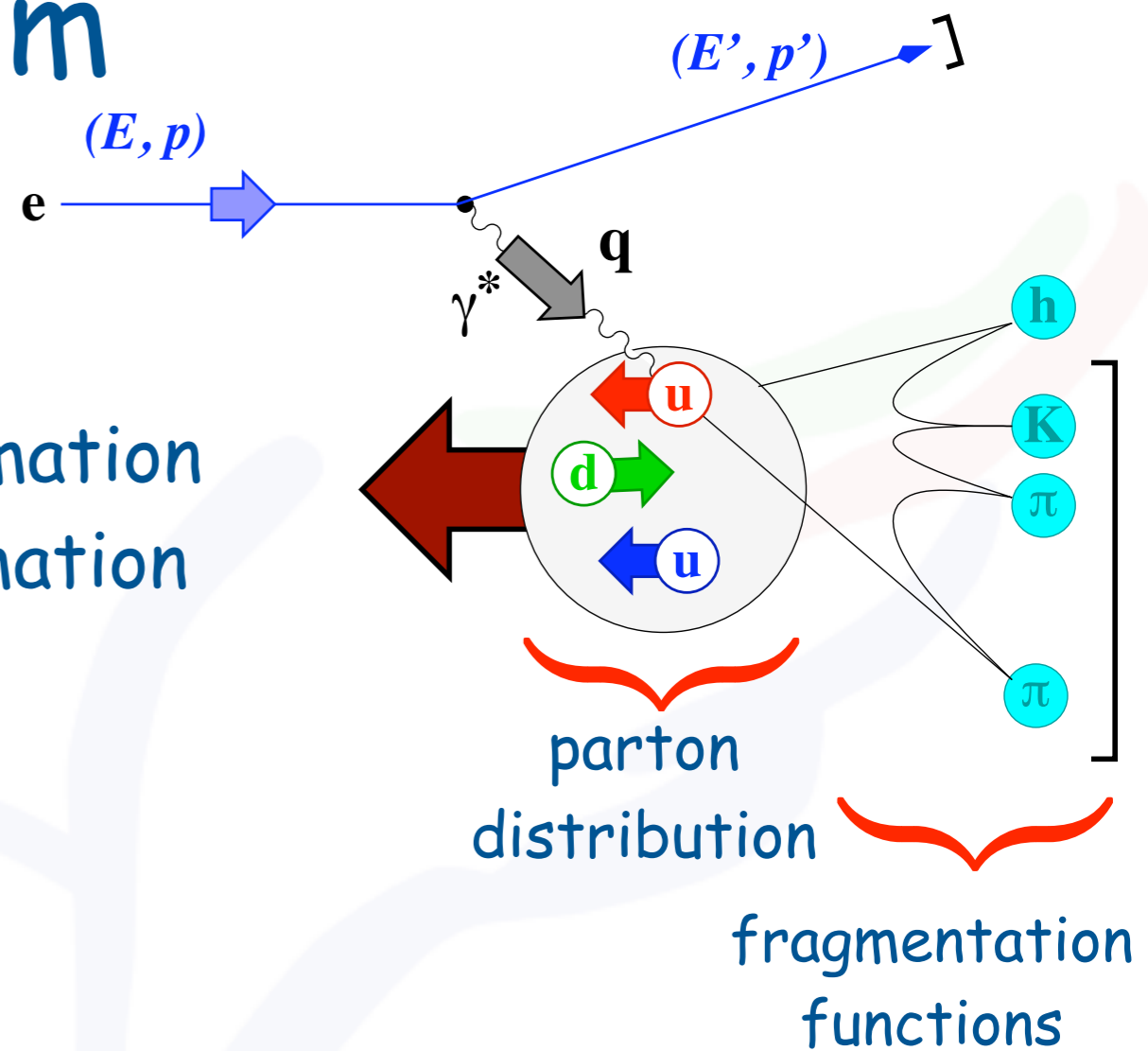


Semi-inclusive DIS off unpolarized targets

--the  hermes perspective--

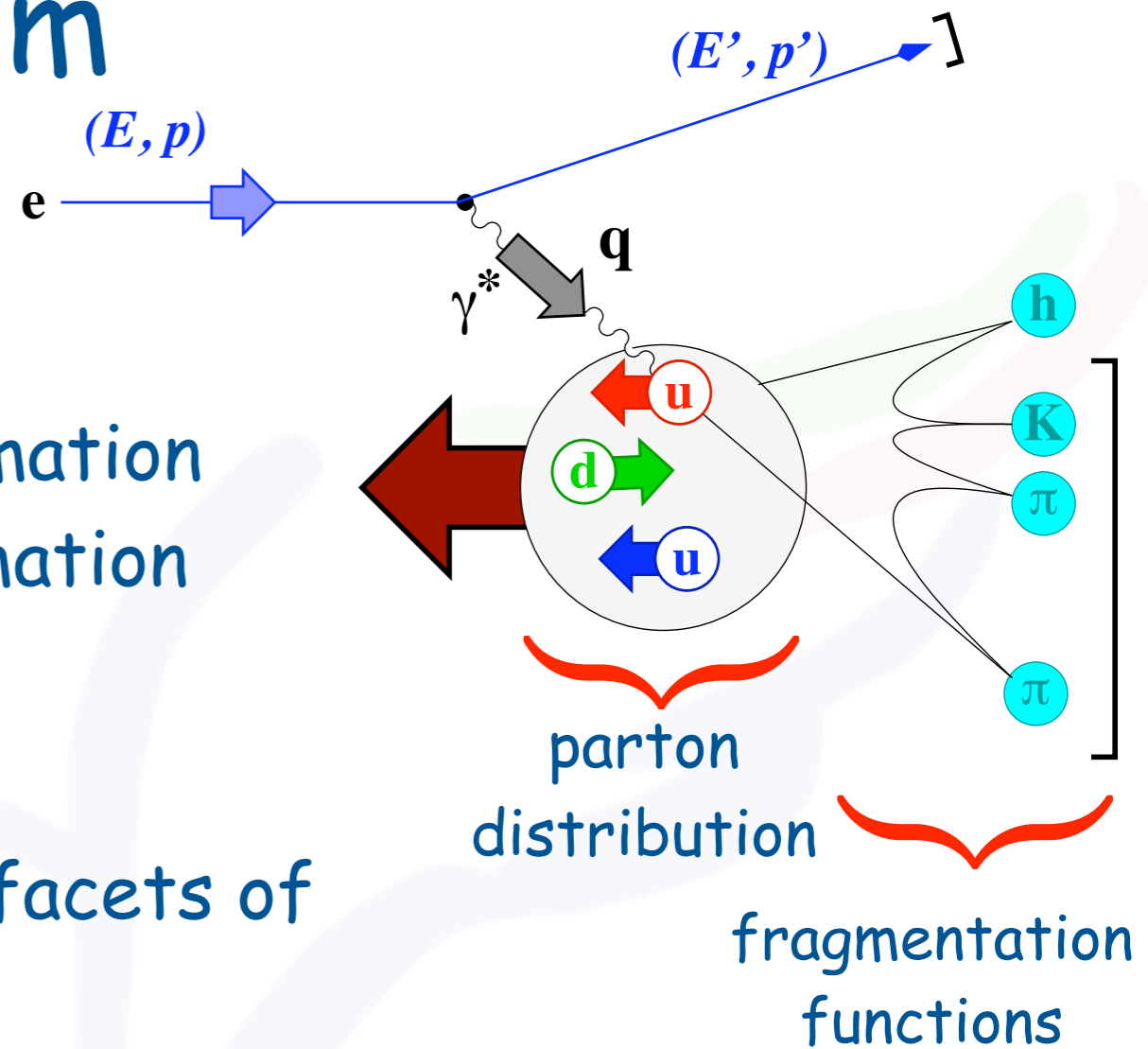
Why study SIDIS from unpolarized targets?

- Semi-inclusive DIS provides information on both hadron structure and formation



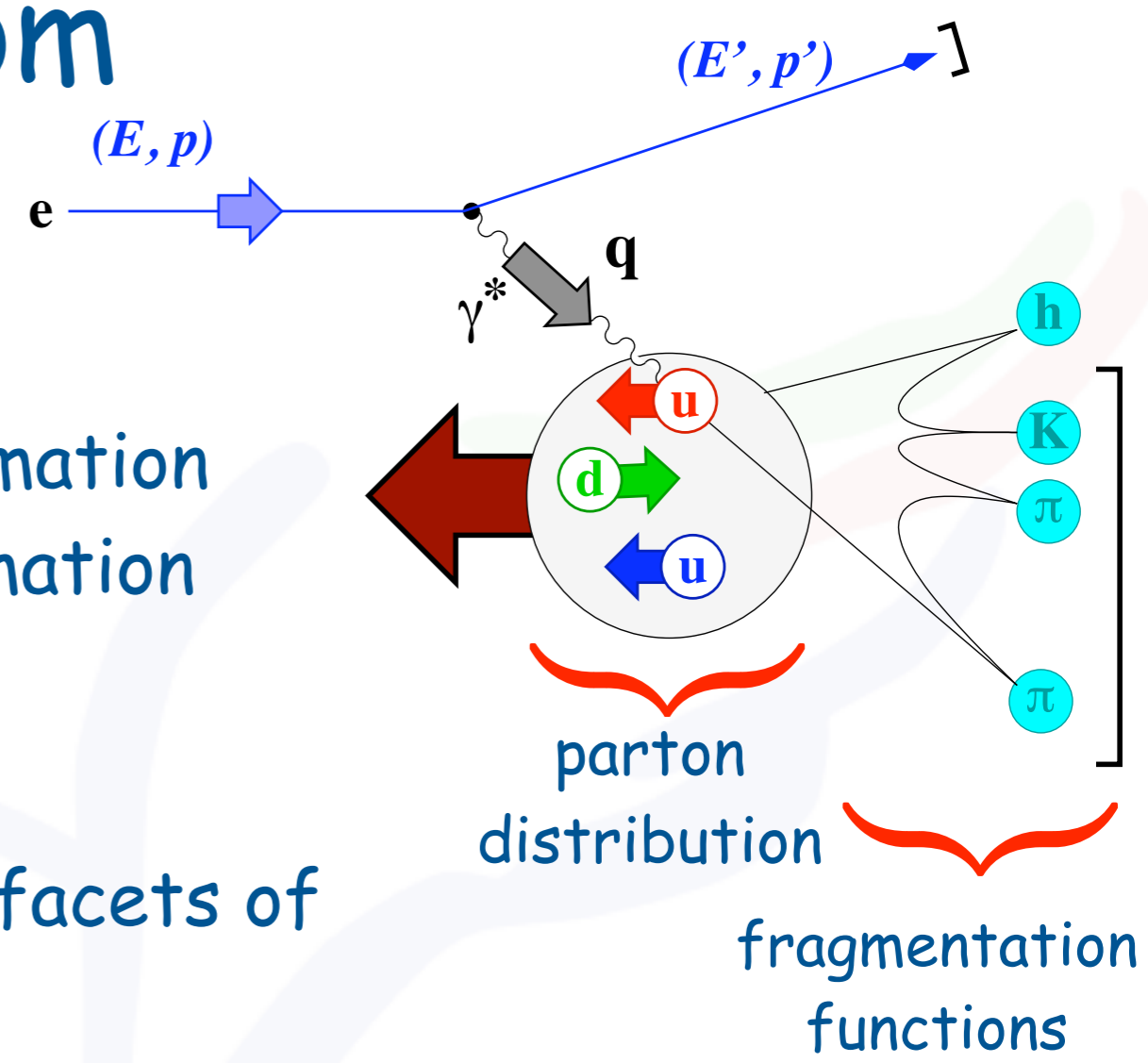
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- (probably) easiest one to study facets of hadron structure, even in 3D



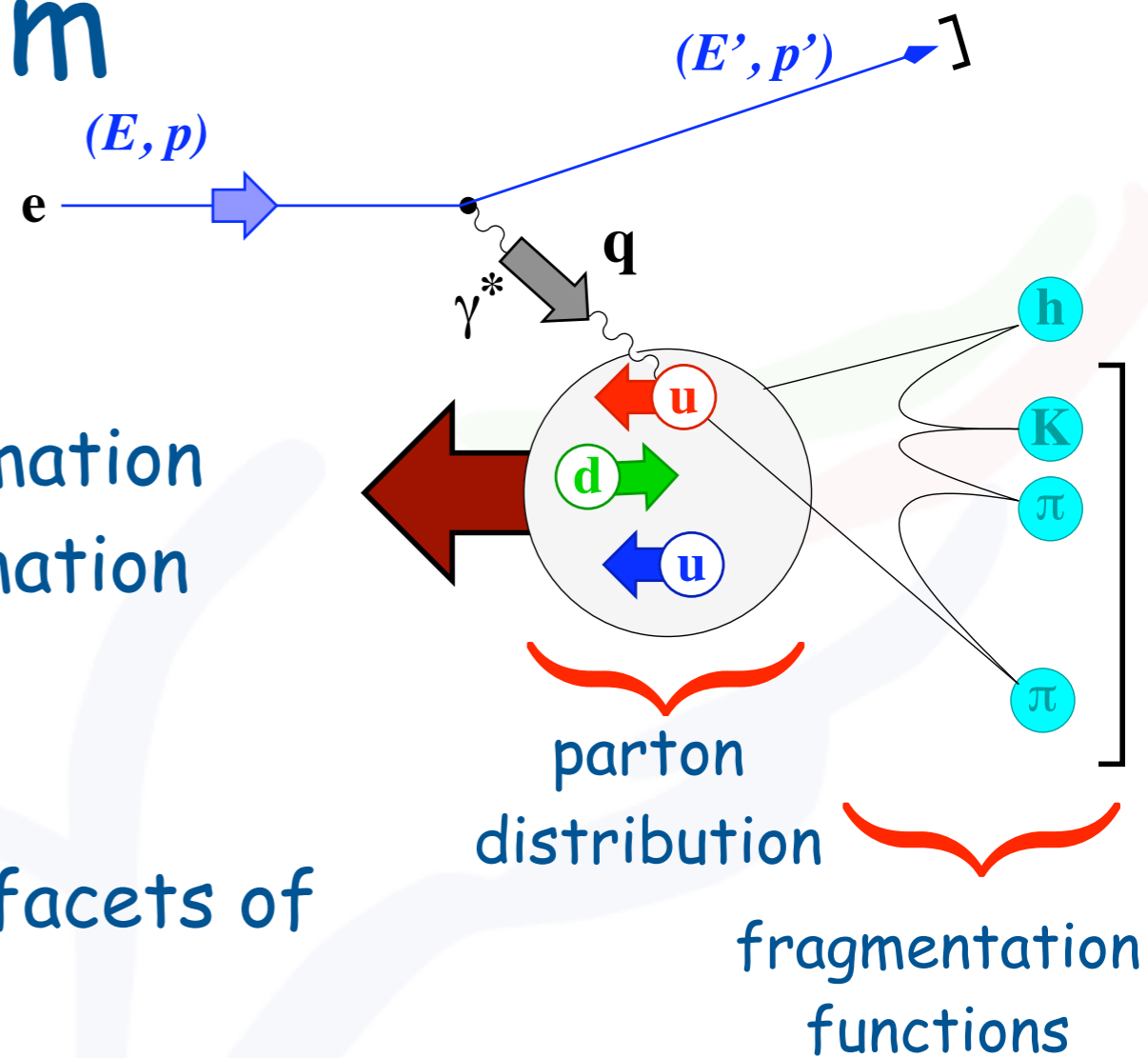
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 - both are ingredients of basically every (spin) asymmetry



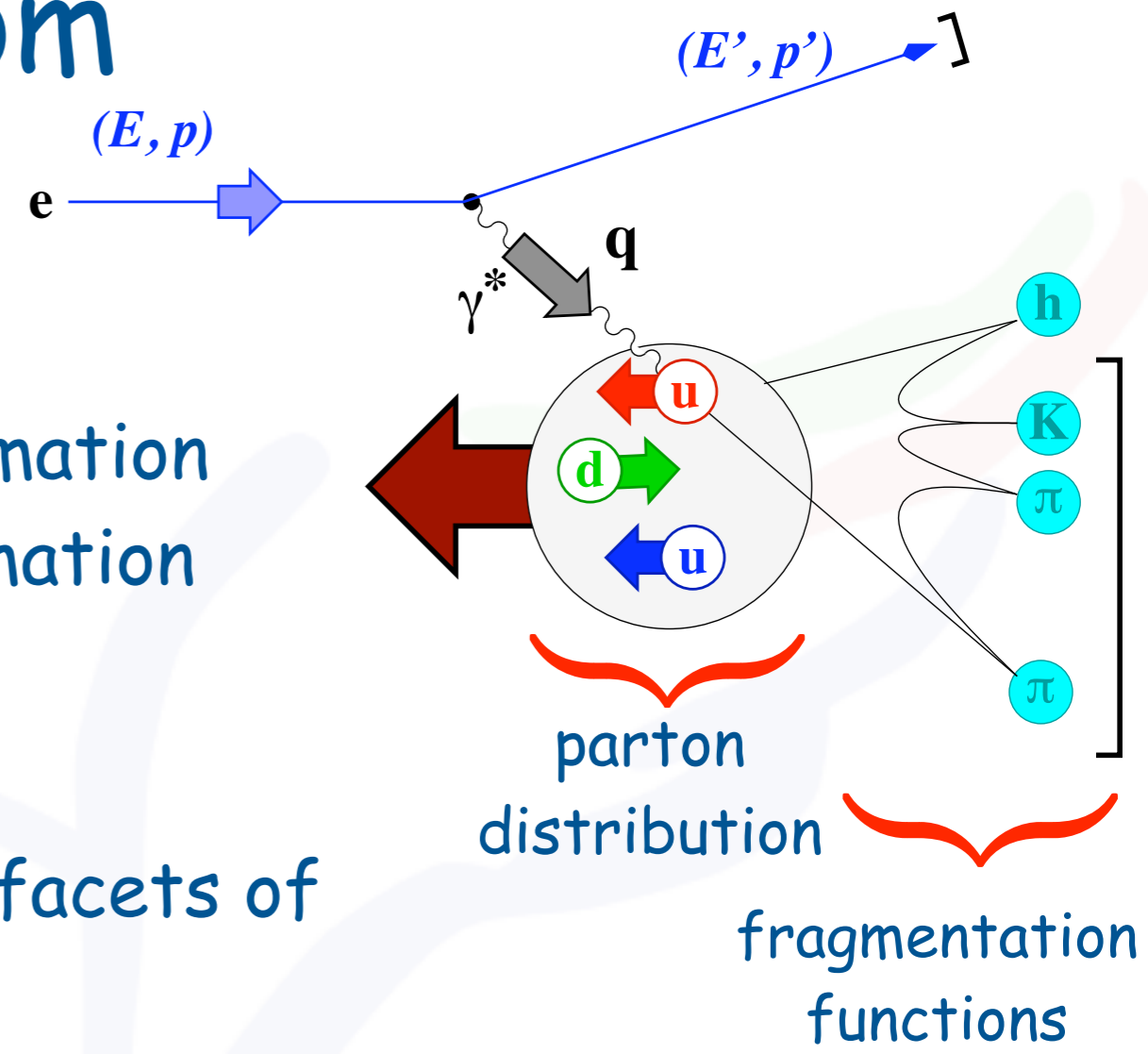
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- in semi-inclusive DIS, f_1 couples to D_1 fragmentation function
- both are ingredients of basically every (spin) asymmetry
- complimentary info on FFs to e^+e^- (e.g., charge separation)
- nuclear targets provide laboratory for hadronization studies



Polarization-averaged cross section

$$F_{XY,Z} = F_{XY,Z}(x, y, z, P_{h\perp})$$

target polarization
↓
↑ beam polarization ↑ virtual-photon polarization

$$\frac{d^5\sigma}{dx dy dz d\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \{F_{UU,T} + \epsilon F_{UU,L}\}$$

$$\gamma = \frac{2Mx}{Q}$$
$$\epsilon = \frac{1 - y - \frac{1}{4}\gamma^2 y^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}\gamma^2 y^2}$$

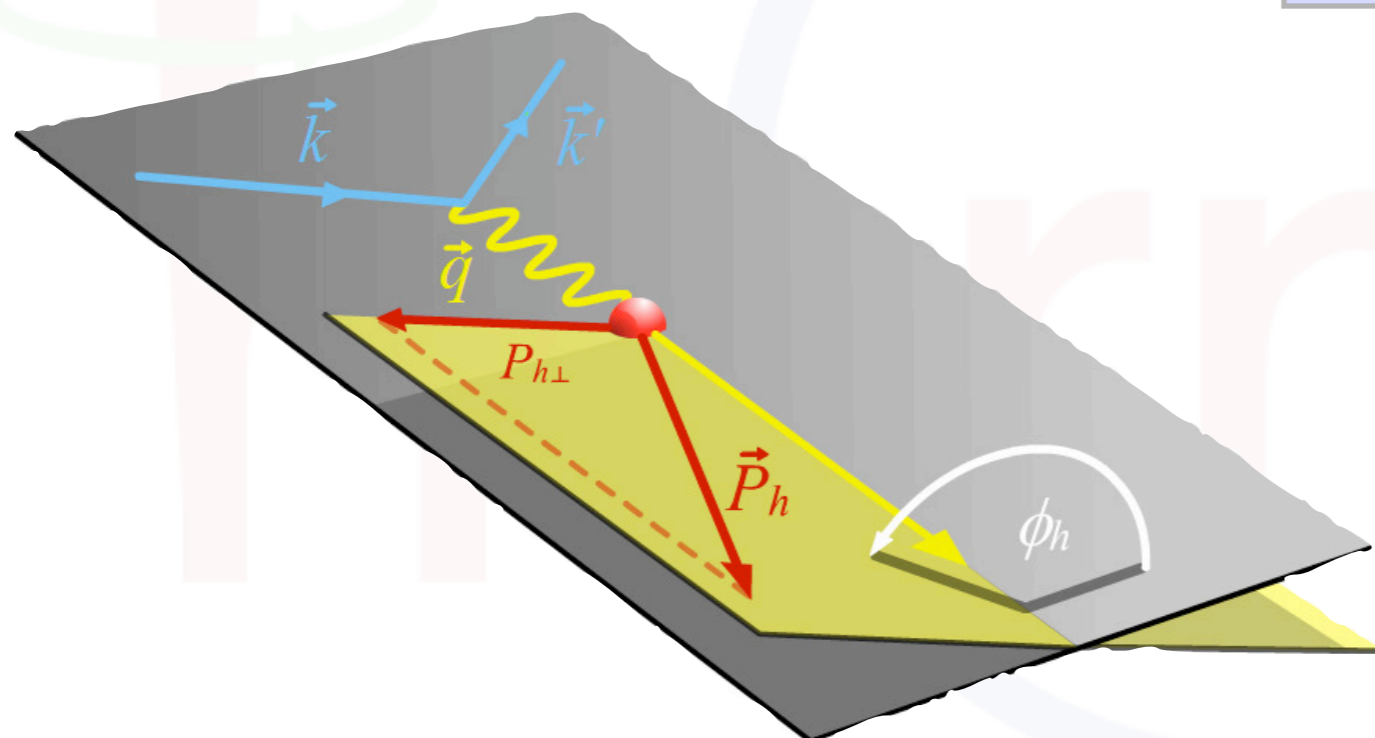
[see, e.g., Bacchetta et al., JHEP 0702 (2007) 093]

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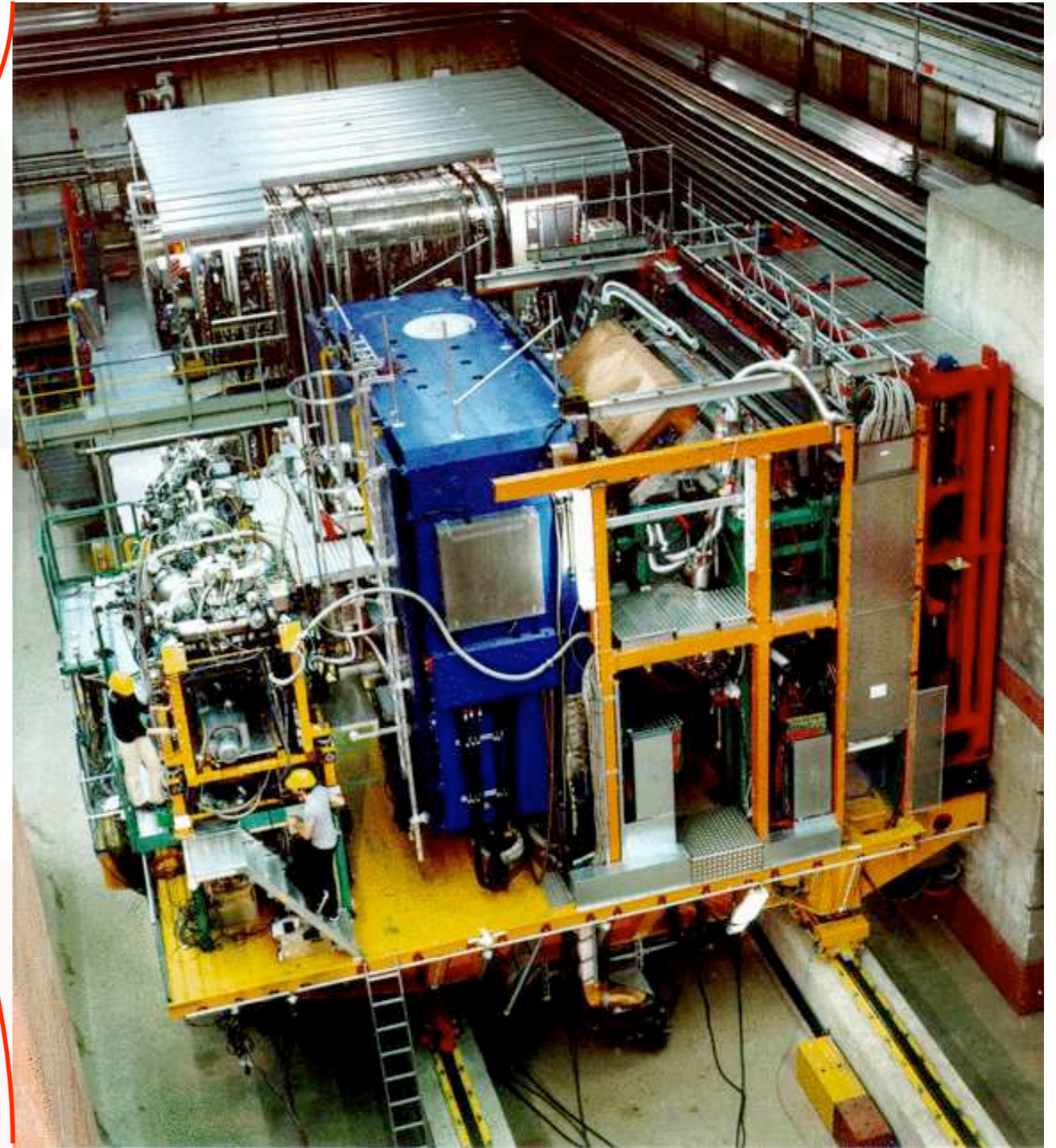
[see, e.g., Bacchetta et al., JHEP 0702 (2007) 093]

Some experimental challenges ...

- pure targets
- large acceptance
- excellent particle identification
- no spin asymmetry \rightarrow few systematics cancel
- efficiencies
- absolute luminosity
- acceptance
- smearing

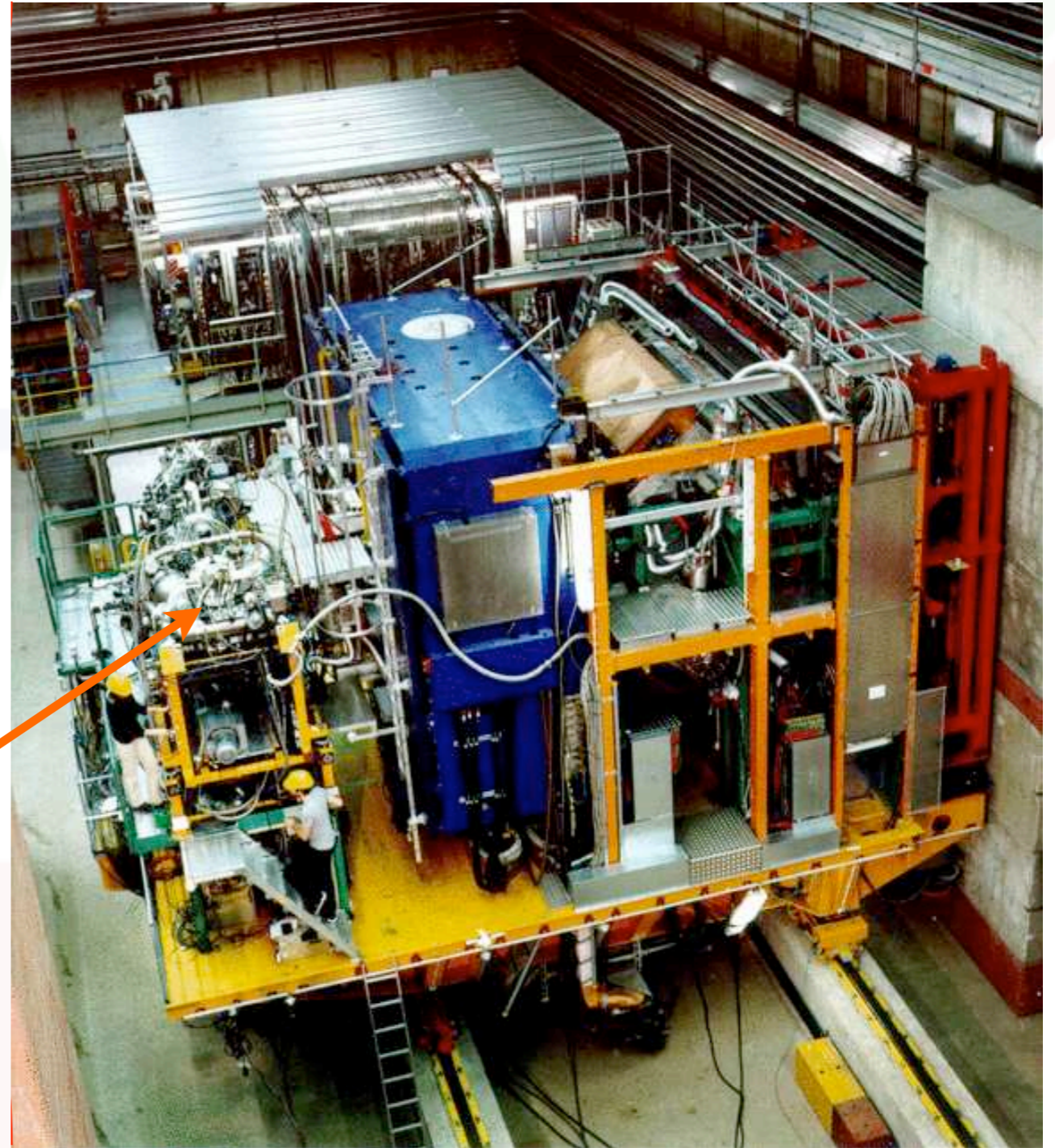
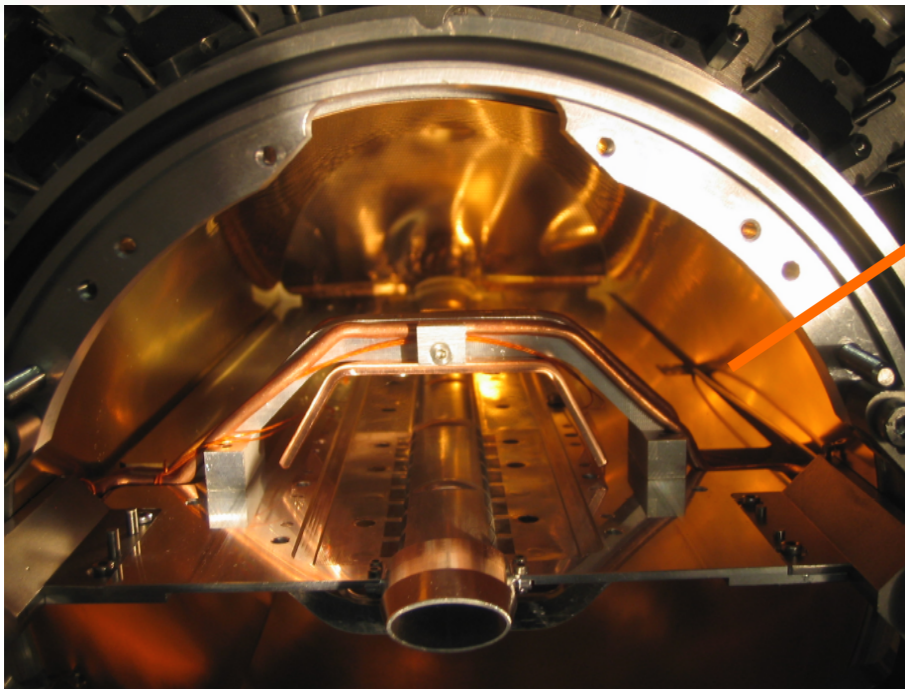
The HERMES Experiment

27.5 GeV e^+/e^- beam of HERA

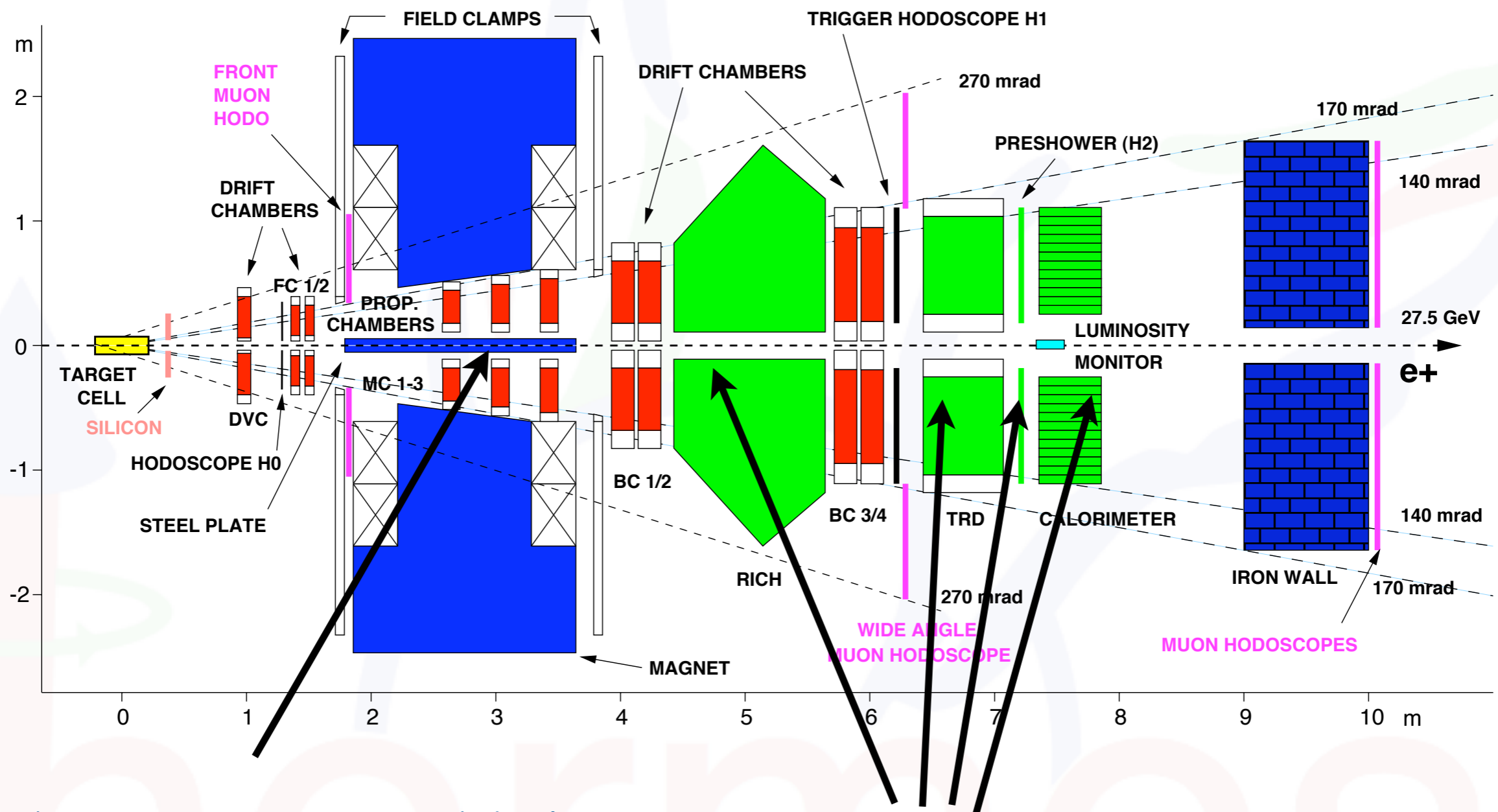


The HERMES Experiment

- pure gas targets
- internal to lepton ring
- unpolarized (^1H ... Xe)
- long. polarized: ^1H , ^2H , ^3He
- transversely polarized: ^1H



... and solutions



two (mirror-symmetric) halves
-> no homogenous azimuthal coverage

Particle ID detectors allow for

- lepton/hadron separation
- RICH: pion/kaon/proton discrimination $2\text{GeV} < p < 15\text{GeV}$

... and solutions ...

$$\frac{d^5 \sigma}{dx dy dz d\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \{F_{UU,T} + \epsilon F_{UU,L}$$
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hadron multiplicity:
normalize to inclusive DIS
cross section

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➡ this talk

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➡ S. Gliske (Saturday)

... geometric acceptance ...

extract acceptance from Monte Carlo simulation?

$$\epsilon(\phi, \Omega) = \frac{\epsilon(\phi, \Omega) \sigma_{UU}(\phi, \Omega)}{\sigma_{UU}(\phi, \Omega)}$$

$$\Omega = x, y, z, \dots$$

simulated acceptance

simulated cross section

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extract acceptance from Monte Carlo simulation?

$$\begin{aligned}\epsilon(\phi, \Omega) &= \frac{\epsilon(\phi, \Omega) \sigma_{UU}(\phi, \Omega)}{\sigma_{UU}(\phi, \Omega)} \\ &\neq \frac{\int d\Omega \sigma_{UU}(\phi, \Omega) \epsilon(\phi, \Omega)}{\int d\Omega \sigma_{UU}(\phi, \Omega)}\end{aligned}$$

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"Aus Differenzen und Summen kürzen nur die Dummen."

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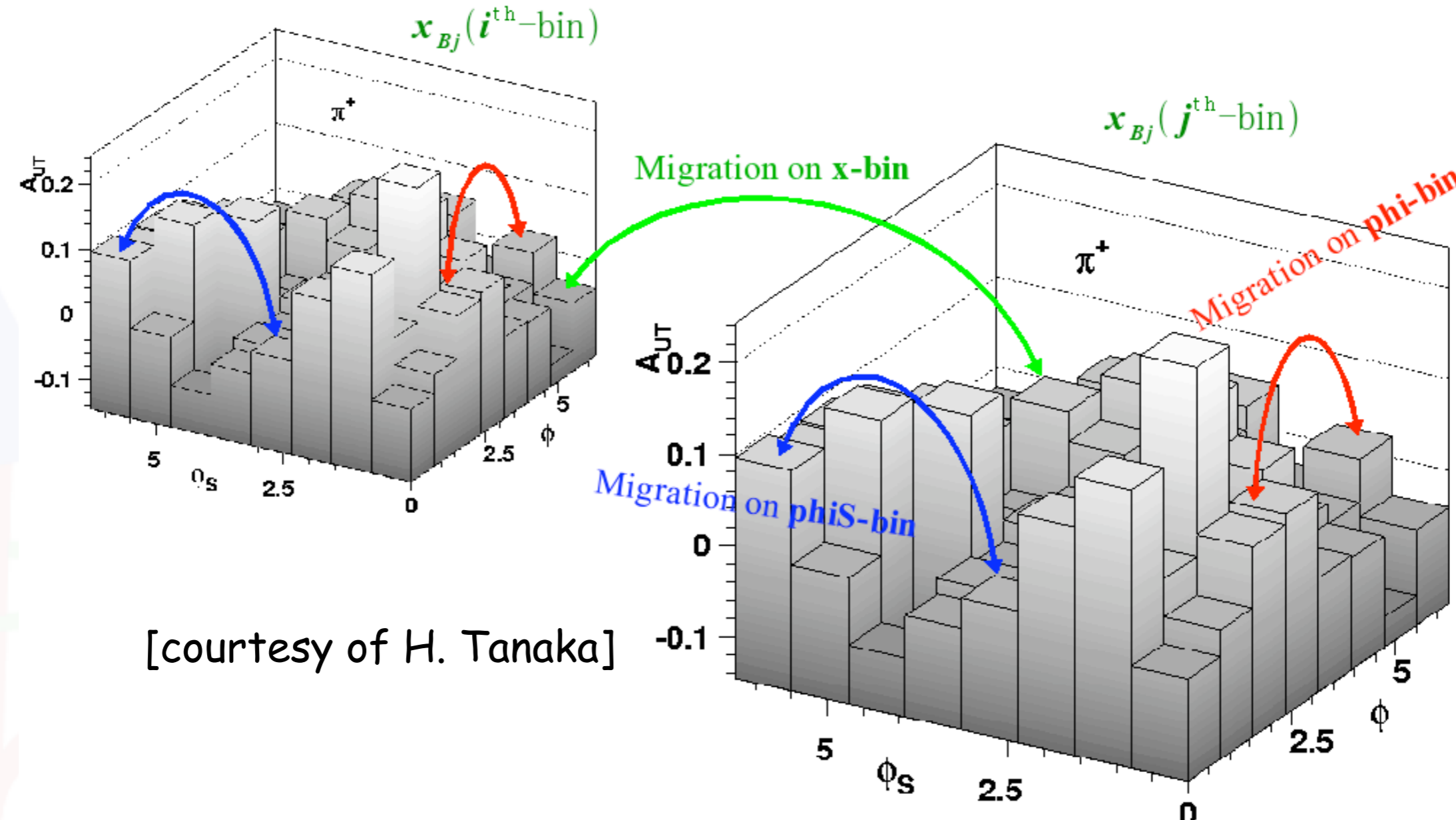
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$$\neq \int d\Omega \epsilon(\phi, \Omega) \equiv \epsilon(\phi)$$

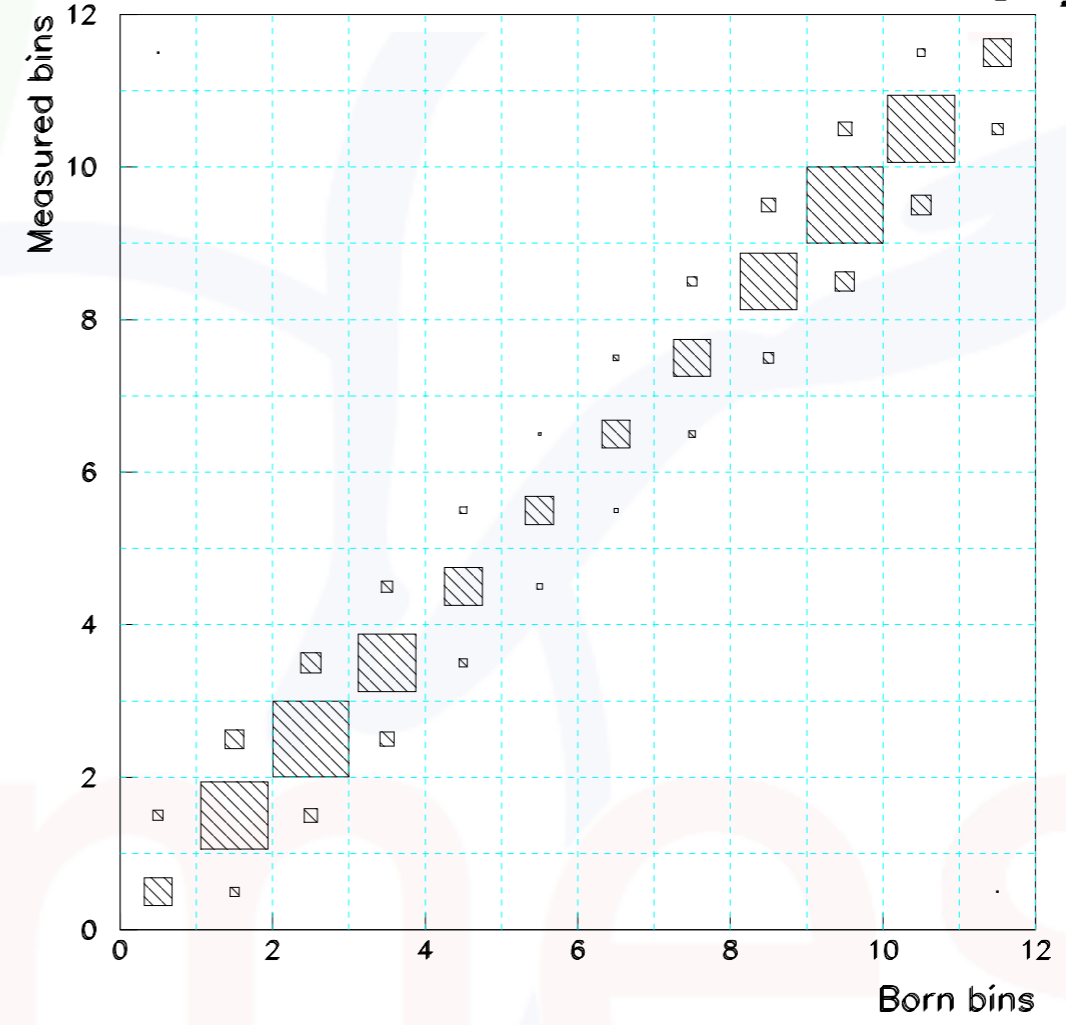
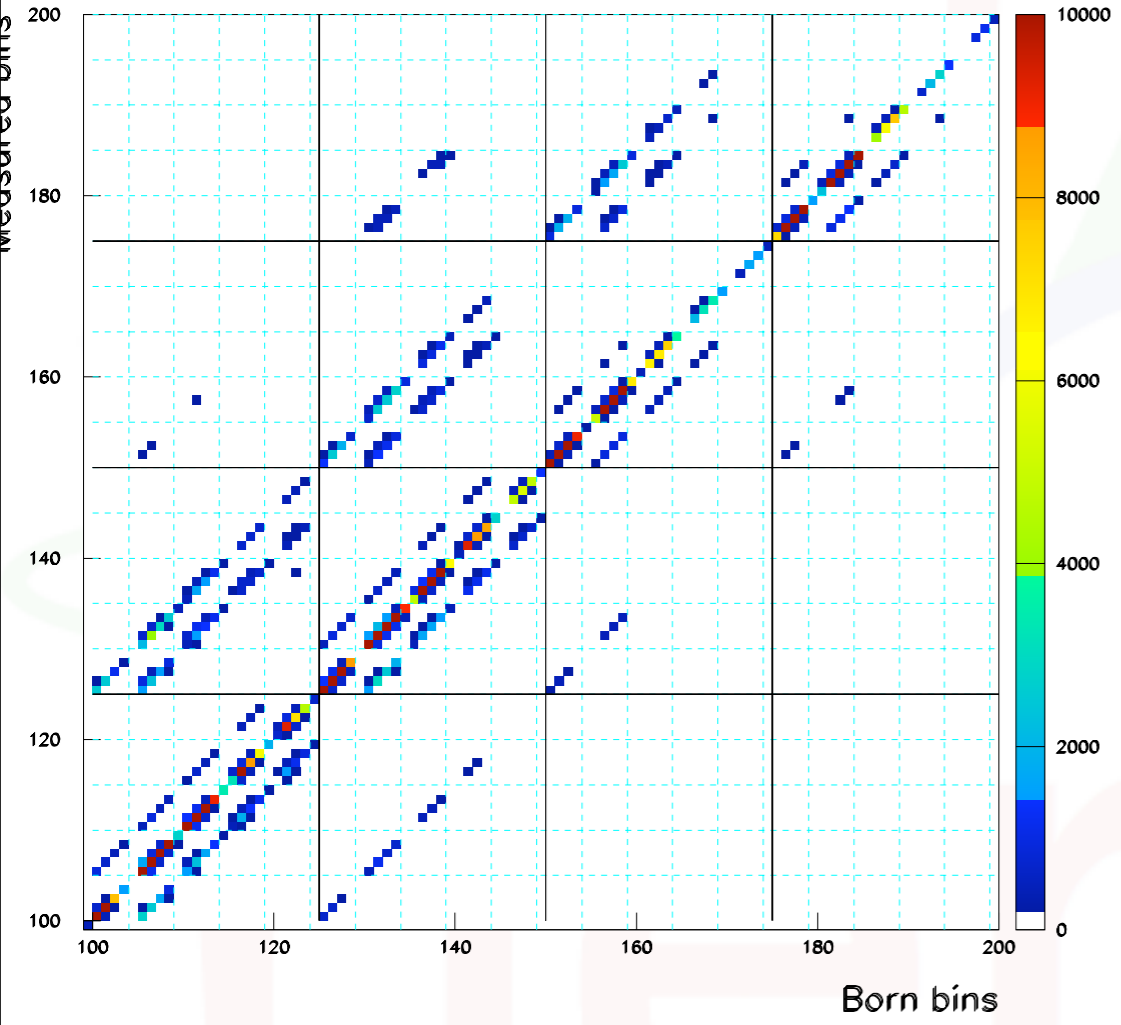
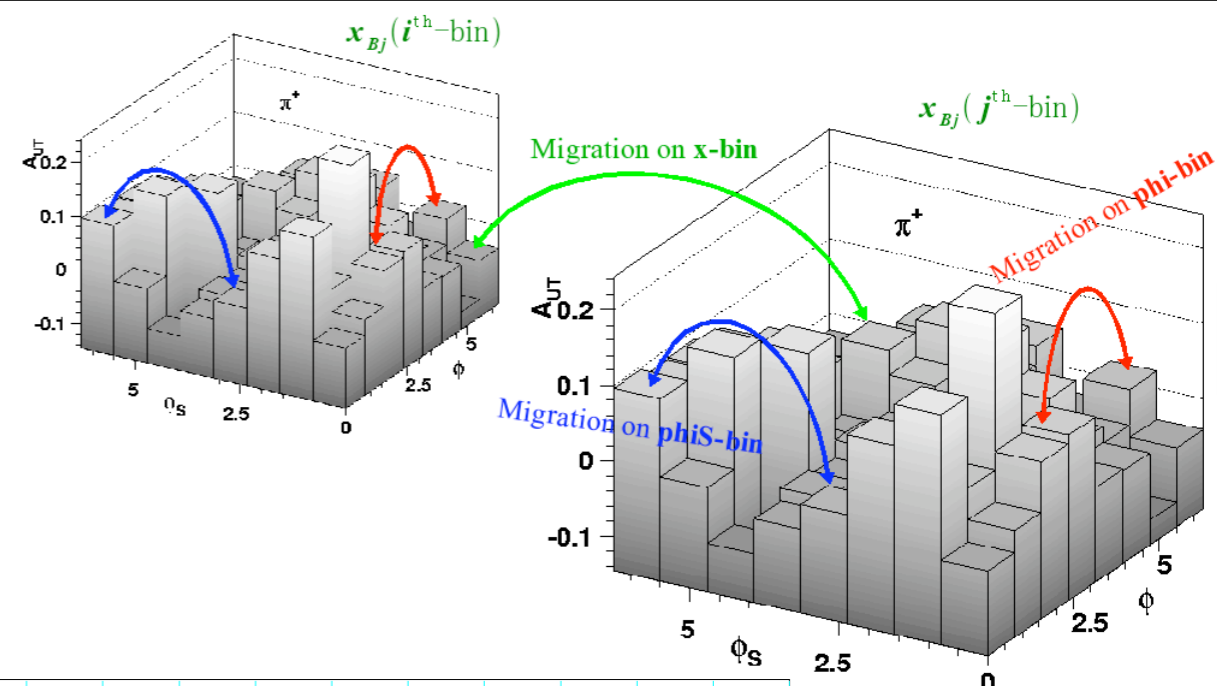
Cross-section model does NOT CANCEL in general when integrating numerator and denominator over (large) ranges in kinematic variables!

... event migration ...



[courtesy of H. Tanaka]

... event migration ...



- migration correlates yields in different bins
- can't be corrected properly in bin-by-bin approach

... event migration -> unfolding

$$\mathcal{Y}^{\text{exp}}(\Omega_i) \propto \sum_{j=1}^N S_{ij} \int_j d\Omega d\sigma(\Omega) + \mathcal{B}(\Omega_i)$$



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- experimental yield in i^{th} bin depends on all Born bins j ...

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- ... and on BG entering kinematic range from outside region

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- smearing matrix S_{ij} embeds information on migration
- determined from Monte Carlo - independent of physics model in limit of infinitesimally small bins and/or flat acceptance/cross-section in every bin
- in real life: dependence on BG and physics model due to finite bin sizes

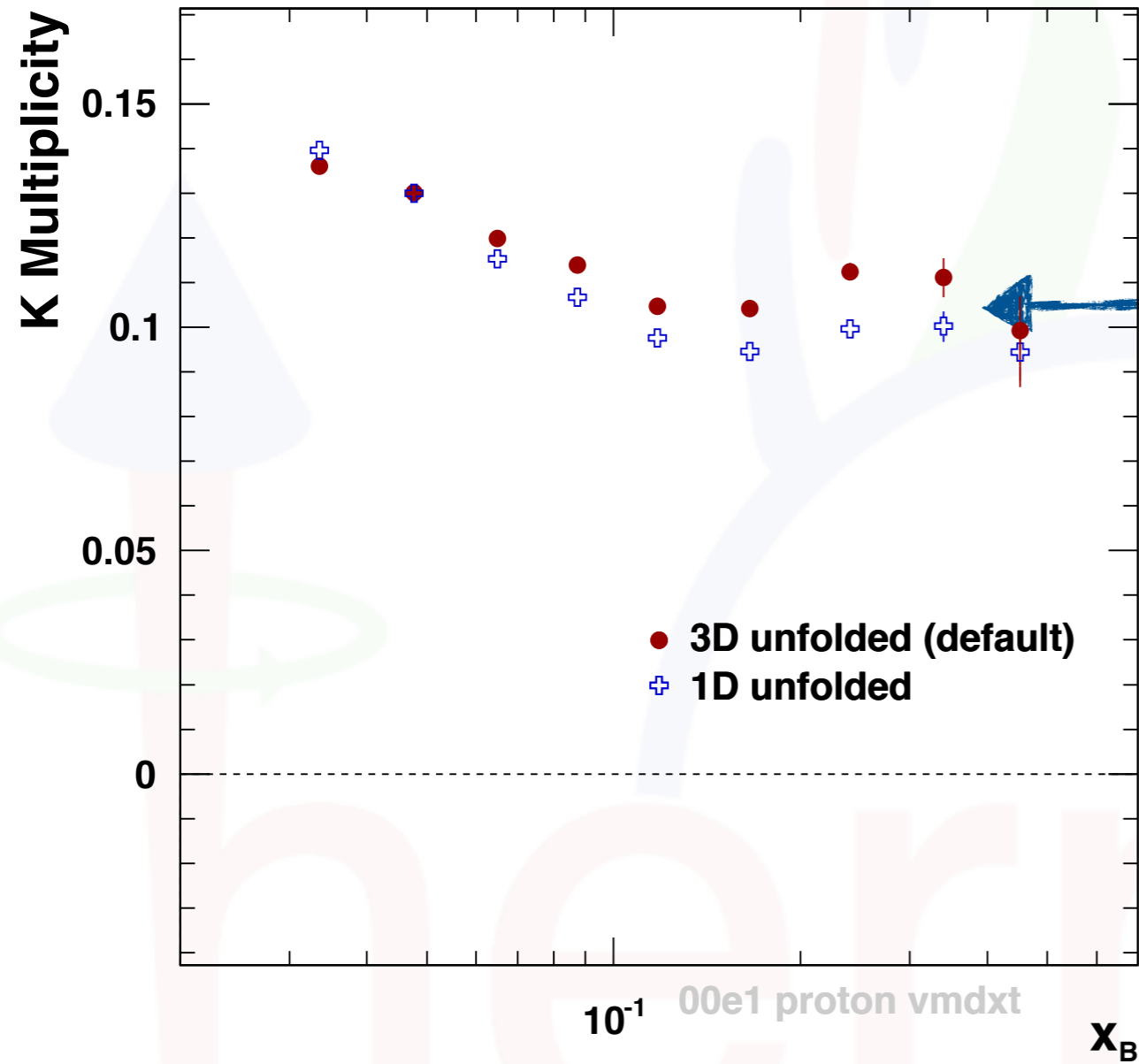
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- inversion of relation gives Born cross section from measured yields

Multi-D vs. 1D unfolding at work

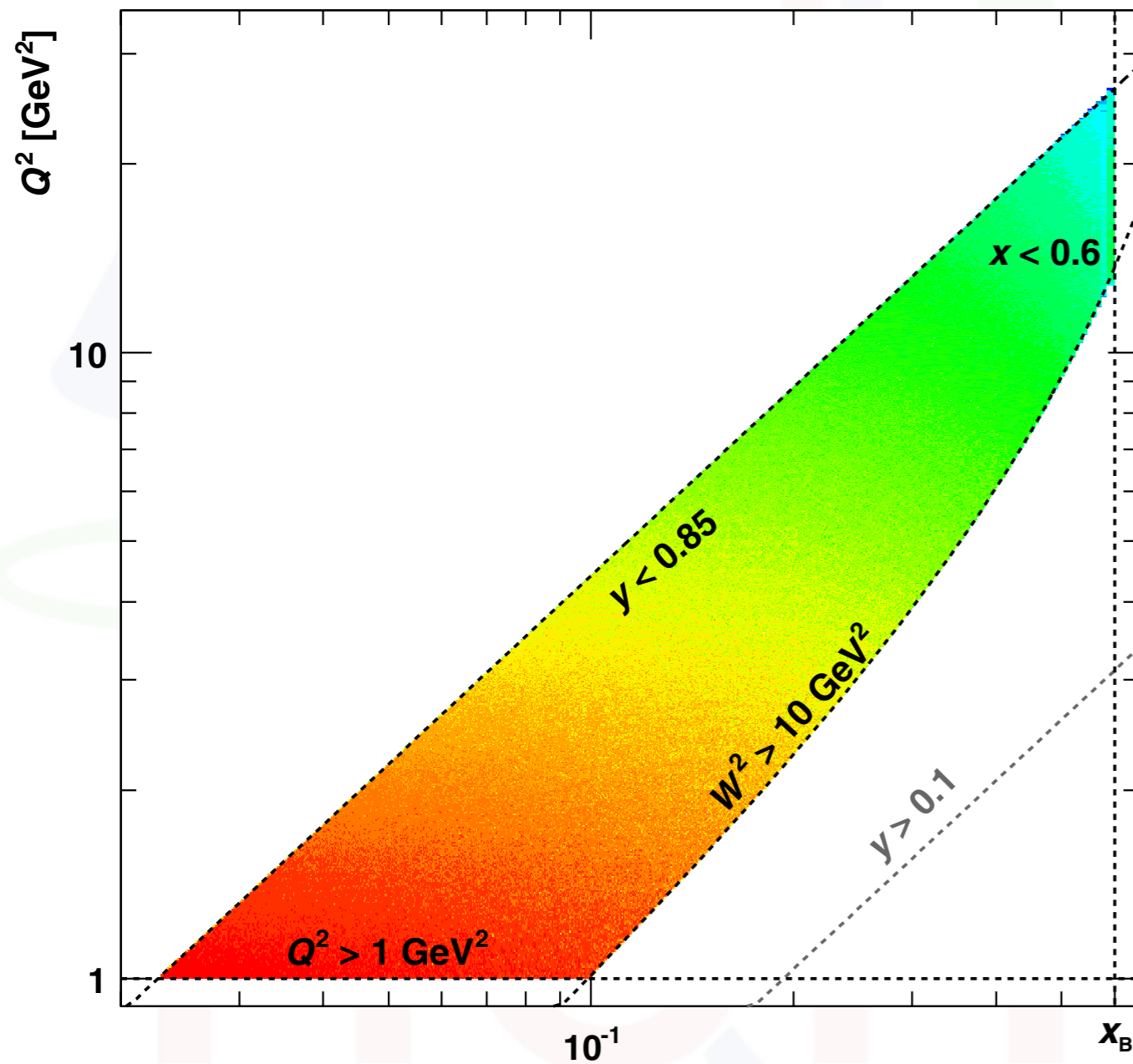
[S.J. Joosten, PhD thesis UIUC (2013)]



Neglecting to unfold in z changes x dependence dramatically

➔ 1D unfolding clearly insufficient

Kinematic range at HERMES



- $0.023 < x < 0.6$
- $0.1 < y < 0.85$
- $0.2 < z < 0.8$
- $W^2 > 10 \text{ GeV}^2$
- $Q^2 > 1 \text{ GeV}^2$

Results I:

charged pions and kaons from proton and deuteron targets

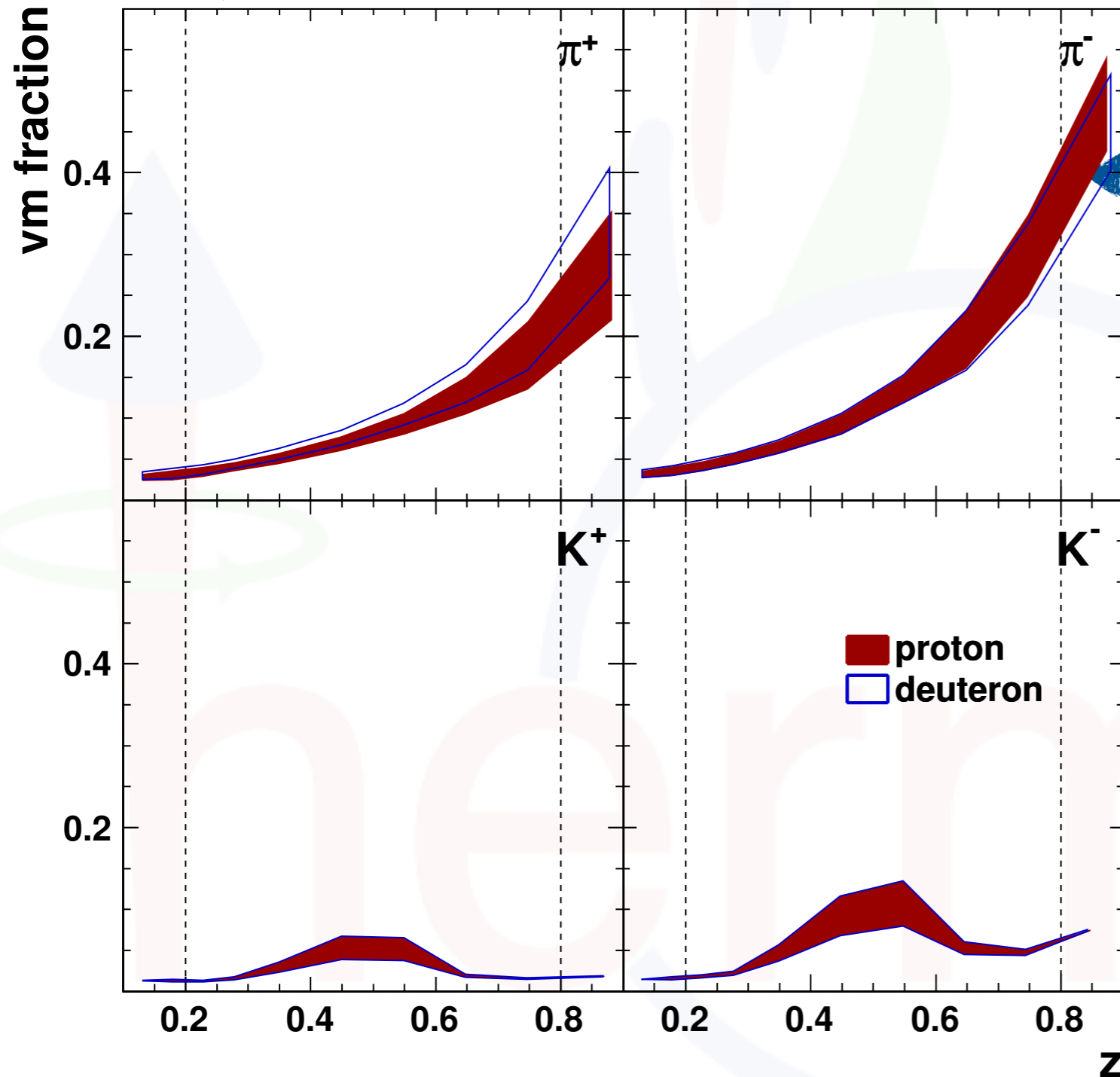
A. Airapetian et al., Phys. Rev. D87 (2013) 074029

<http://www-hermes.desy.de/multiplicities>

Influence from exclusive VM

for instance: $ep \rightarrow ep \rho^0 \rightarrow ep \pi^+ \pi^-$

[Airapetian et al., PRD 87 (2013) 074029]

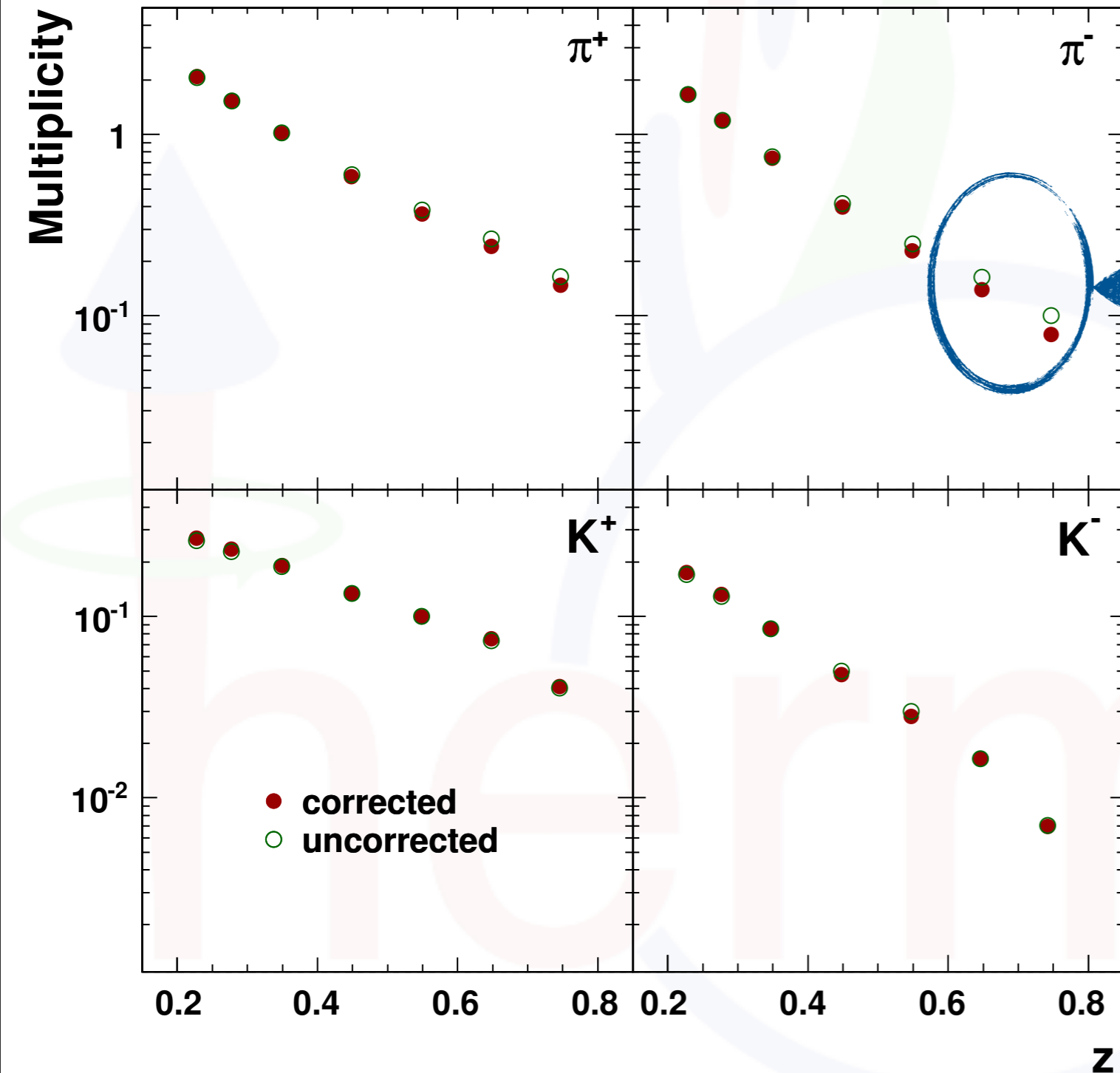


partially large contribution from exclusive VM production, in particular at high z

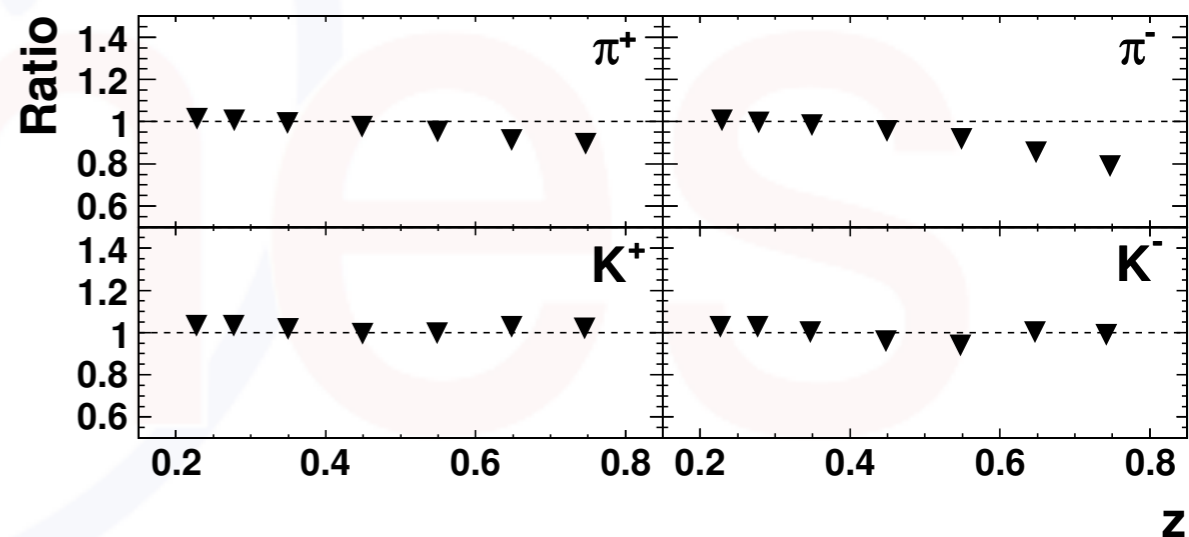
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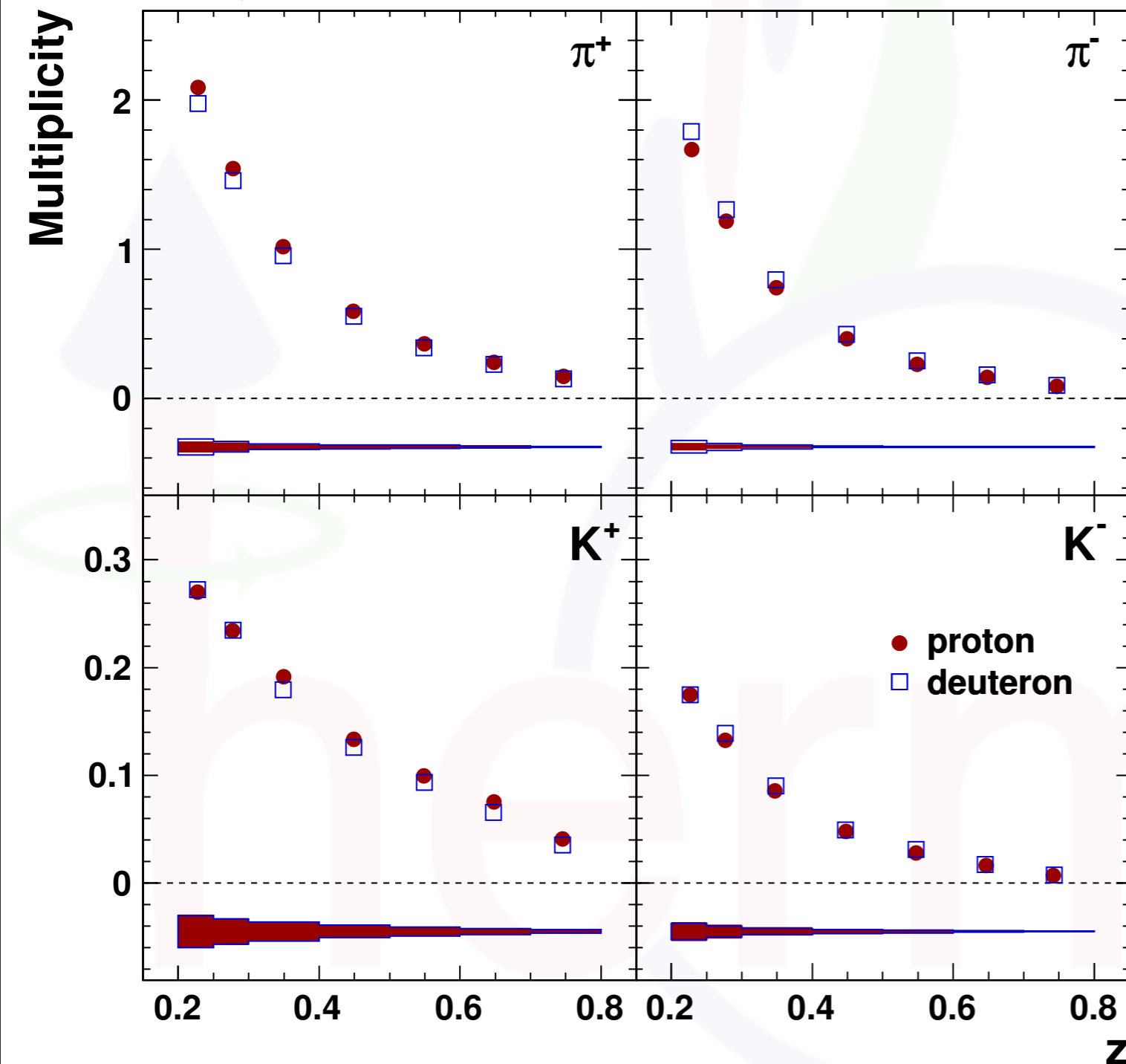
multiplicities before
and after subtraction
of contributions from
exclusively produced
VMs



Multiplicities: z projection

most exhaustive data set on ($P_{h\perp}$ -integrated) electro-production of charged identified mesons on nucleons

[Airapetian et al., PRD 87 (2013) 074029]



➔ slight differences between proton and deuteron targets: reflection of valence structure of target and produced meson, e.g.

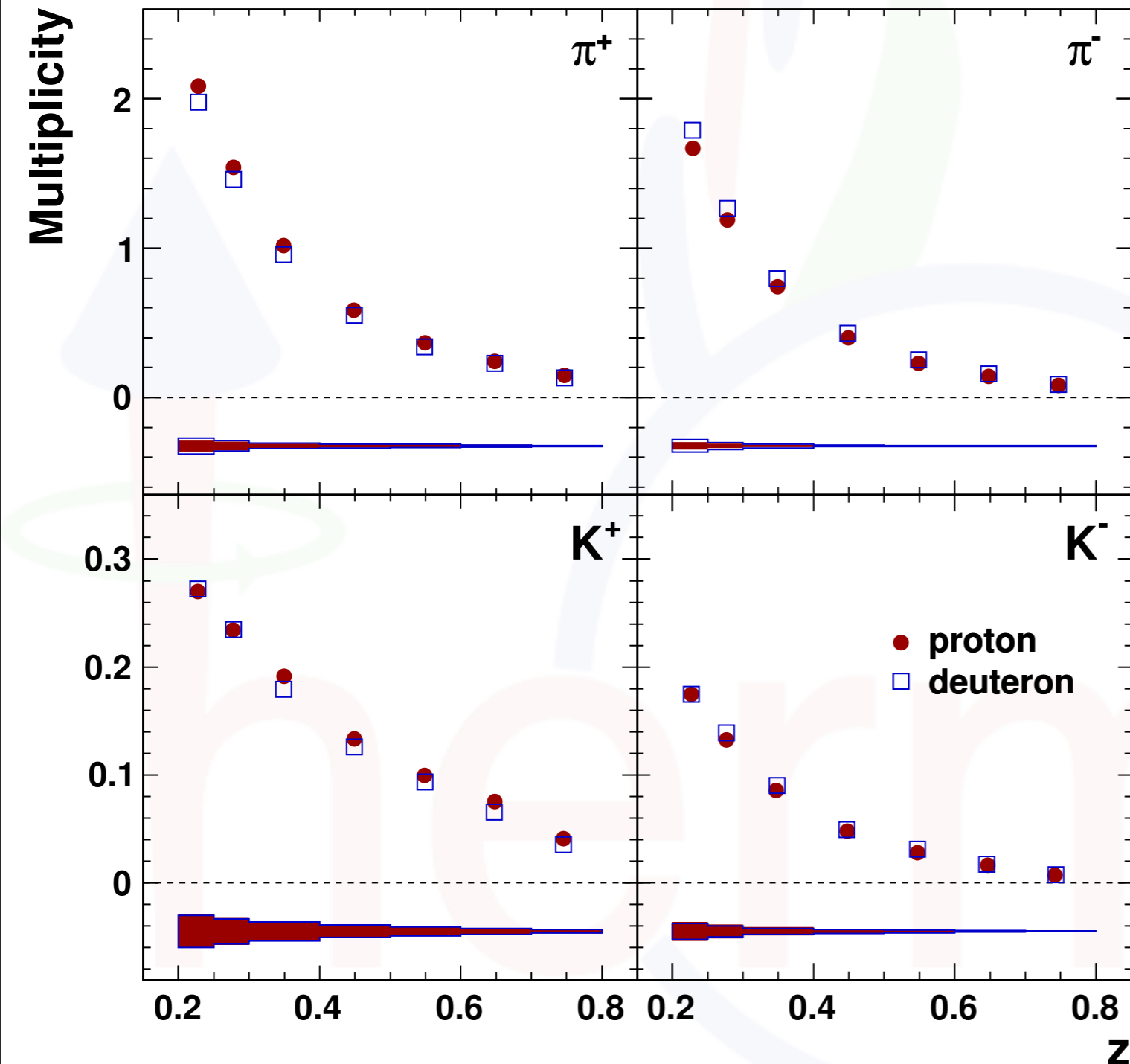
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➔ K^- pure "sea object" hence suppressed and hardly any difference for proton and deuteron

Multiplicities: z projection

proton target:
(deuteron similar)

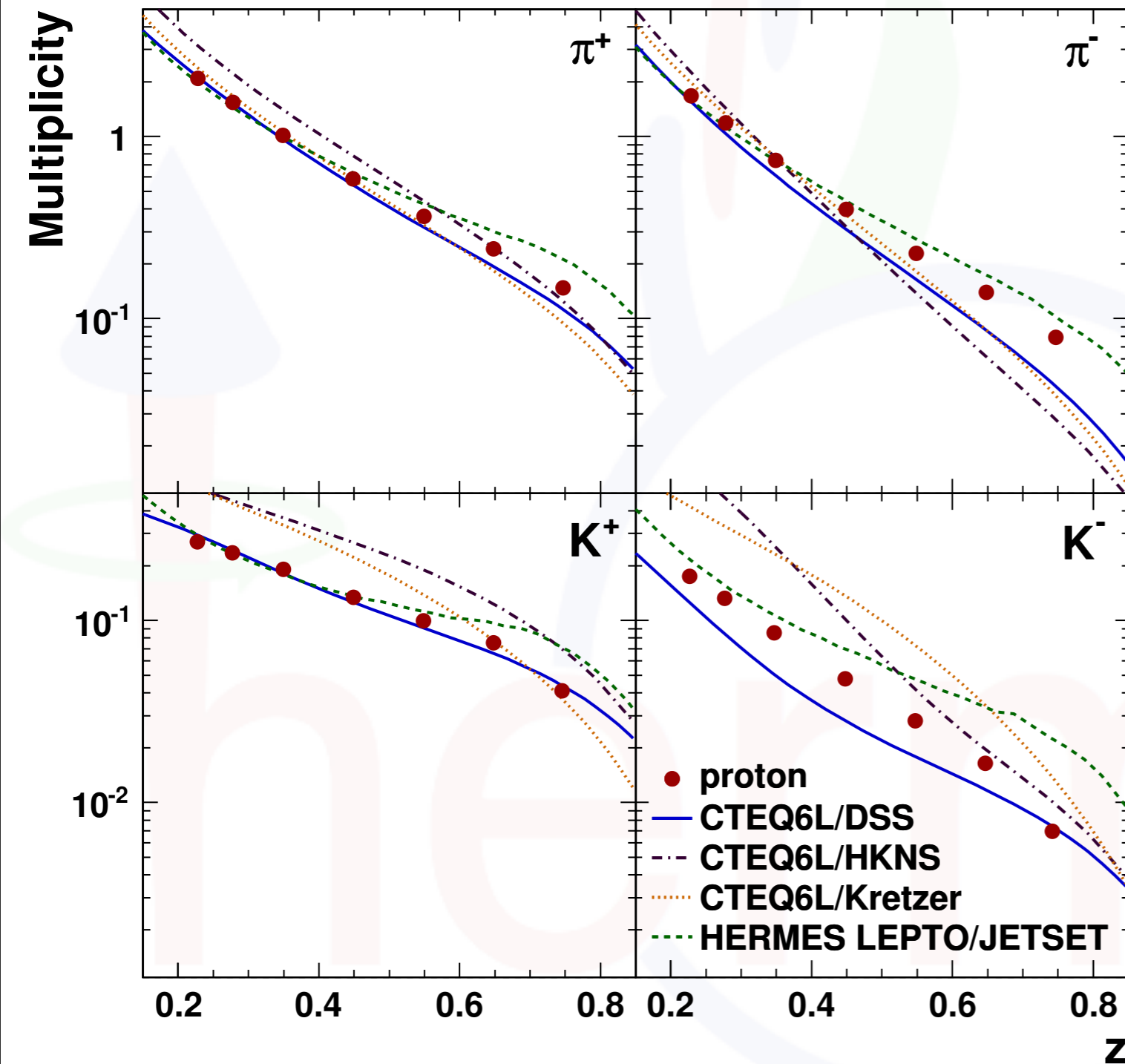
positive hadrons in general
better described than
negative ones

➔ better understanding of
favored fragmentation?

➔ best described by
HERMES Jetset tune and
DSS FF set

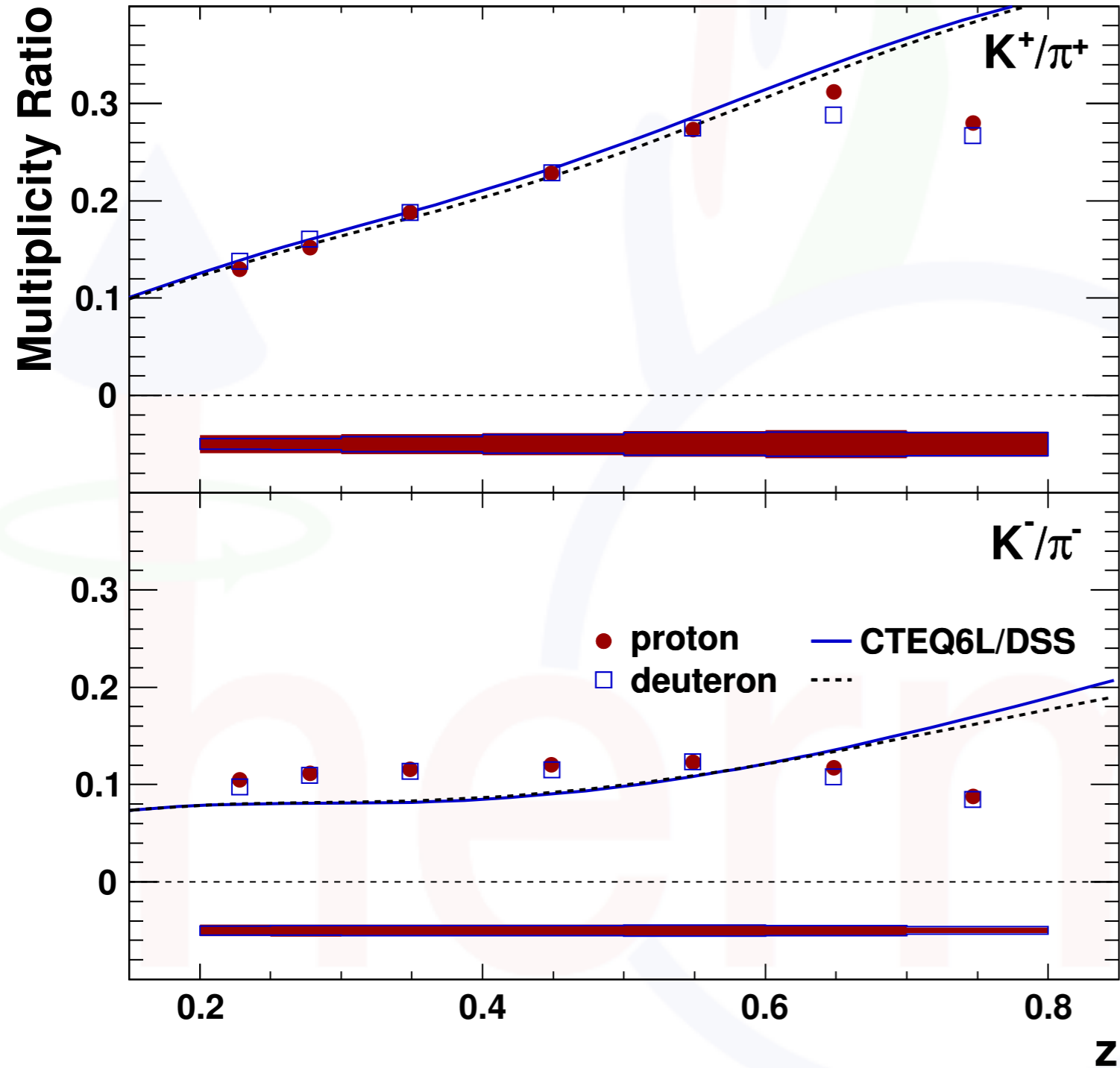
kaons best described by
DSS FF set, though all with
problems in describing K^-

[Airapetian et al., PRD 87 (2013) 074029]



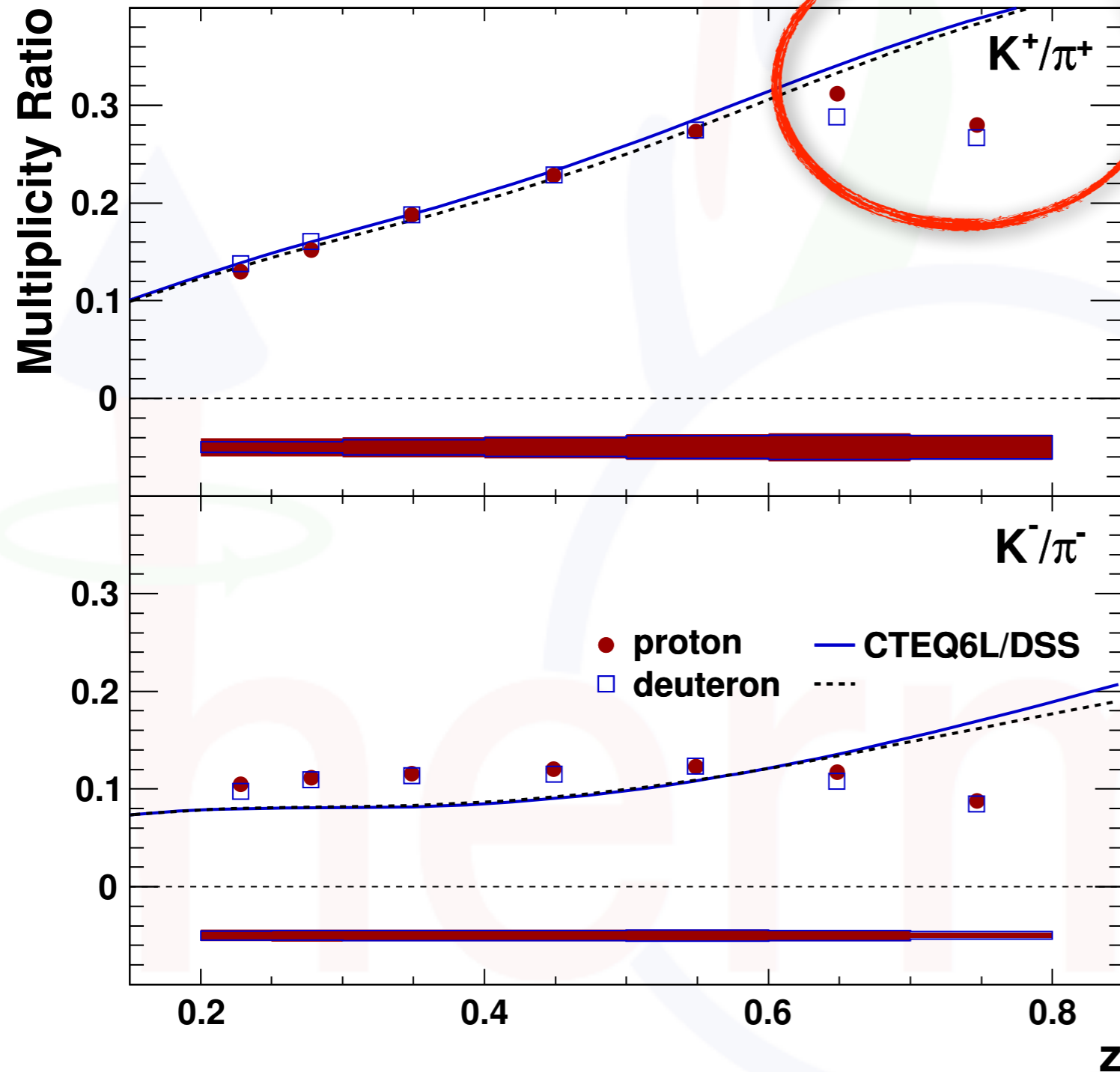
Multiplicity ratio: z projection

[<http://www-hermes.desy.de/multiplicities>]



Multiplicity ratio: z projection

[<http://www-hermes.desy.de/multiplicities>]



at large z mainly favored fragmentation:

➔ dominated by up quarks

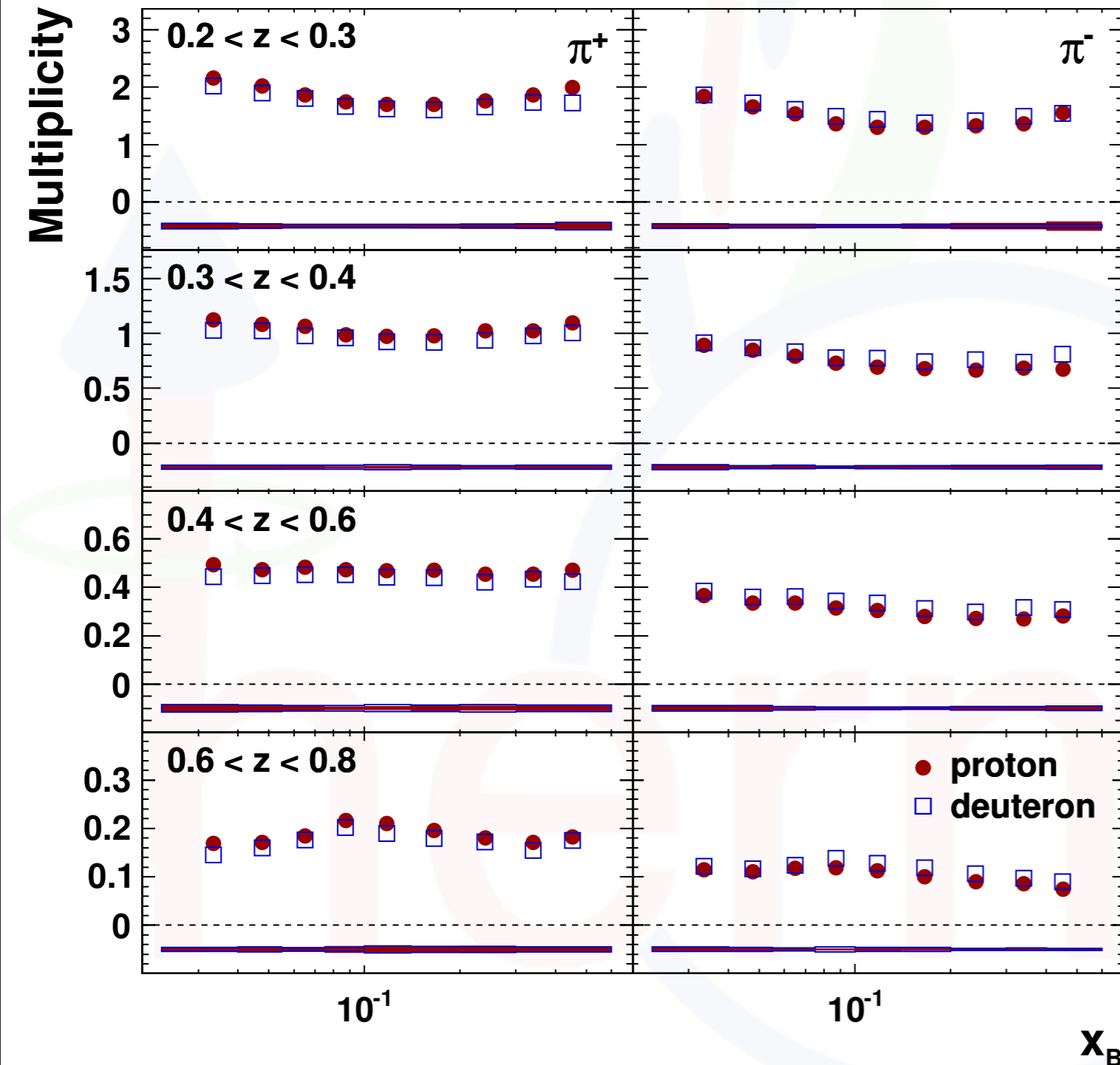
➔ kaon requires strangeness production

➔ strangeness suppression of about 0.3 (apparently stronger than modeled in DSS FF set)

➔ in rough agreement with typical ansatz of 1/3

Multiplicities: x-z projection

[Airapetian et al., PRD 87 (2013) 074029]

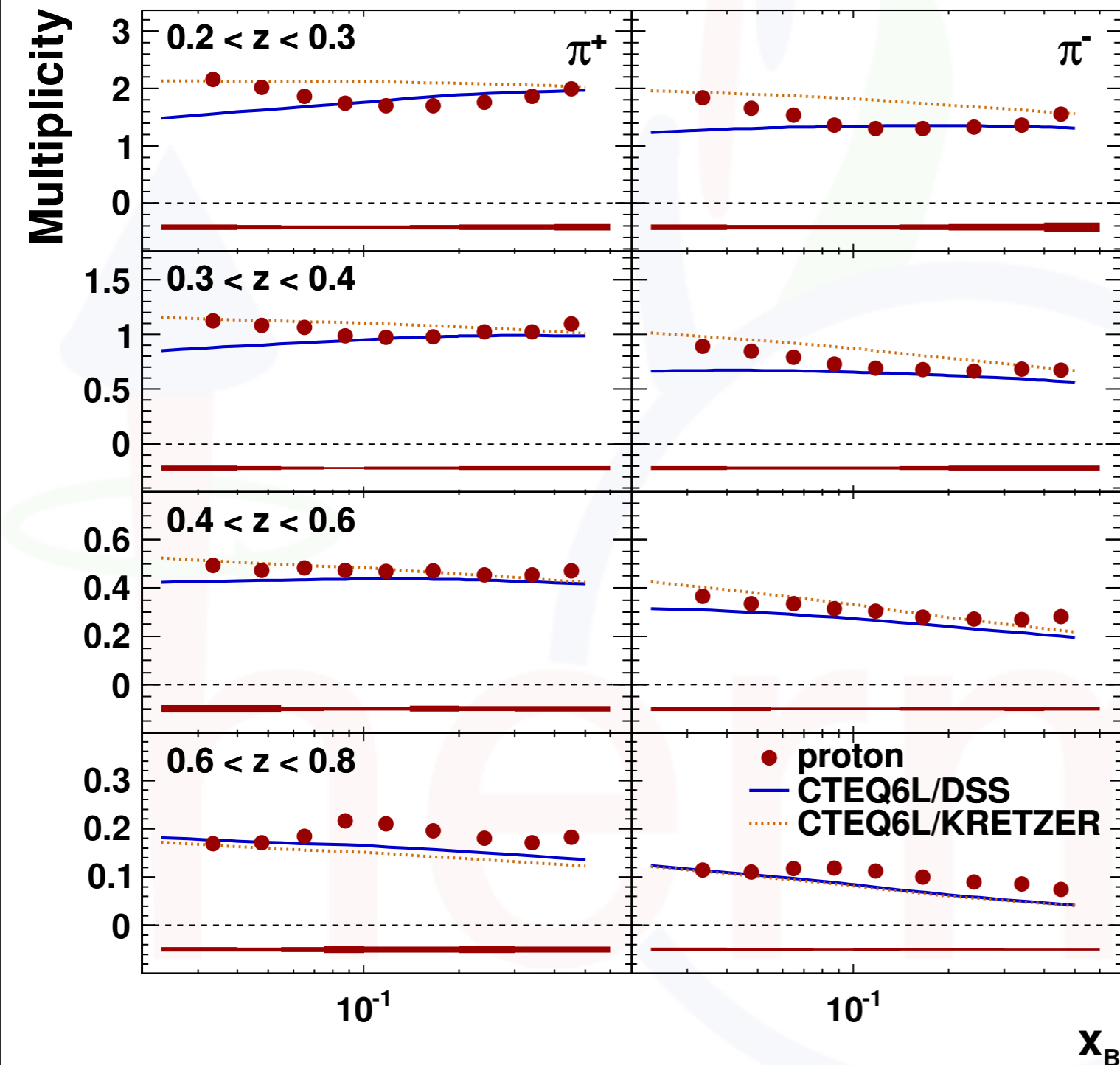


➔ weaker dependence on x

$$\sum_q \frac{e_q^2 f_1^q(x)}{\sum_{q'} e_{q'}^2 f_1^{q'}(x)} D_1^{q \rightarrow \pi}(z)$$

Multiplicities: x-z projection

[Airapetian et al., PRD 87 (2013) 074029]



➔ weaker dependence on x

➔ remaining dependence from $f_1 - D_1$ convolution over quark flavors

$$\sum_q \frac{e_q^2 f_1^q(x)}{\sum_{q'} e_{q'}^2 f_1^{q'}(x)} D_1^{q \rightarrow \pi}(z)$$

Strange-quark distribution

- use isoscalar probe and target to extract strange-quark distribution
- only need K^+K^- multiplicities on deuteron

$$S(x) \int \mathcal{D}_S^K(z) dz \simeq Q(x) \left[5 \frac{d^2 N^K(x)}{d^2 N^{\text{DIS}}(x)} - \int \mathcal{D}_Q^K(z) dz \right]$$

$$S(x) = s(x) + \bar{s}(x)$$

$$Q(x) = u(x) + \bar{u}(x) + d(x) + \bar{d}(x)$$

$$\mathcal{D}_S^K = D_1^{s \rightarrow K^+} + D_1^{\bar{s} \rightarrow K^+} + D_1^{s \rightarrow K^-} + D_1^{\bar{s} \rightarrow K^-}$$

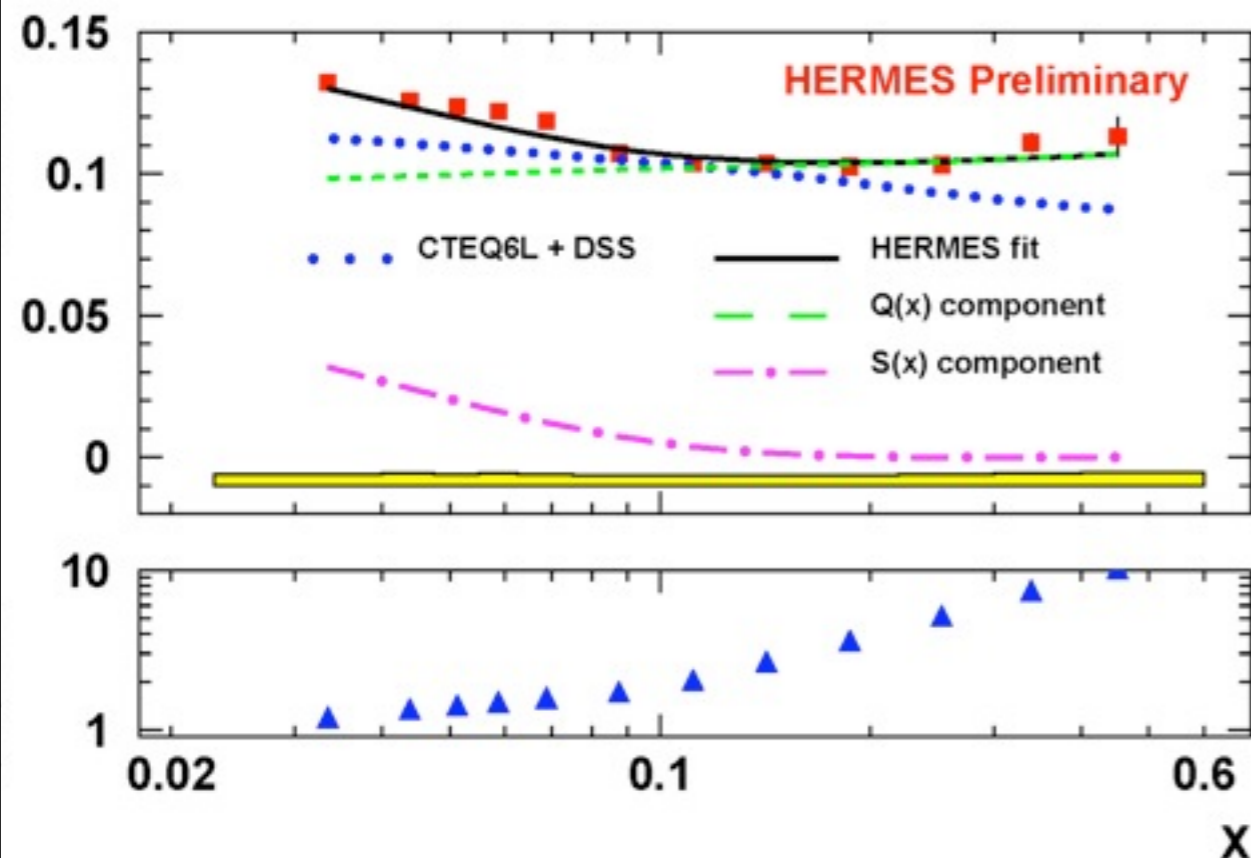
$$\mathcal{D}_Q^K = 4D_1^{u \rightarrow K^+} + 4D_1^{\bar{u} \rightarrow K^+} + D_1^{d \rightarrow K^+} + D_1^{\bar{d} \rightarrow K^+} + \dots$$

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- assume vanishing strangeness at high x to extract non-strange fragmentation



$$S(x) = s(x) + \bar{s}(x)$$

$$Q(x) = u(x) + \bar{u}(x) + d(x) + \bar{d}(x)$$

$$\mathcal{D}_S^K = \mathcal{D}_1^{s \rightarrow K^+} + \mathcal{D}_1^{\bar{s} \rightarrow K^+} + \mathcal{D}_1^{s \rightarrow K^-} + \mathcal{D}_1^{\bar{s} \rightarrow K^-}$$

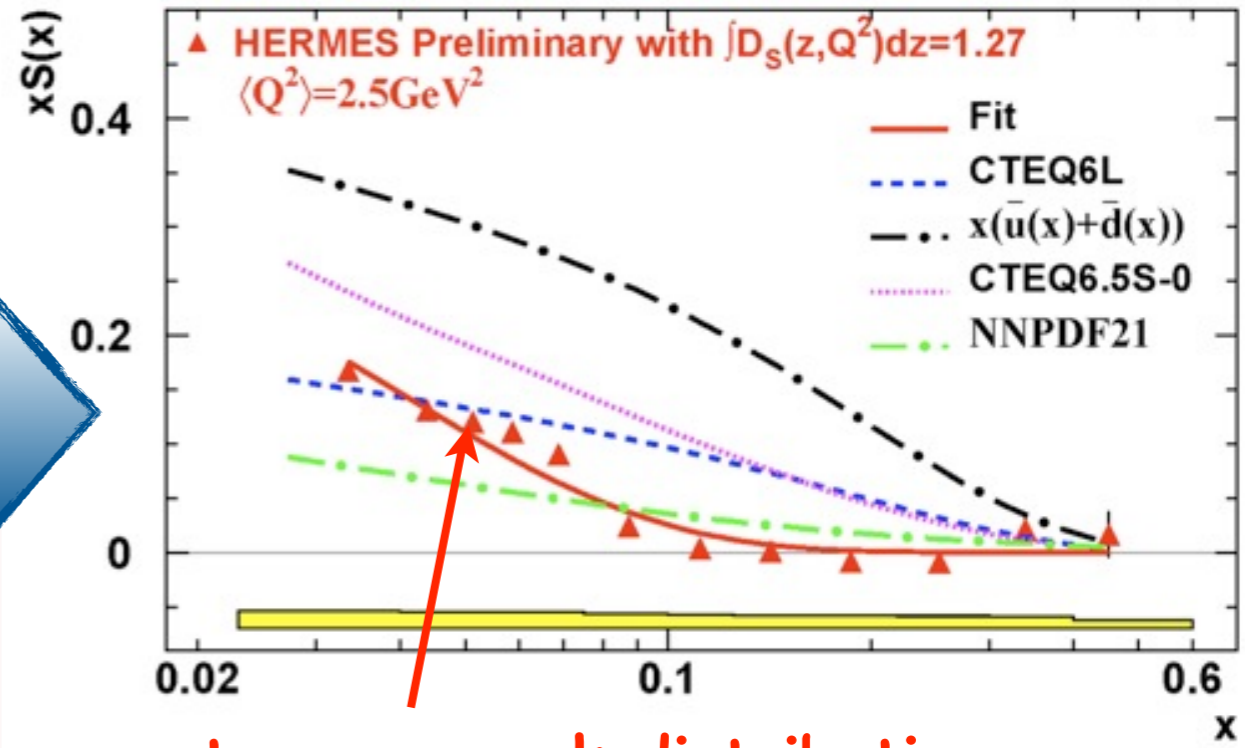
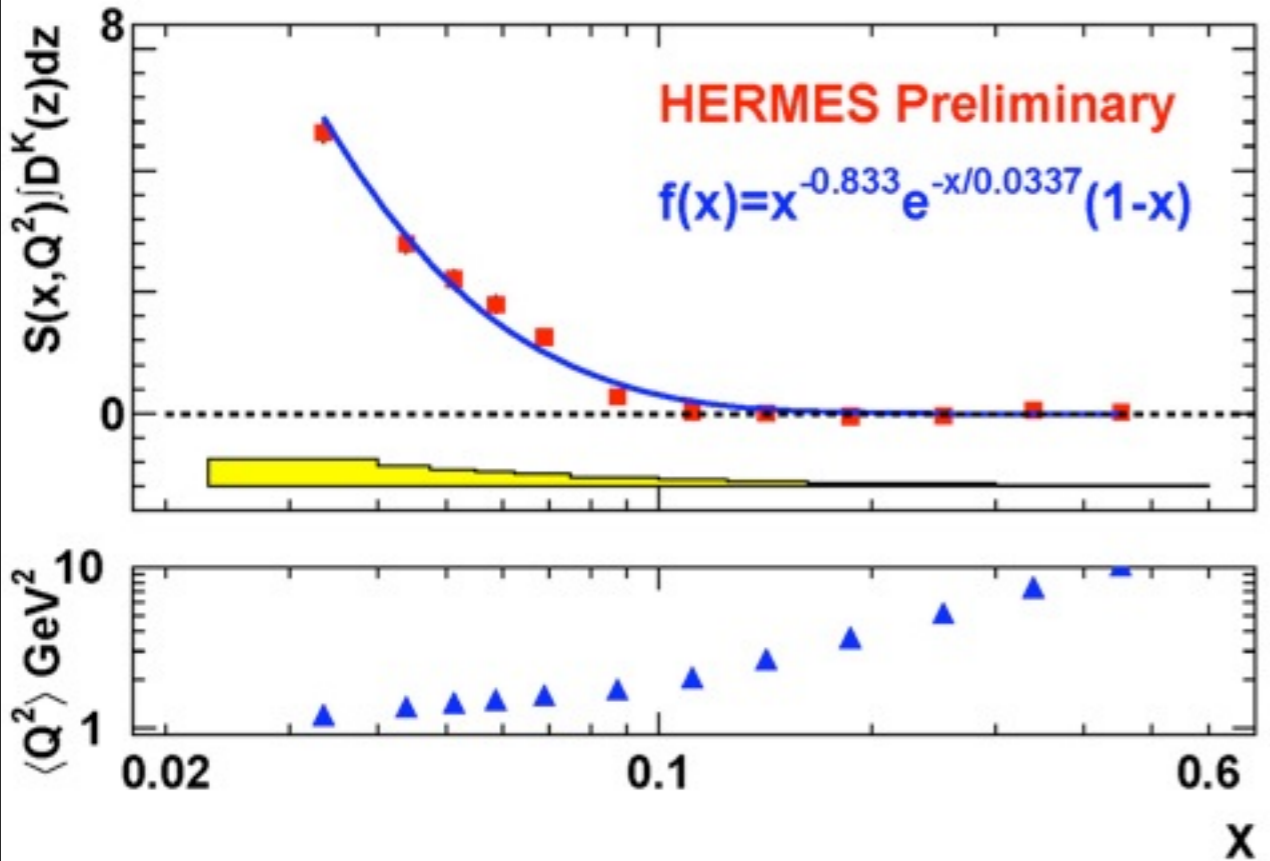
$$\mathcal{D}_Q^K = 4\mathcal{D}_1^{u \rightarrow K^+} + 4\mathcal{D}_1^{\bar{u} \rightarrow K^+} + \mathcal{D}_1^{d \rightarrow K^+} + \mathcal{D}_1^{\bar{d} \rightarrow K^+} + \dots$$

Strange-quark distribution

- use isoscalar probe and target to extract strange-quark distribution
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- assume vanishing strangeness at high x to extract non-strange fragmentation

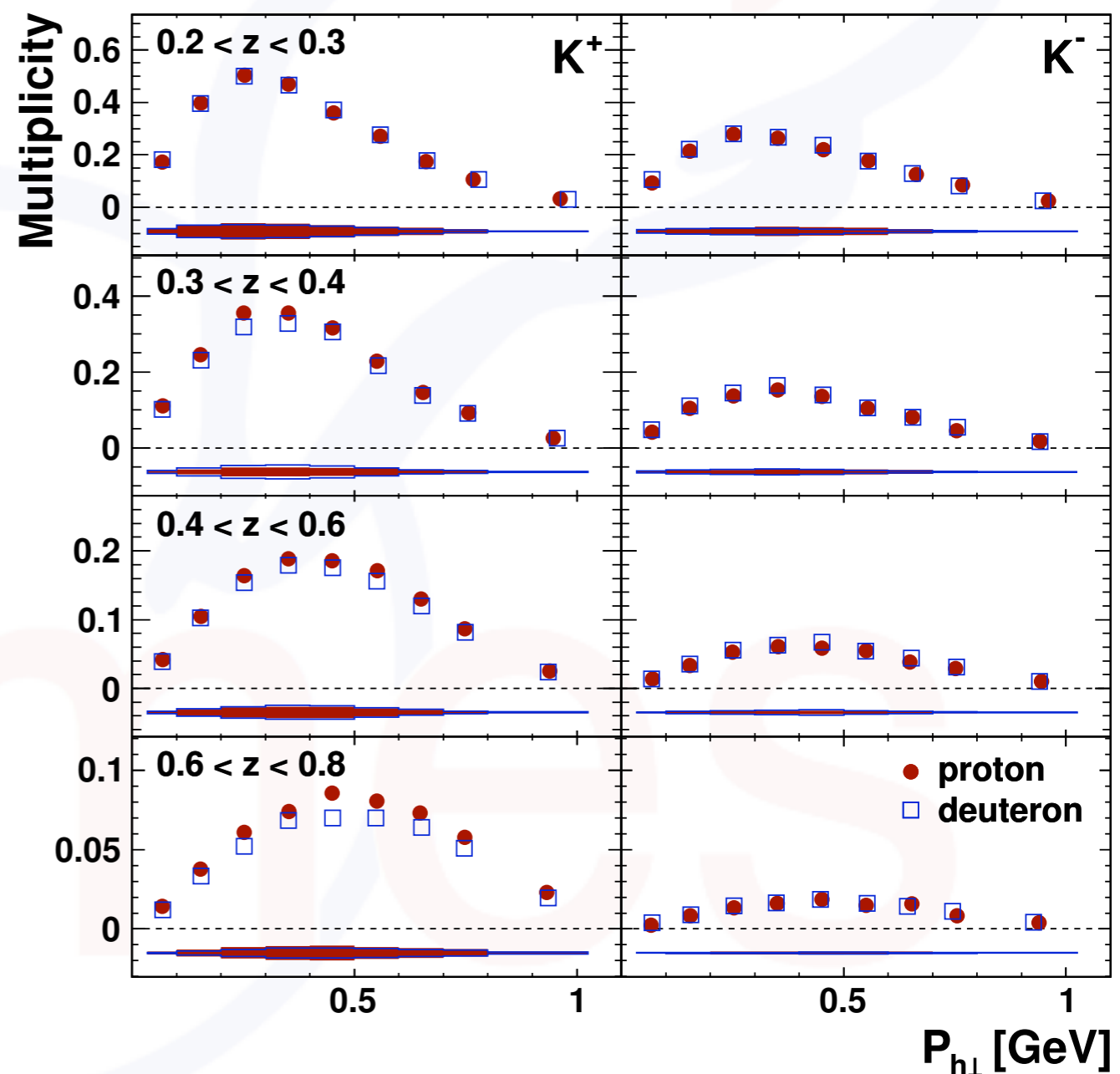
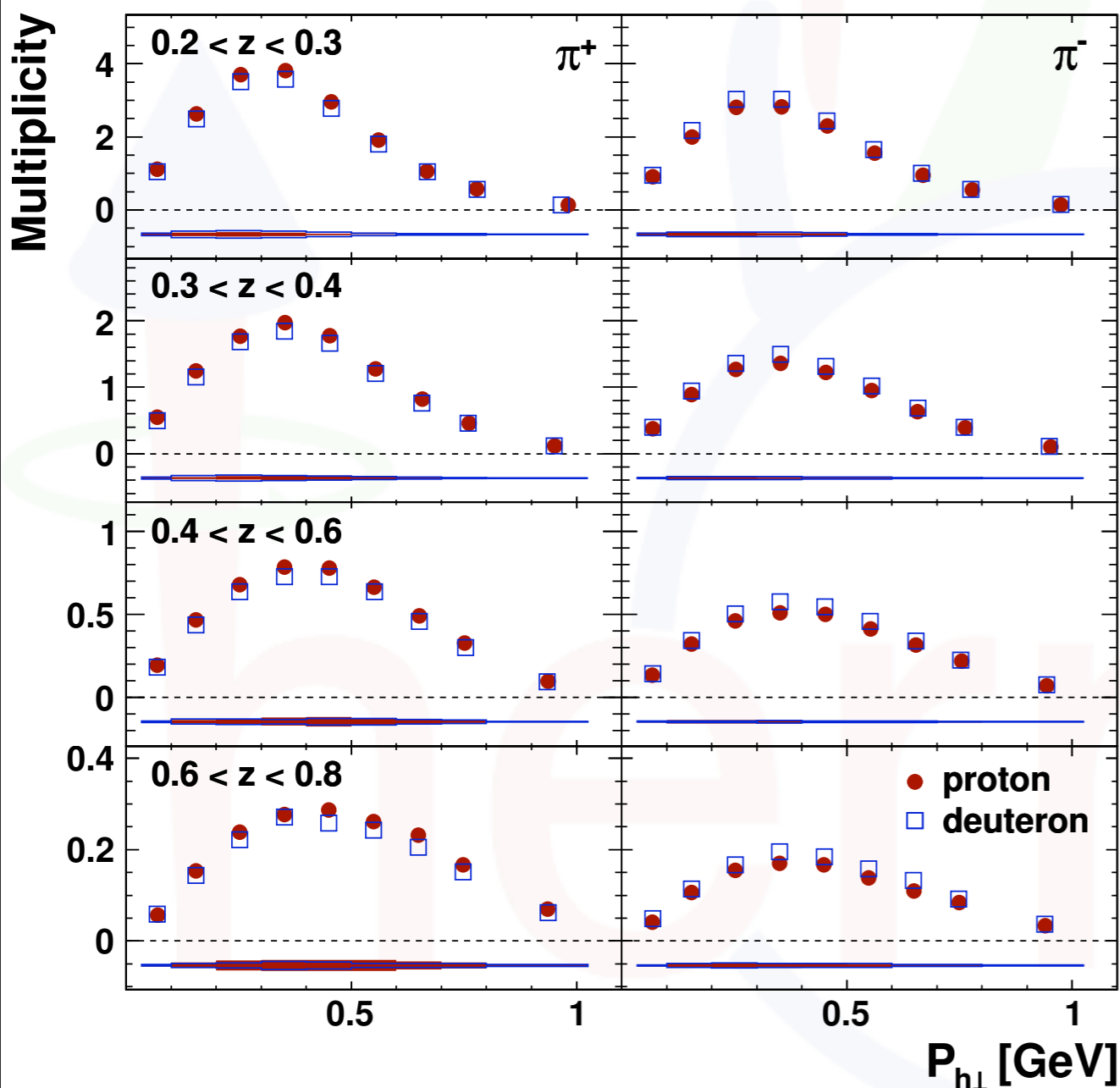


strange-quark distribution
 softer than (maybe) expected

Transverse momentum dependence

- multi-dimensional analysis allows going beyond collinear factorization
- flavor information on transverse momenta via target variation and hadron ID

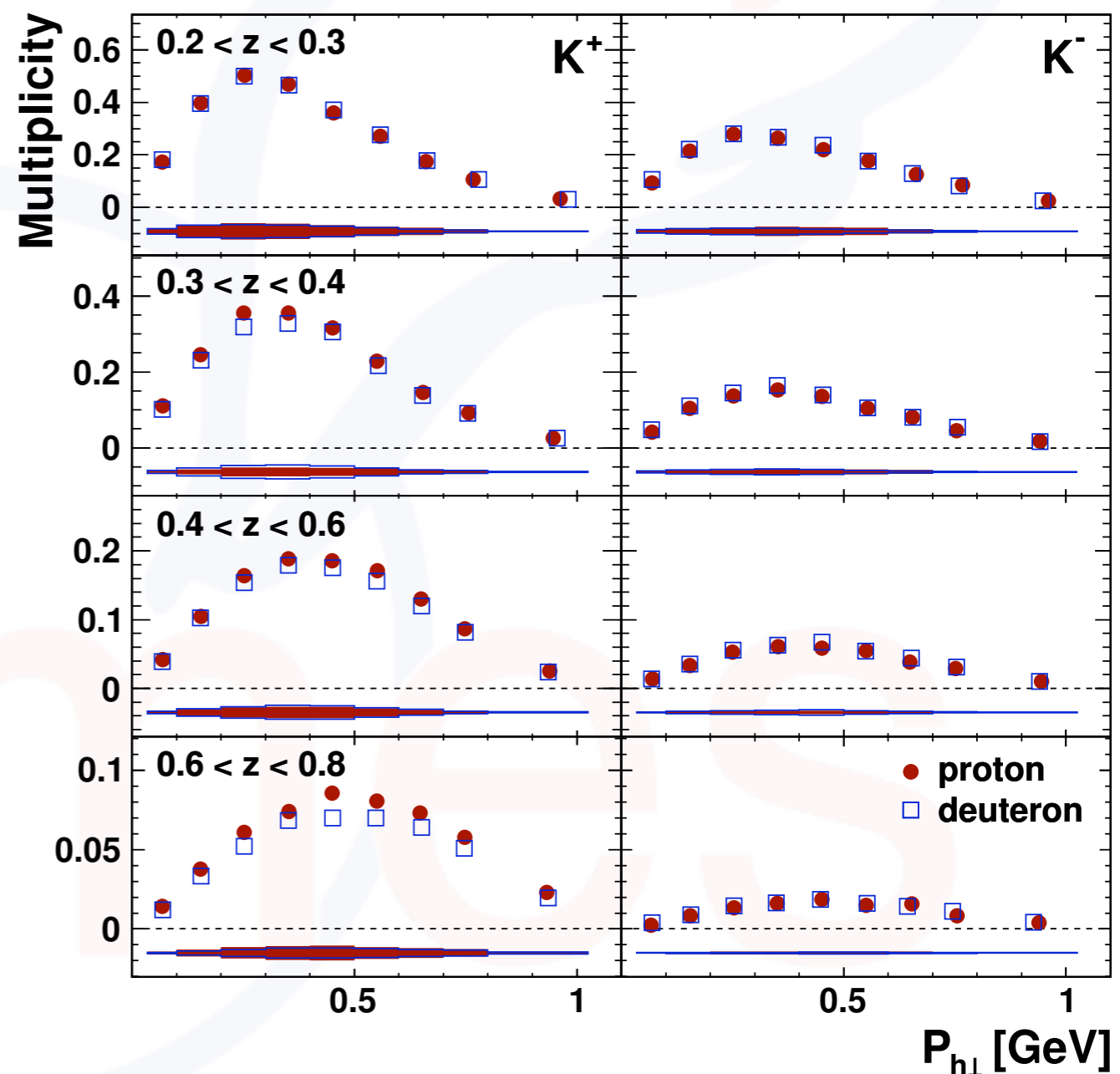
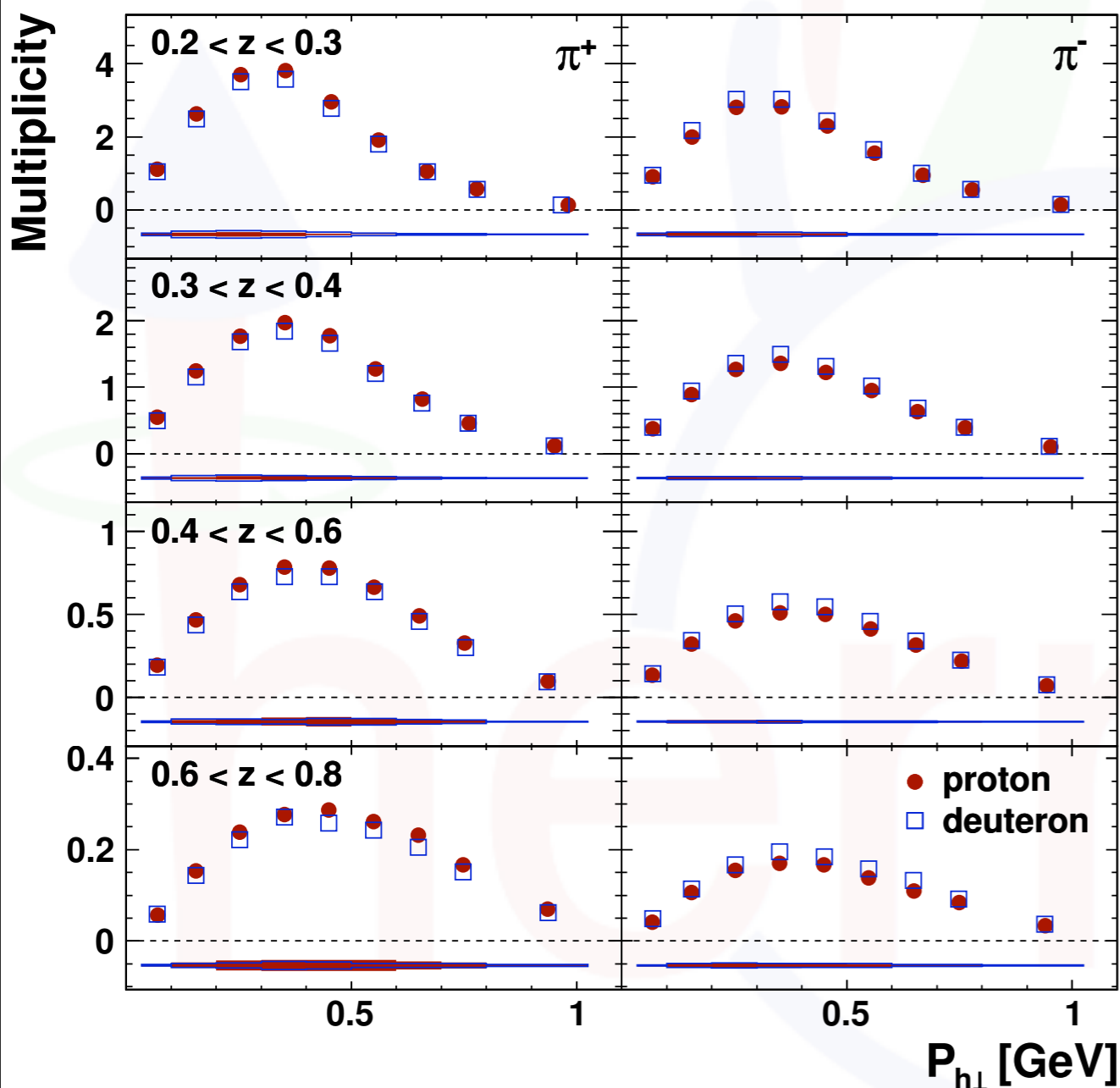
[Airapetian et al., PRD 87 (2013) 074029]



Transverse momentum dependence

- multi-dimensional analysis allows going beyond collinear factorization
 - flavor information on transverse momenta via target variation and hadron ID
- ➔ A. Signori (Friday)

[Airapetian et al., PRD 87 (2013) 074029]



Results II:

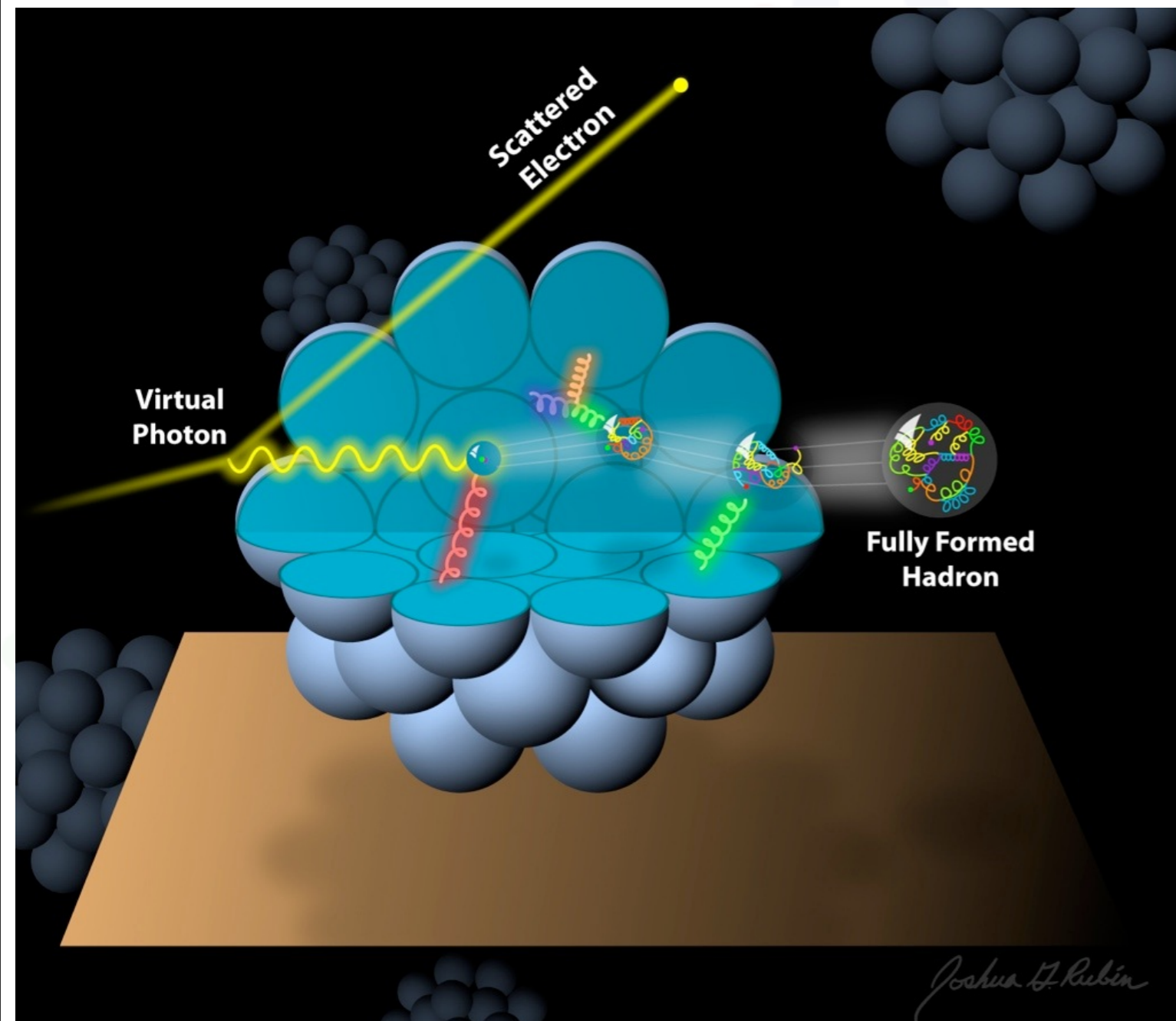
multiplicity ratios - nuclear attenuation

A. Airapetian et al., Nucl. Phys. B 780 (2007) 1-27

A. Airapetian et al., EPJ A 47 (2011) 113

<http://inspirebeta.net/record/918944/>

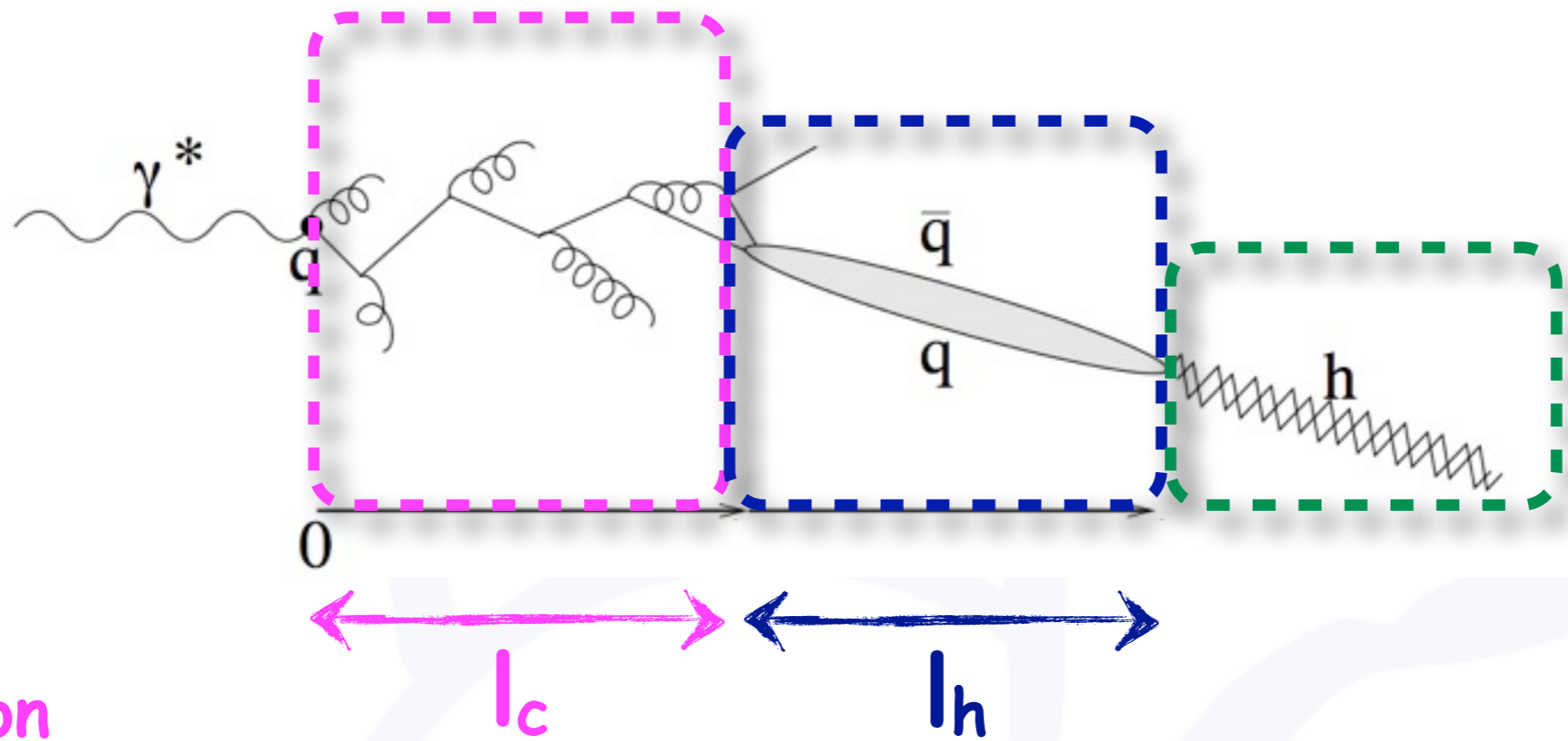
Nuclei: a hadronization laboratory



[J. Rubin]

- partons in nuclear medium:
 - PDFs modified (e.g, EMC effect)
 - gluon radiation and rescattering effects
- (pre)hadron in nuclear medium:
 - rescattering
 - absorption

Nuclei: a hadronization laboratory



- parton

- pre-hadron

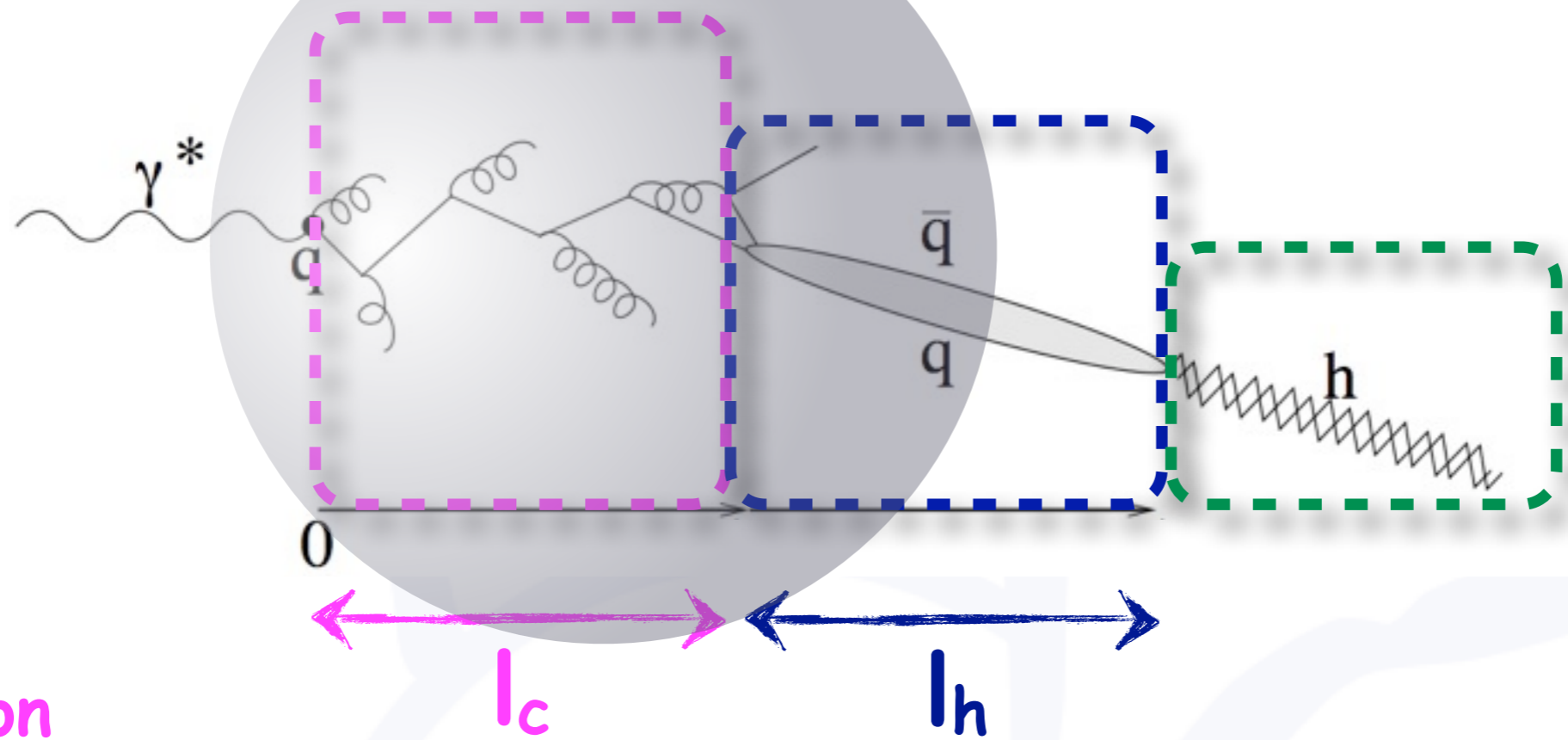
- colorless

- quantum numbers of final hadron

- final state hadron

- differences predicted for partonic and (pre-)hadronic interactions

Nuclei: a hadronization laboratory



- parton

- pre-hadron

- colorless

- quantum numbers of final hadron

- final state hadron

- differences predicted for partonic and (pre-)hadronic interactions

- depends on formation lengths (1-10fm) = O(nucleus size)

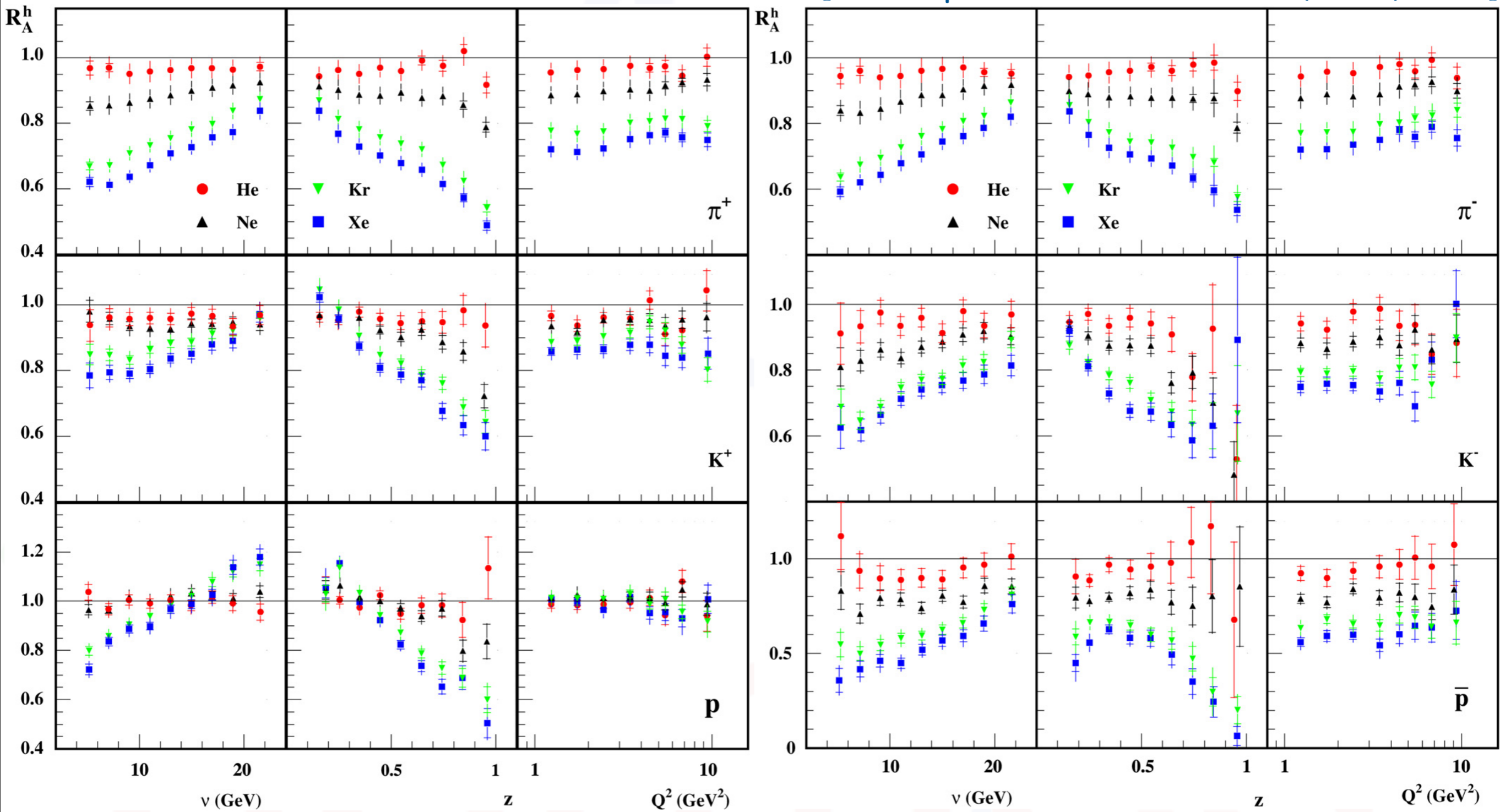
Multiplicity ratios

$$R_A^h(\nu, Q^2, z, p_t^2) = \frac{\left(\frac{N^h(\nu, Q^2, z, p_t^2)}{N^e(\nu, Q^2)} \right)_A}{\left(\frac{N^h(\nu, Q^2, z, p_t)}{N^e(\nu, Q^2)} \right)_D}$$

- nuclear targets: (He,) Ne, Kr, Xe compared to D
- ratio \implies approximate cancellation of:
 - QED radiative effects (RADGEN)
 - limited geometric and kinematic acceptance of spectrometer
 - detector resolution
- multi-dimensional extraction

Nuclear attenuation

[A. Airapetian et al., NPB 780 (2007) 1-27]

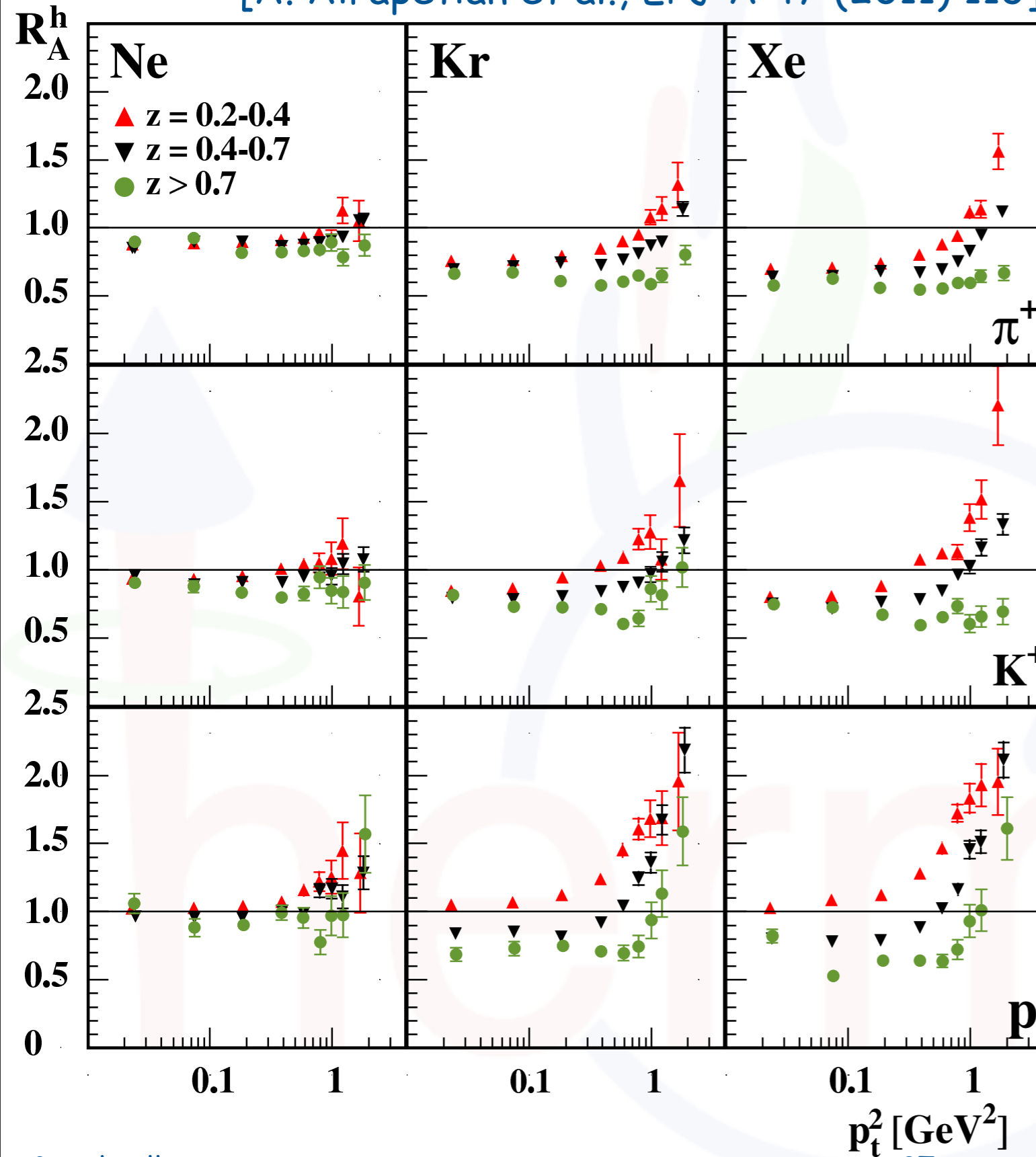


- strong mass dependence: attenuation mainly increases with A
- quite different behavior for protons

Nuclear attenuation

[A. Airapetian et al., EPJ A 47 (2011) 113]

$$R_A^h \equiv \frac{\mathcal{M}_A^h}{\mathcal{M}_d^h}$$

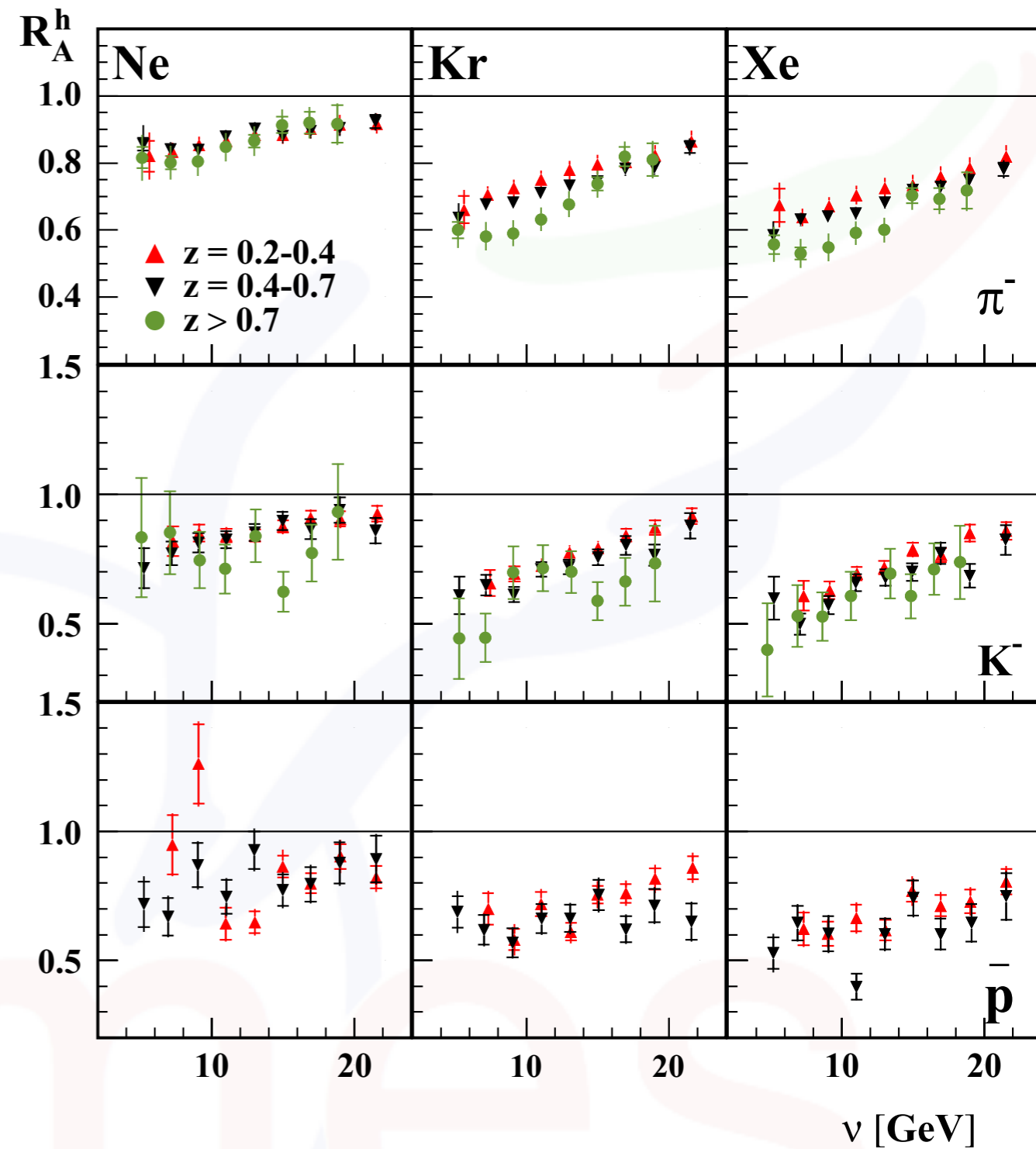
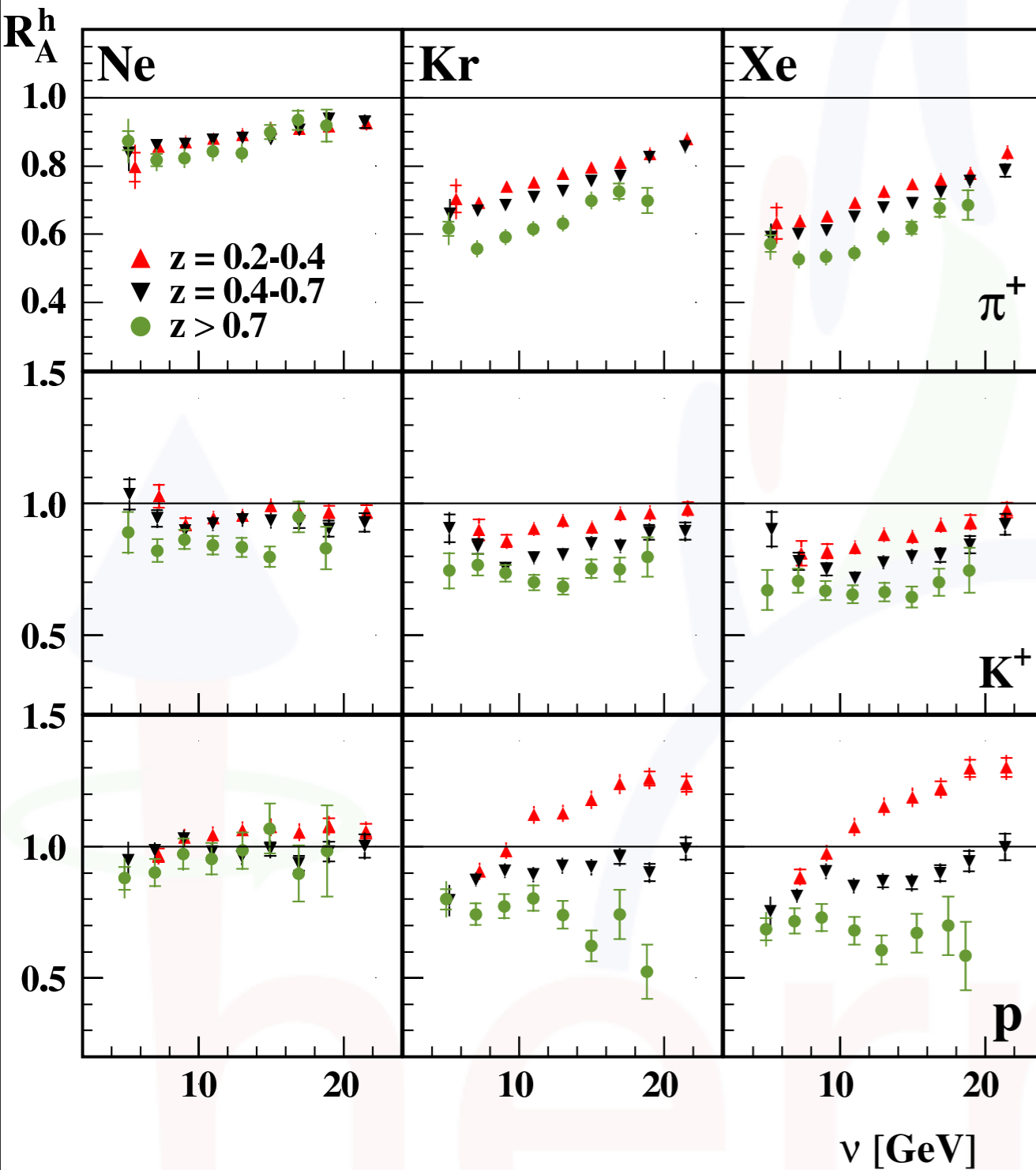


strong p_T dependence of nuclear attenuation (e.g., Cronin effect - enhancement at large p_T)

except maybe at large z for pions and kaons (little energy loss dictates few interactions)

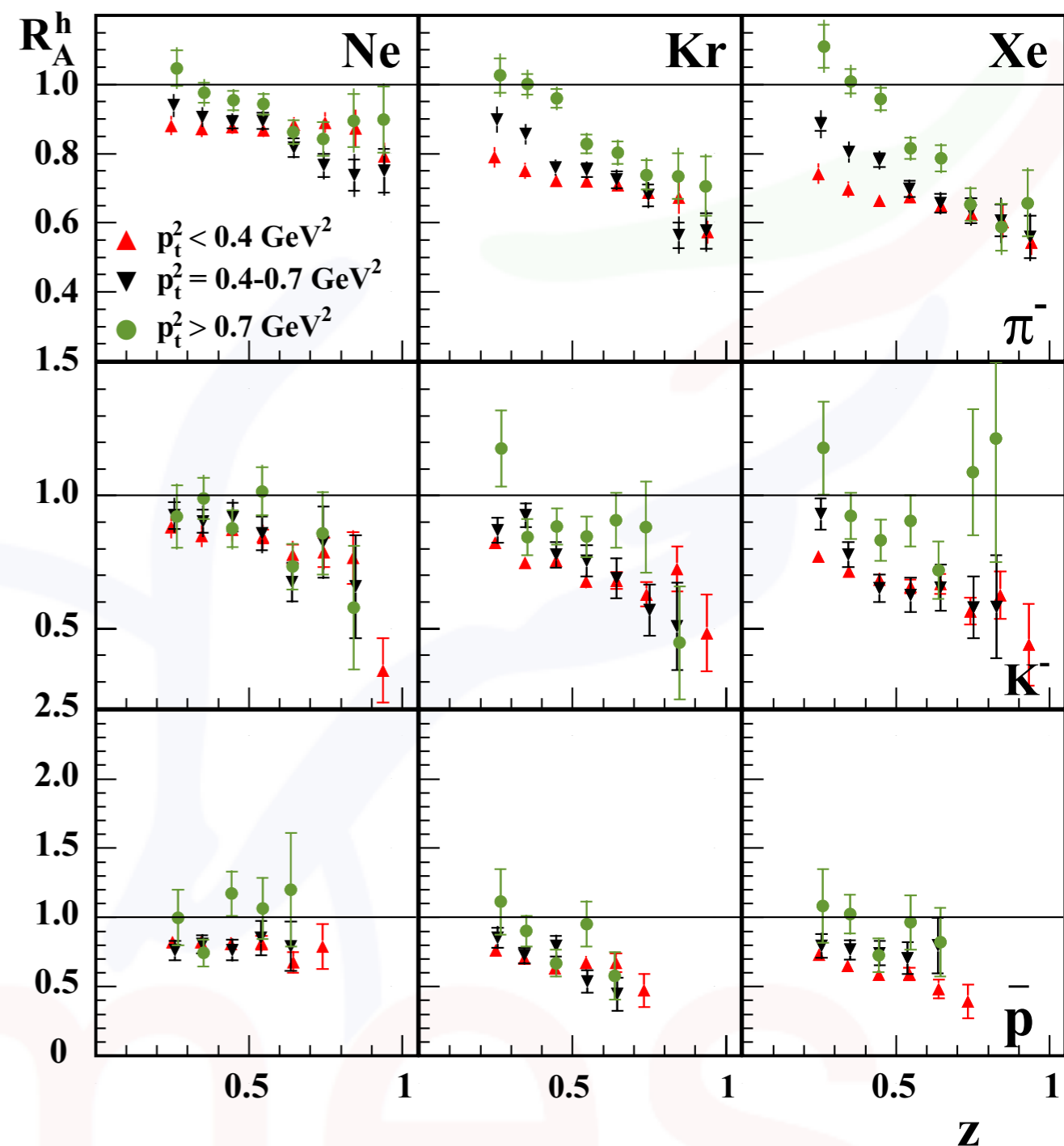
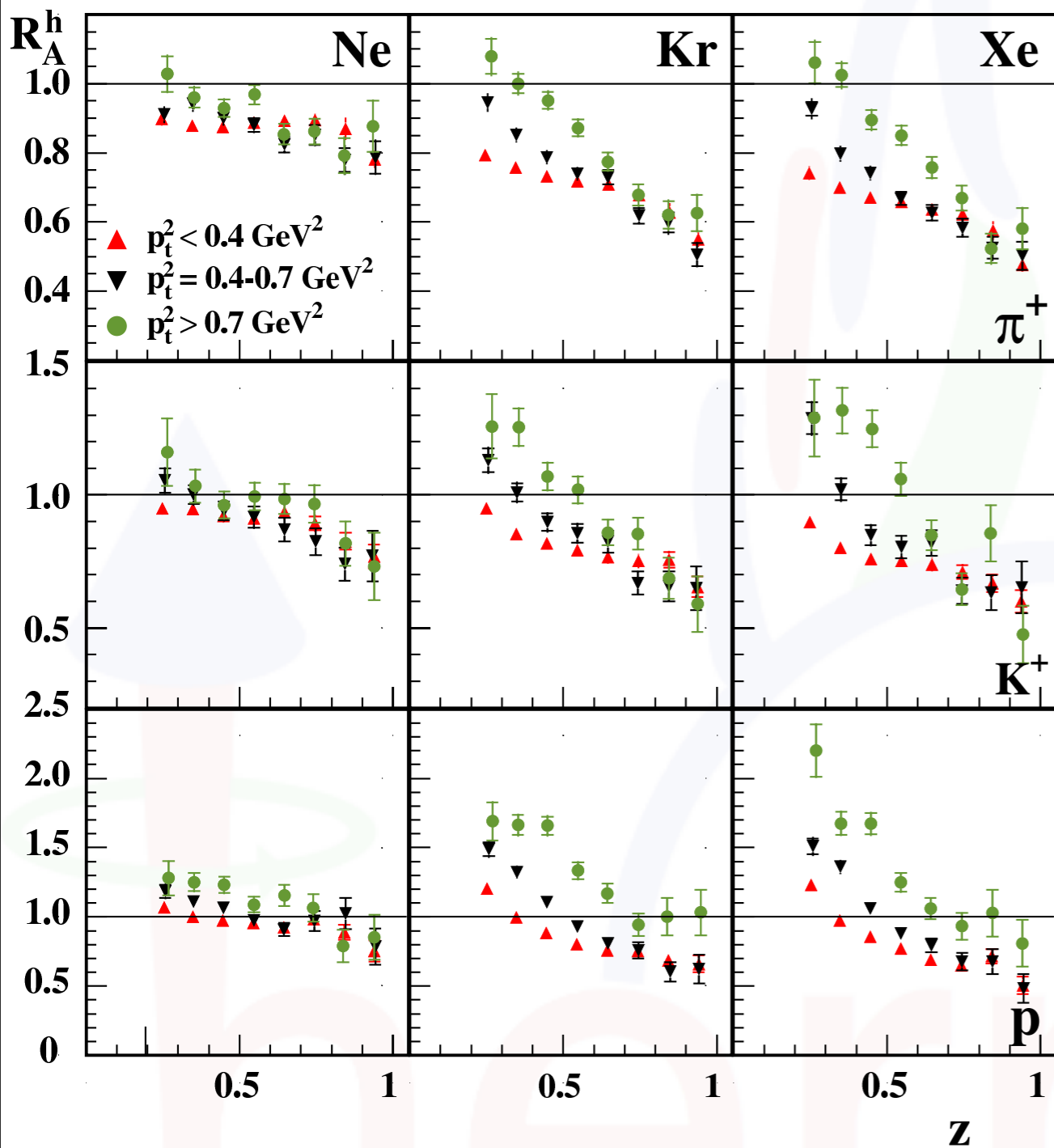
larger effect for protons

Nuclear attenuation



- mostly decrease of attenuation with increasing ν
- enhancement of proton multiplicities at low z and high ν

Nuclear attenuation



- strong z dependence of attenuation
- amplified by transverse momentum and target mass (i.e., size)

Conclusions

- HERMES managed step from spin-asymmetry experiment to unpolarized-target experiment
- largest data set on charged-separated identified meson lepto-production
- multi-dimensional analysis and various targets allow study of correlations and flavor dependences
- large attenuation effects at HERMES energies, mainly increasing with nucleus size (except protons) with correlated kinematic dependences
- nuclear environment can play significant role in TMD effects
- don't forget longitudinal photons