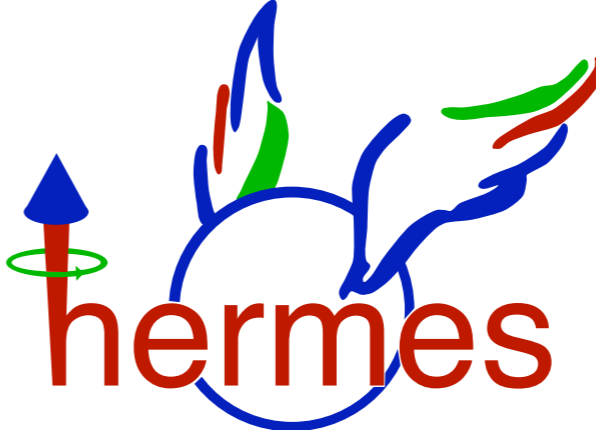
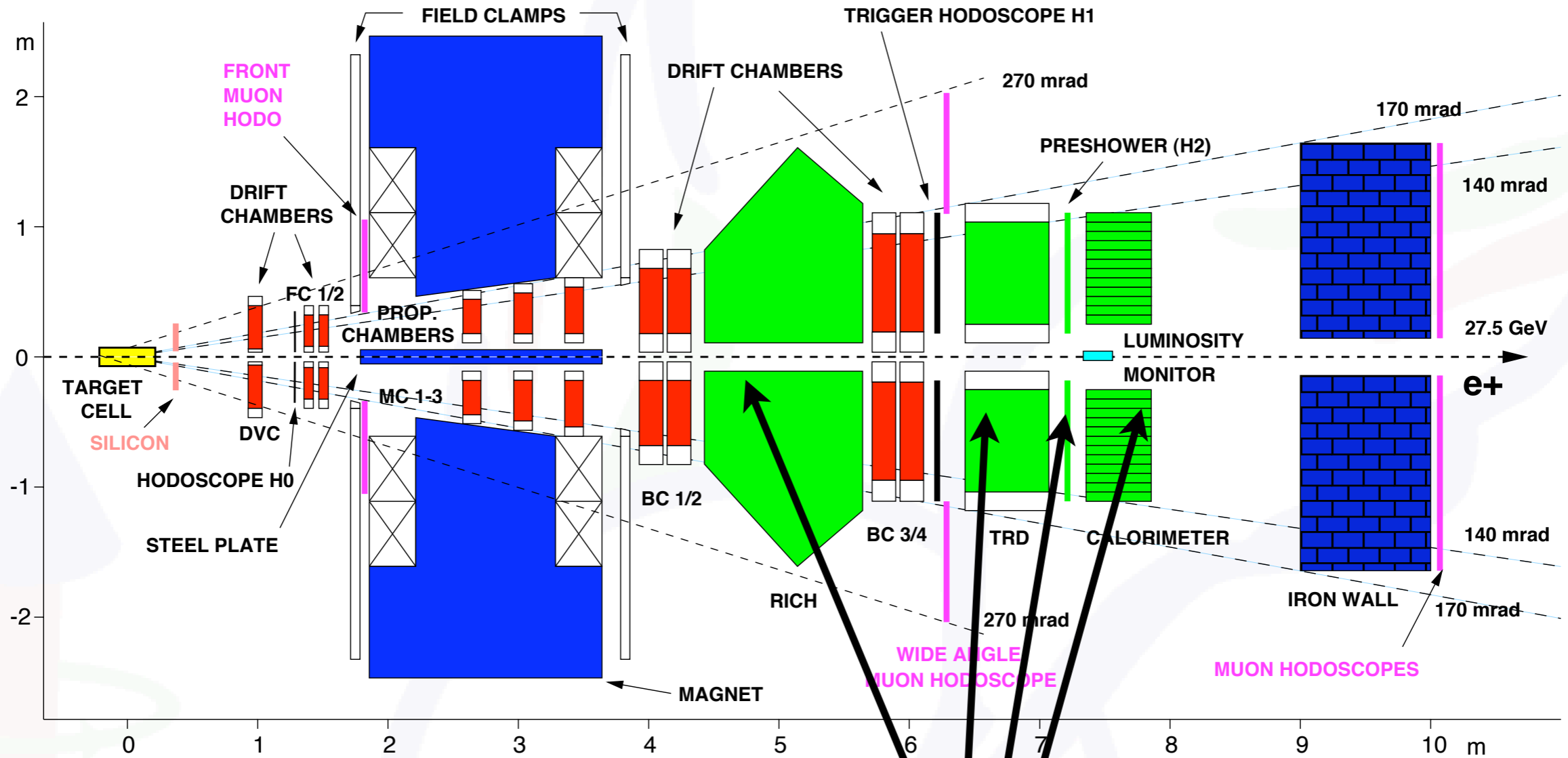




Overview of TMD results

from  hermes

HERMES schematically



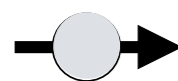

- pure gas targets internal to HERA 27.6 GeV lepton ring
- unpolarized ($^1\text{H} \dots \text{Xe}$)
- long. polarized: ^1H , ^2H , ^3He
- transversely polarized: ^1H



Particle ID detectors allow for

- lepton/hadron separation
- RICH: pion/kaon/proton discrimination $2\text{GeV} < p < 15\text{GeV}$

TMDs - probabilistic interpretation

Proton goes out of the screen/ photon goes into the screen

  nucleon with transverse or longitudinal spin

  parton with transverse or longitudinal spin

 parton transverse momentum

$$f_1 = \text{[Diagram: circle with red dot]}$$

$$g_1 = \text{[Diagram: circle with black dot and red dot]} - \text{[Diagram: circle with black dot and red cross]}$$

$$h_1 = \text{[Diagram: circle with red dot and right arrow]} - \text{[Diagram: circle with red dot and left arrow]}$$

$$f_{1T}^\perp = \text{[Diagram: circle with blue down arrow and red dot]} - \text{[Diagram: circle with blue up arrow and red dot]}$$

$$h_1^\perp = \text{[Diagram: circle with blue down arrow, red dot, and red right arrow]} - \text{[Diagram: circle with blue up arrow, red dot, and red right arrow]}$$

$$g_{1T} = \text{[Diagram: circle with red dot and blue right arrow]} - \text{[Diagram: circle with red dot and blue left arrow]}$$

$$h_{1L}^\perp = \text{[Diagram: circle with black dot, red dot, and blue right arrow]} - \text{[Diagram: circle with black dot, red dot, and blue left arrow]}$$

$$h_{1T}^\perp = \text{[Diagram: circle with red dot, blue right arrow, and blue right arrow]} - \text{[Diagram: circle with red dot, blue left arrow, and blue left arrow]}$$

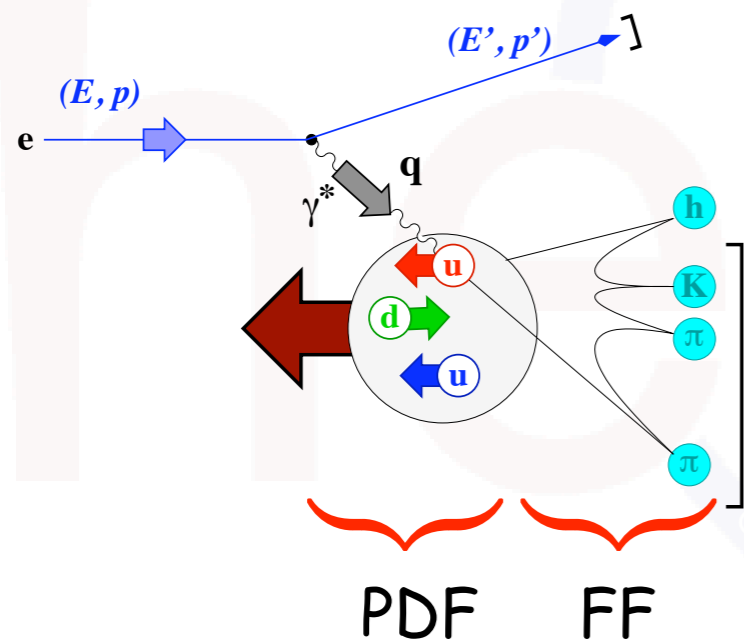
[courtesy of A. Bacchetta]

Probing TMDs in semi-inclusive DIS

quark pol.

	U	L	T
nucleon pol.	U		h_1^\perp
	L	g_{1L}	h_{1L}^\perp
	T	f_{1T}^\perp	h_1, h_{1T}^\perp

in SIDIS*) couple PDFs to:



*) semi-inclusive DIS with unpolarized final state

Probing TMDs in semi-inclusive DIS

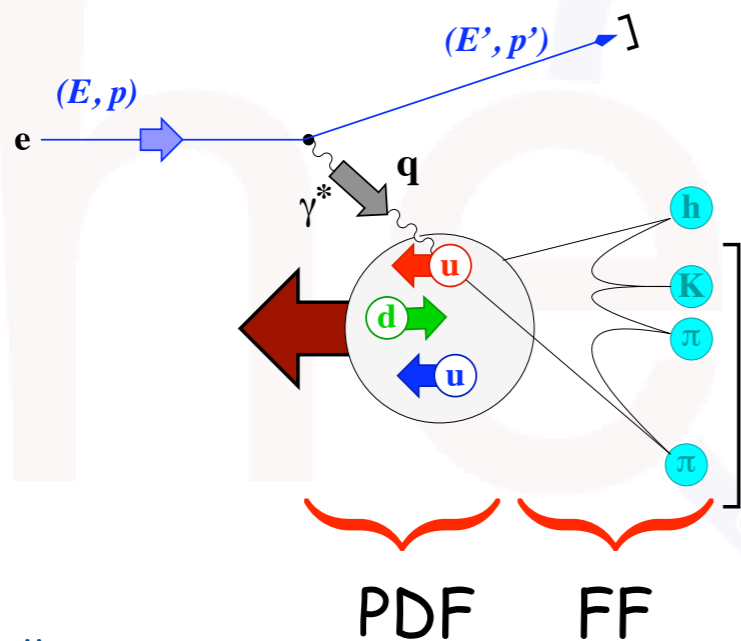
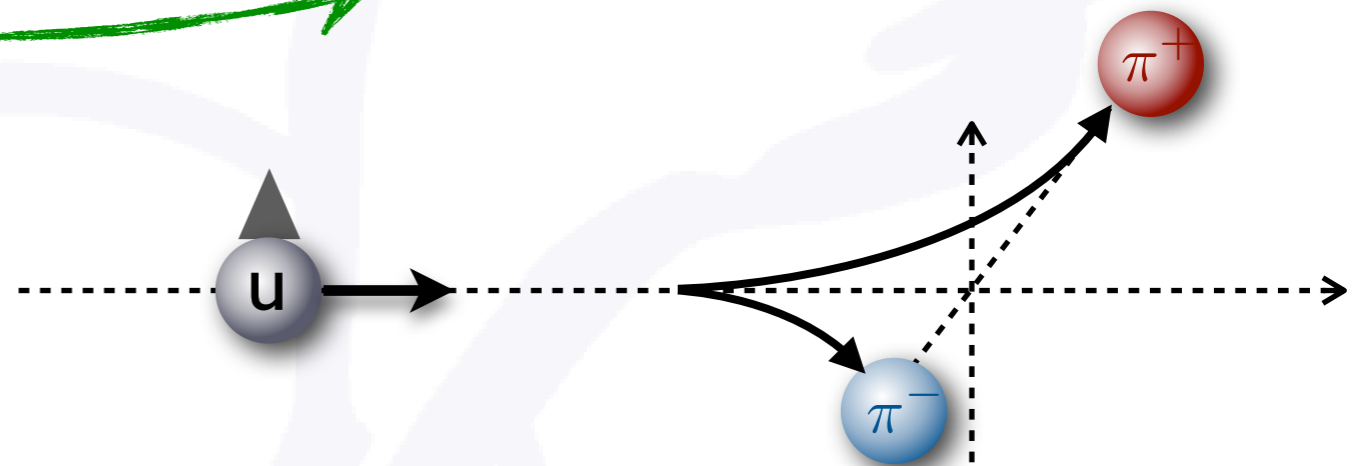
nucleon pol.

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Collins FF: $H_1^{\perp, q \rightarrow h}$



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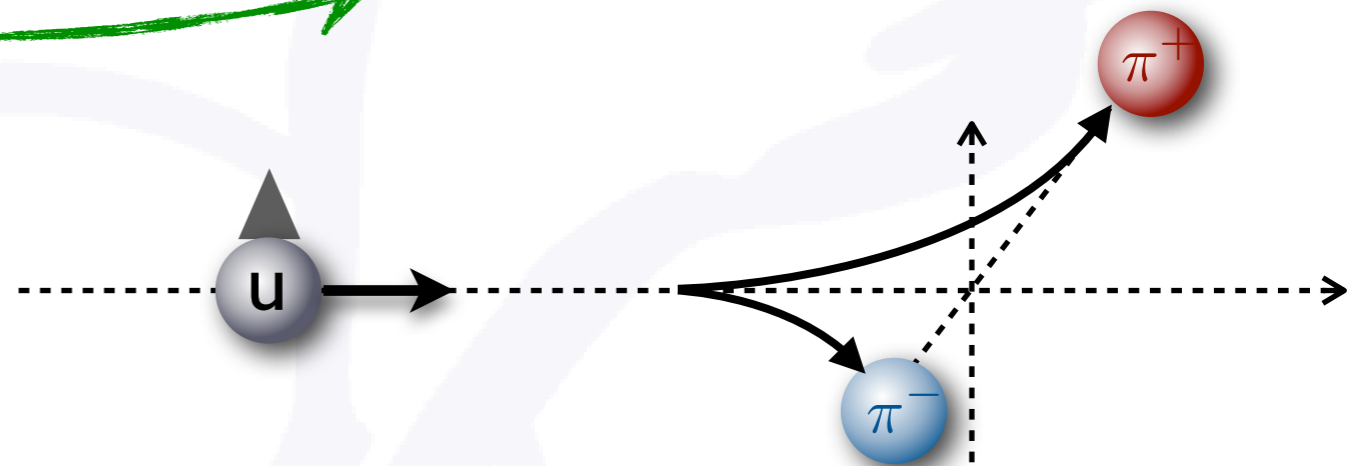
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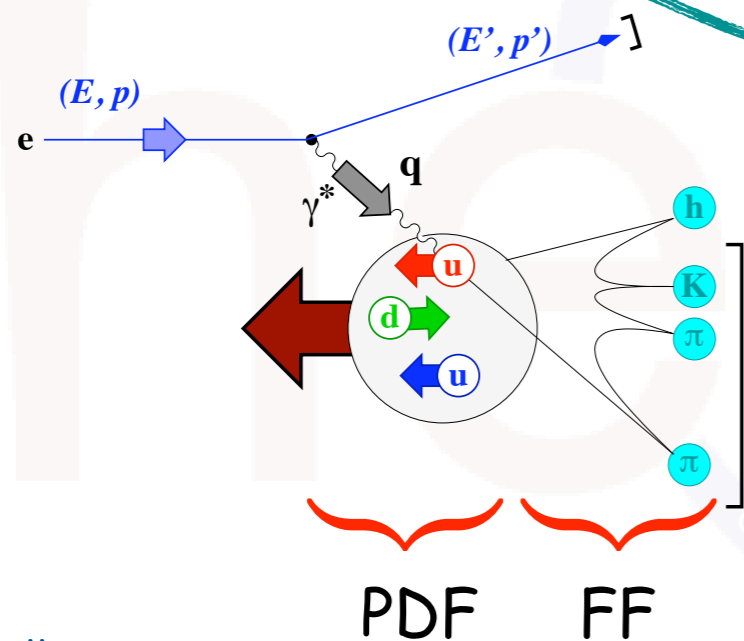
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ordinary FF: $D_1^{q \rightarrow h}$



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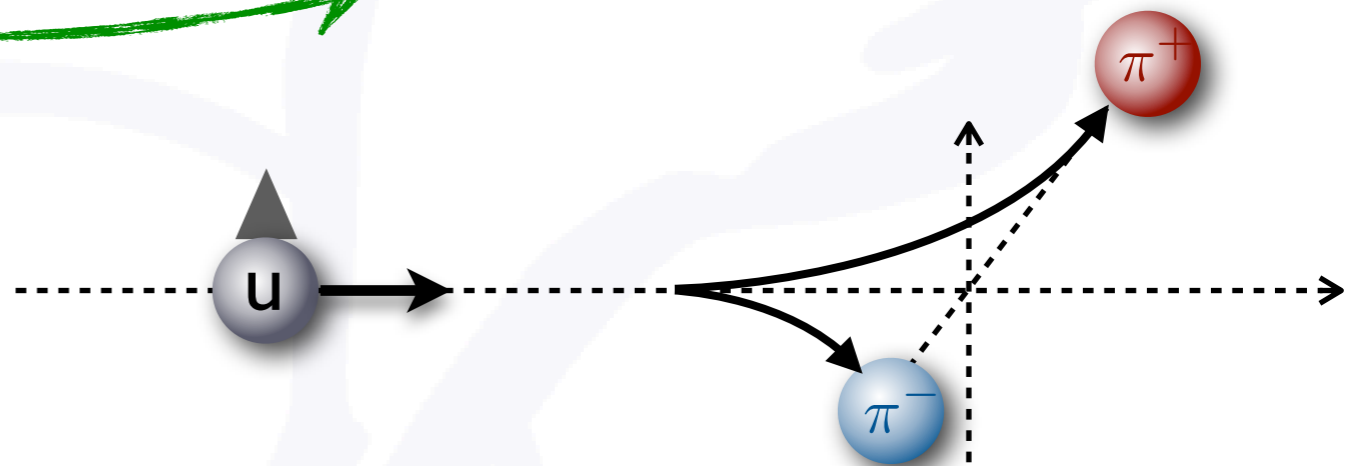
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	U	L	T
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L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

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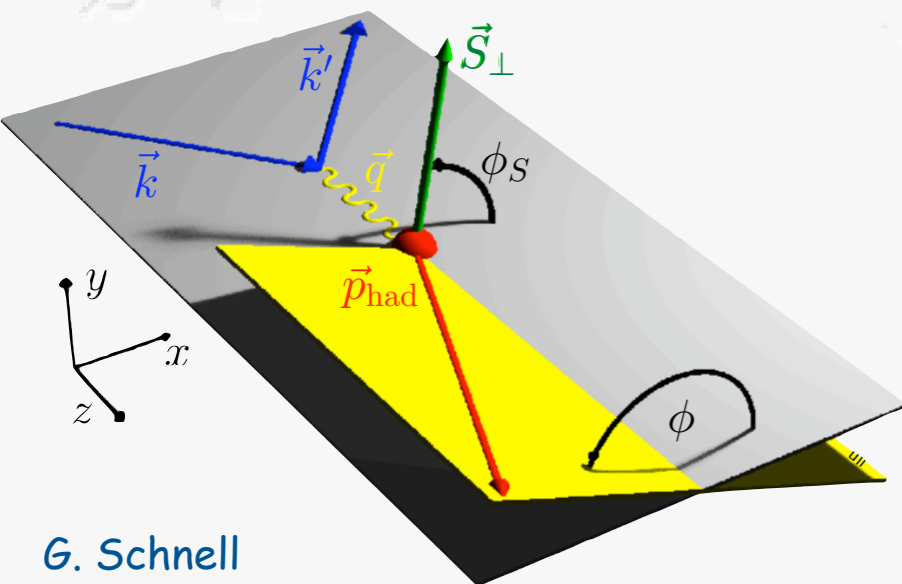
gives rise to characteristic azimuthal dependences

*) semi-inclusive DIS with unpolarized final state

1-Hadron production ($ep \rightarrow ehX$)

$$\begin{aligned}
 d\sigma = & d\sigma_{UU}^0 + \cos 2\phi d\sigma_{UU}^1 + \frac{1}{Q} \cos \phi d\sigma_{UU}^2 + \lambda_e \frac{1}{Q} \sin \phi d\sigma_{LU}^3 \\
 & + S_L \left\{ \sin 2\phi d\sigma_{UL}^4 + \frac{1}{Q} \sin \phi d\sigma_{UL}^5 + \lambda_e \left[d\sigma_{LL}^6 + \frac{1}{Q} \cos \phi d\sigma_{LL}^7 \right] \right\} \\
 & + S_T \left\{ \sin(\phi - \phi_S) d\sigma_{UT}^8 + \sin(\phi + \phi_S) d\sigma_{UT}^9 + \sin(3\phi - \phi_S) d\sigma_{UT}^{10} \frac{1}{Q} \right. \\
 & \quad \left. + \frac{1}{Q} (\sin(2\phi - \phi_S) d\sigma_{UT}^{11} + \sin \phi_S d\sigma_{UT}^{12}) \right. \\
 & \quad \left. + \lambda_e \left[\cos(\phi - \phi_S) d\sigma_{LT}^{13} + \frac{1}{Q} (\cos \phi_S d\sigma_{LT}^{14} + \cos(2\phi - \phi_S) d\sigma_{LT}^{15}) \right] \right\}
 \end{aligned}$$

σ_{XY}
 ↙ ↘
Beam Target
Polarization



Mulders and Tangermann, Nucl. Phys. B 461 (1996) 197

Boer and Mulders, Phys. Rev. D 57 (1998) 5780

Bacchetta et al., Phys. Lett. B 595 (2004) 309

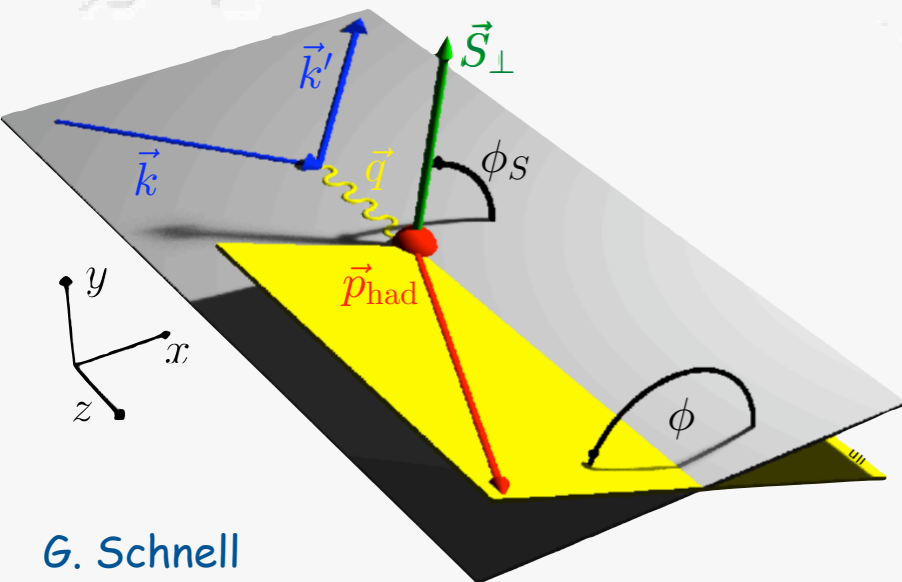
Bacchetta et al., JHEP 0702 (2007) 093

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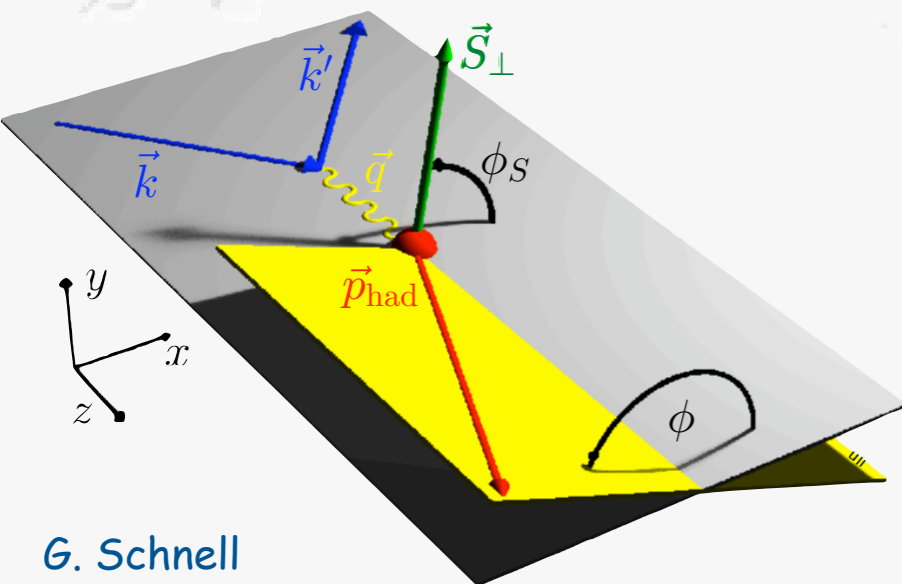
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Hadron multiplicities in DIS

$$\frac{d^5\sigma}{dx dy dz d\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \{F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1-\epsilon)} F_{UU}^{\cos\phi_h} \cos\phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h\}$$

$$F_{XY,Z} = F_{XY,Z}(x, y, z, P_{h\perp})$$

target polarization \downarrow
 beam polarization \uparrow virtual-photon polarization \uparrow

[see, e.g., Bacchetta et al., JHEP 0702 (2007) 093]

$$\gamma = \frac{2Mx}{Q}$$

$$\epsilon = \frac{1 - y - \frac{1}{4}\gamma^2 y^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}\gamma^2 y^2}$$

Hadron multiplicities in DIS

hadron multiplicity:
normalize to inclusive DIS
cross section

$$\frac{d^2 \sigma^{\text{incl. DIS}}}{dxdy} \propto F_T + \epsilon F_L$$

$$\frac{d^4 \mathcal{M}^h(x, y, z, P_{h\perp}^2)}{dxdydzdP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \frac{F_{UU,T} + \epsilon F_{UU,L}}{F_T + \epsilon F_L}$$

$$\approx \frac{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^{q \rightarrow h}(z, K_T^2)}{\sum_q e_q^2 f_1^q(x)}$$

$$\frac{d^5 \sigma}{dxdydzd\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \left\{ F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1-\epsilon)} F_{UU}^{\cos \phi_h} \cos \phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h \right\}$$

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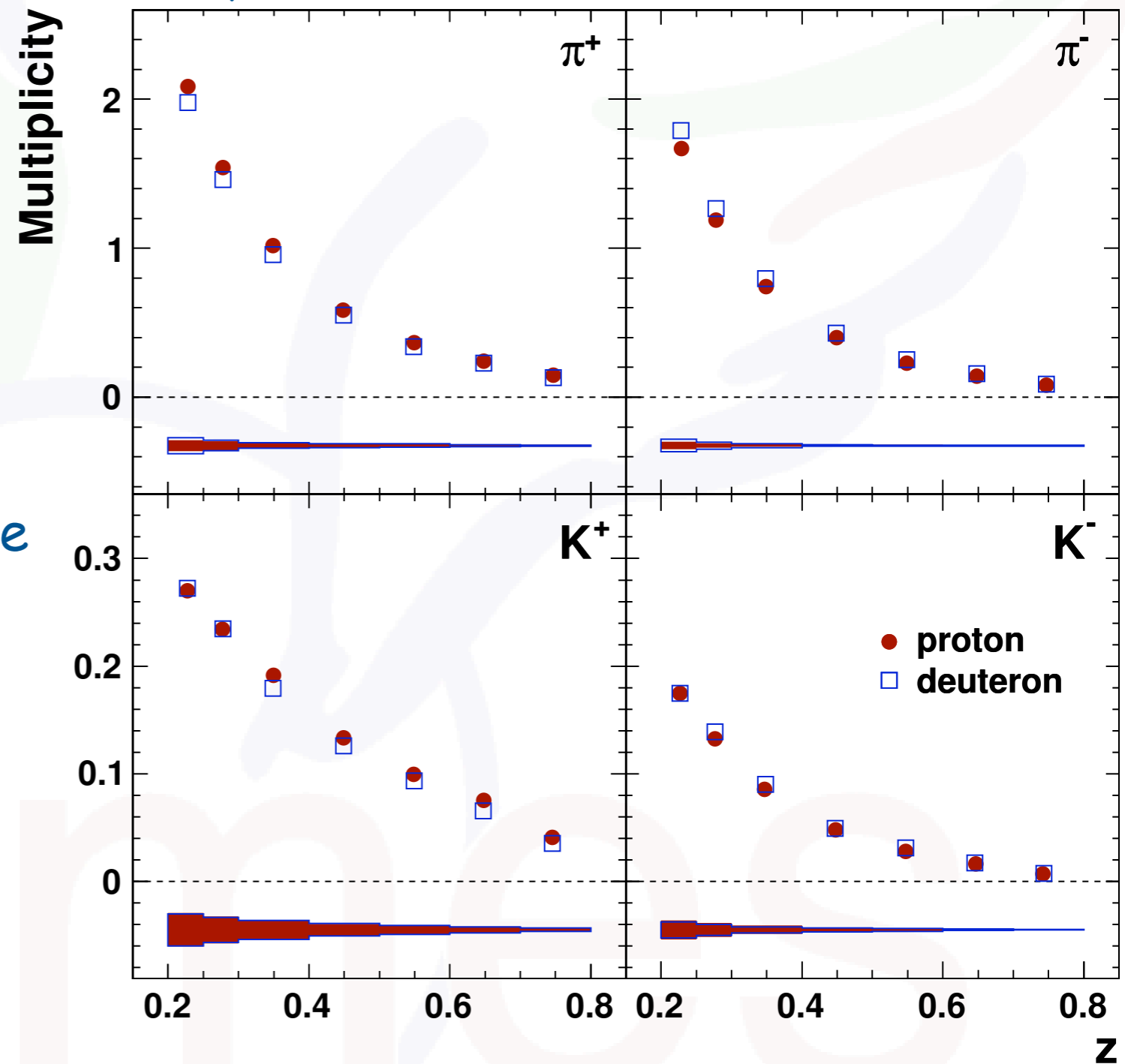
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Multiplicities @ HERMES

- extensive data set on pure proton and deuteron targets for identified charged mesons <http://www-hermes.desy.de/multiplicities>
- extracted in a multi-dimensional unfolding procedure

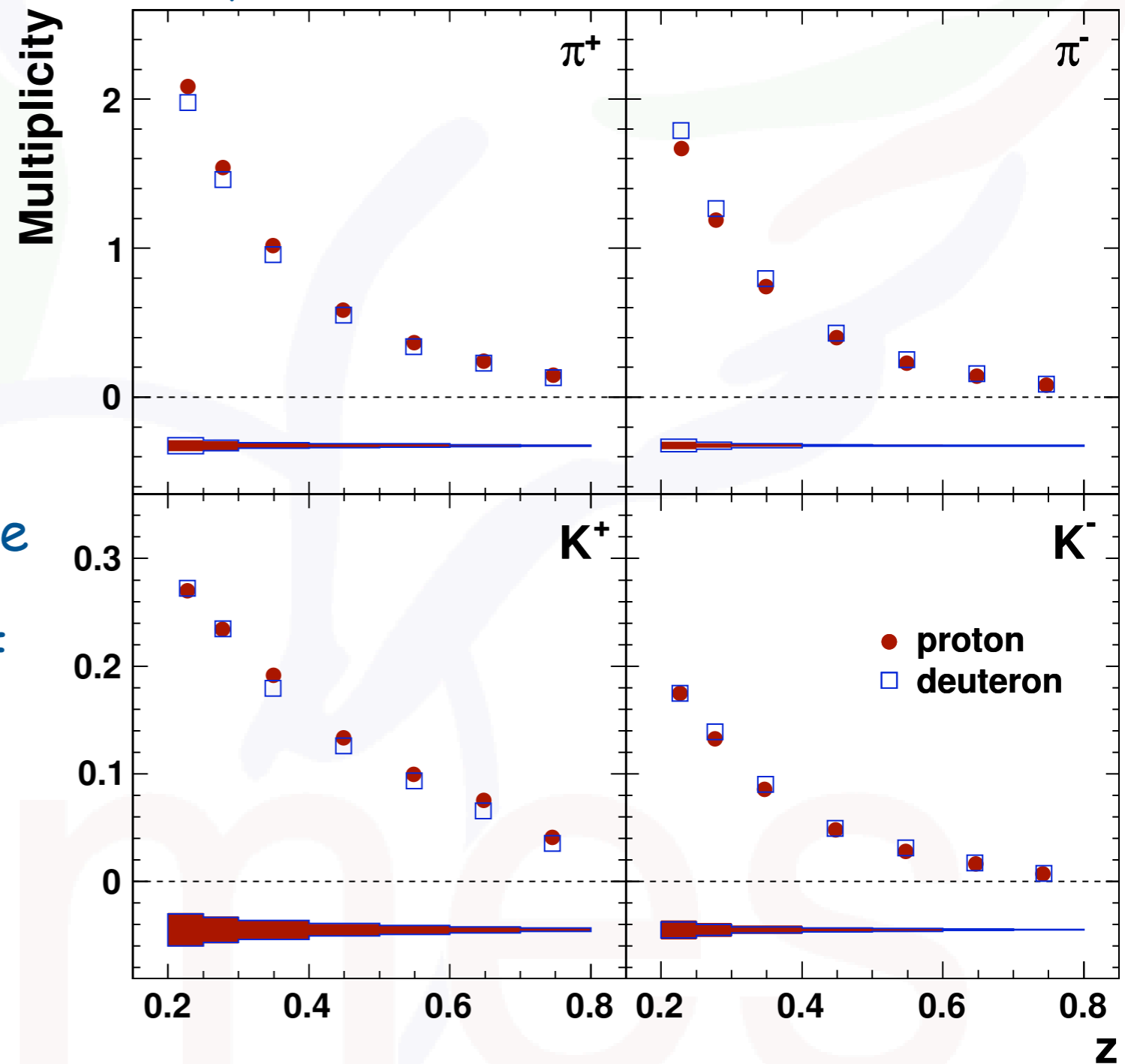
[Airapetian et al., PRD 87 (2013) 074029]



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[Airapetian et al., PRD 87 (2013) 074029]

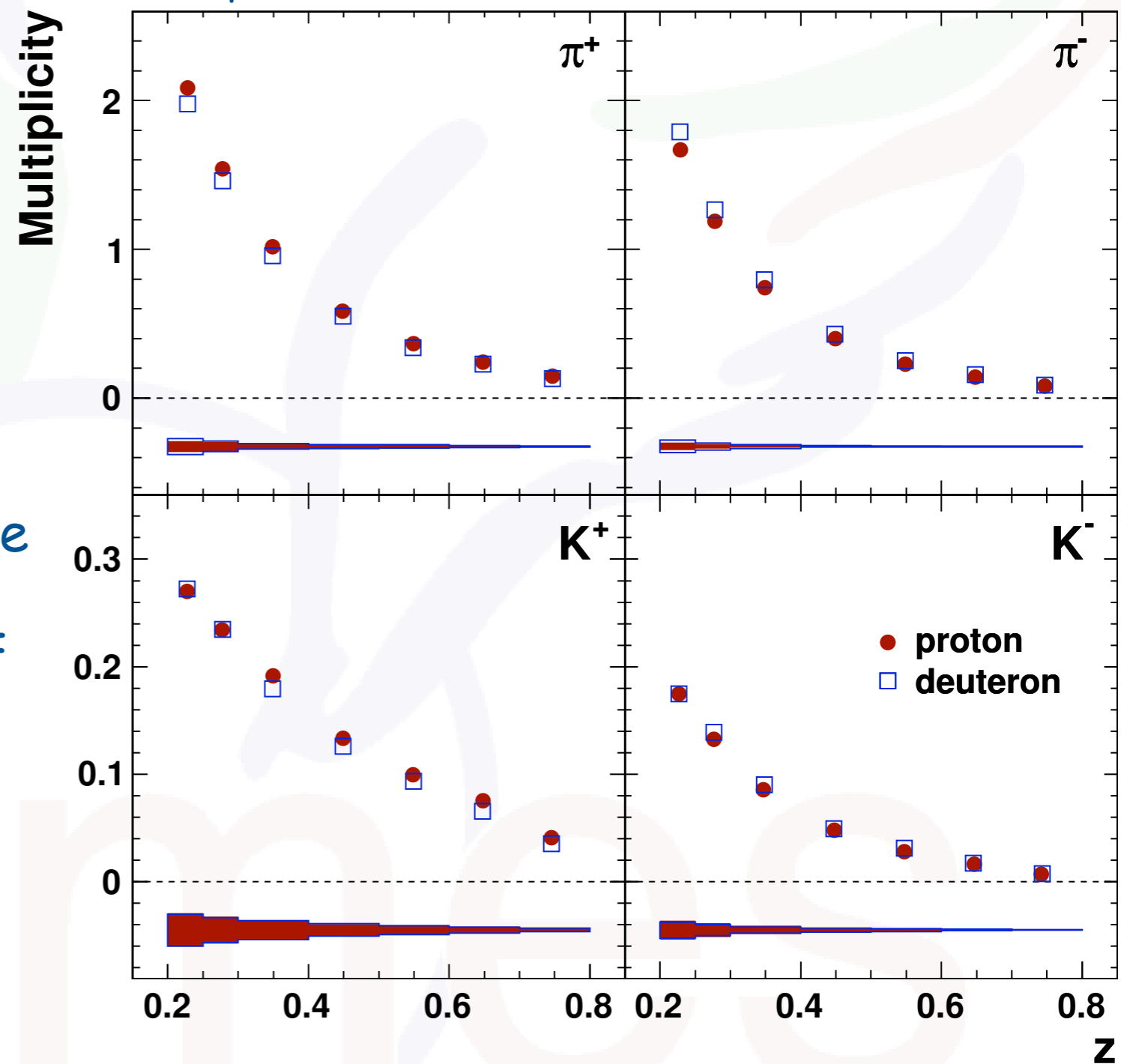


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➔ Helicity VI session

[Airapetian et al., PRD 87 (2013) 074029]

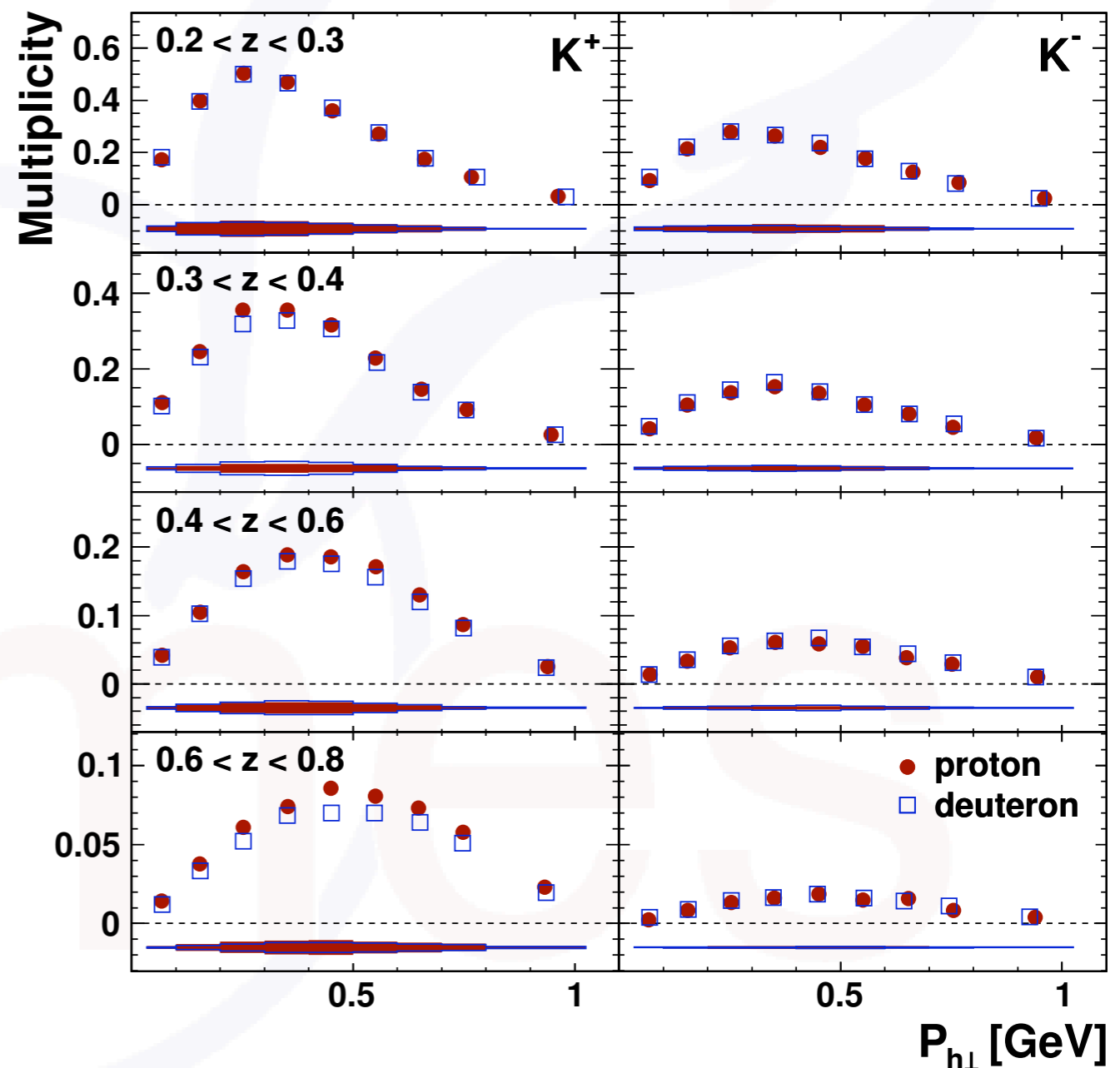
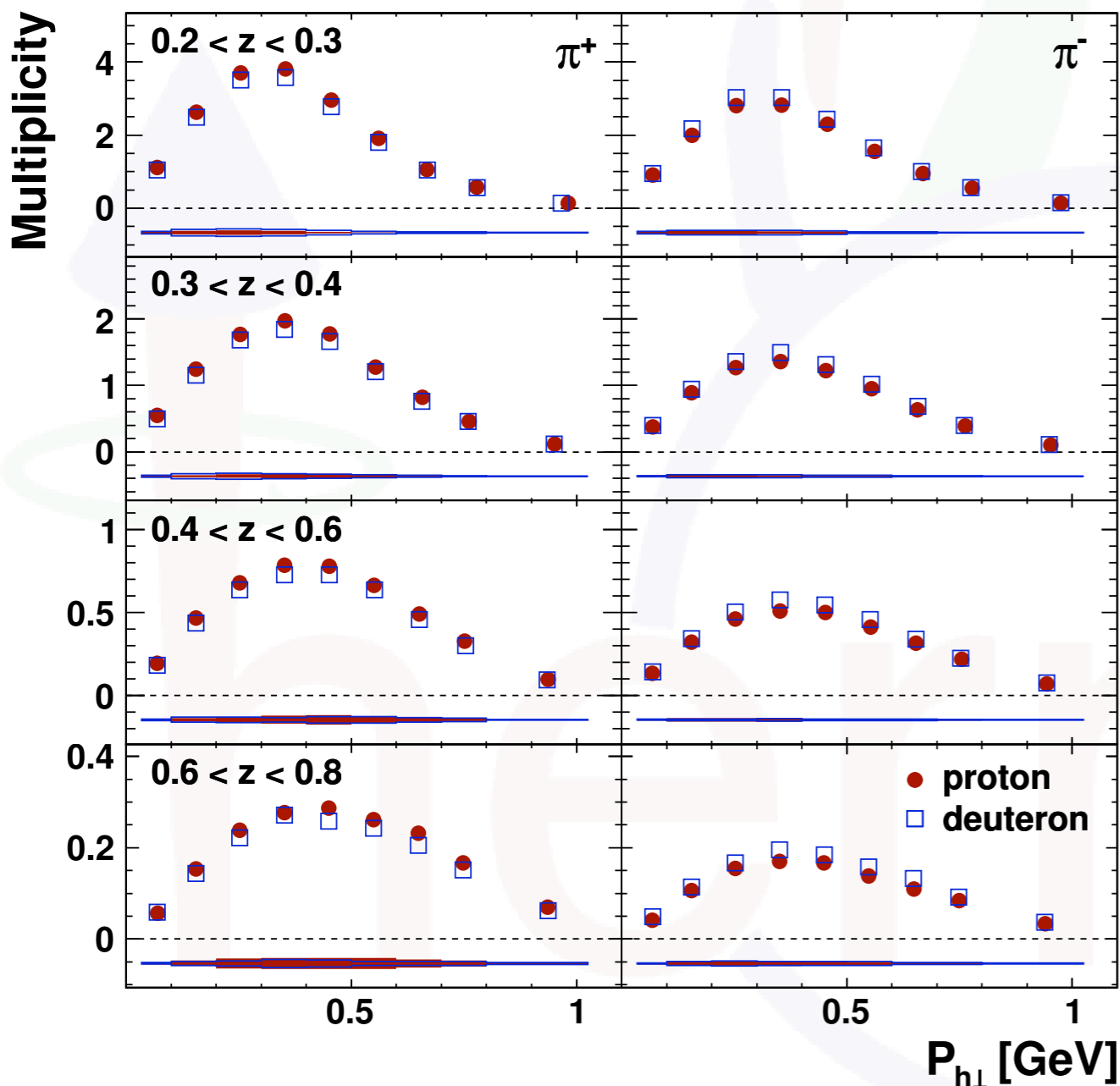


● proton
□ deuteron

Transverse momentum dependence

- multi-dimensional analysis allows going beyond collinear factorization
- flavor information on transverse momenta via target variation and hadron ID

[Airapetian et al., PRD 87 (2013) 074029]

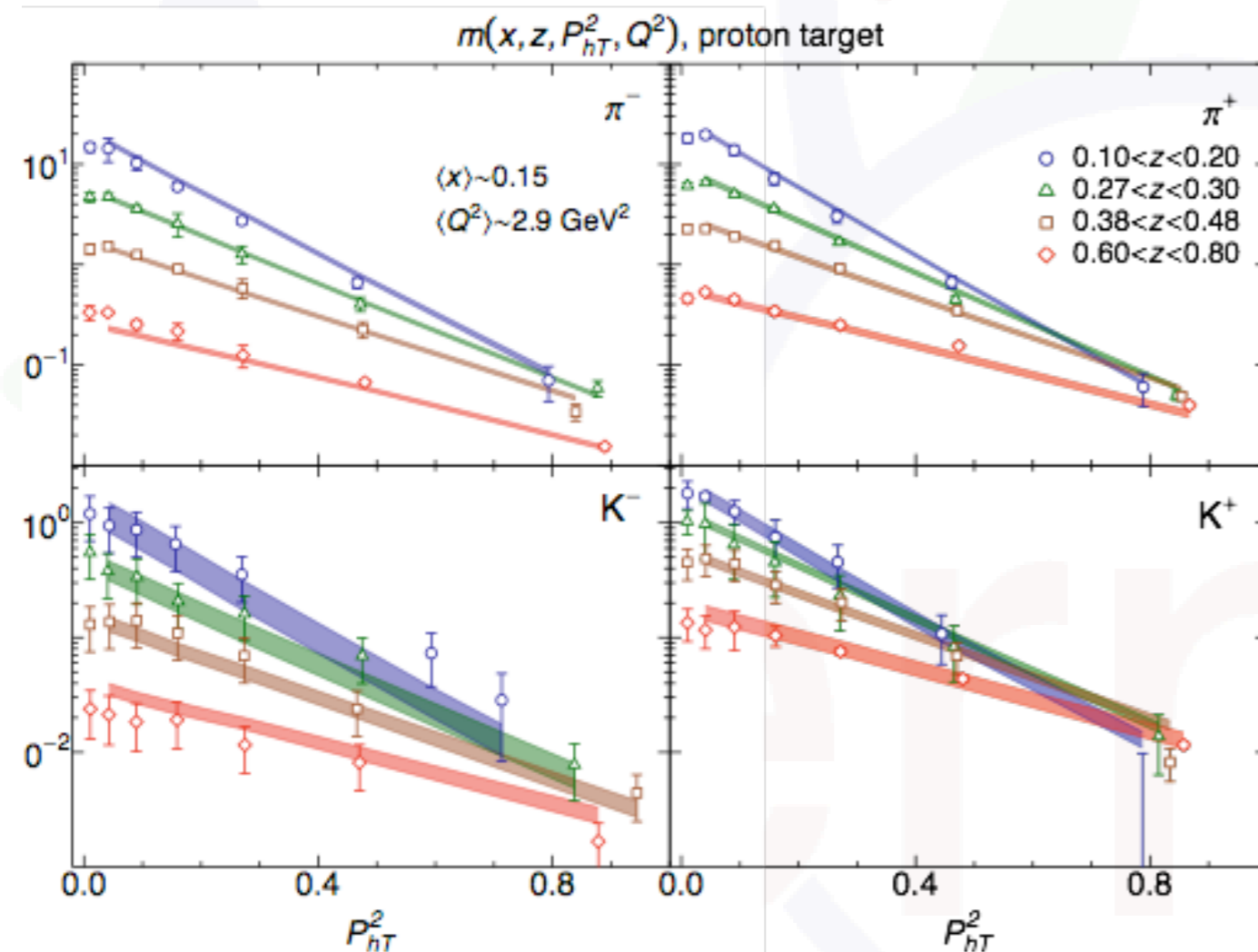


FF TMD flavor dependence

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

- fit to HERMES multiplicity data:

$$m_N^h(x, z, \mathbf{P}_{hT}^2; Q^2) = \frac{\pi}{\sum_q e_q^2 f_1^q(x; Q^2)} \sum_q e_q^2 f_1^q(x; Q^2) D_1^{q \rightarrow h}(z; Q^2) \frac{e^{-\mathbf{P}_{hT}^2 / \langle \mathbf{P}_{hT,q}^2 \rangle}}{\pi \langle \mathbf{P}_{hT,q}^2 \rangle}$$



$$f_1^q(x, \mathbf{k}_\perp^2; Q^2) = f_1^q(x; Q^2) \frac{e^{-\mathbf{k}_\perp^2 / \langle \mathbf{k}_{\perp,q}^2 \rangle}}{\pi \langle \mathbf{k}_{\perp,q}^2 \rangle}$$

$$D_1^{q \rightarrow h}(z, \mathbf{P}_\perp^2; Q^2) = D_1^{q \rightarrow h}(z; Q^2) \frac{e^{-\mathbf{P}_\perp^2 / \langle \mathbf{P}_{\perp,q \rightarrow h}^2 \rangle}}{\pi \langle \mathbf{P}_{\perp,q \rightarrow h}^2 \rangle}$$

$$\langle \mathbf{P}_{hT,q}^2 \rangle = z^2 \langle \mathbf{k}_{\perp,q}^2 \rangle + \langle \mathbf{P}_{\perp,q \rightarrow h}^2 \rangle$$

[A. Signori, A. Bacchetta, M. Radici and GS, JHEP 11(2013)194]

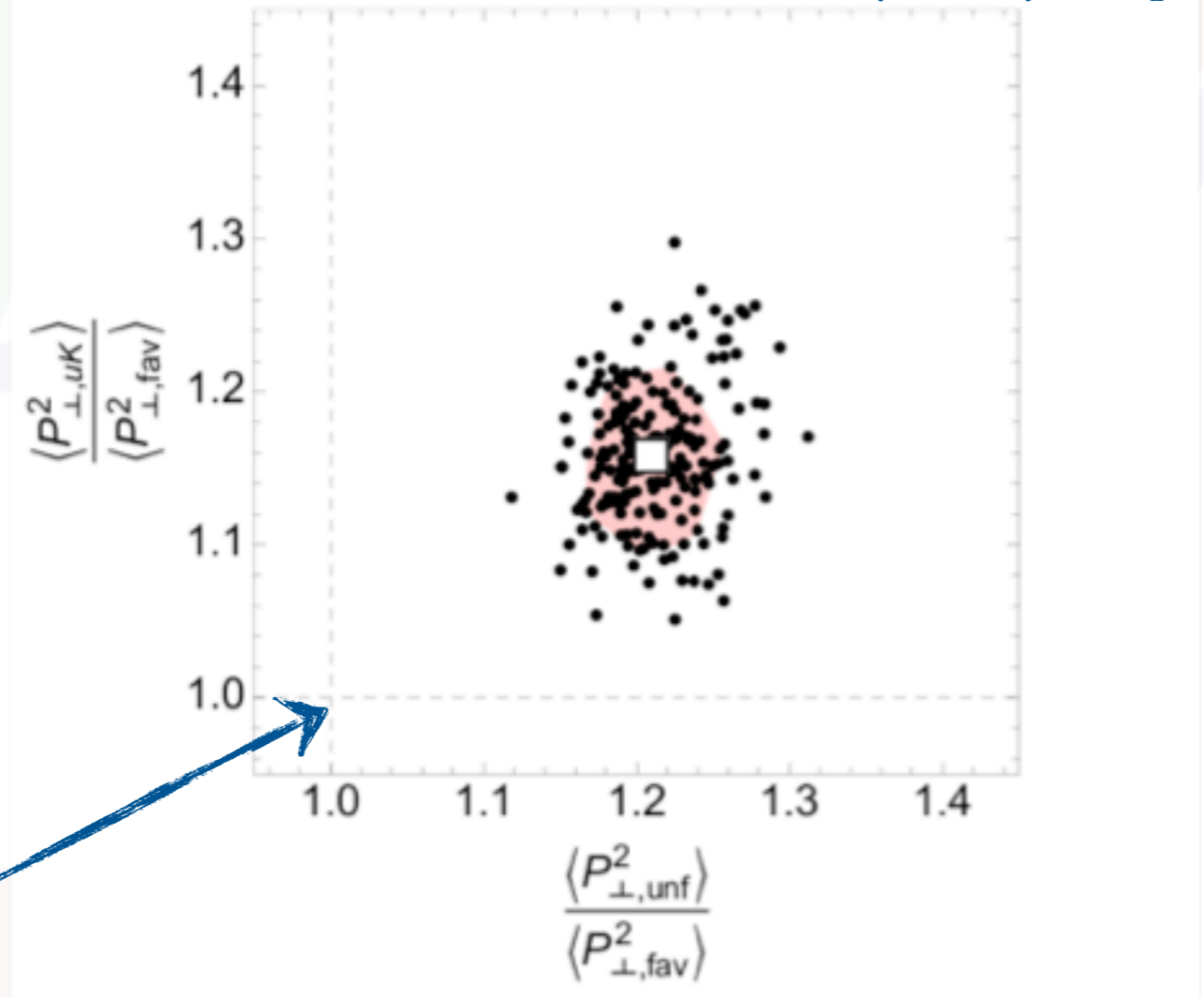
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$q \rightarrow \pi$ favored width
 $<$
 $q \rightarrow K$ favored width



point of
no flavor dep.

$q \rightarrow \pi$ favored width $<$ unfavored

Transversely polarized quarks?

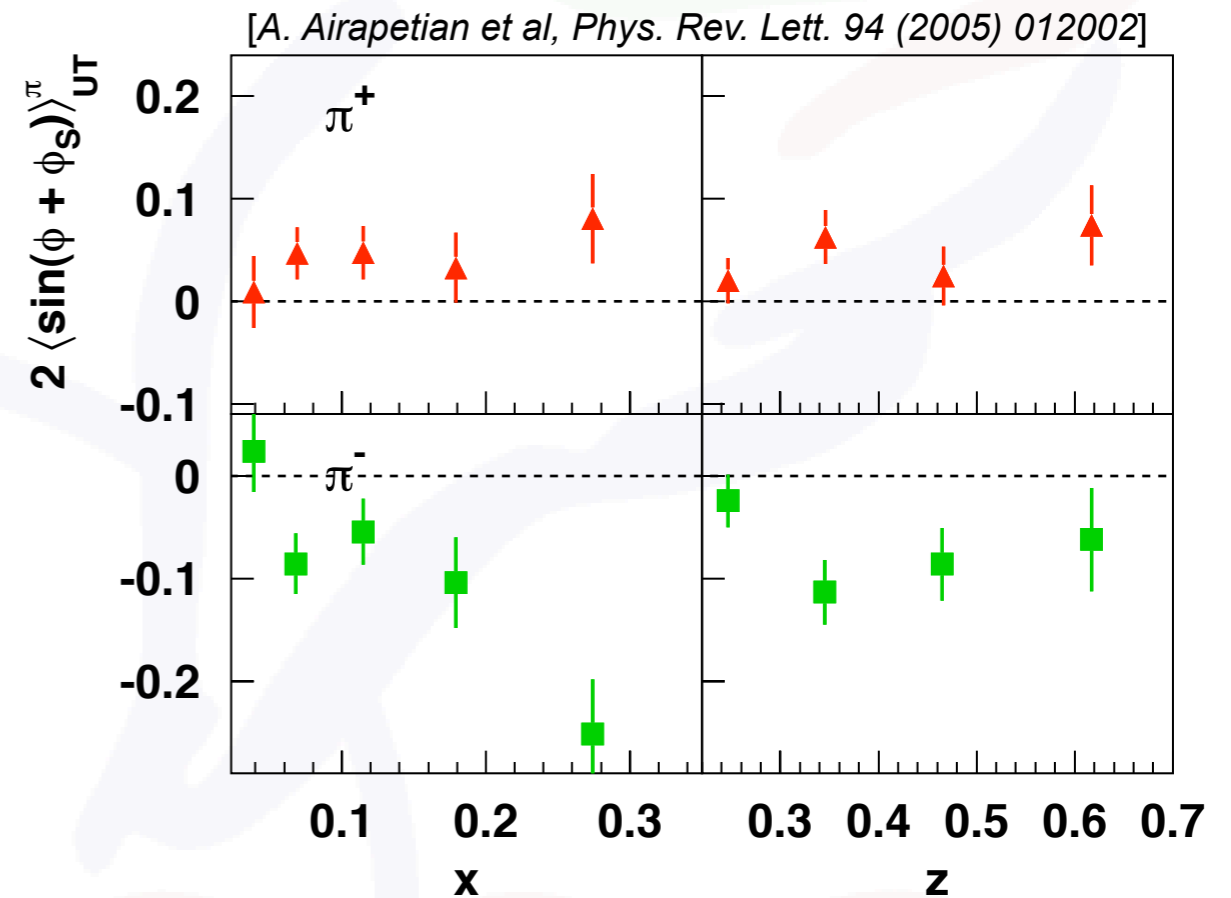
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- transverse polarization of quarks leads to large effects!



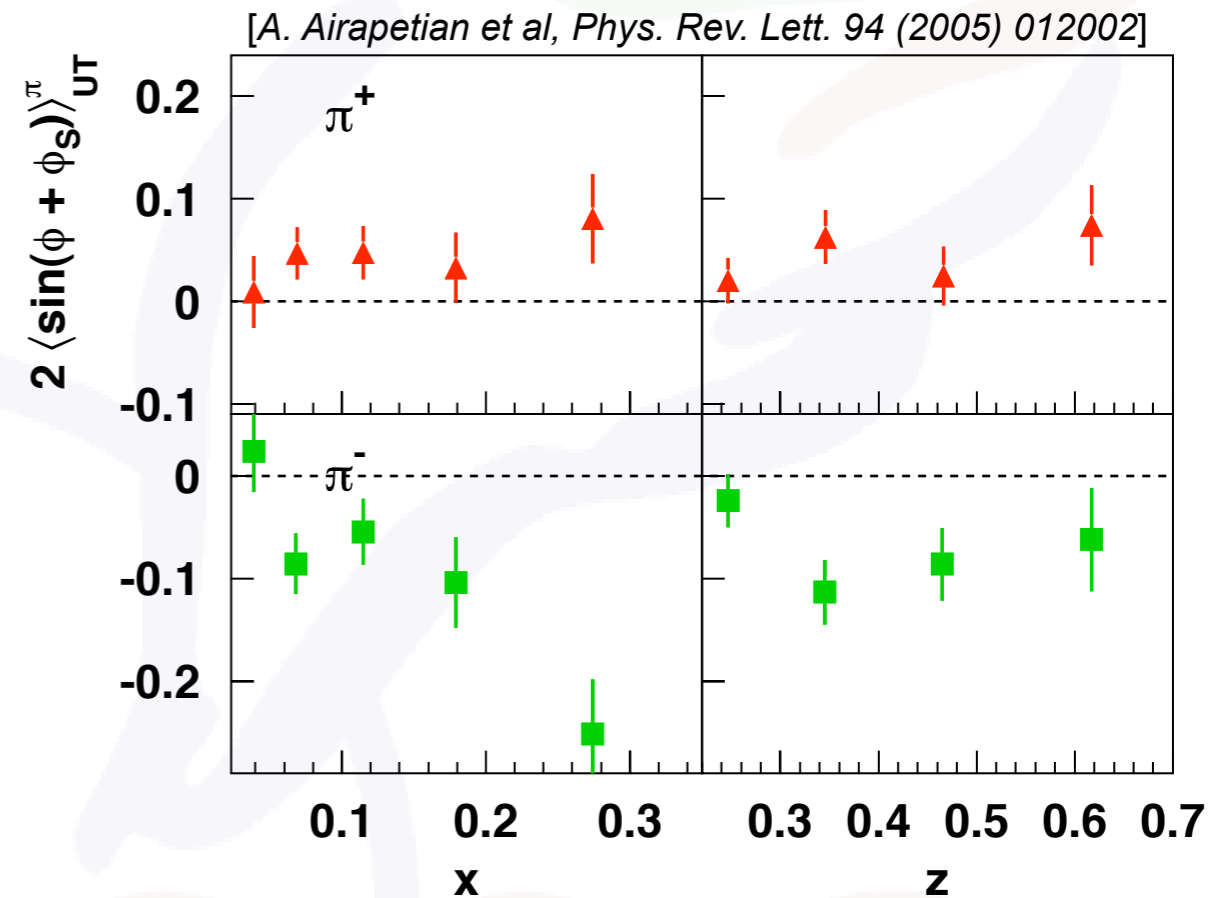
2005: First evidence from HERMES
SIDIS on proton

Non-zero transversity
Non-zero Collins function

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- transverse polarization of quarks leads to large effects!
- opposite in sign for charged pions



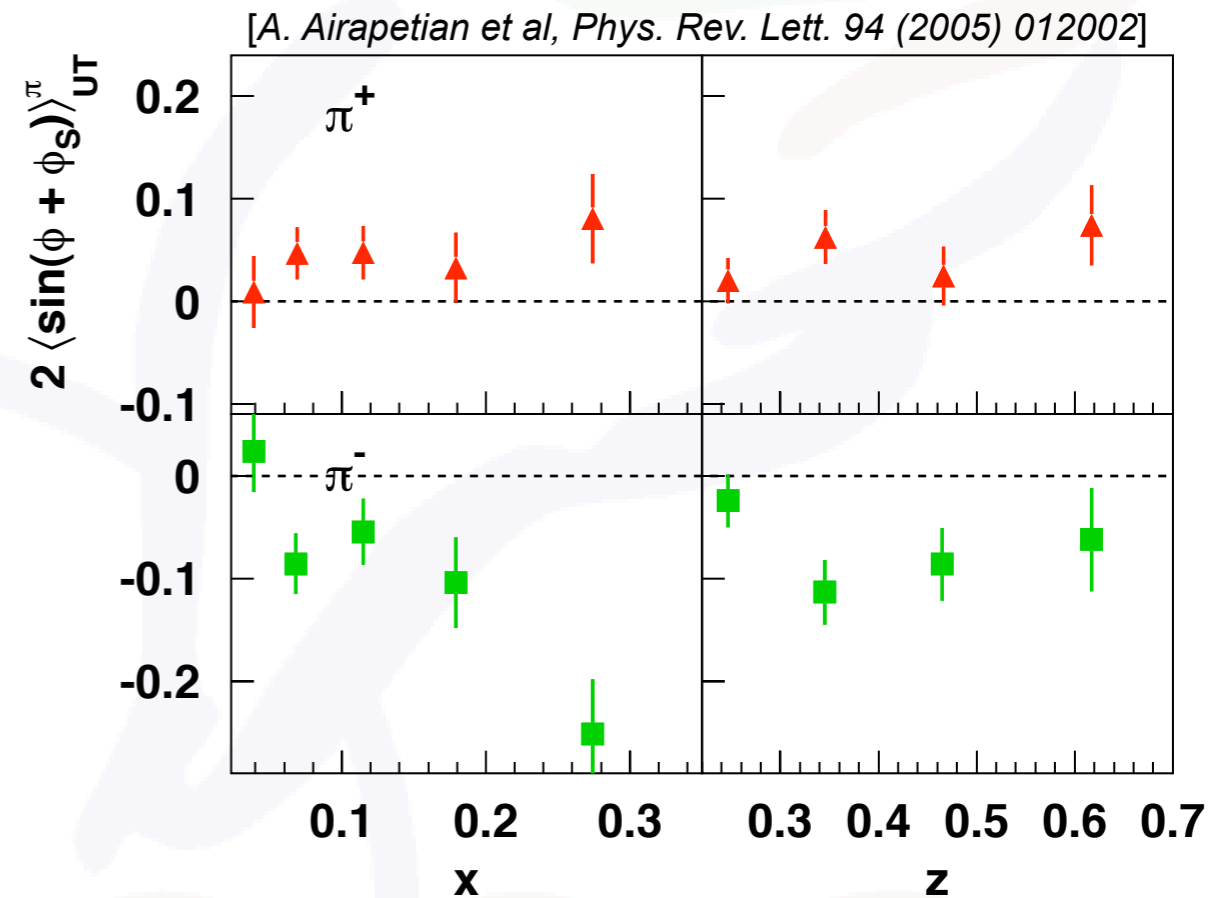
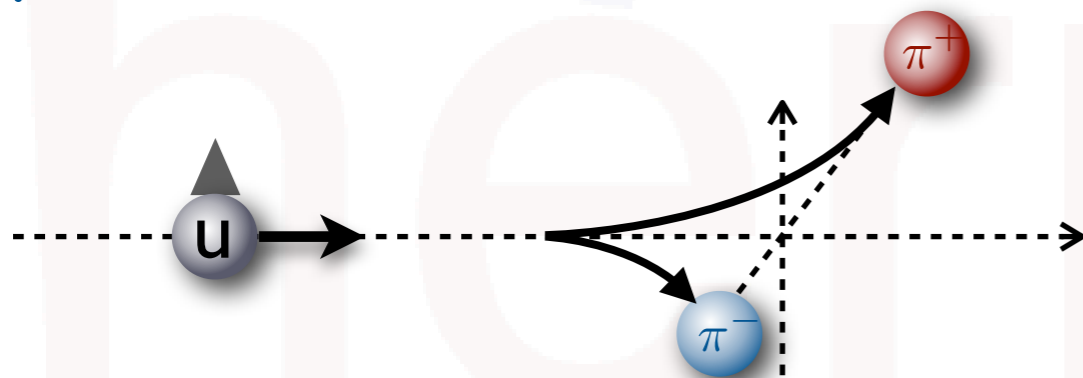
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- transverse polarization of quarks leads to large effects!
- opposite in sign for charged pions
- disfavored Collins FF large and opposite in sign to favored one



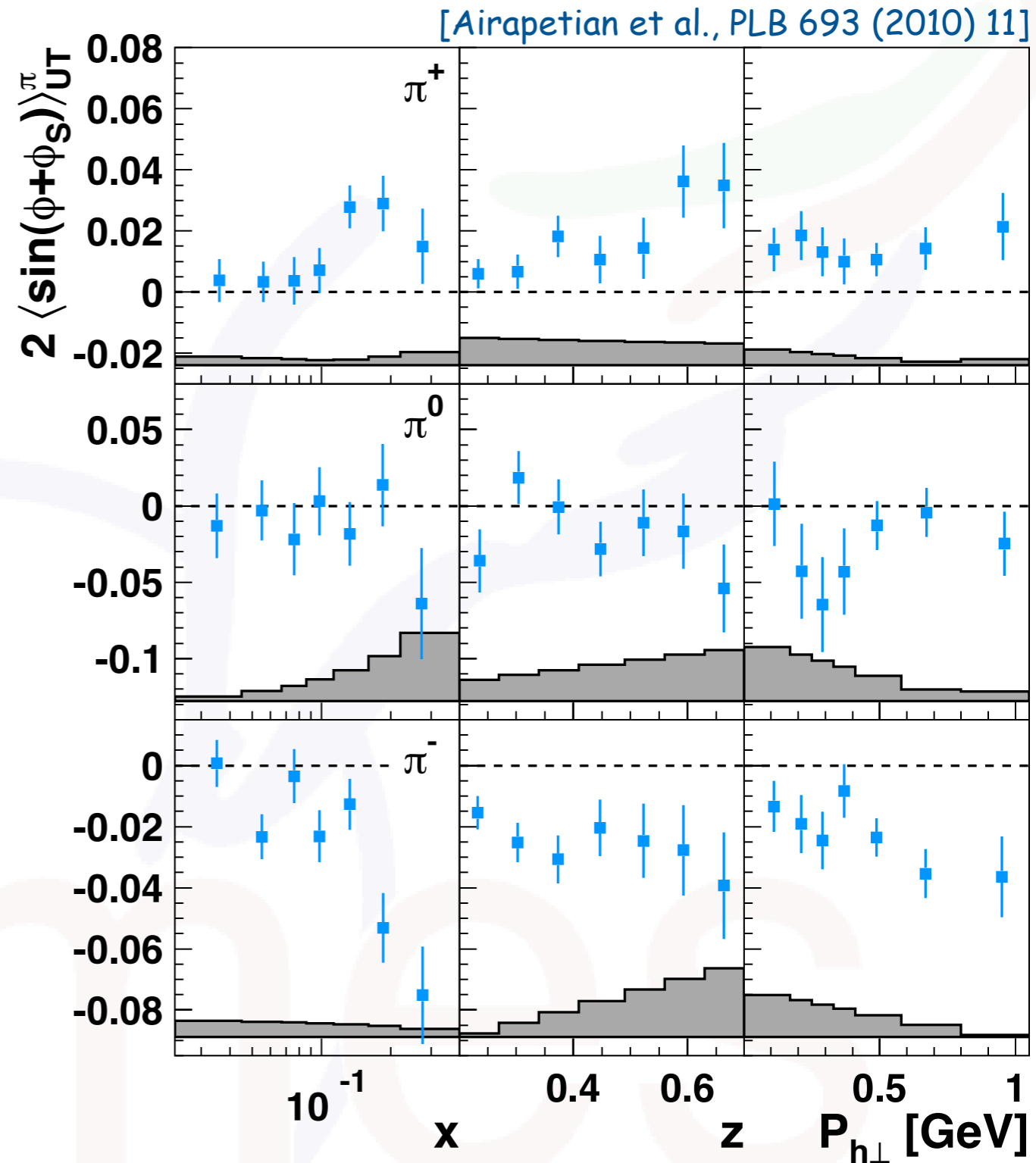
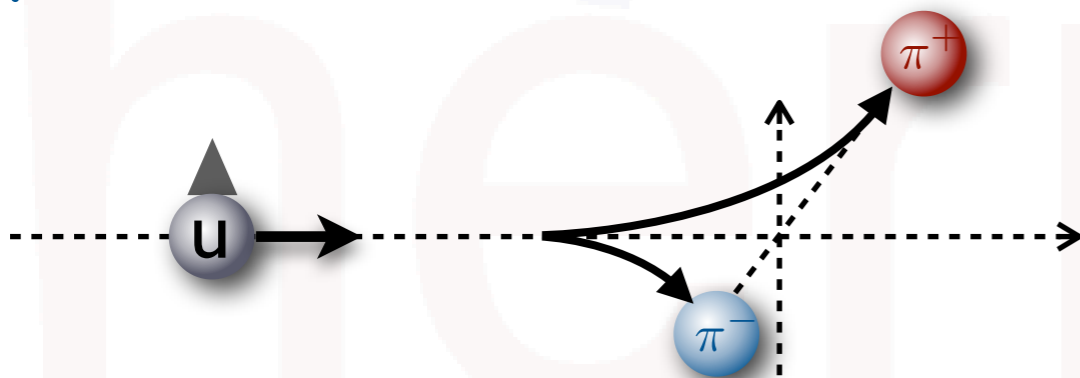
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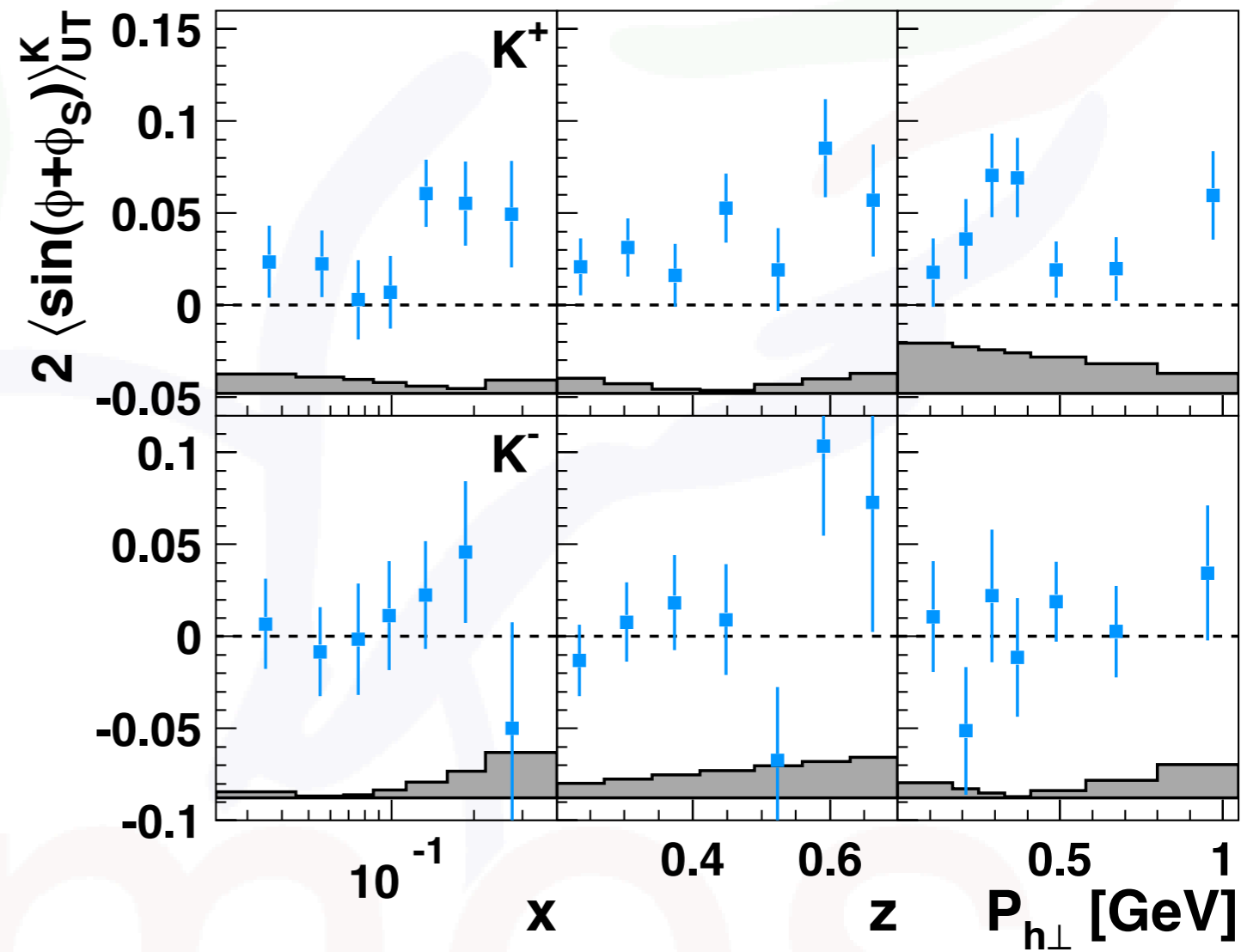
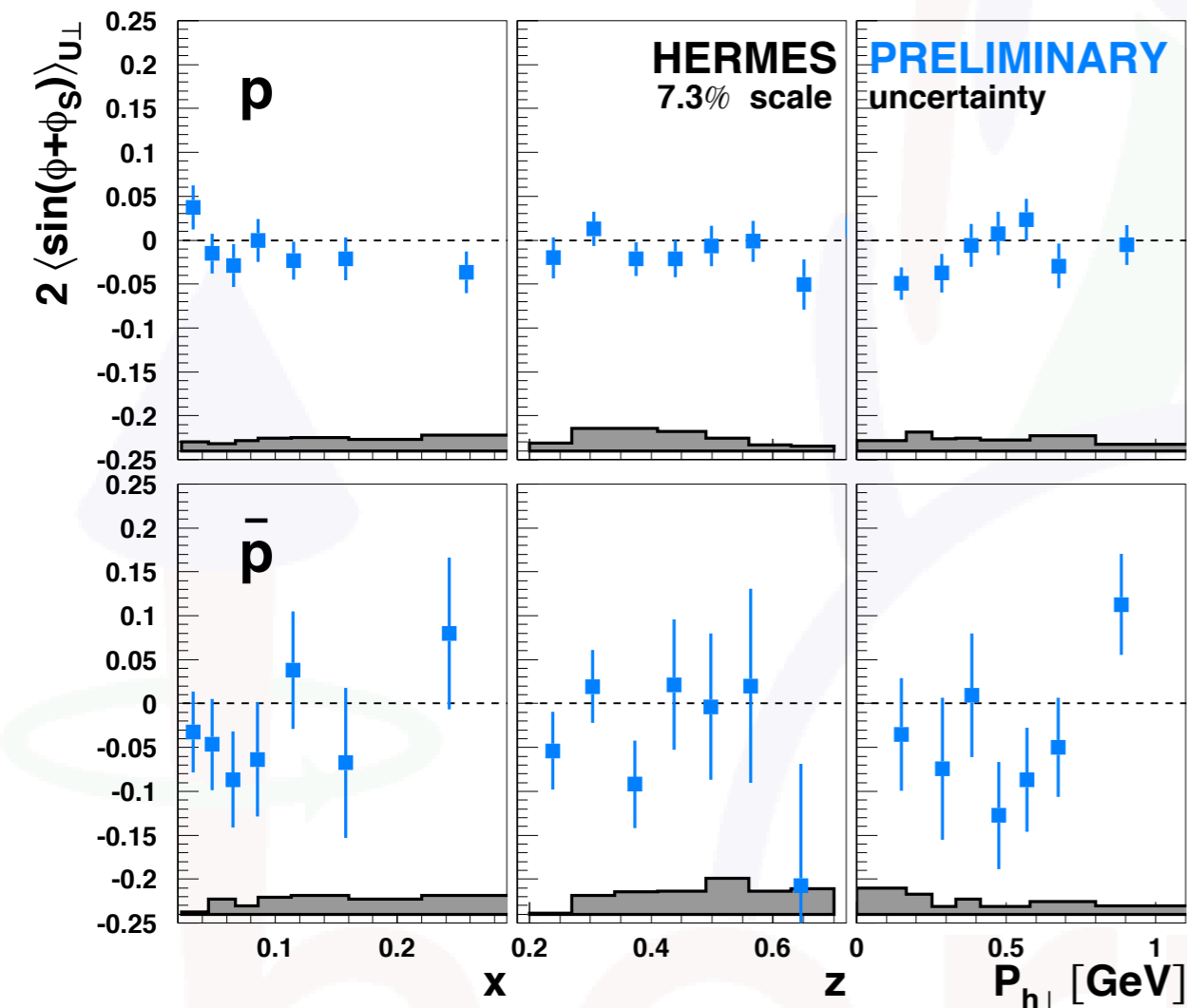
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Collins effect for kaons and (anti) protons

[Airapetian et al., PLB 693 (2010) 11]

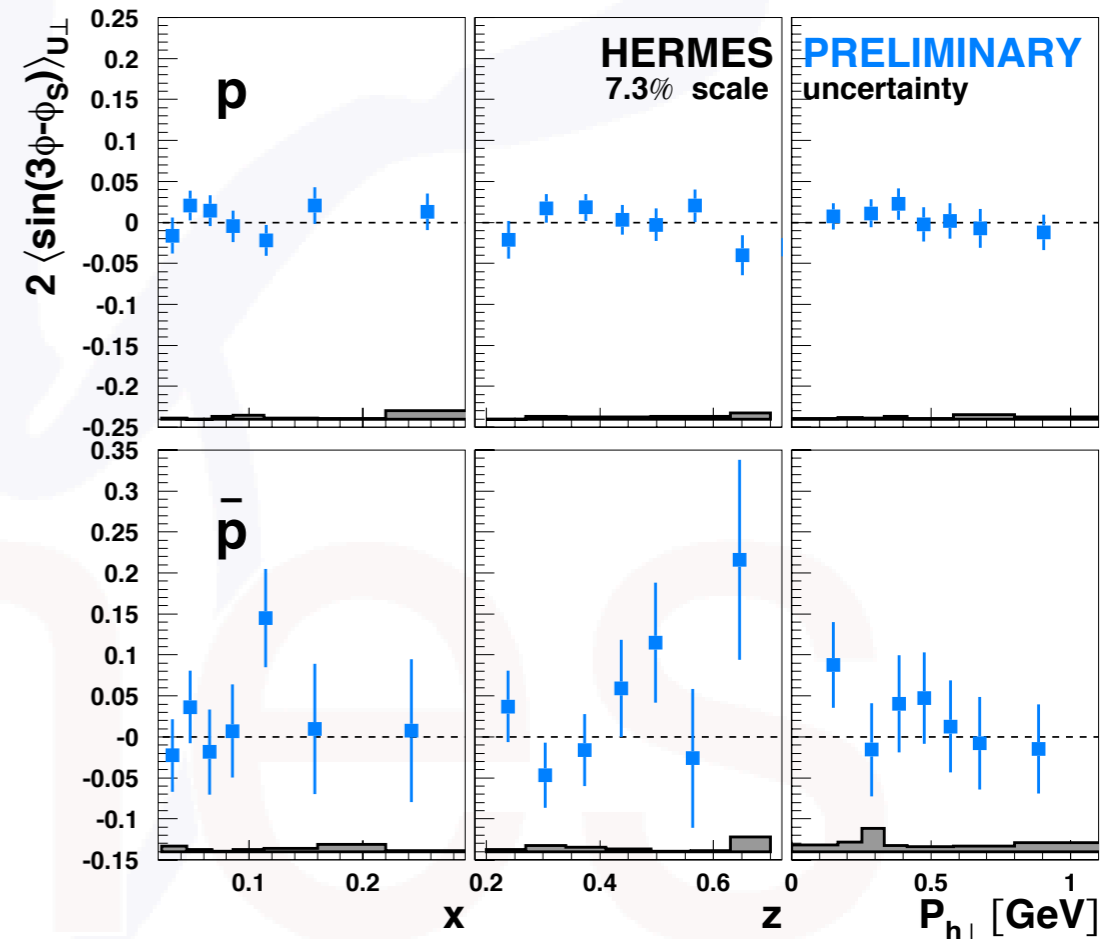
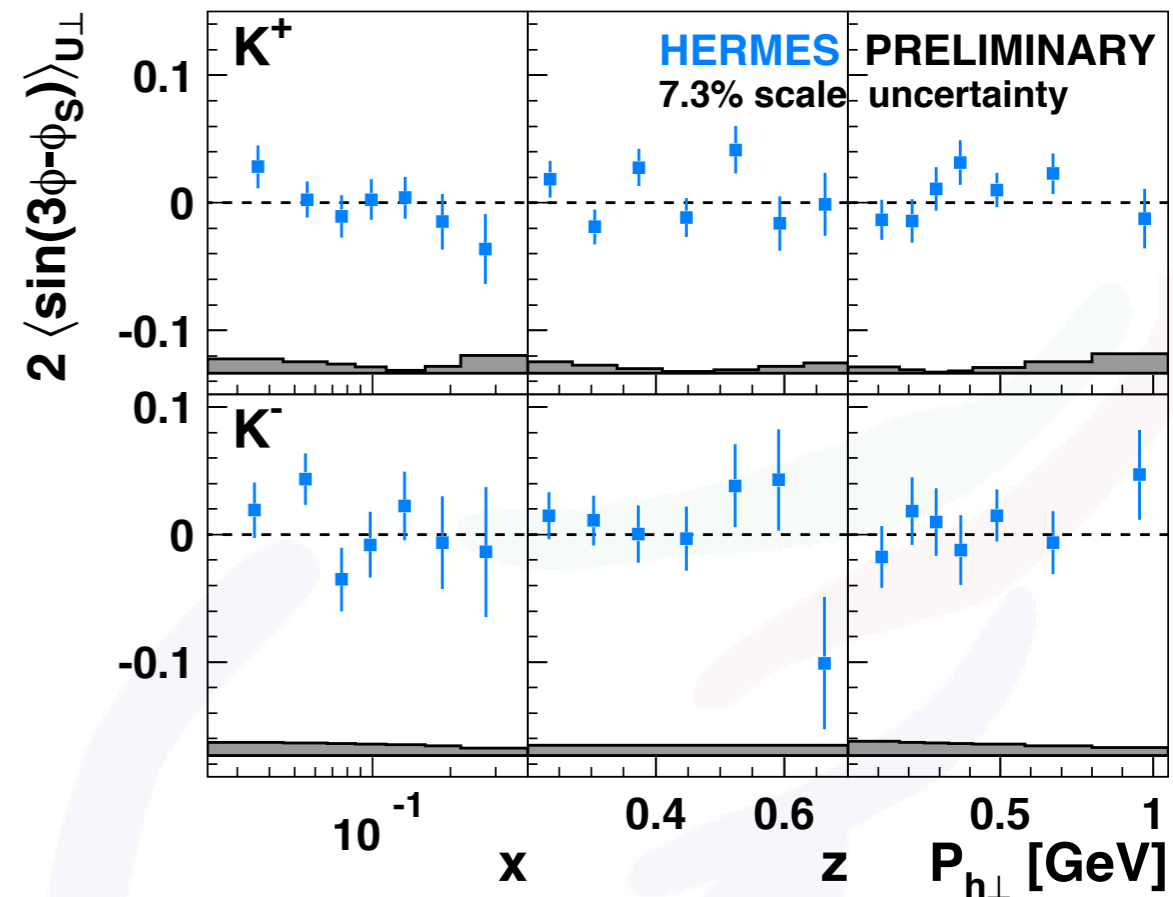
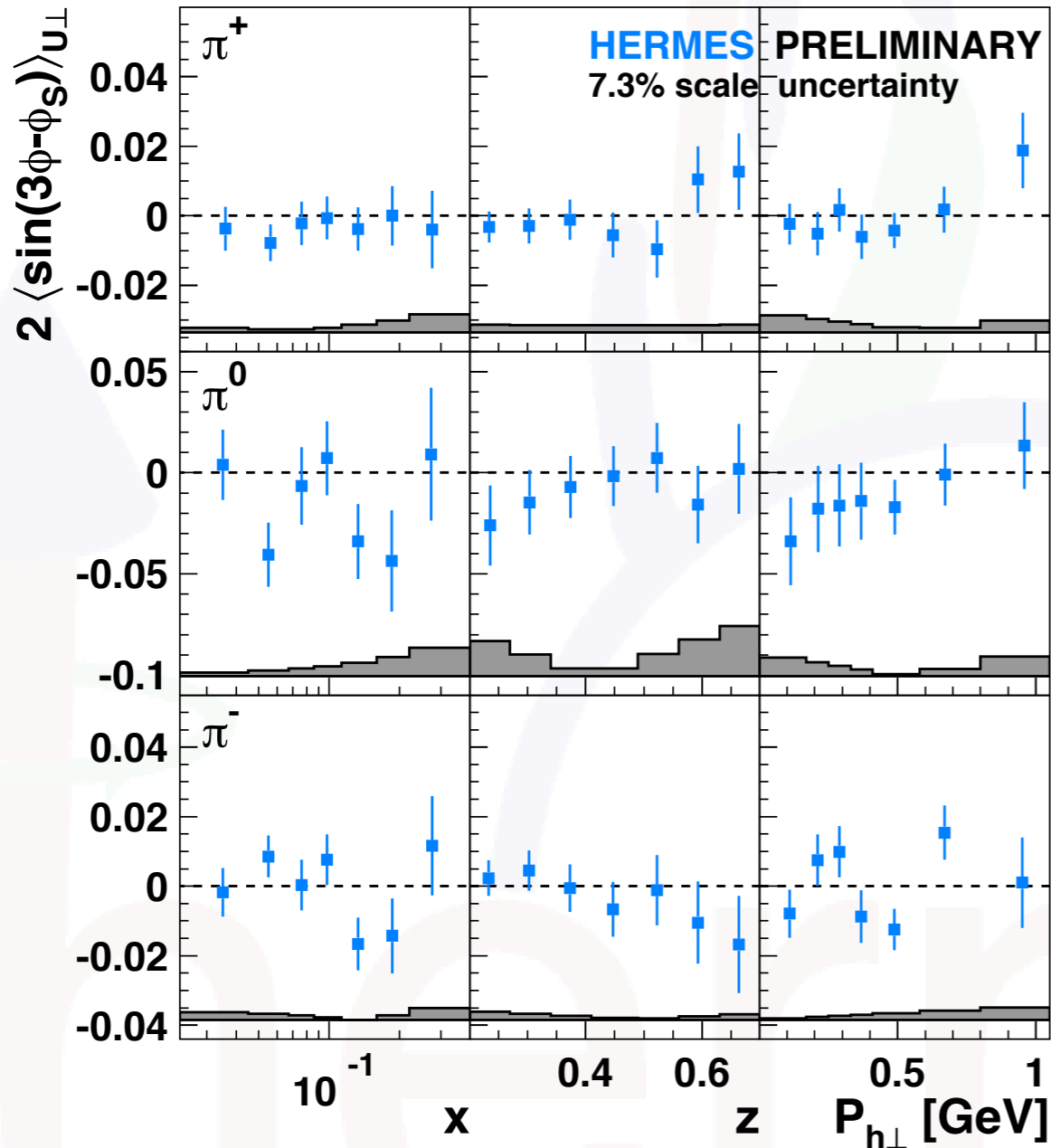
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L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp



- positive Collins SSA amplitude for positive kaons
 - consistent with zero for negative kaons and (anti)protons
- ➔ vanishing sea-quark transversity and baryon Collins effect?

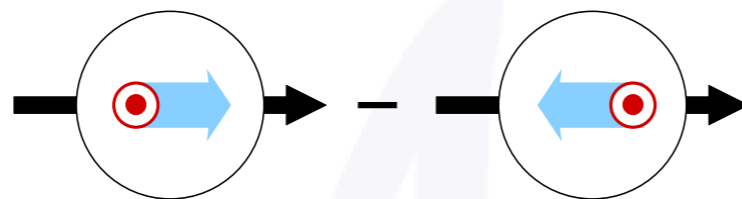
Pretzelosity?

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp



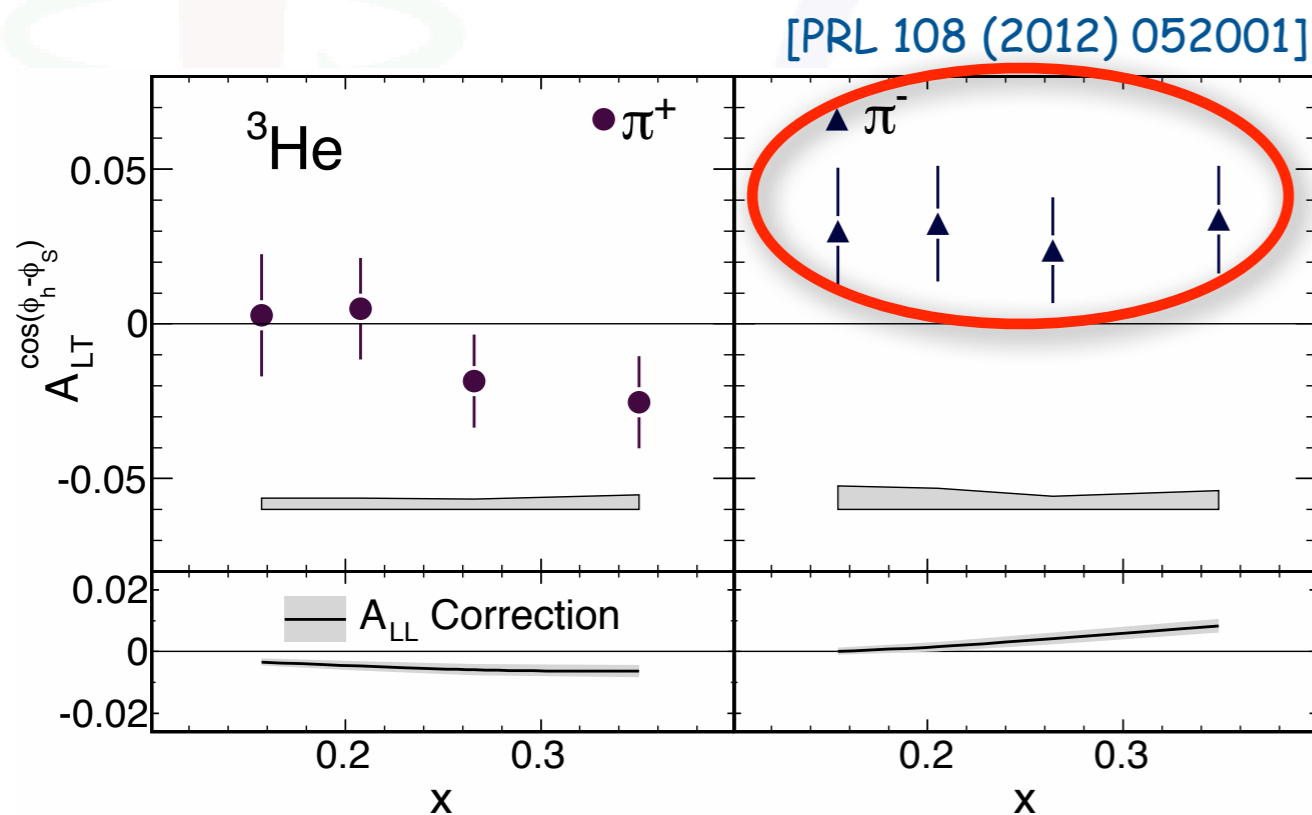
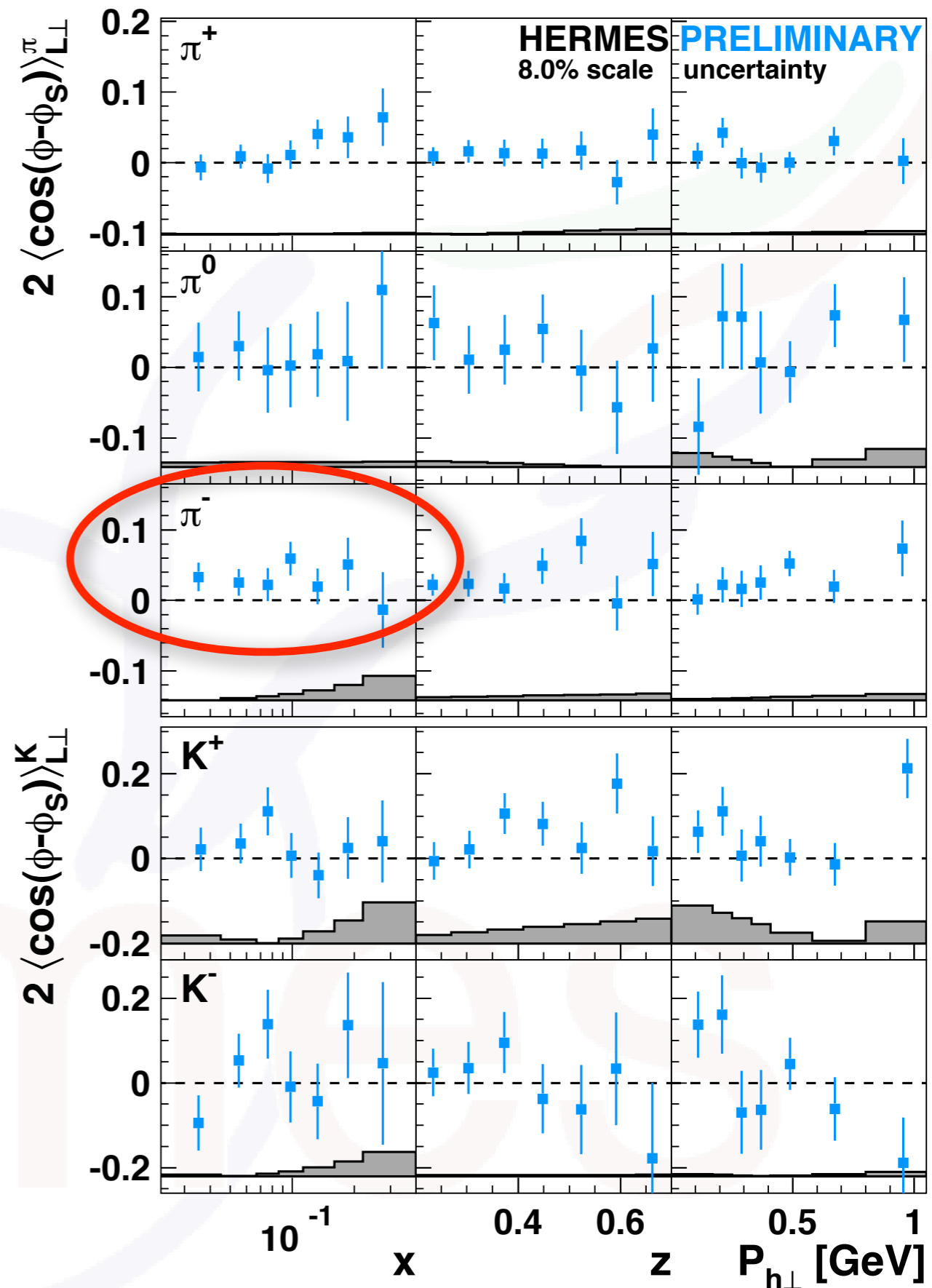
● consistent with zero; but suppressed by two powers of $P_{h\perp}$ (compared to, e.g., Collins)

	U	L	T
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L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

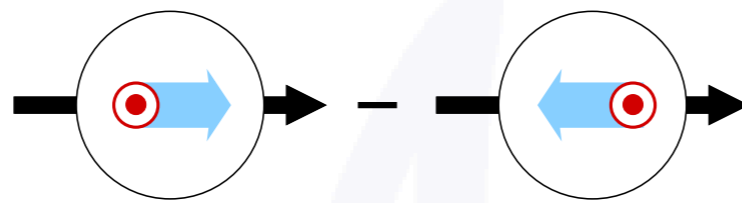


Worm-Gear

- chiral even
- first direct evidence for worm-gear g_{1T} on
 - ^3He target at JLab
 - H target at HERMES

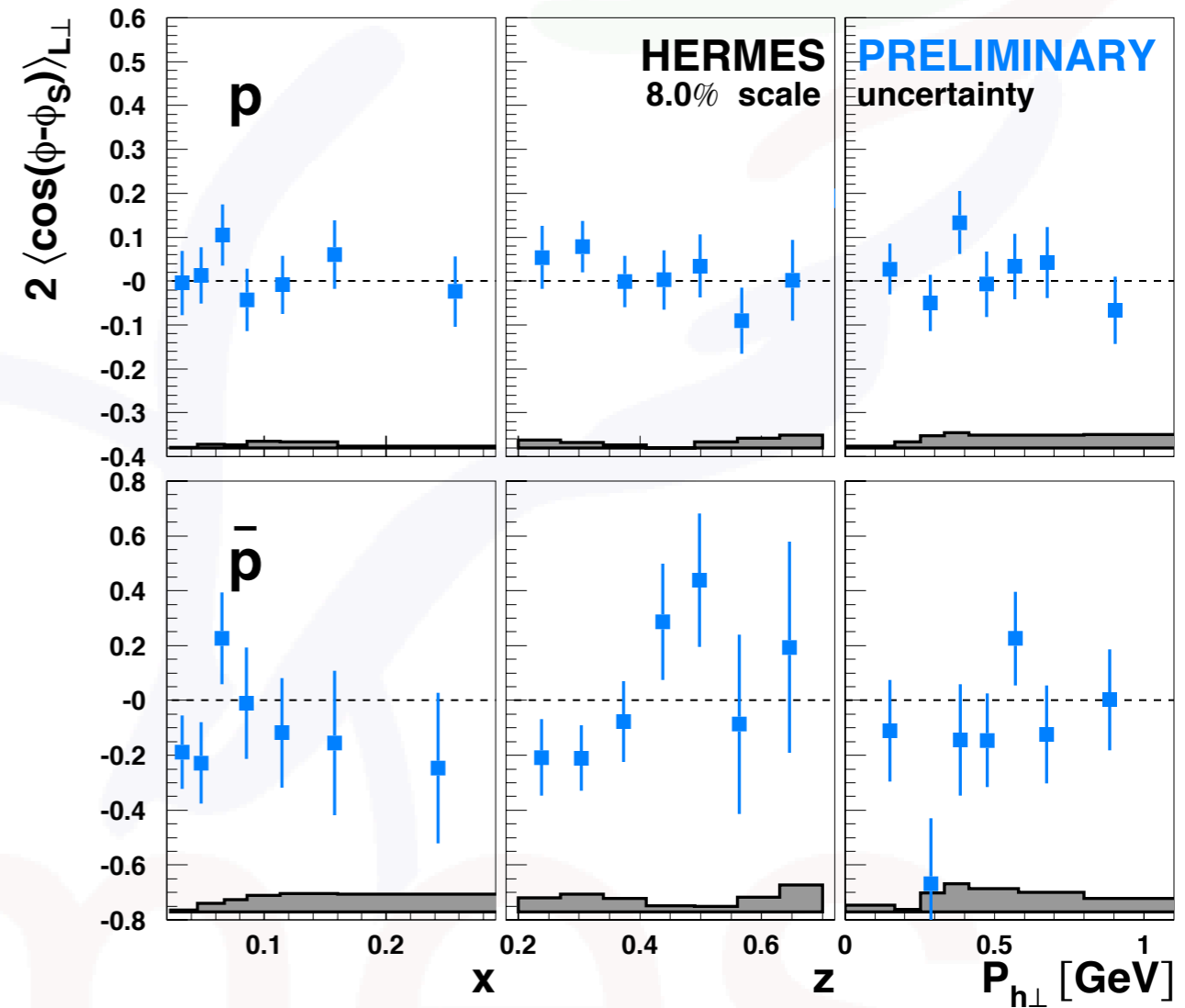


Worm-Gear



	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

- chiral even
- first direct evidence for worm-gear g_{1T} on
 - ^3He target at JLab
 - H target at HERMES
- results for protons and anti-protons consistent with zero

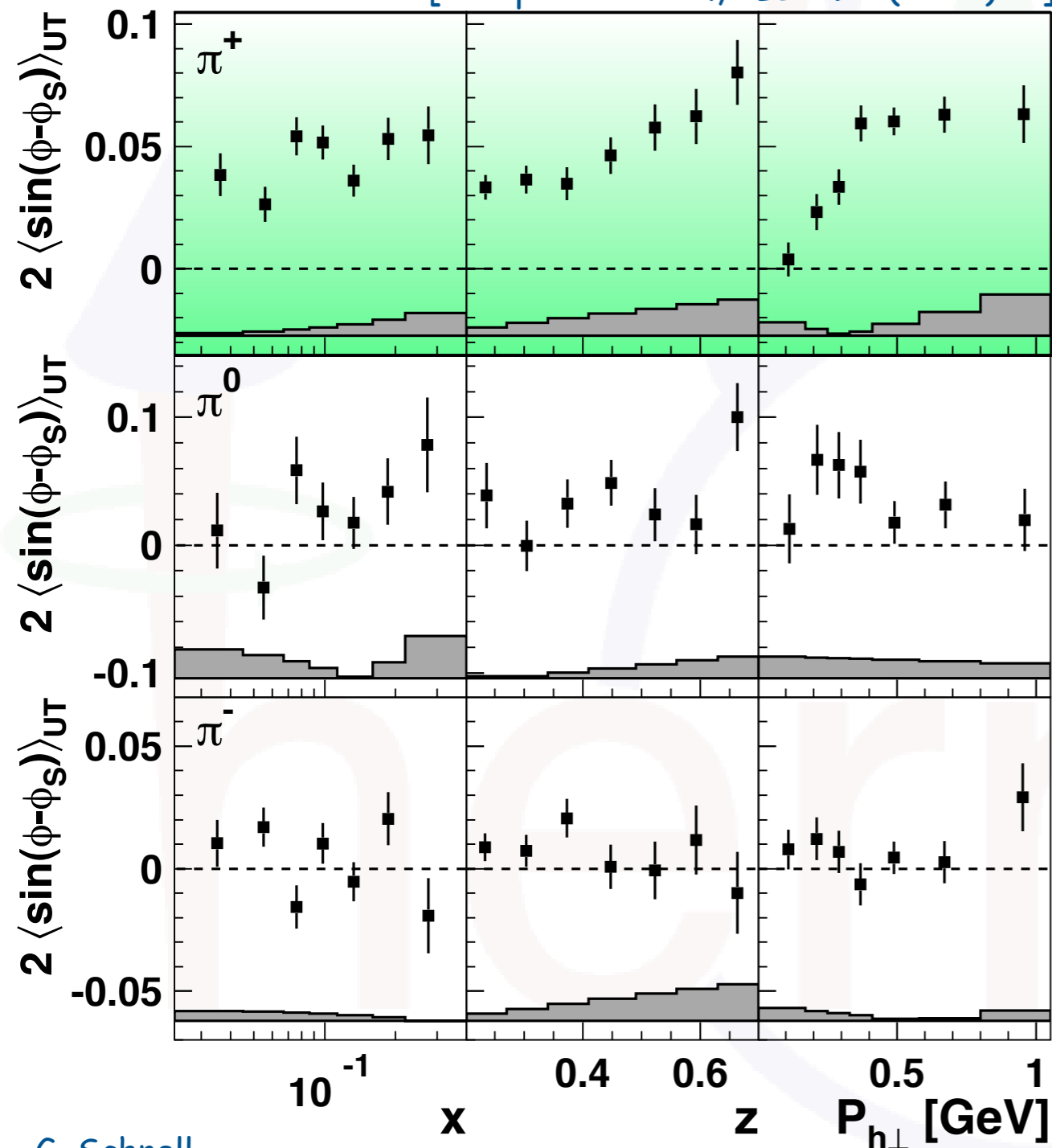


Sivers amplitudes for pions

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

$$2\langle \sin(\phi - \phi_S) \rangle_{UT} = - \frac{\sum_q e_q^2 f_{1T}^{\perp,q}(x, p_T^2) \otimes_{\mathcal{W}} D_1^q(z, k_T^2)}{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^q(z, k_T^2)}$$

[Airapetian et al., PLB 693 (2010) 11]

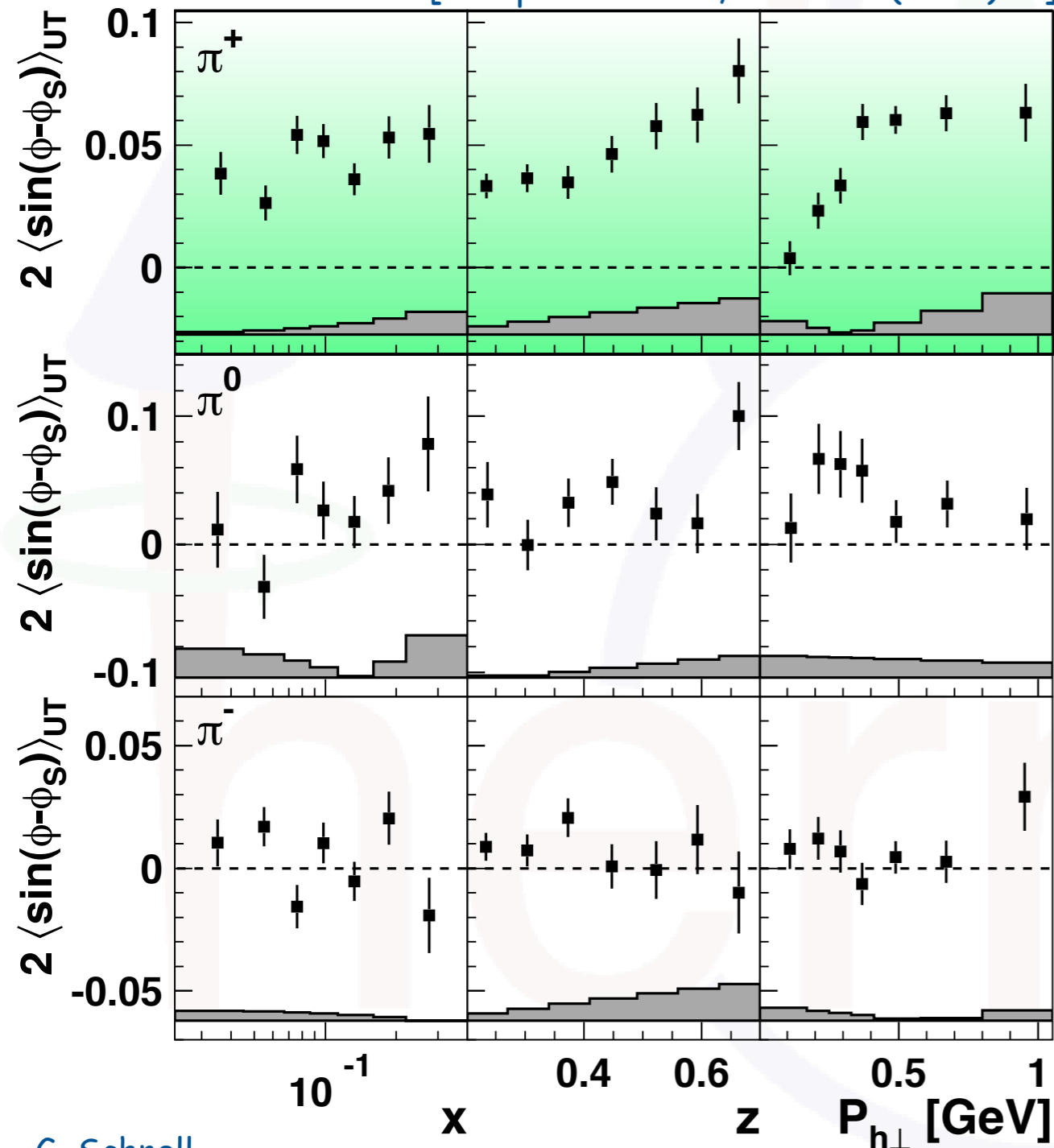


Sivers amplitudes for pions

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

$$2\langle \sin(\phi - \phi_S) \rangle_{UT} = - \frac{\sum_q e_q^2 f_{1T}^{\perp,q}(x, p_T^2) \otimes_{\mathcal{W}} D_1^q(z, k_T^2)}{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^q(z, k_T^2)}$$

[Airapetian et al., PLB 693 (2010) 11]



π^+ dominated by u-quark scattering:

$$\simeq - \frac{f_{1T}^{\perp,u}(x, p_T^2) \otimes_{\mathcal{W}} D_1^{u \rightarrow \pi^+}(z, k_T^2)}{f_1^u(x, p_T^2) \otimes D_1^{u \rightarrow \pi^+}(z, k_T^2)}$$

➡ u-quark Sivers DF < 0

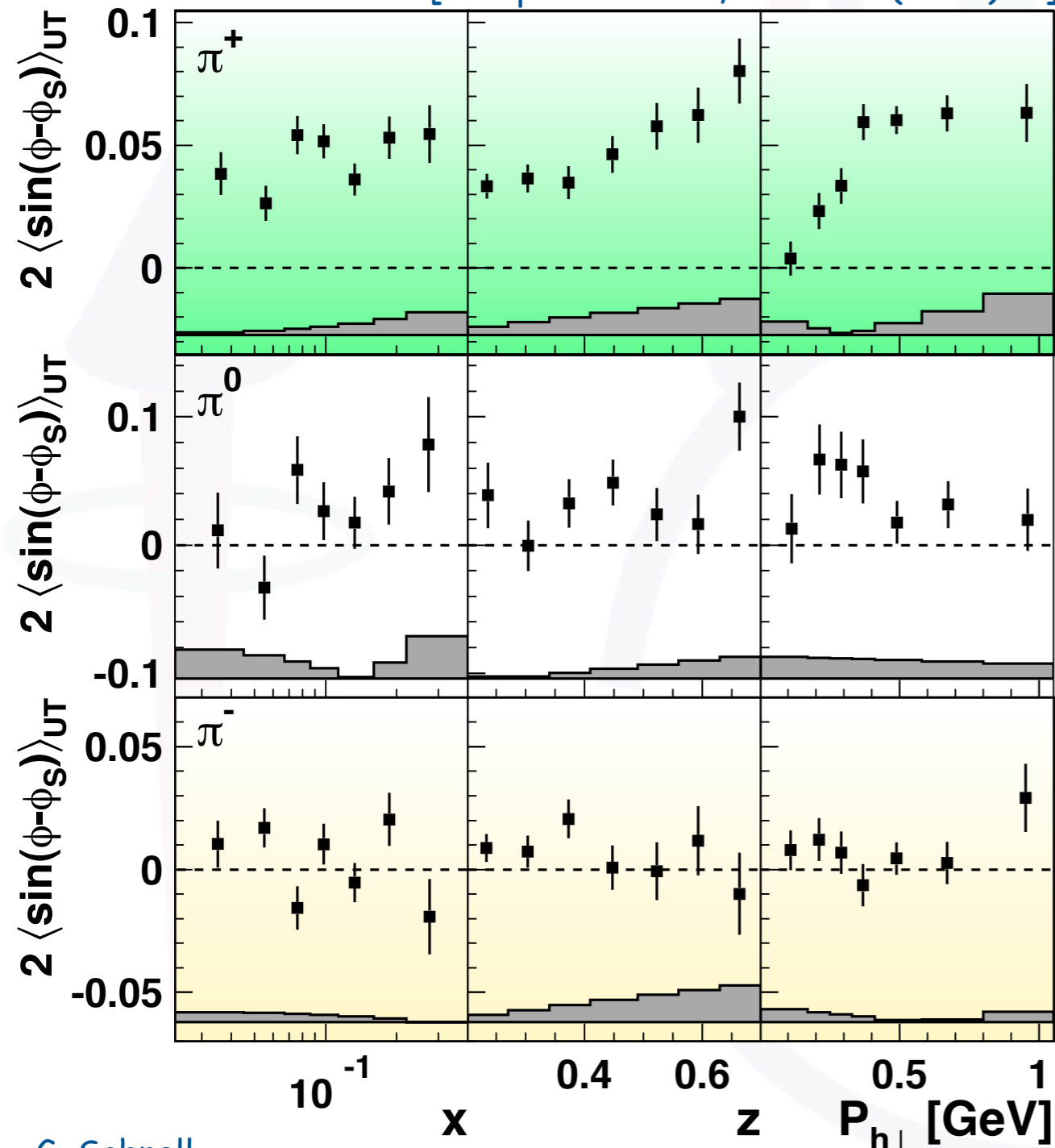
Sivers amplitudes for pions

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

$$2\langle \sin(\phi - \phi_S) \rangle_{UT} =$$

$$= - \frac{\sum_q e_q^2 f_{1T}^{\perp,q}(x, p_T^2) \otimes_{\mathcal{W}} D_1^q(z, k_T^2)}{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^q(z, k_T^2)}$$

[Airapetian et al., PLB 693 (2010) 11]



π^+ dominated by u-quark scattering:

$$\simeq - \frac{f_{1T}^{\perp,u}(x, p_T^2) \otimes_{\mathcal{W}} D_1^{u \rightarrow \pi^+}(z, k_T^2)}{f_1^u(x, p_T^2) \otimes D_1^{u \rightarrow \pi^+}(z, k_T^2)}$$

➡ u-quark Sivers DF < 0

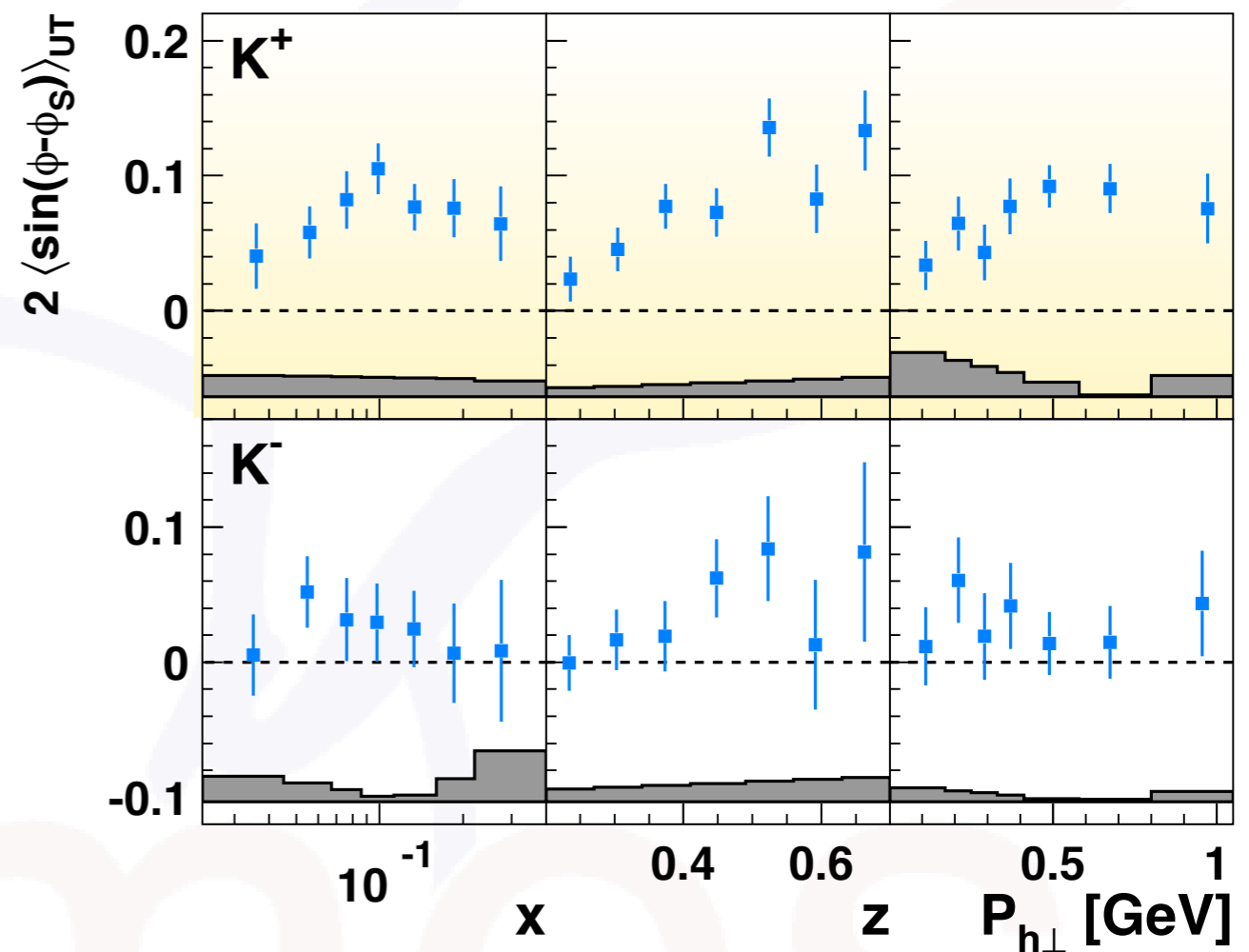
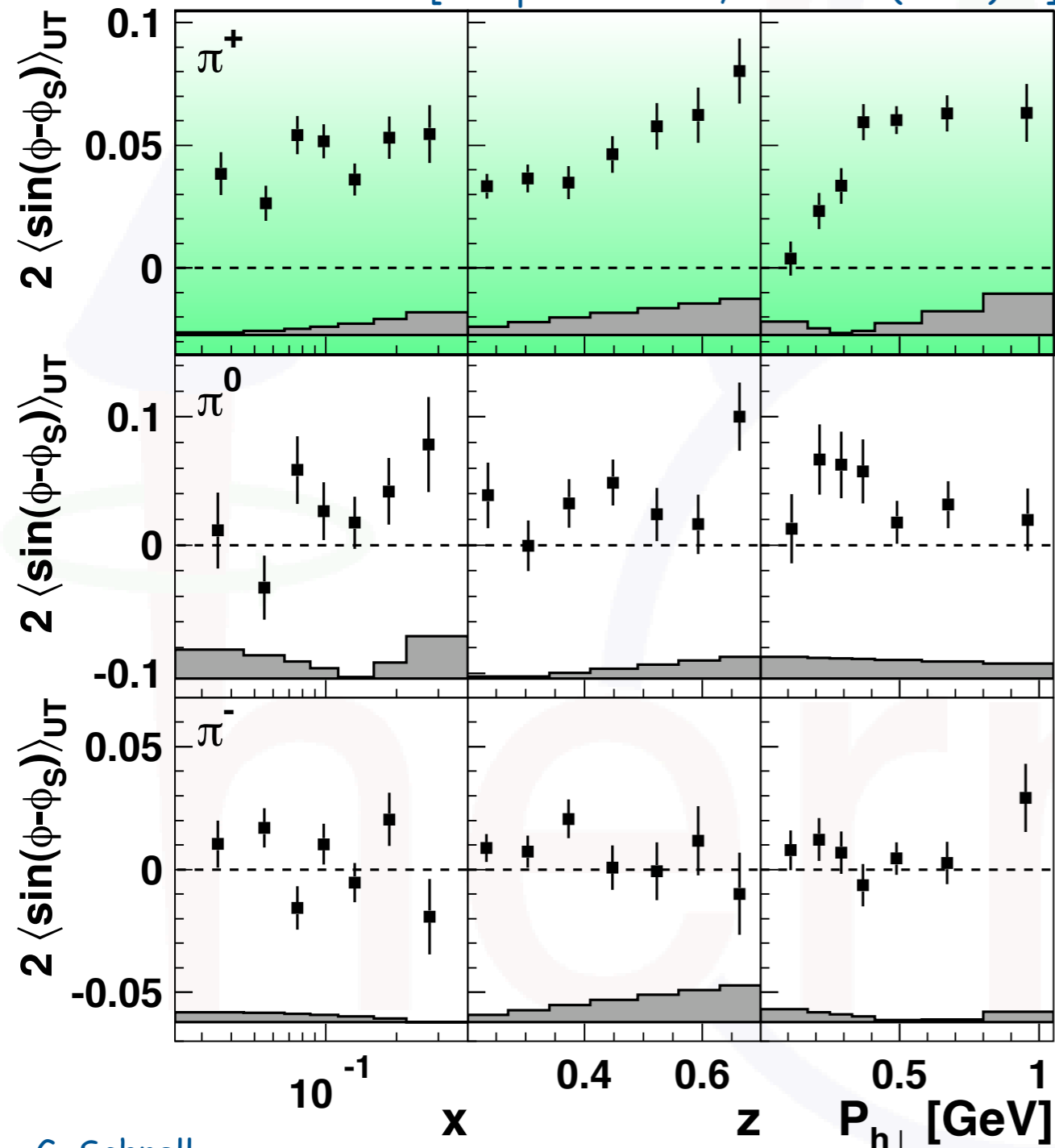
➡ d-quark Sivers DF > 0
(cancellation for π^-)

Sivers amplitudes for mesons

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

$$2\langle \sin(\phi - \phi_S) \rangle_{UT} = - \frac{\sum_q e_q^2 f_{1T}^{\perp,q}(x, p_T^2) \otimes_{\mathcal{W}} D_1^q(z, k_T^2)}{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^q(z, k_T^2)}$$

[Airapetian et al., PLB 693 (2010) 11]



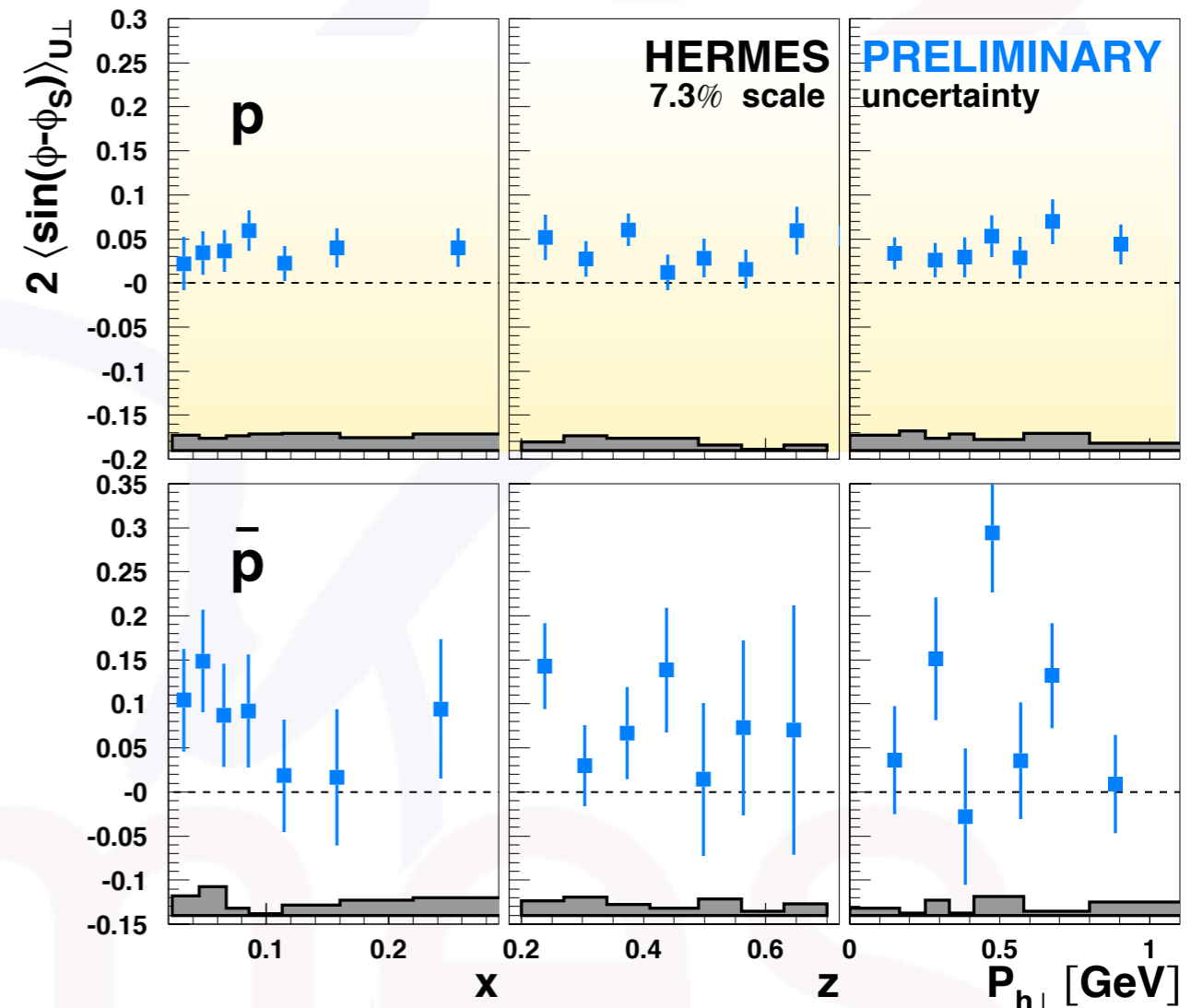
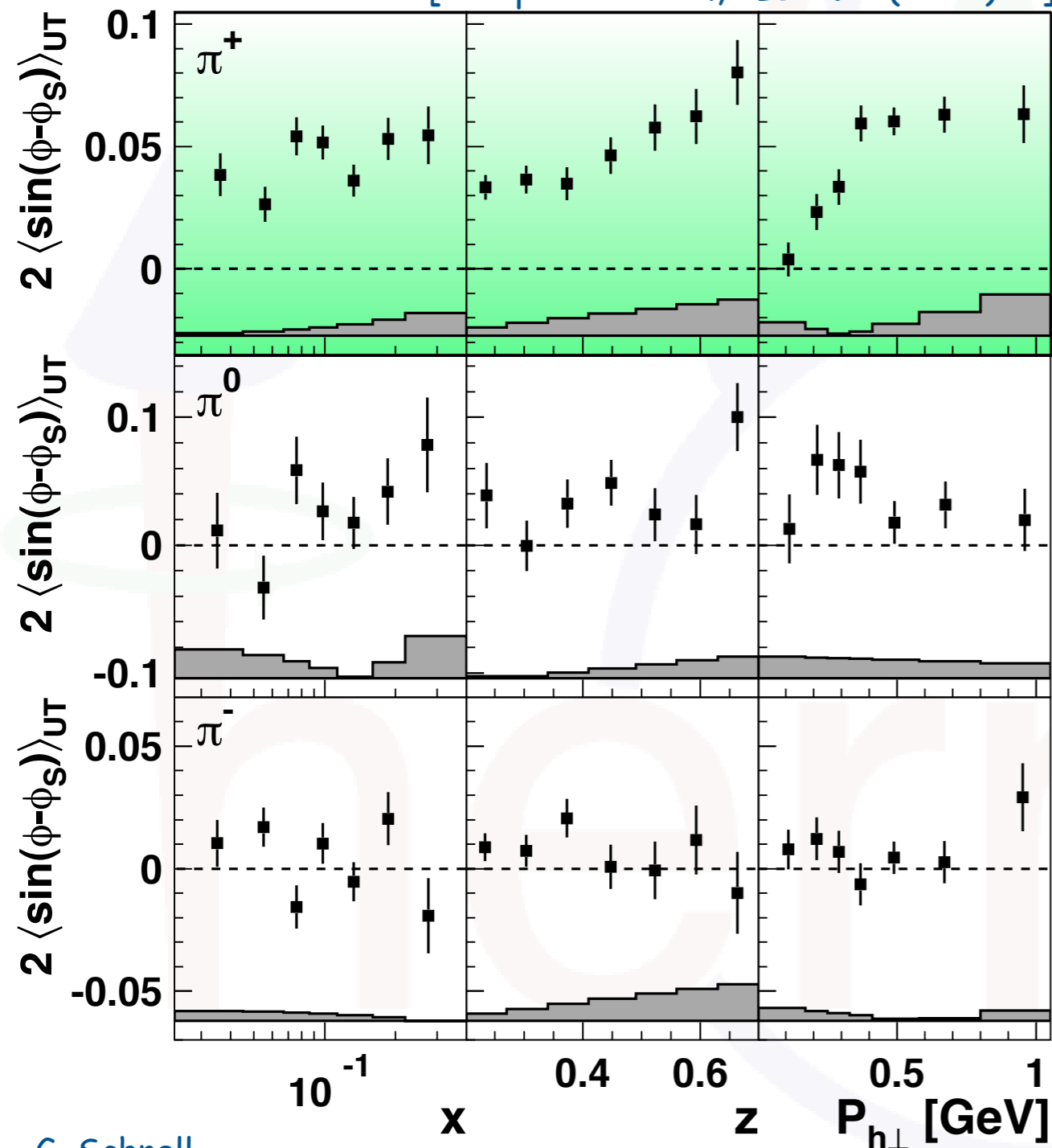
➡ larger amplitudes for positive kaons vs. pions

Sivers amplitudes for baryons

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

$$2\langle \sin(\phi - \phi_S) \rangle_{UT} = - \frac{\sum_q e_q^2 f_{1T}^{\perp,q}(x, p_T^2) \otimes_{\mathcal{W}} D_1^q(z, k_T^2)}{\sum_q e_q^2 f_1^q(x, p_T^2) \otimes D_1^q(z, k_T^2)}$$

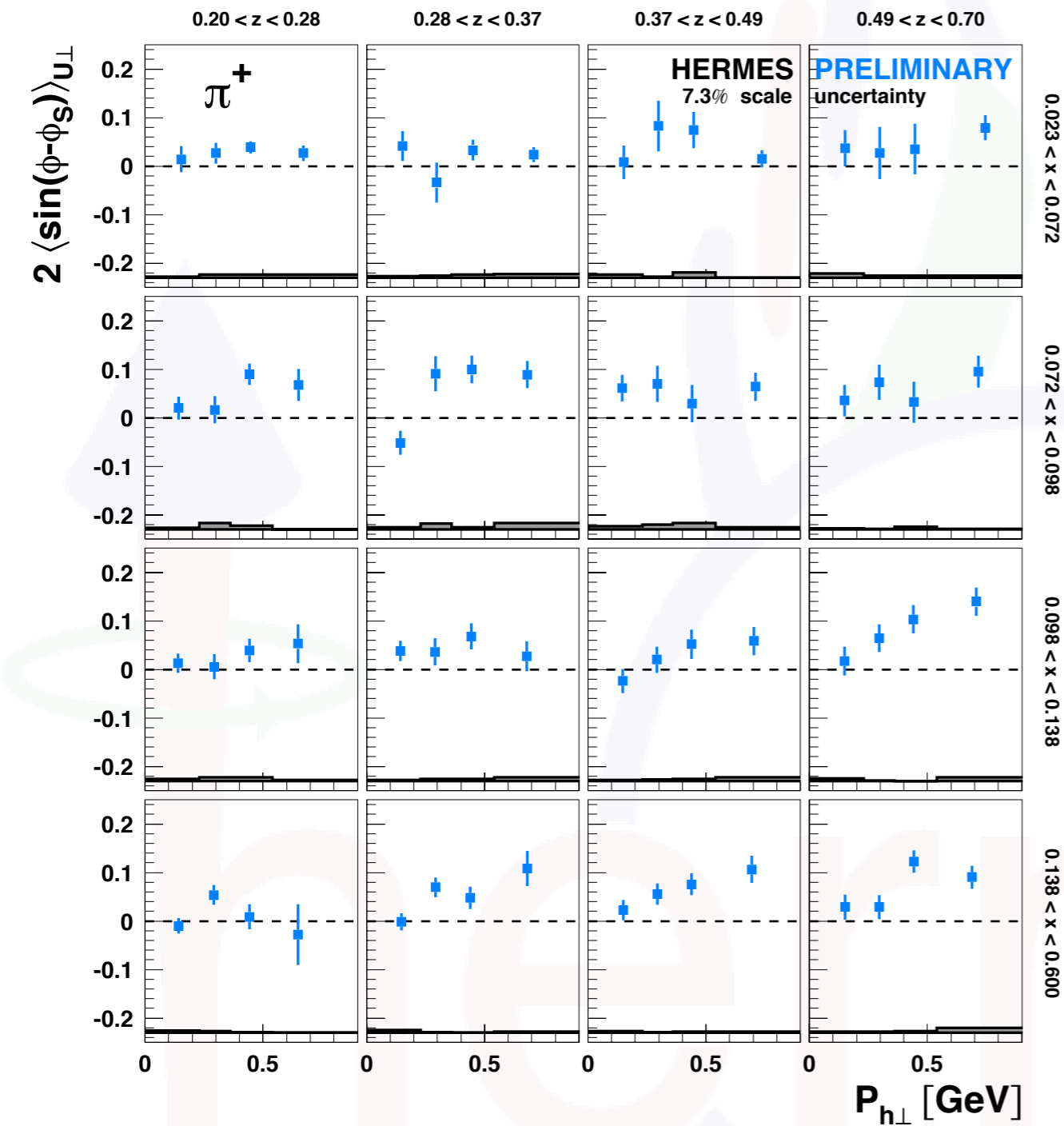
[Airapetian et al., PLB 693 (2010) 11]



similar amplitudes for positive pions and protons \rightarrow u-quark dominance (and not a FF effect)?

Sivers amplitudes - 3d binning

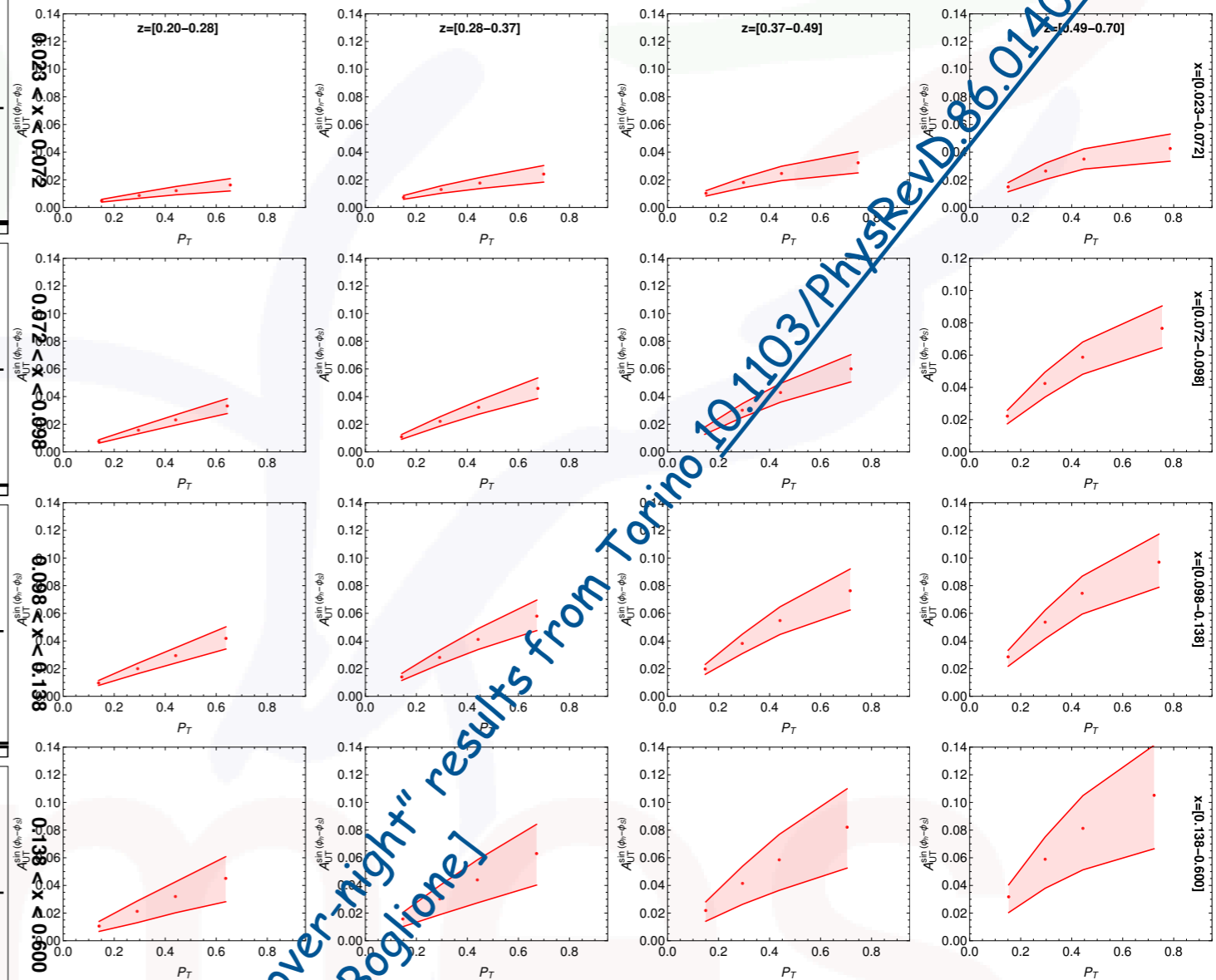
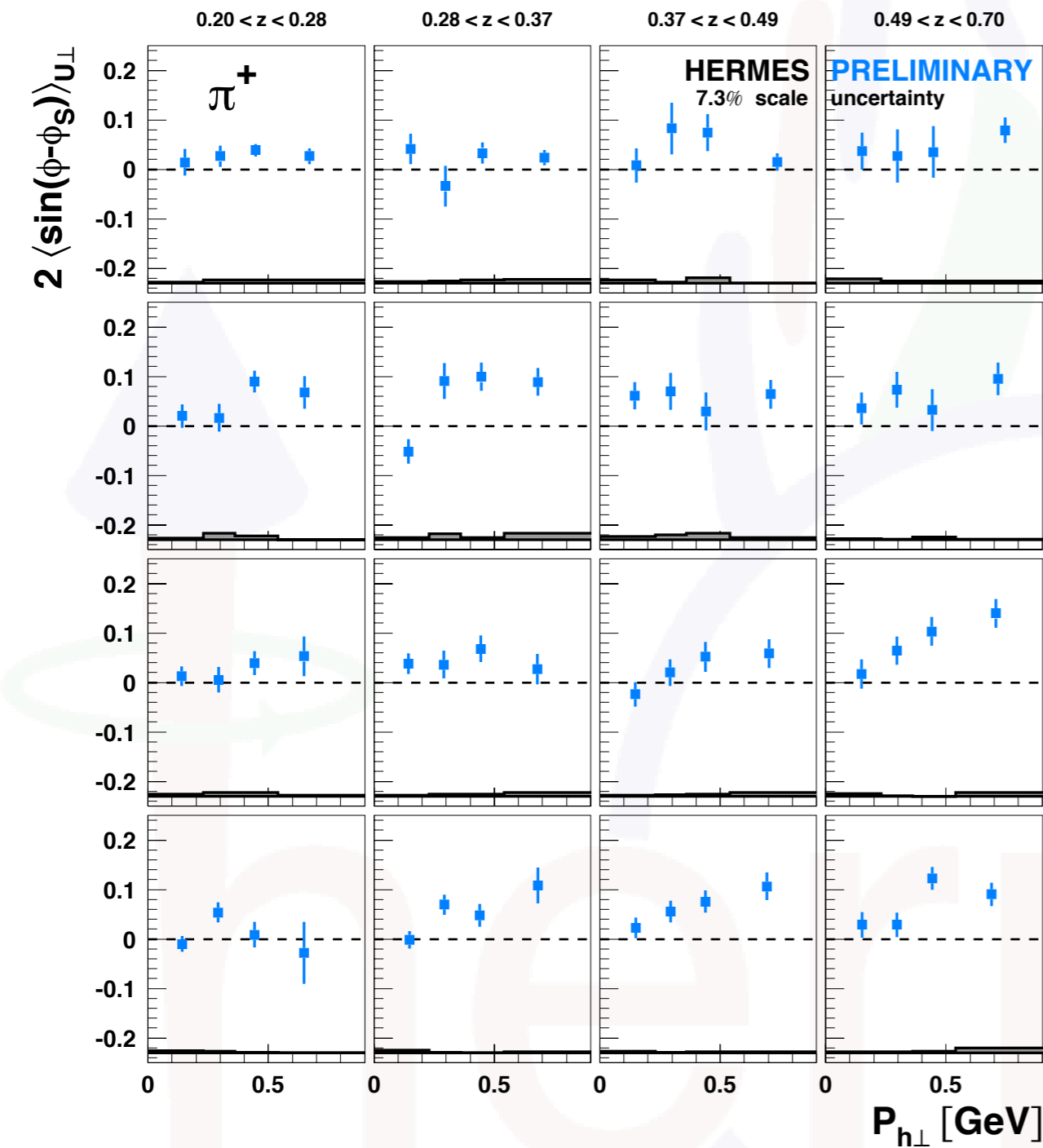
	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp



- 3d analysis: 4x4x4 bins in (x, z, P_{h_\perp})
- disentangle correlations
- isolate phase-space region with strong signal strength
- allows more detailed comparison with calculations (e.g., unofficial "over-night" results from Torino [10.1103/PhysRevD.86.014028](https://arxiv.org/abs/10.1103/PhysRevD.86.014028) fit - courtesy M. Boglione)

Sivers amplitudes - 3d binning

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

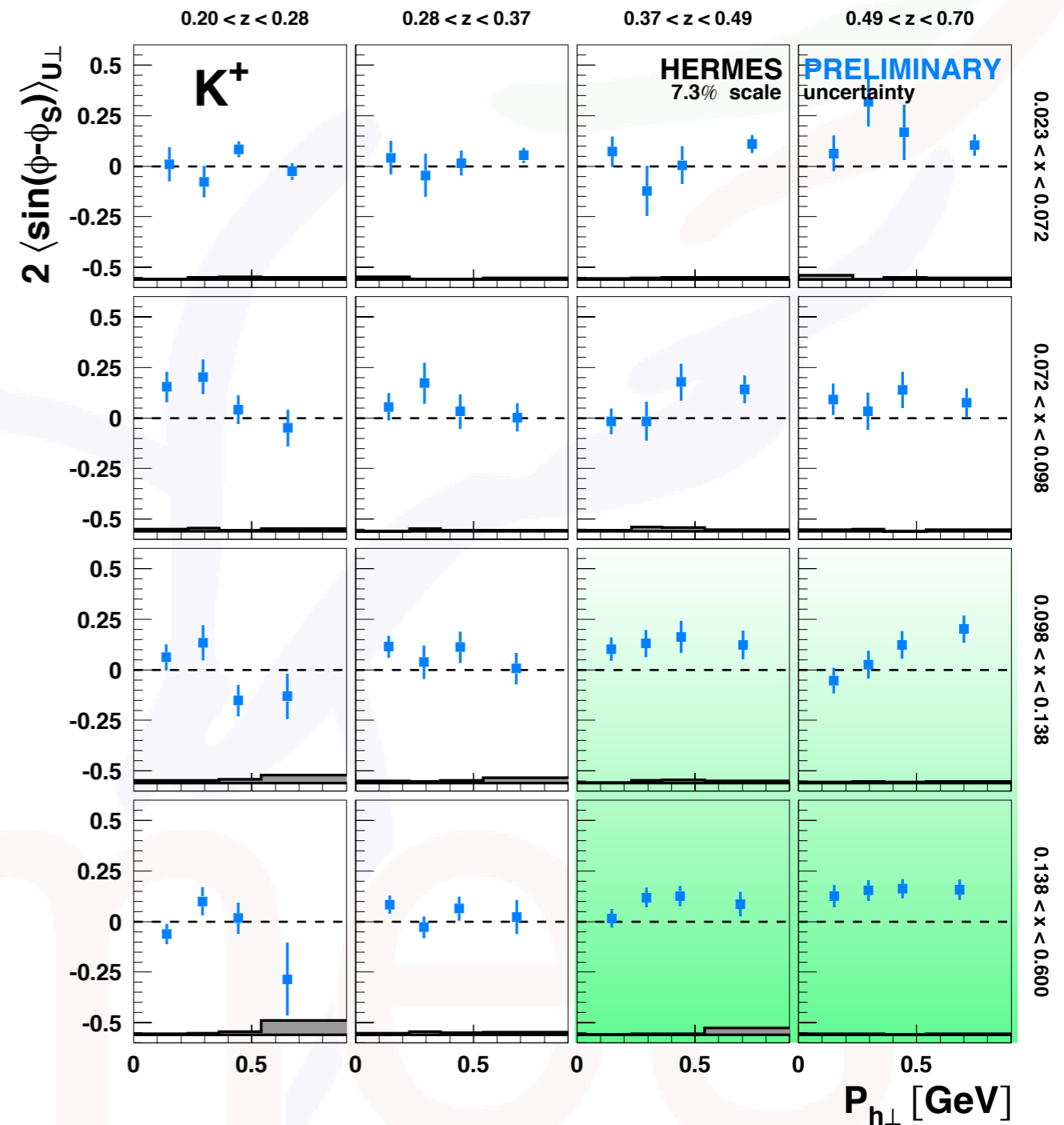


[unofficial "over right" results from Torino 10.1103/PhysRevD.86.014028 fit
- courtesy M. Boglione]

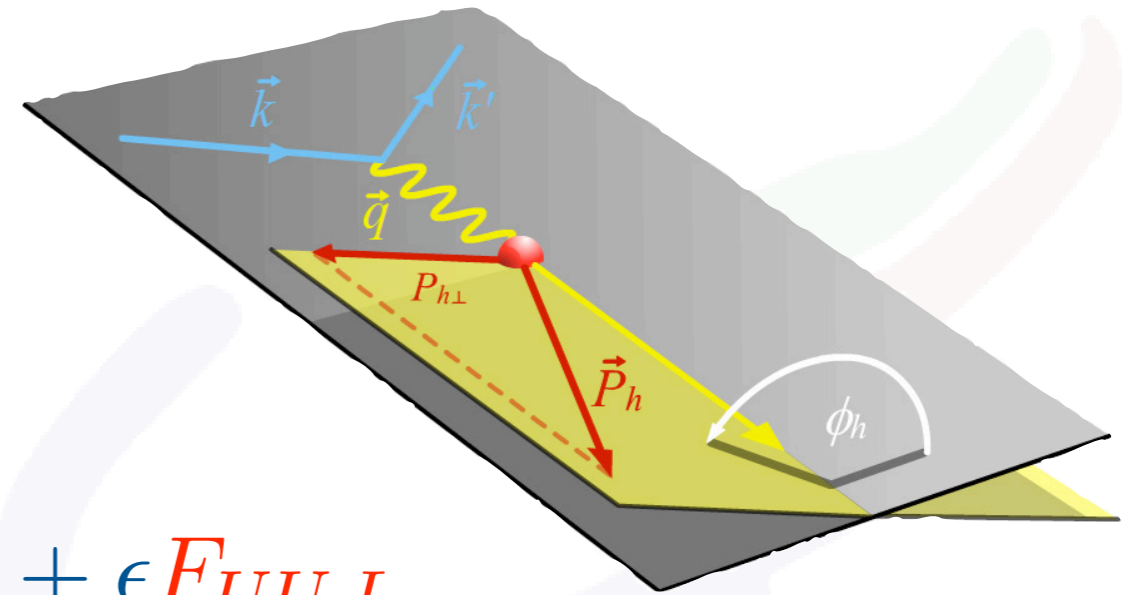
Sivers amplitudes - 3d binning

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

- large K^+ amplitudes $O(20\%)$ seen at large values of (x, z)
- region of purest "u-quark probe"



Cross section without polarization



$$\frac{d^5\sigma}{dx dy dz d\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \{F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1-\epsilon)} F_{UU}^{\cos\phi_h} \cos\phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h\}$$

$$F_{XY,Z} = F_{XY,Z}(x, y, z, P_{h\perp})$$

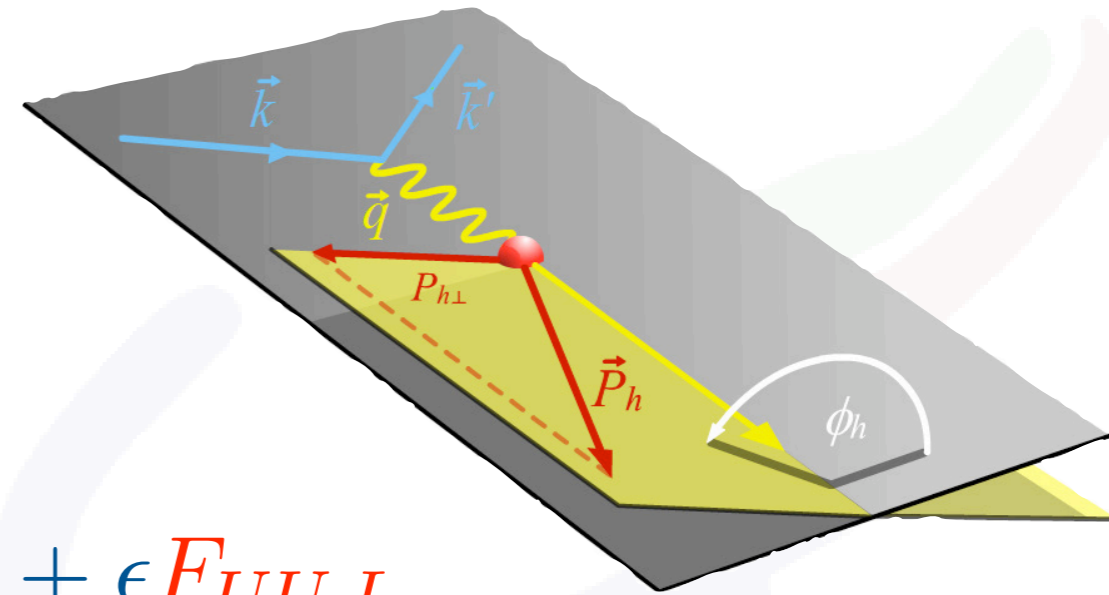
target polarization \downarrow
 \uparrow beam polarization \uparrow virtual-photon polarization

$$\gamma = \frac{2Mx}{Q}$$

$$\epsilon = \frac{1 - y - \frac{1}{4}\gamma^2 y^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}\gamma^2 y^2}$$

[see, e.g., Bacchetta et al., JHEP 0702 (2007) 093]

Cross section without polarization



$$\frac{d^5\sigma}{dx dy dz d\phi_h dP_{h\perp}^2} \propto \left(1 + \frac{\gamma^2}{2x}\right) \{F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1-\epsilon)} F_{UU}^{\cos\phi_h} \cos\phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h\}$$

leading twist
 $F_{UU}^{\cos 2\phi_h} \propto C \left[\frac{2(\hat{P}_{h\perp} \cdot \vec{k}_T)(\hat{P}_{h\perp} \cdot \vec{p}_T) - \vec{k}_T \cdot \vec{p}_T}{MM_h} h_1^\perp H_1^\perp \right]$

next to leading twist
 $F_{UU}^{\cos\phi_h} \propto \frac{2M}{Q} C \left[\frac{\hat{P}_{h\perp} \cdot \vec{p}_T}{M_h} x h_1^\perp H_1^\perp - \frac{\hat{P}_{h\perp} \cdot \vec{k}_T}{M} x f_1 D_1 + \dots \right]$

BOER-MULDERS EFFECT (points to the blue box in the leading twist equation)

CAHN EFFECT (points to the red box in the next to leading twist equation)

Interaction dependent terms neglected (points to the ellipsis in the next to leading twist equation)

$$\gamma = \frac{2Mx}{Q}$$

$$\epsilon = \frac{1 - y - \frac{1}{4}\gamma^2 y^2}{1 - y + \frac{1}{2}y^2 + \frac{1}{4}\gamma^2 y^2}$$

[see, e.g., Bacchetta et al., JHEP 0702 (2007) 093]

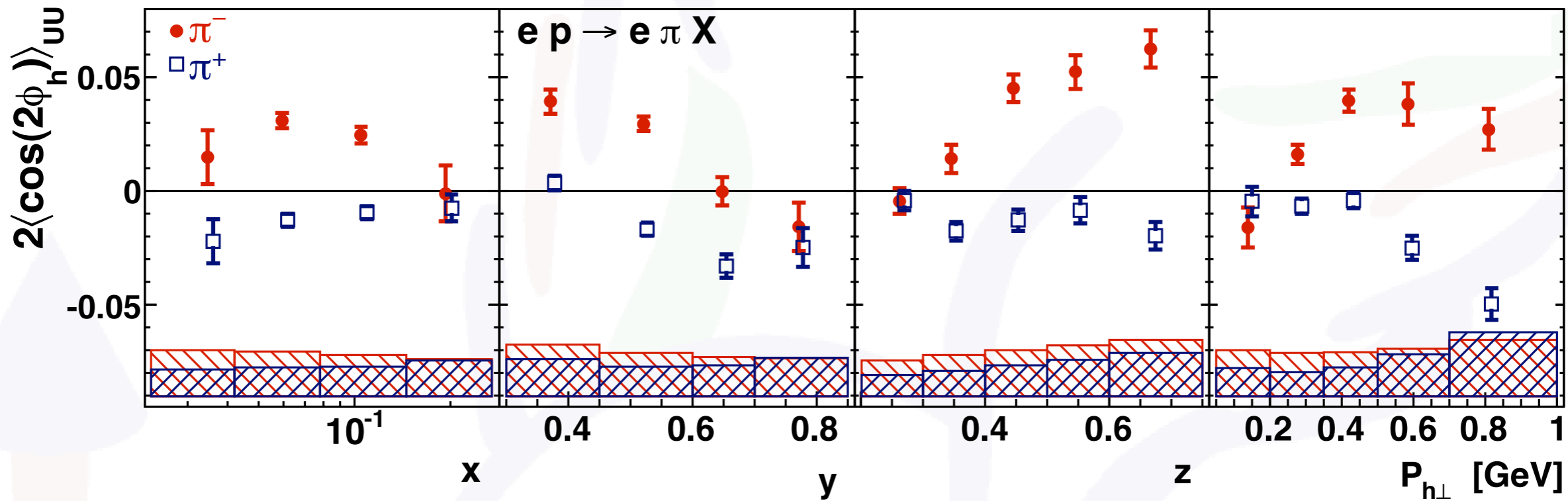
(Implicit sum over quark flavours)

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp



signs of Boer-Mulders

[Airapetian et al., PRD 87 (2013) 012010]

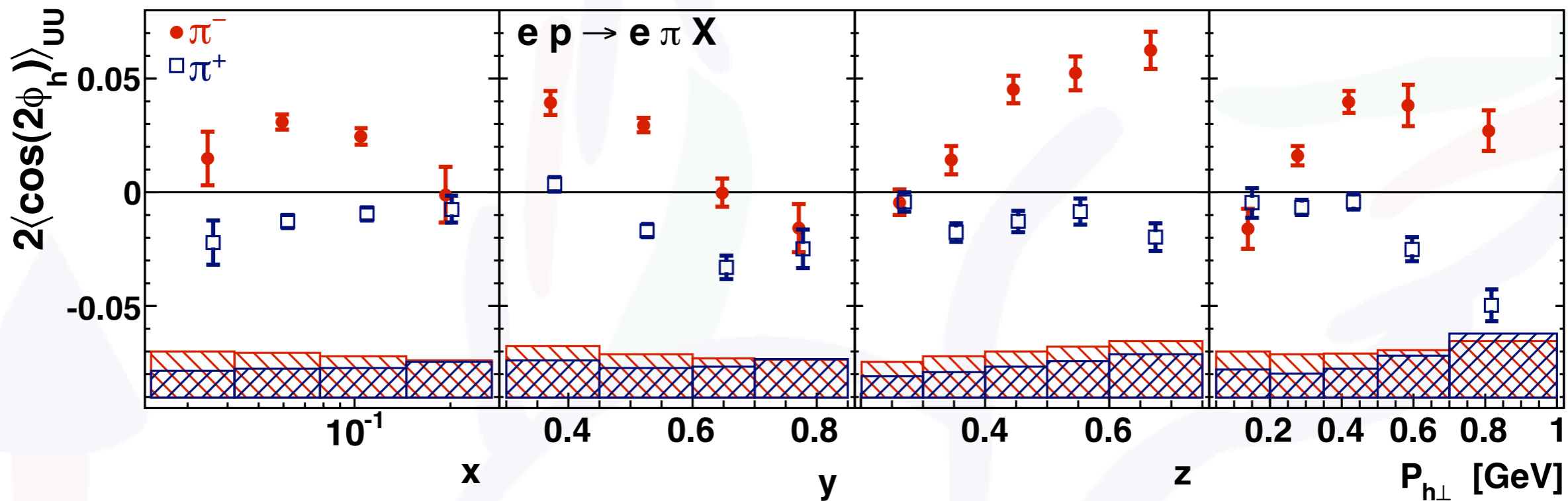


	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp



signs of Boer-Mulders

[Airapetian et al., PRD 87 (2013) 012010]



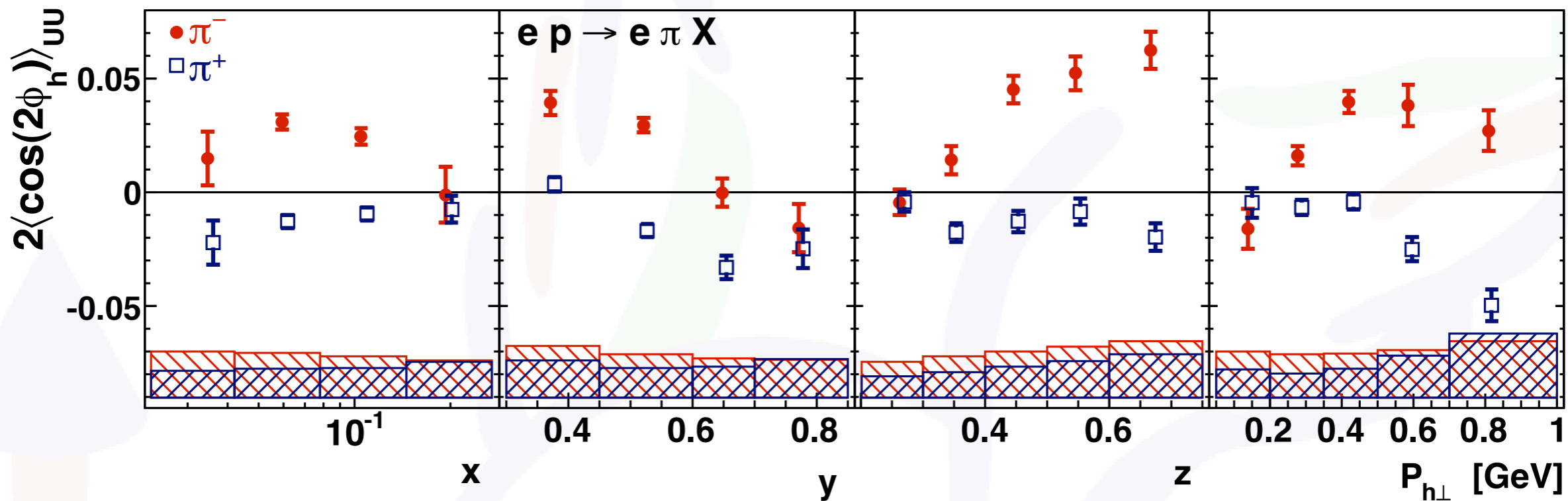
● modulations are not zero!

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp



signs of Boer-Mulders

[Airapetian et al., PRD 87 (2013) 012010]



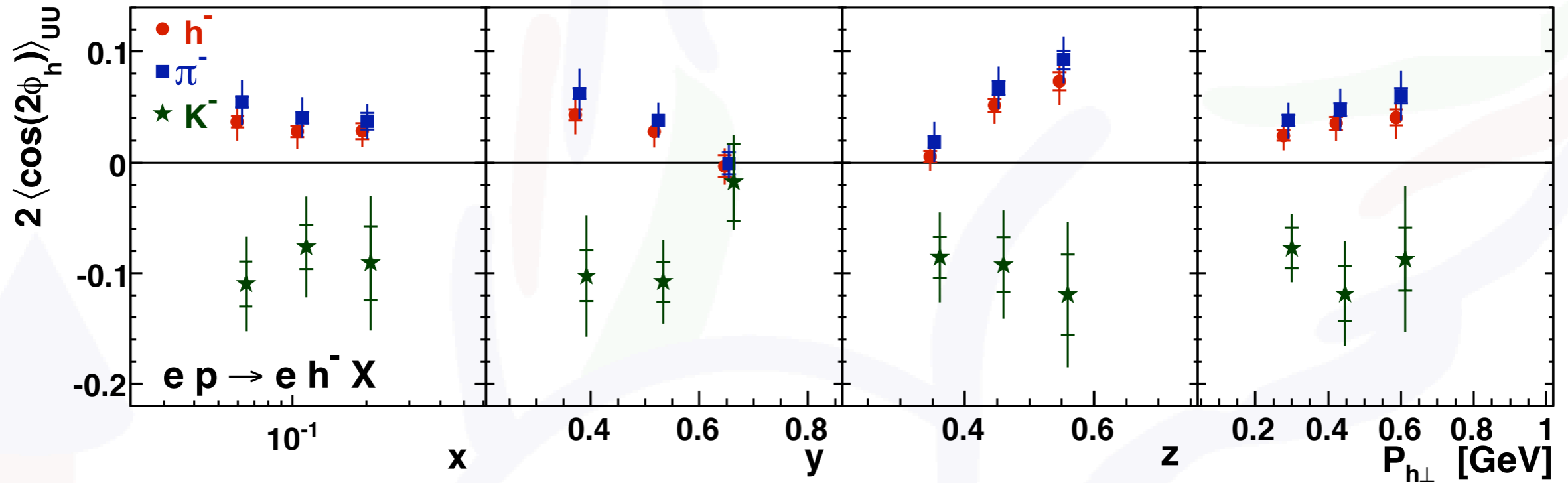
- modulations are not zero!
- opposite sign for charged pions with larger magnitude for π^-

signs of Boer-Mulders



	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp

[Airapetian et al., PRD 87 (2013) 012010]



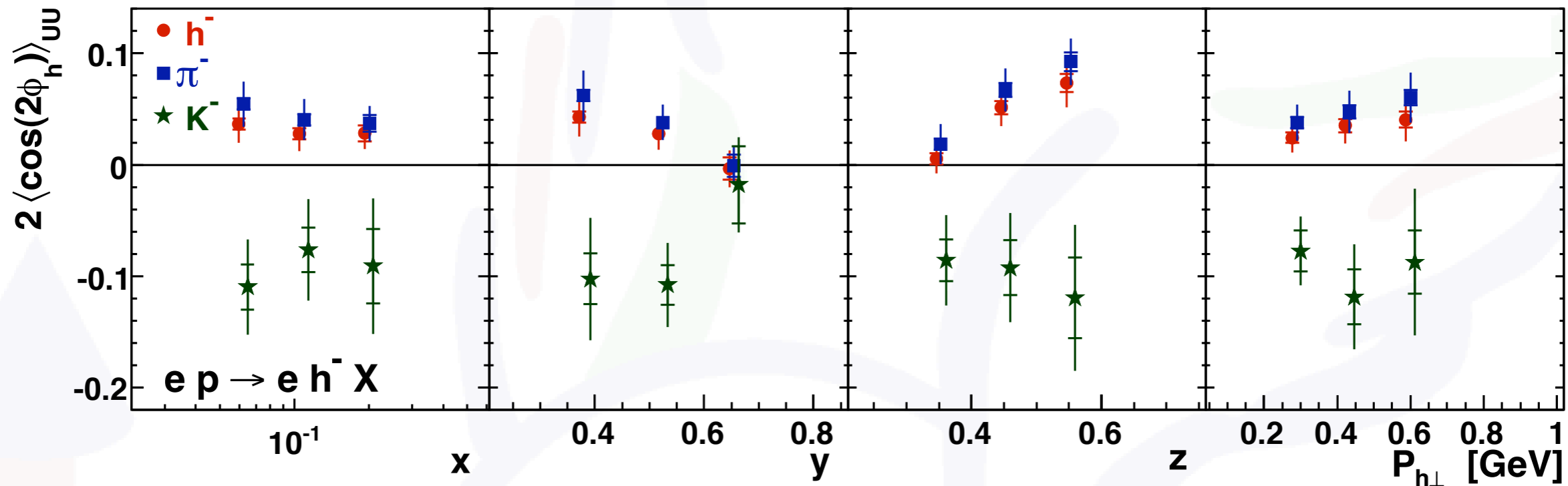
- modulations are not zero!
- opposite sign for charged pions with larger magnitude for π^-
- intriguing behavior for kaons

	U	L	T
U	f_1		h_1^\perp
L		g_{1L}	h_{1L}^\perp
T	f_{1T}^\perp	g_{1T}	h_1, h_{1T}^\perp



signs of Boer-Mulders

[Airapetian et al., PRD 87 (2013) 012010]



- modulations are not zero!
- opposite sign for charged pions with larger magnitude for π^-
- intriguing behavior for kaons
- available in multidimensional binning:
<http://www-hermes.desy.de/cosnphi/>

Cahn effect?

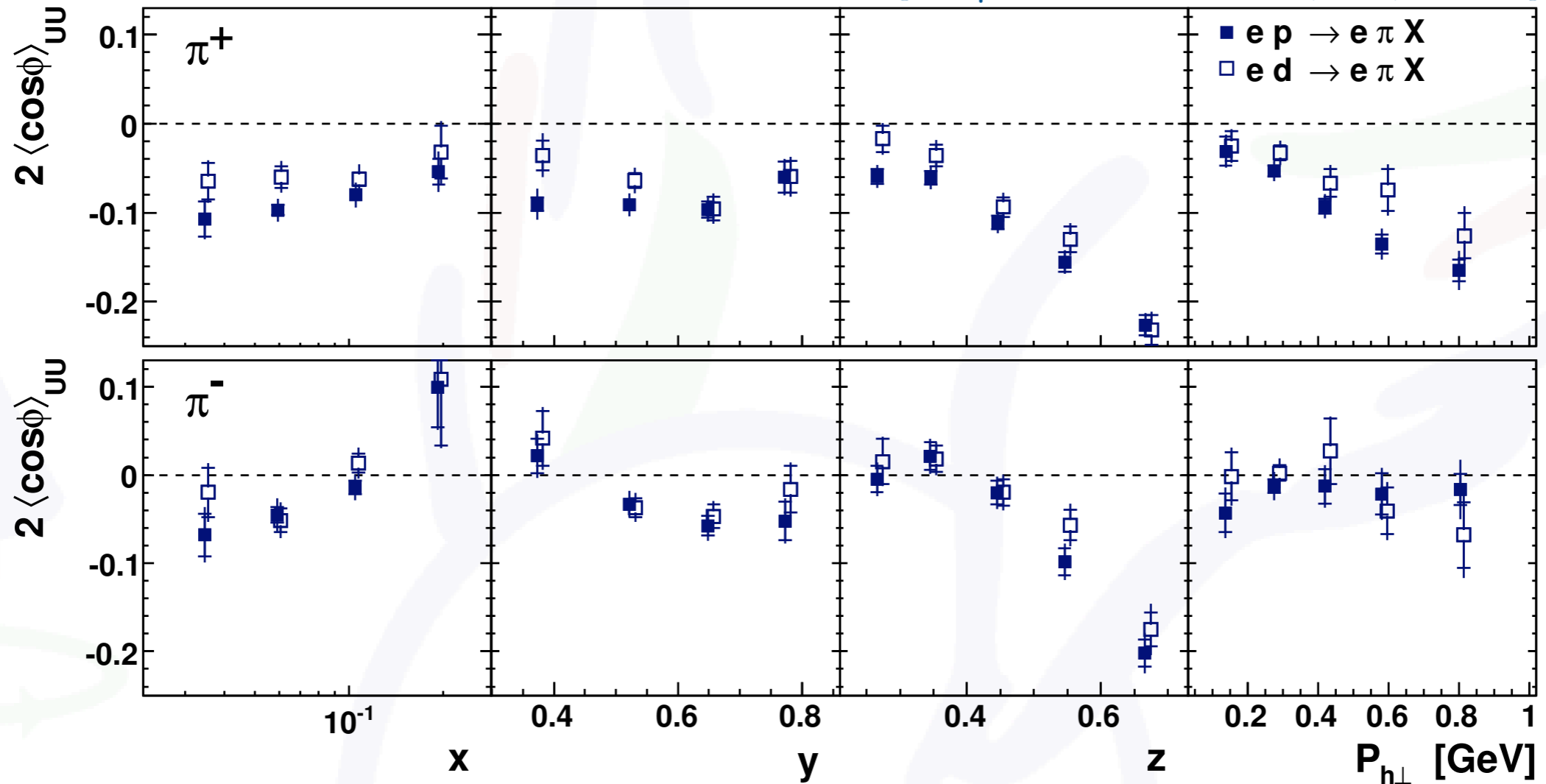
next to leading twist

$$F_{UU}^{\cos\phi_h} \propto \frac{2M}{Q} C$$

$$C \left[\underbrace{-\frac{\hat{P}_{h\perp} \cdot \vec{p}_T}{M_h} x h_1^\perp H_1^\perp}_{\text{BOER-MULDERS EFFECT}} - \underbrace{\frac{\hat{P}_{h\perp} \cdot \vec{k}_T}{M} x f_1 D_1}_{\text{CAHN EFFECT}} + \dots \right]$$

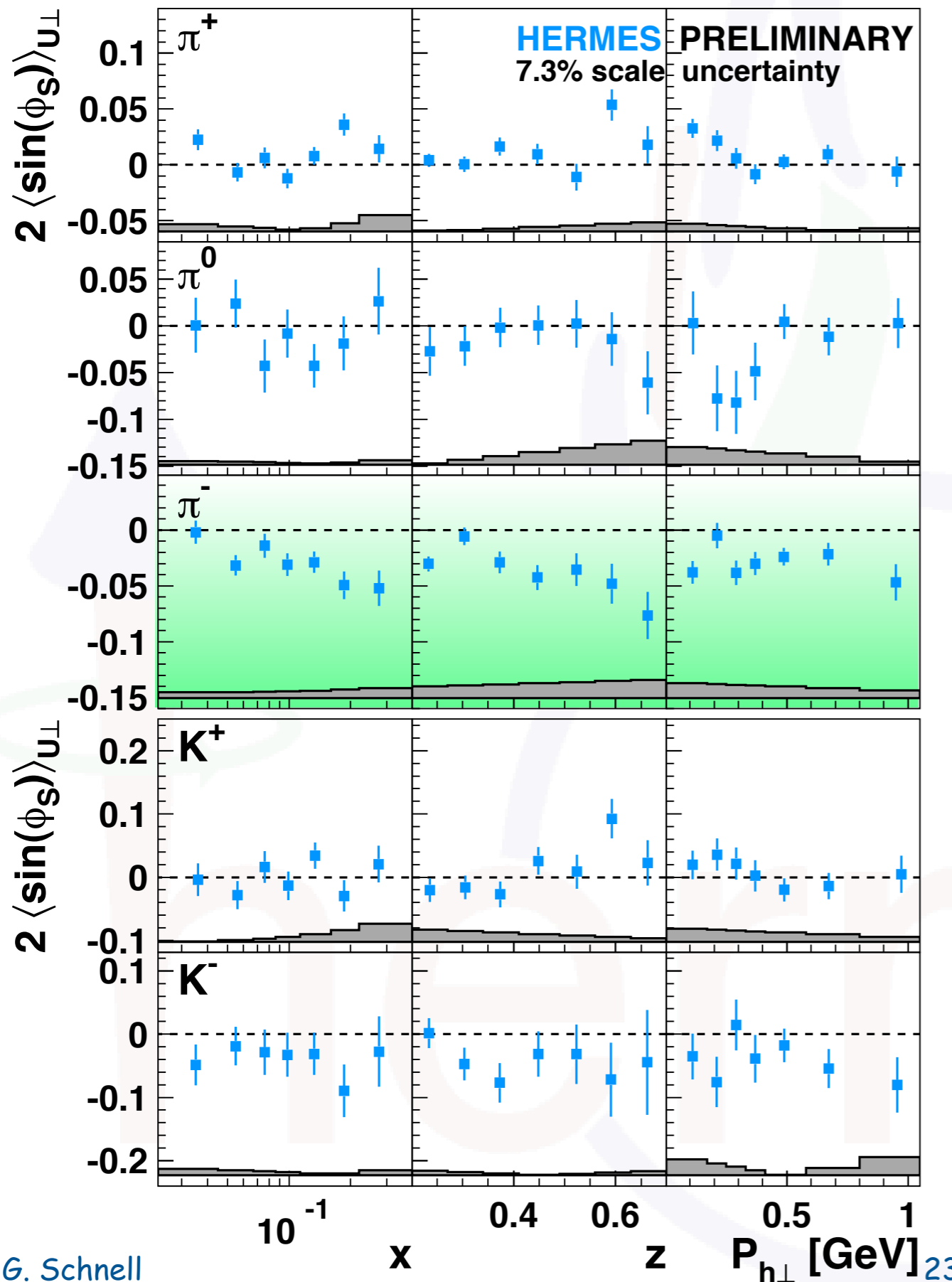
Interaction dependent terms neglected

[Airapetian et al., PRD 87 (2013) 012010]

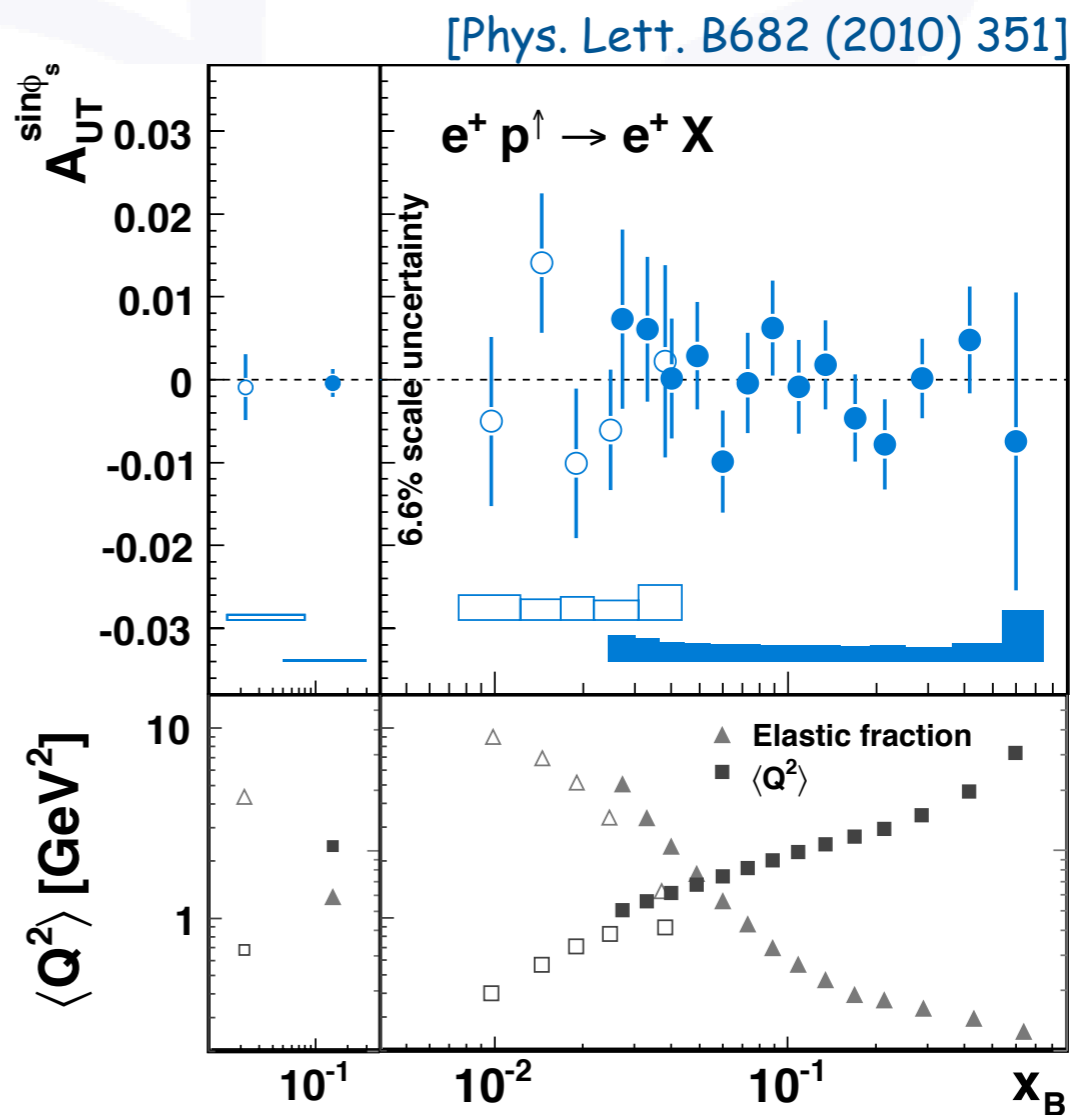


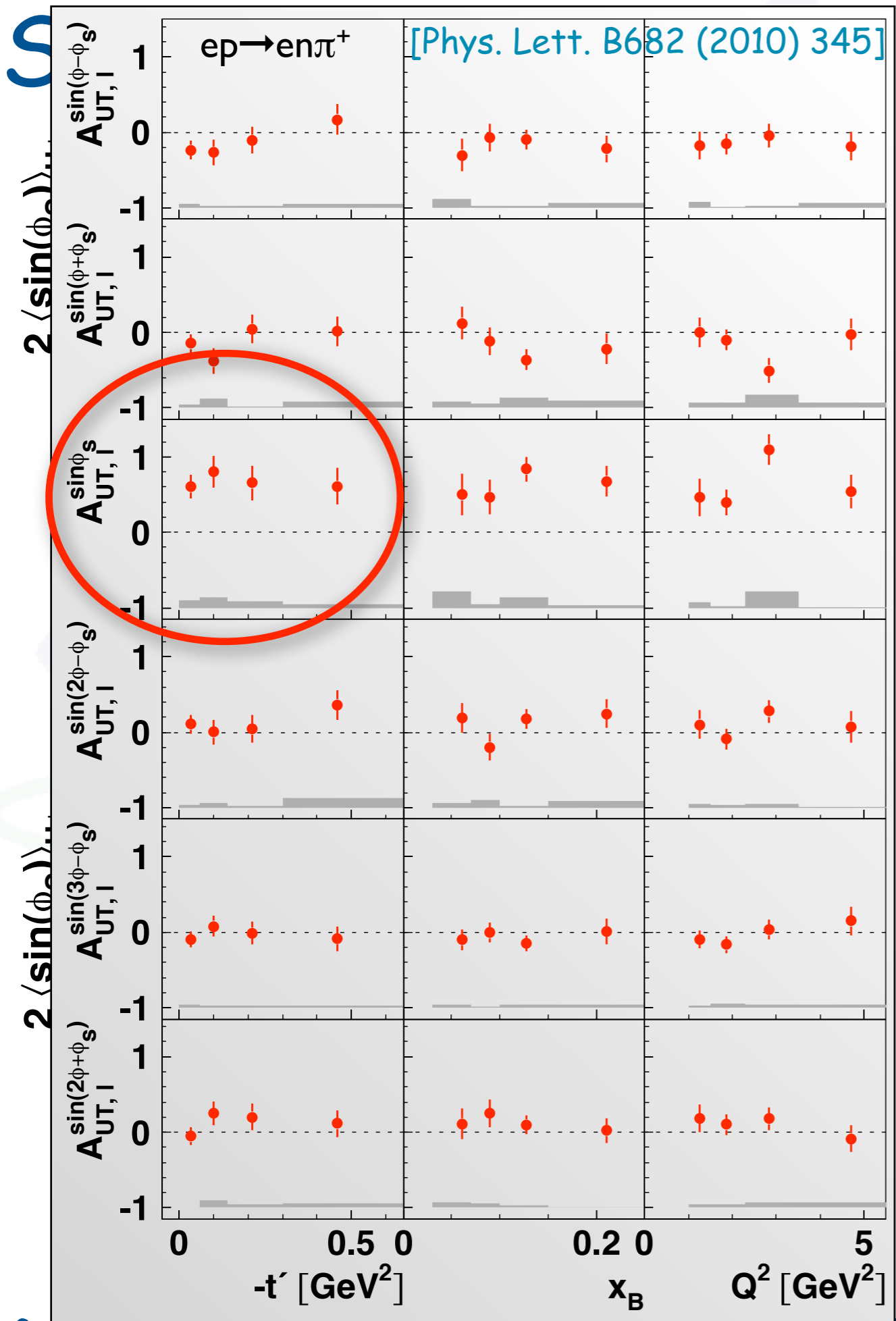
- no dependence on hadron charge **expected** for Cahn effect
- ➔ flavor dependence of transverse momentum
- ➔ sign of Boer-Mulders in $\cos\phi$ modulation (indeed, overall pattern resembles B-M modulations)
- ➔ additional "genuine" twist-3 contributions?

Subleading twist II - $\langle \sin(\phi_s) \rangle_{UT}$



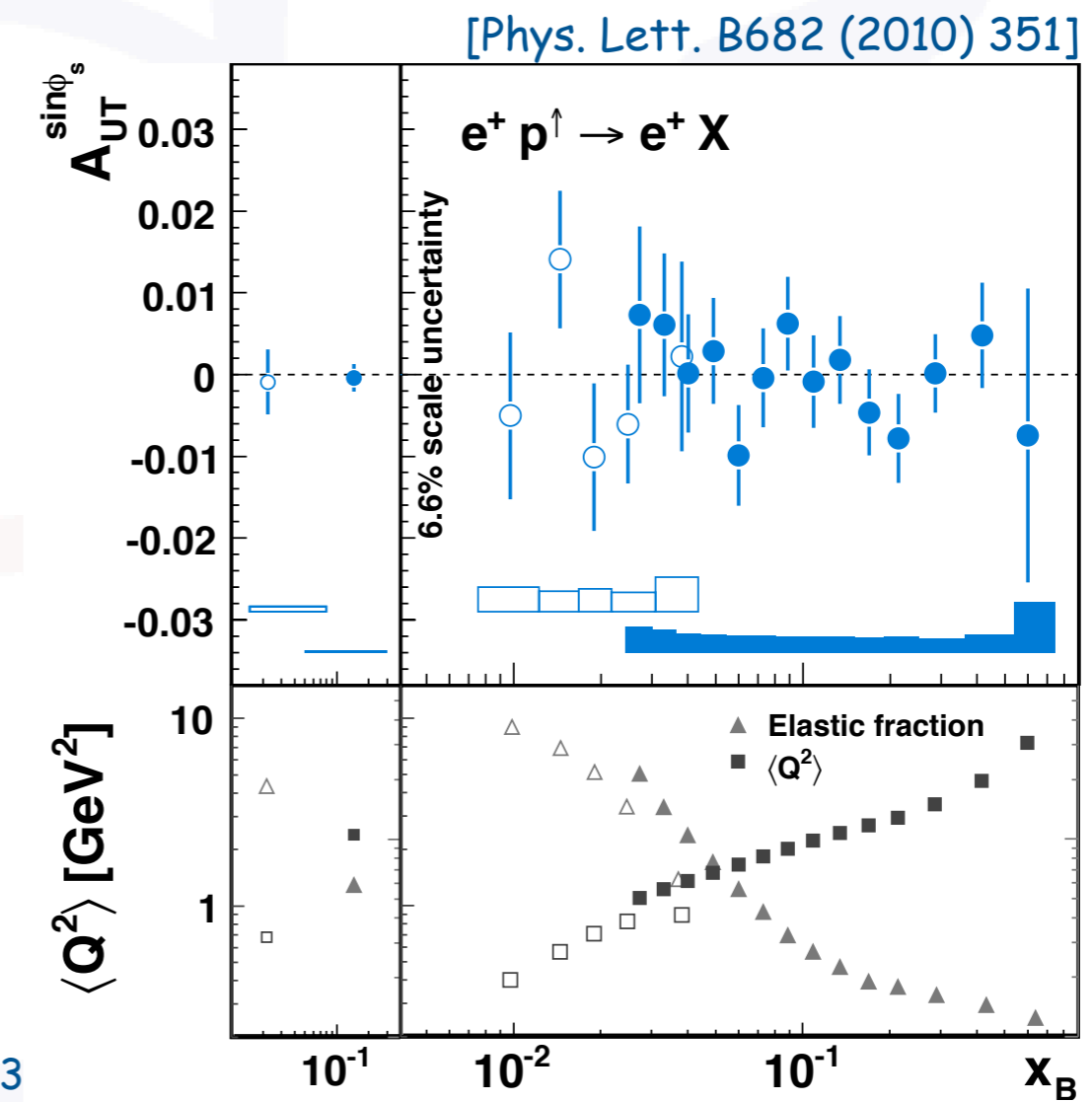
- significant non-zero signal observed for negatively charged mesons
- must vanish after integration over $P_{h\perp}$ and z , and summation over all hadrons



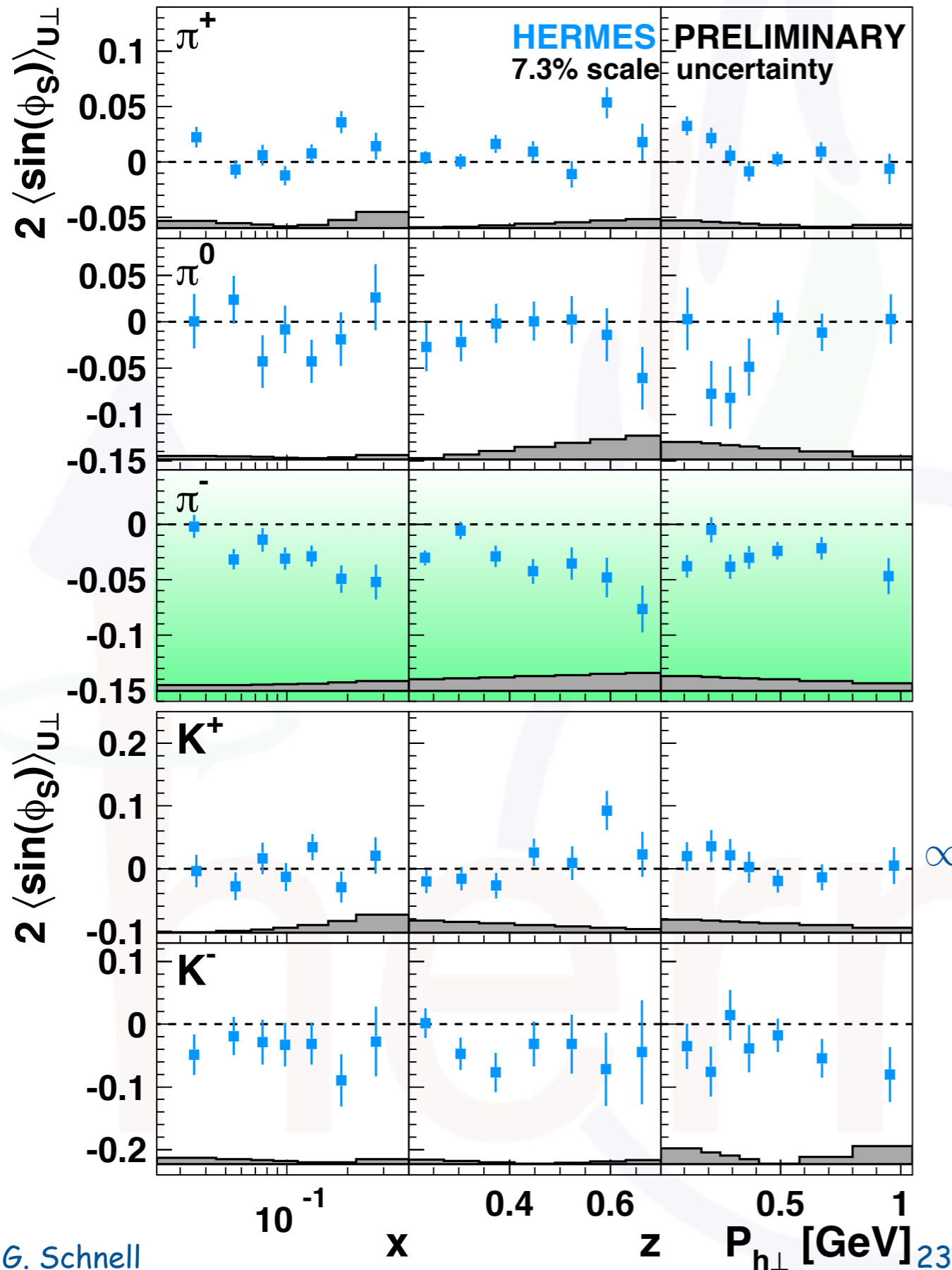


- $\langle \sin(\phi_s) \rangle_{UT}$

- significant non-zero signal observed for negatively charged mesons
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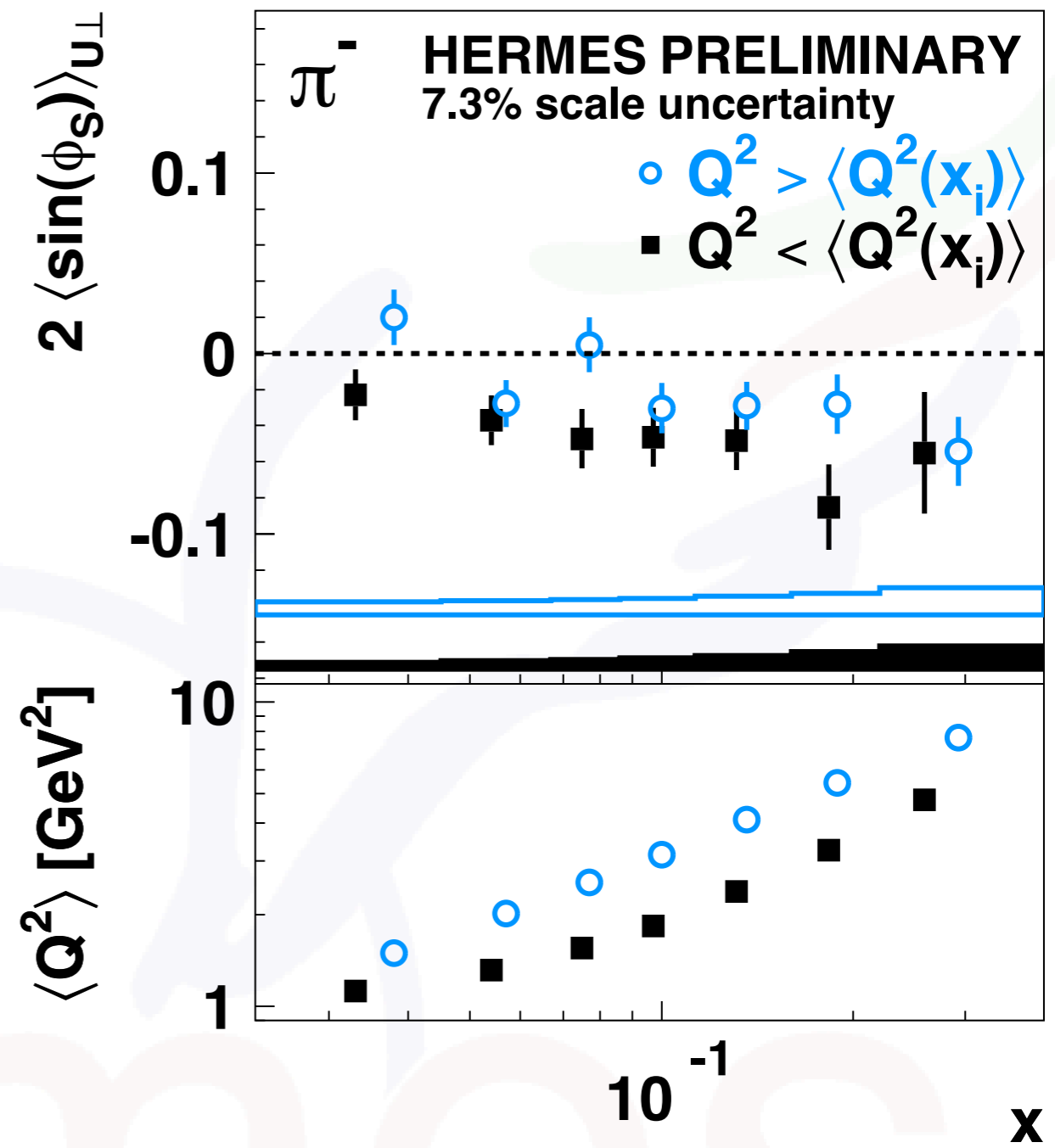
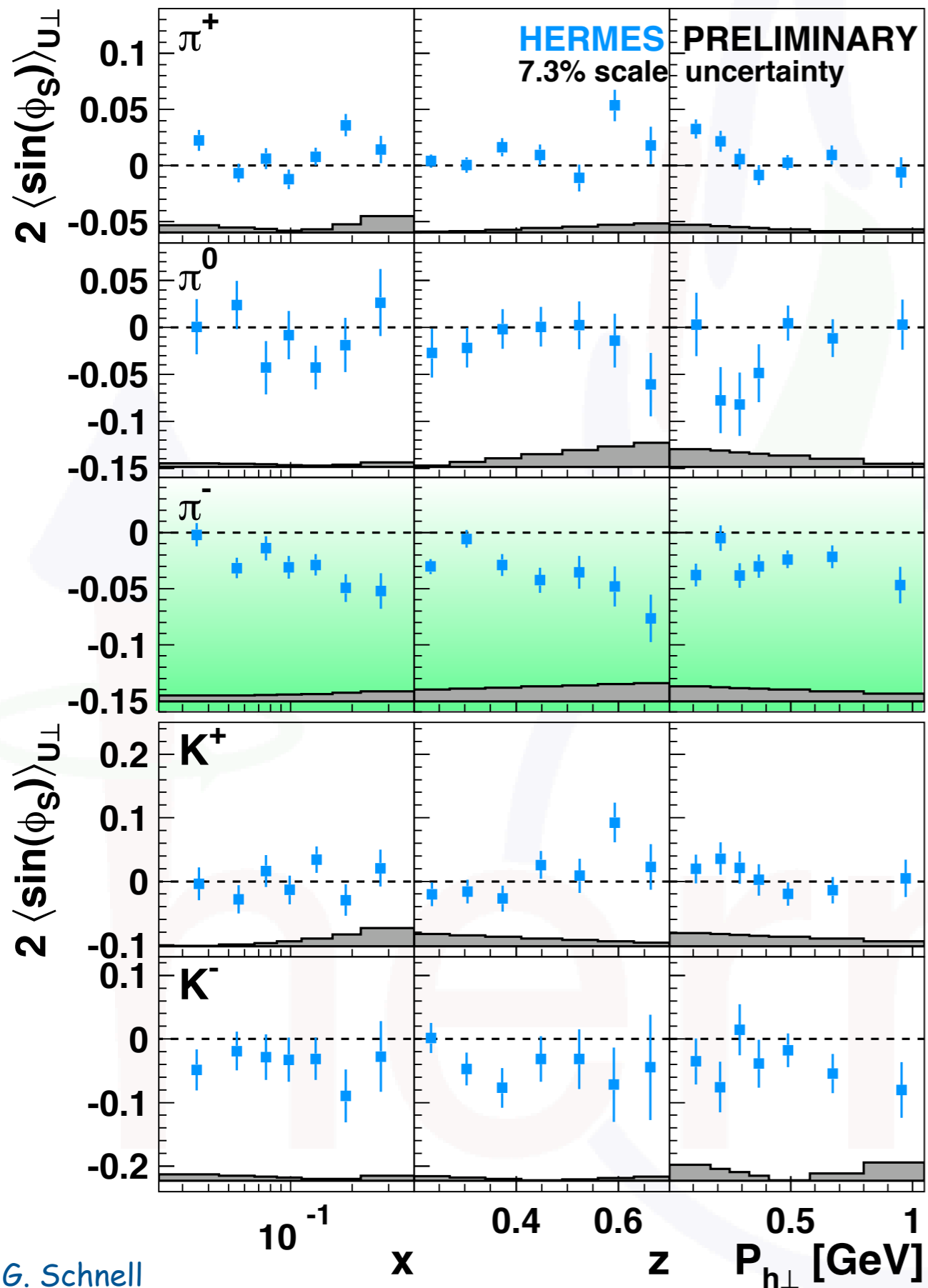
Subleading twist II - $\langle \sin(\phi_s) \rangle_{U\perp}$



- significant non-zero signal observed for negatively charged mesons
- must vanish after integration over $P_{h\perp}$ and z , and summation over all hadrons
- various terms related to transversity, worm-gear, Sivers etc.:

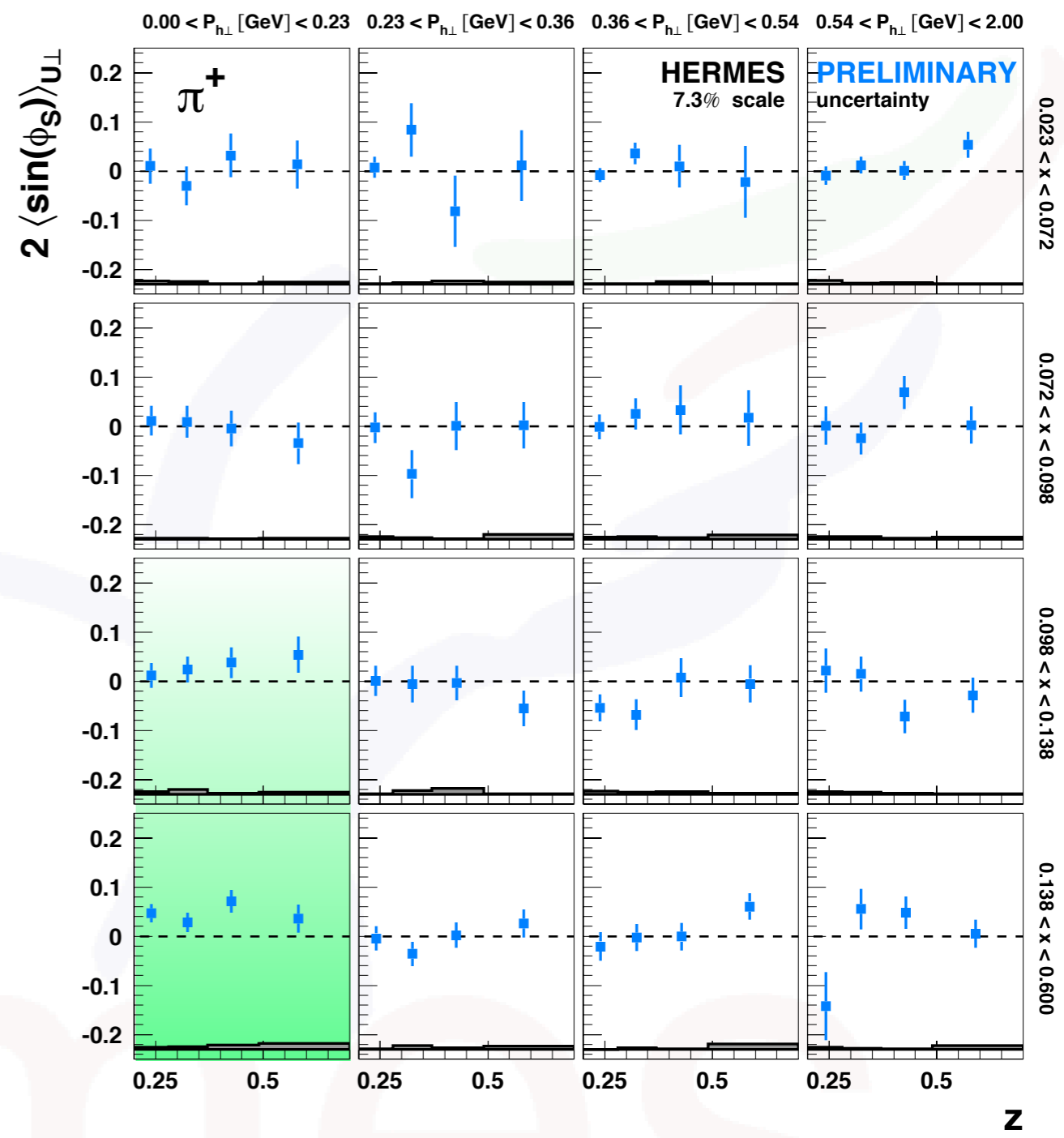
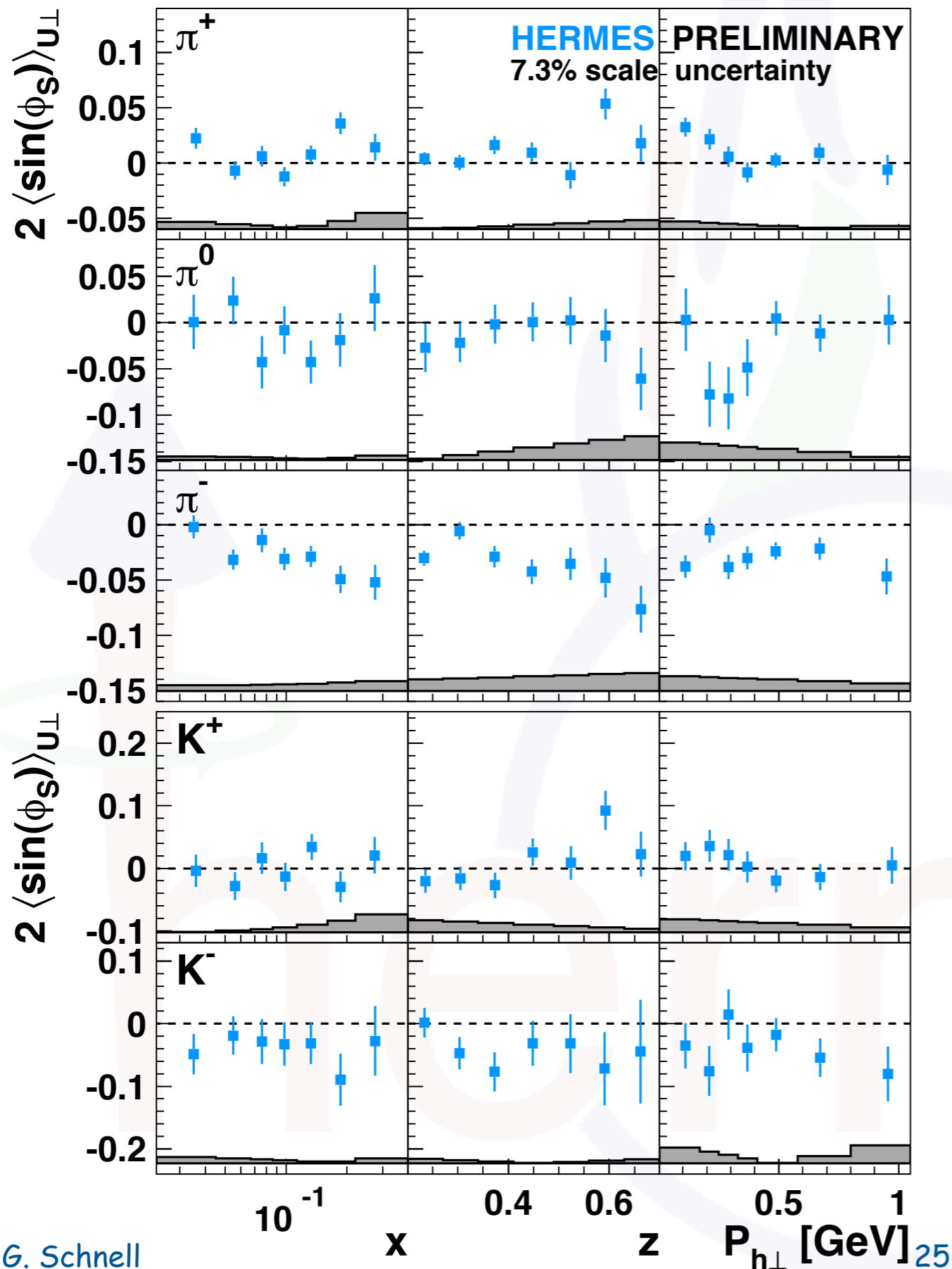
$$\propto \left(x f_{\text{T}}^{\perp} D_1 - \frac{M_{\text{h}}}{M} h_1 \frac{\tilde{\text{H}}}{z} \right) - \mathcal{W}(p_{\text{T}}, k_{\text{T}}, P_{h\perp}) \left[\left(x h_{\text{T}} H_1^{\perp} + \frac{M_{\text{h}}}{M} g_{1\text{T}} \frac{\tilde{\text{G}}^{\perp}}{z} \right) - \left(x h_{\text{T}}^{\perp} H_1^{\perp} - \frac{M_{\text{h}}}{M} f_{1\text{T}}^{\perp} \frac{\tilde{\text{D}}^{\perp}}{z} \right) \right]$$

Subleading twist II - $\langle \sin(\phi_s) \rangle_{UT}$



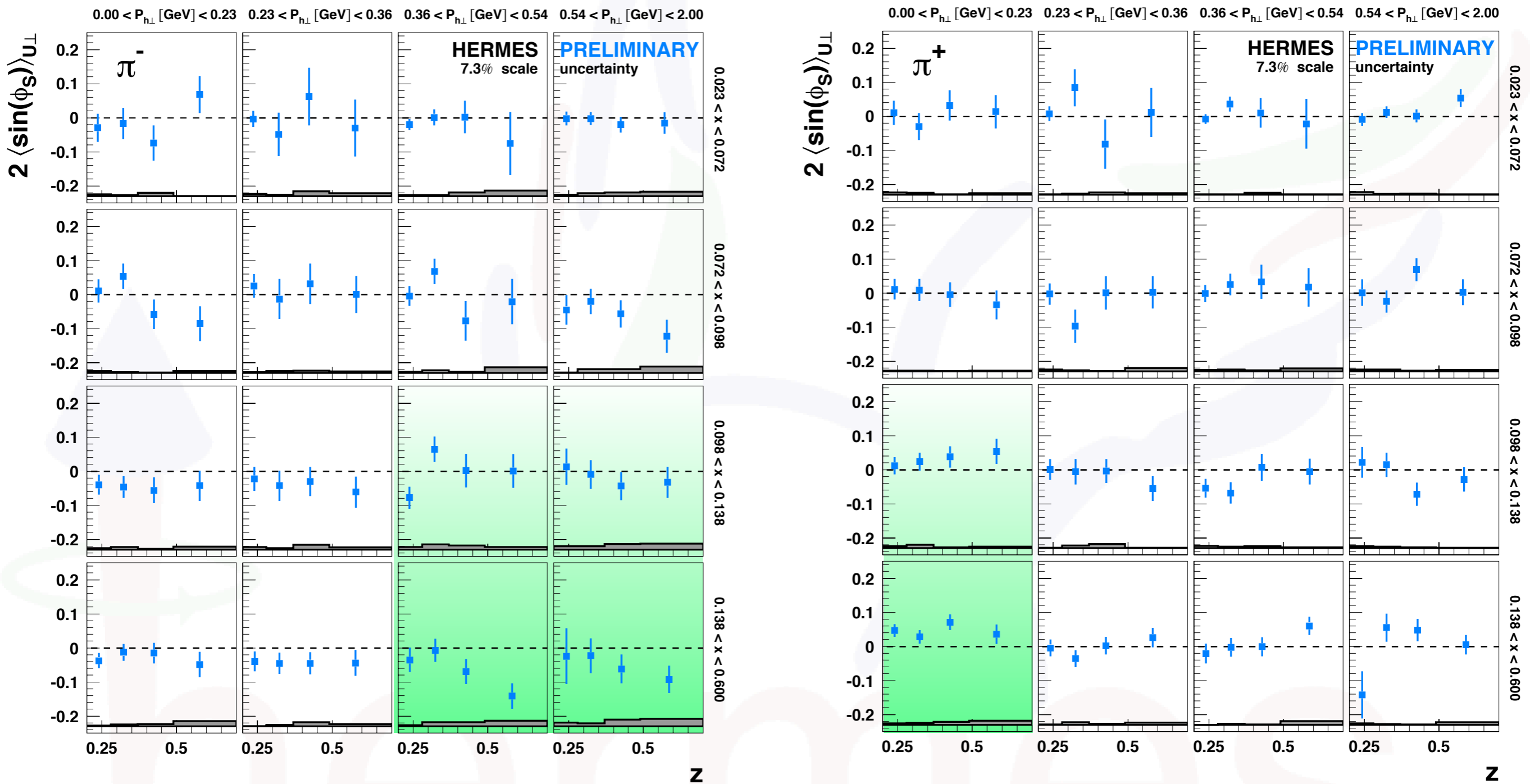
● hint of Q^2 dependence seen in signal for negative pions

Subleading twist II - $\langle \sin(\phi_S) \rangle_{UT}$



● positive amplitudes at low $P_{h\perp}$
also for positive pions

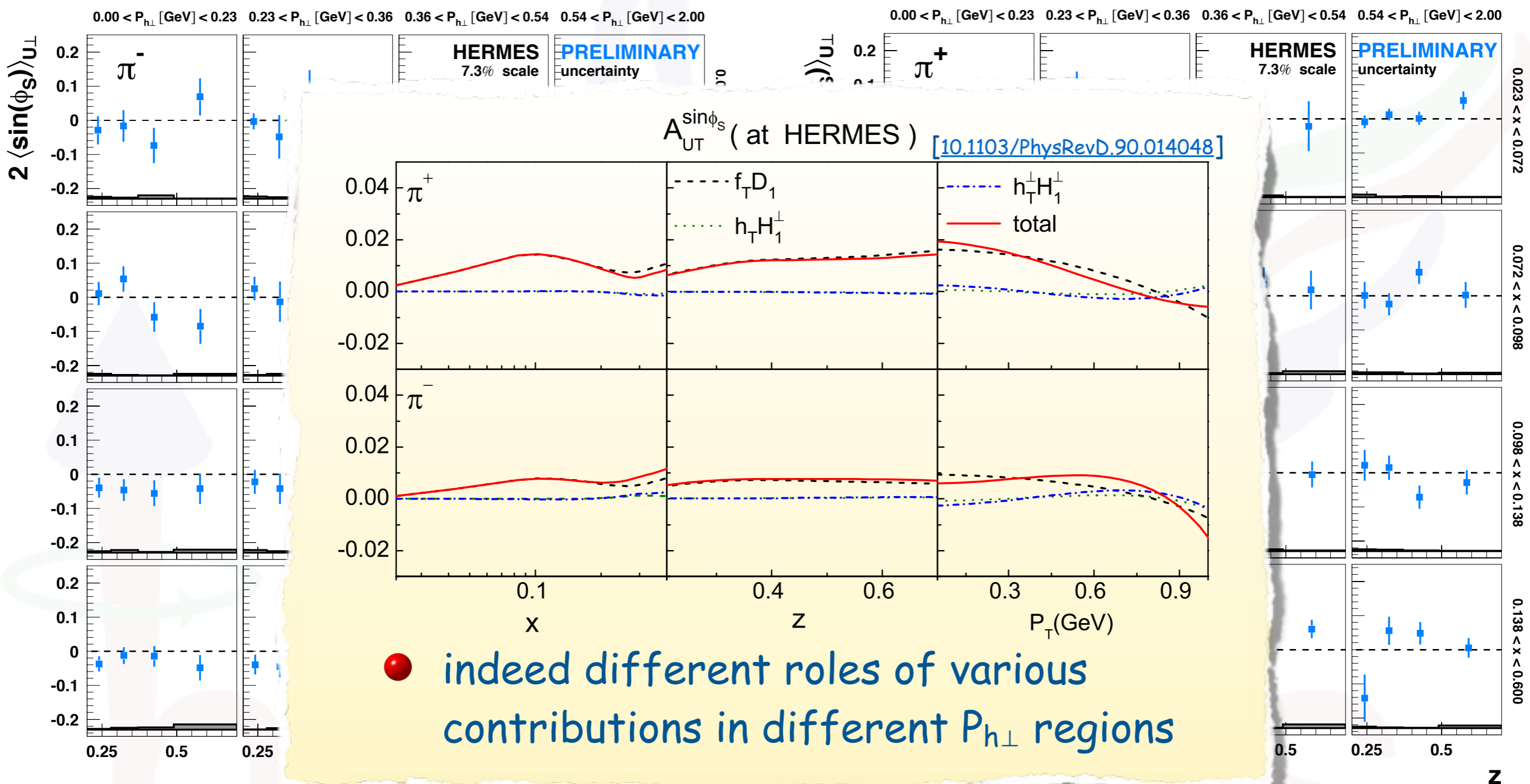
Subleading twist II - $\langle \sin(\phi_S) \rangle_{UT}$



● nonzero amplitudes mainly at large $P_{h\perp}$ in case of negative pions

● positive amplitudes at low $P_{h\perp}$ also for positive pions

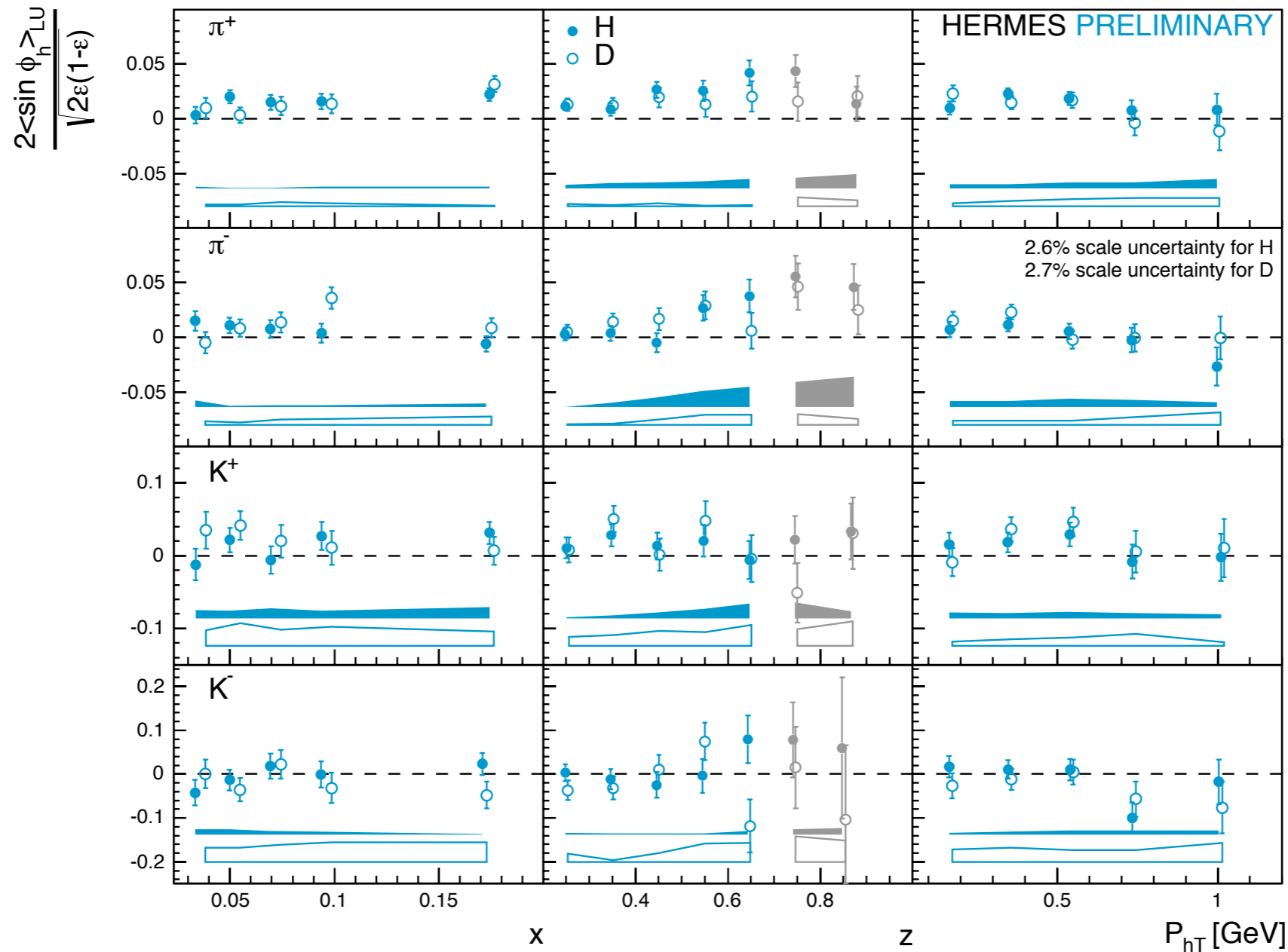
Subleading twist II - $\langle \sin(\phi_S) \rangle_{UT}$



- nonzero amplitudes mainly at large $P_{h\perp}$ in case of negative pions
- positive amplitudes at low $P_{h\perp}$ also for positive pions

Subleading twist III - $\langle \sin(\phi) \rangle_{LU}$

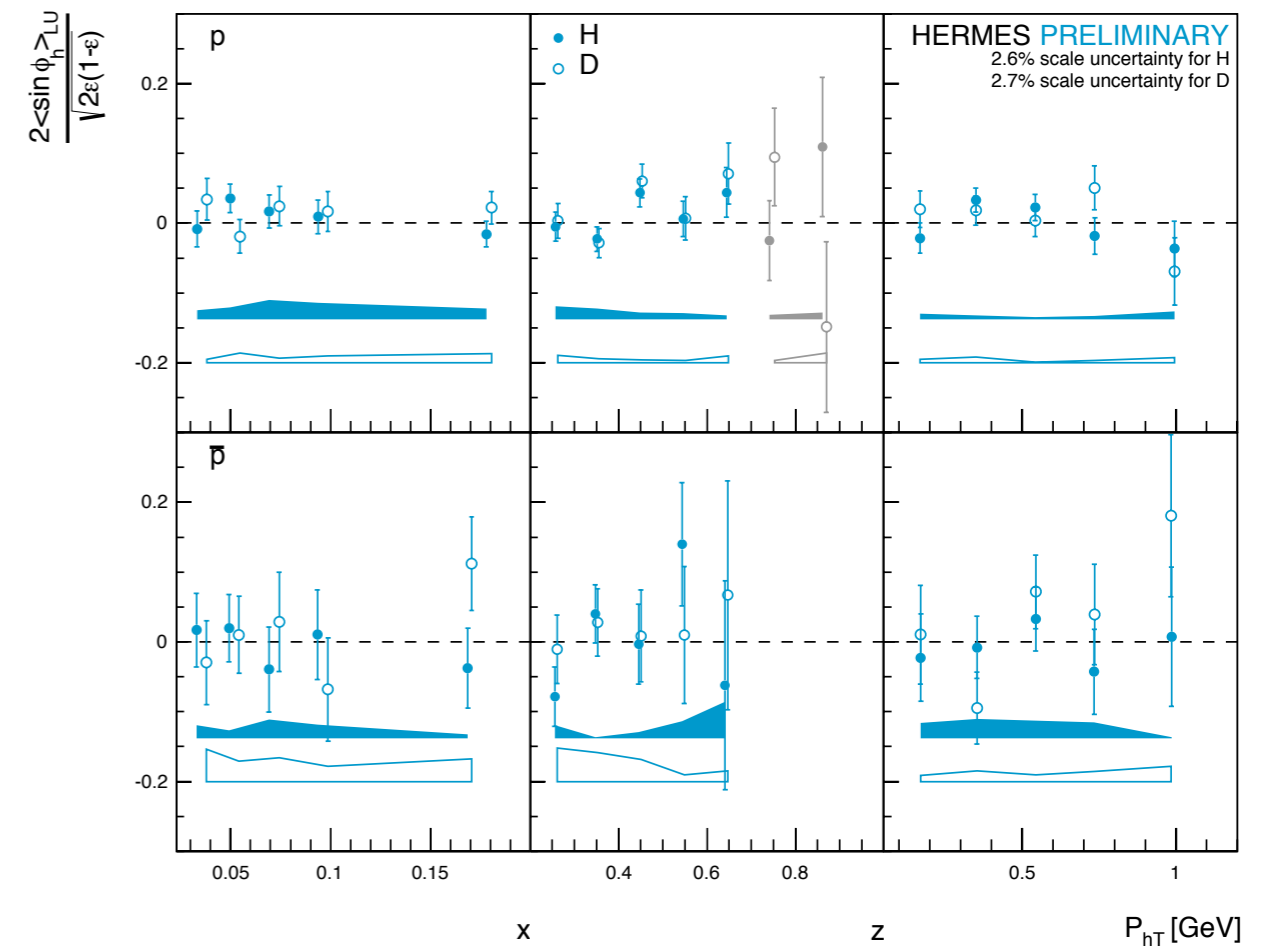
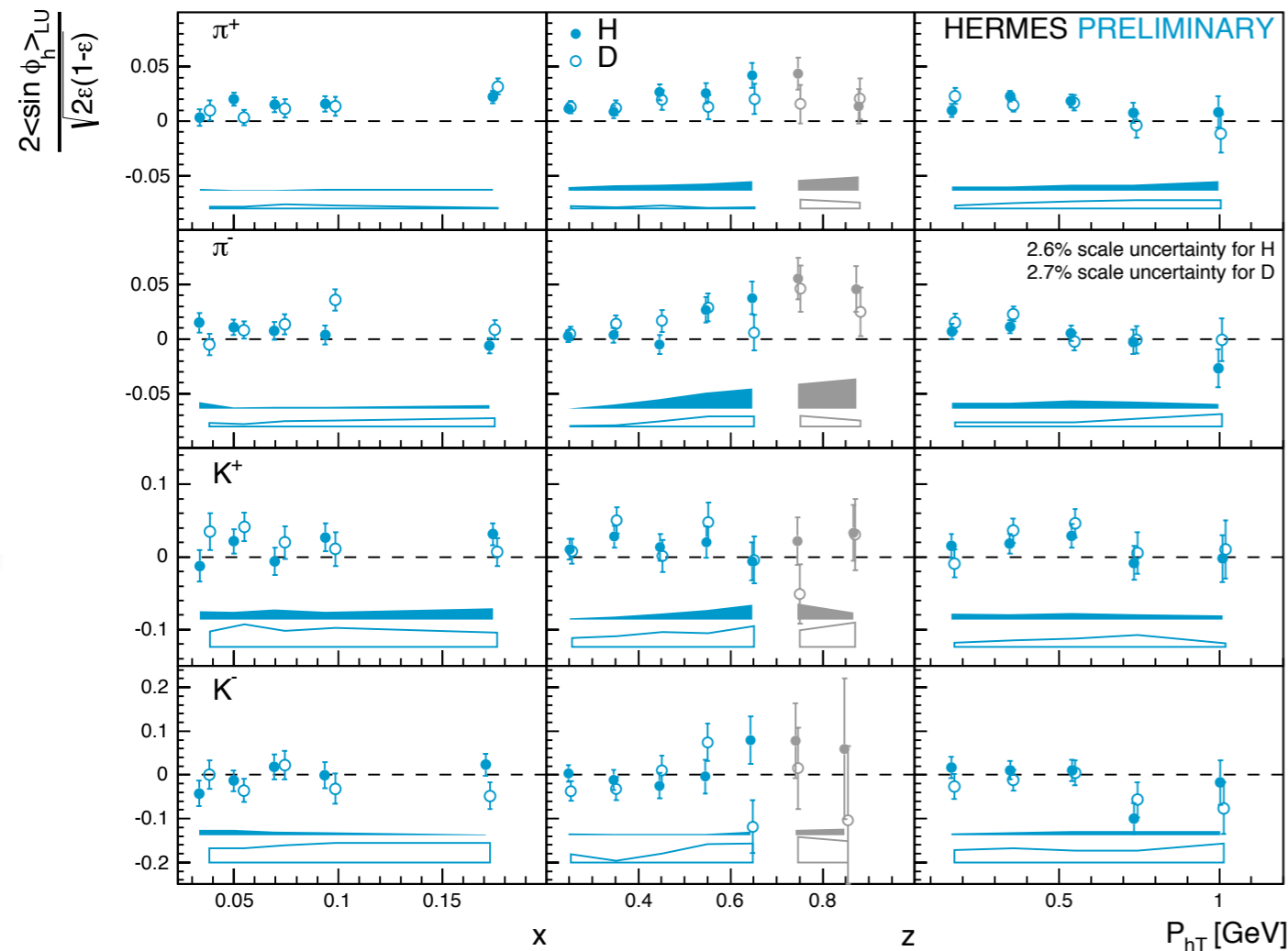
$$\frac{M_h}{Mz} h_1^\perp E \oplus xg^\perp D_1 \oplus \frac{M_h}{Mz} f_1 G^\perp \oplus \color{red} xeH_1^\perp$$



- significant positive amplitudes for (in particular positive) pions

Subleading twist III - $\langle \sin(\phi) \rangle_{LU}$

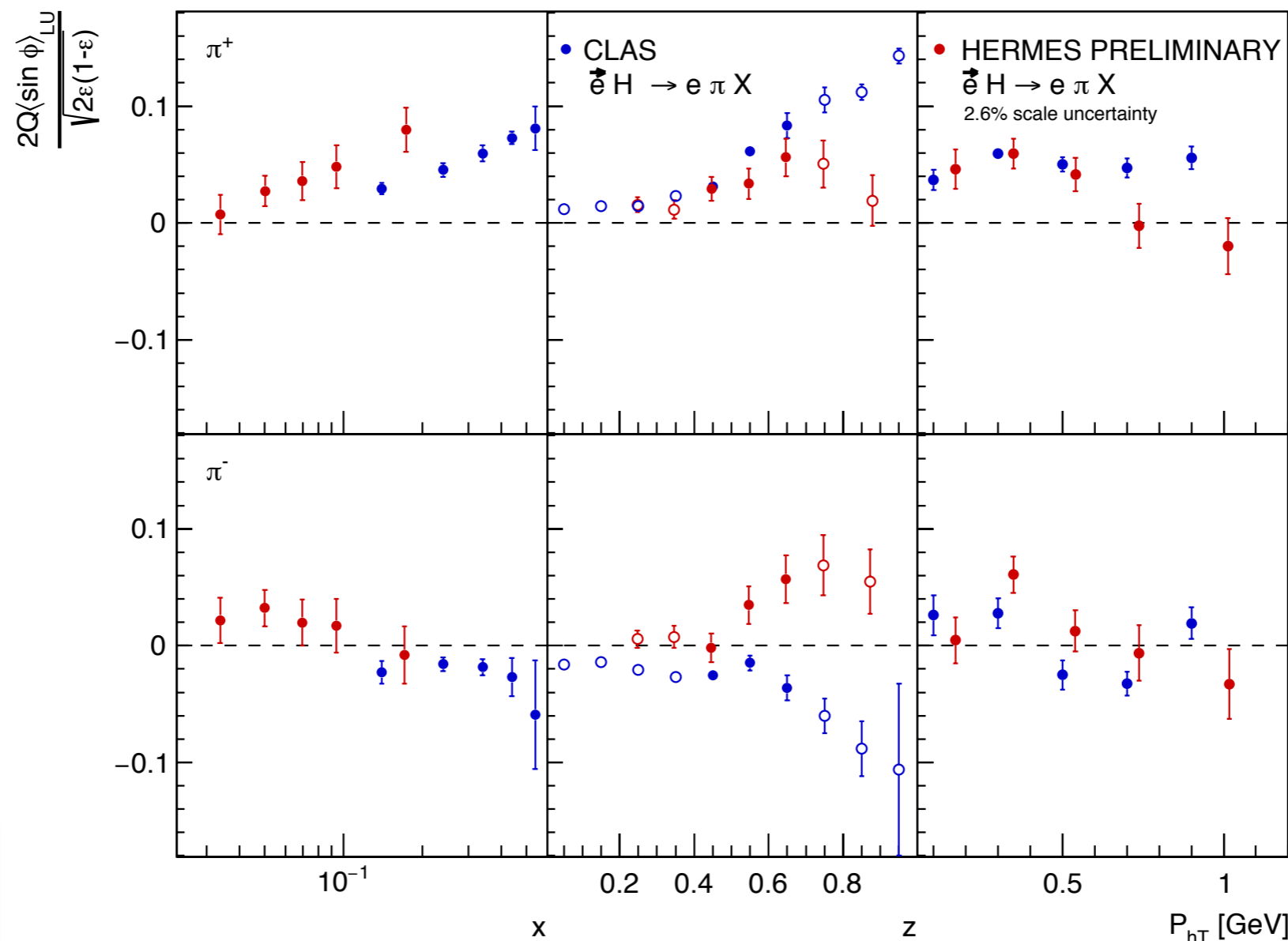
$$\frac{M_h}{M_z} h_1^\perp E \oplus x g^\perp D_1 \oplus \frac{M_h}{M_z} f_1 G^\perp \oplus x e H_1^\perp$$



- mostly consistent w/ zero for other hadrons (except maybe K^+)

Subleading twist III - $\langle \sin(\phi) \rangle_{LU}$

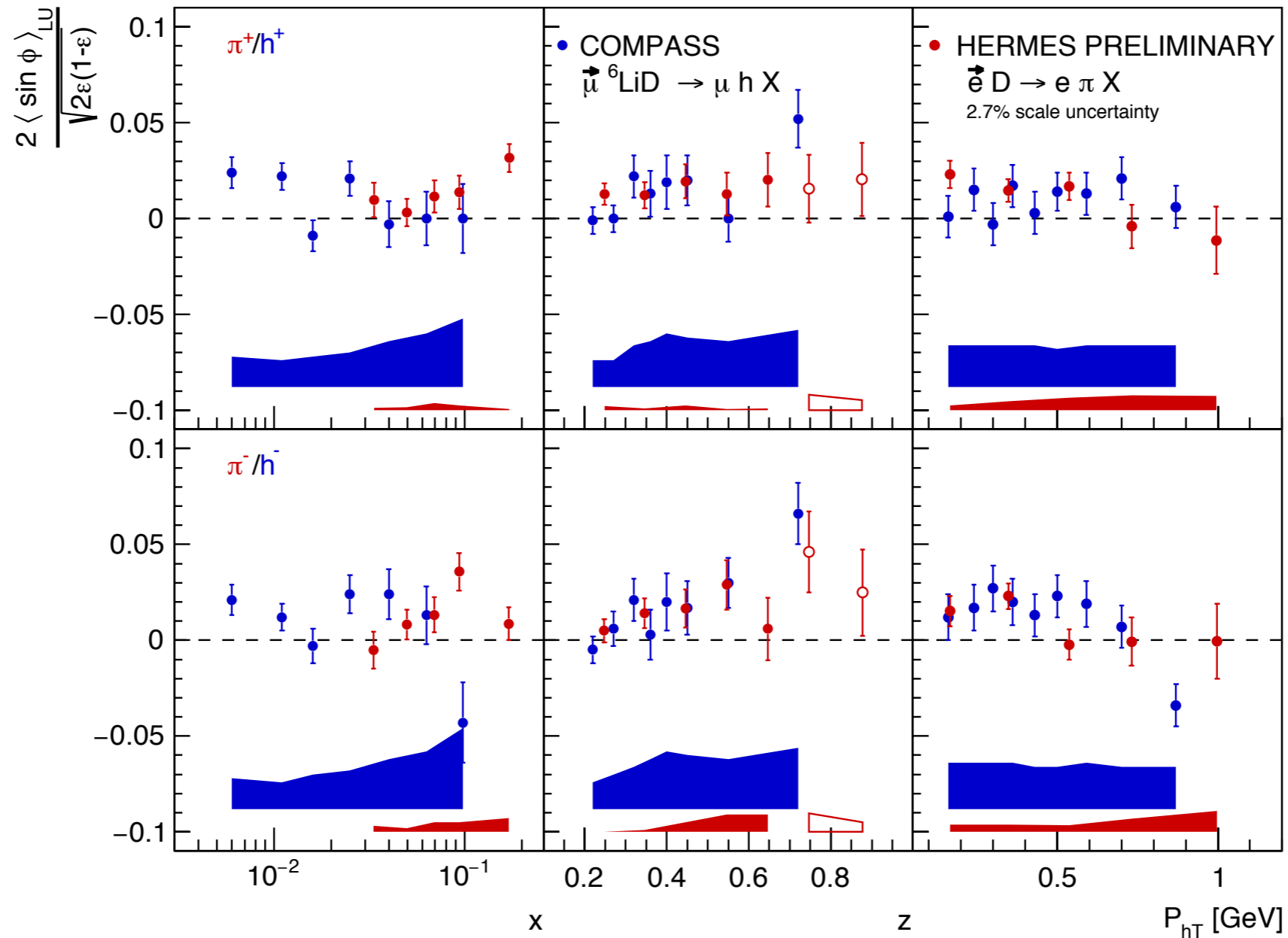
$$\frac{M_h}{M_z} h_1^\perp E \oplus x g^\perp D_1 \oplus \frac{M_h}{M_z} f_1 G^\perp \oplus x e H_1^\perp$$



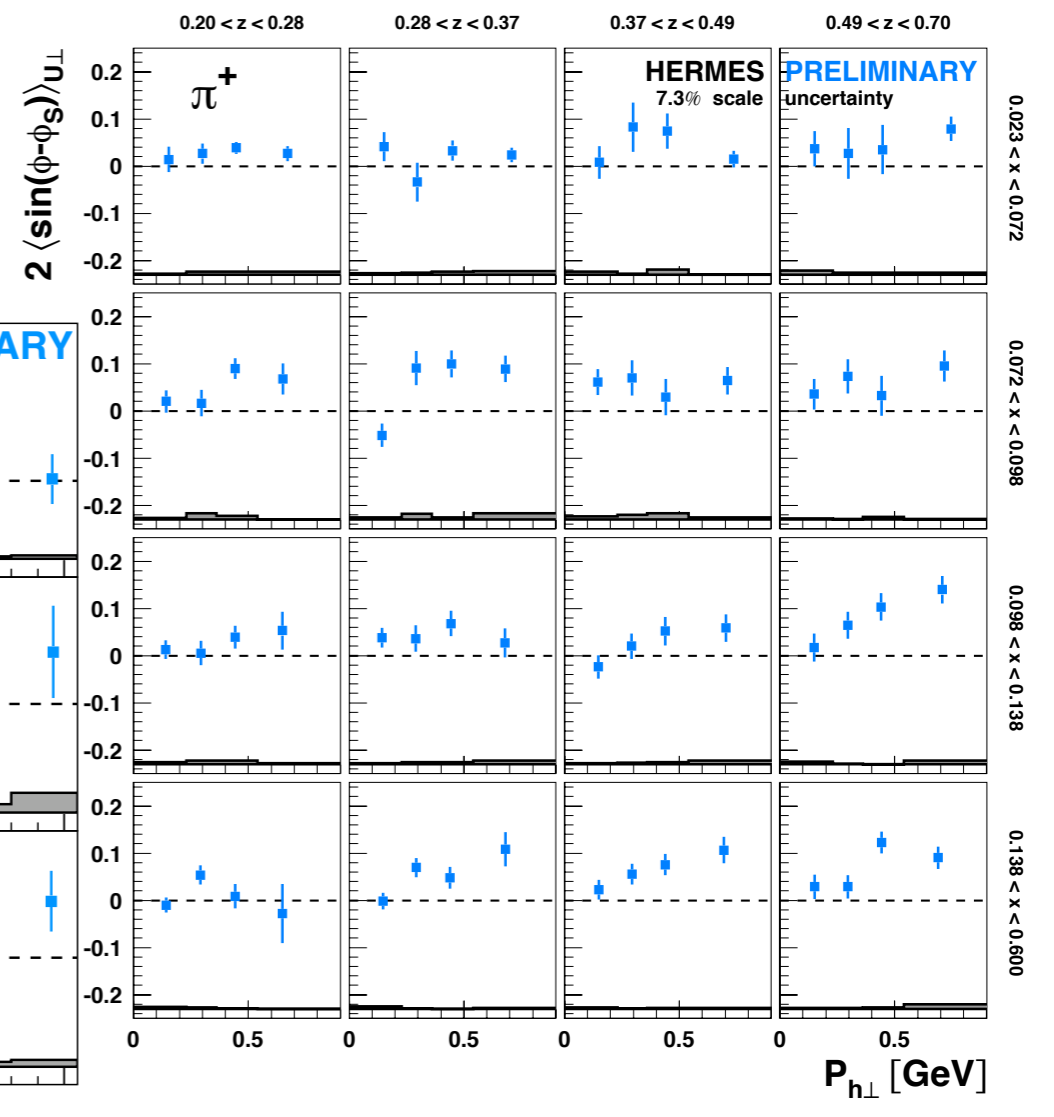
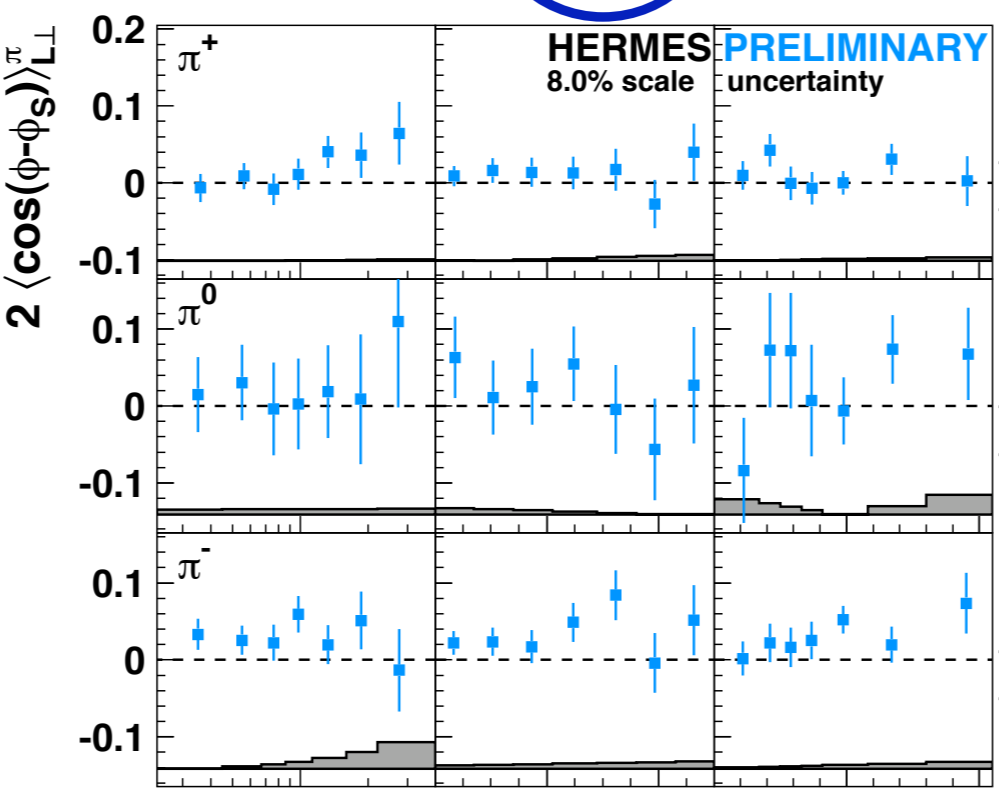
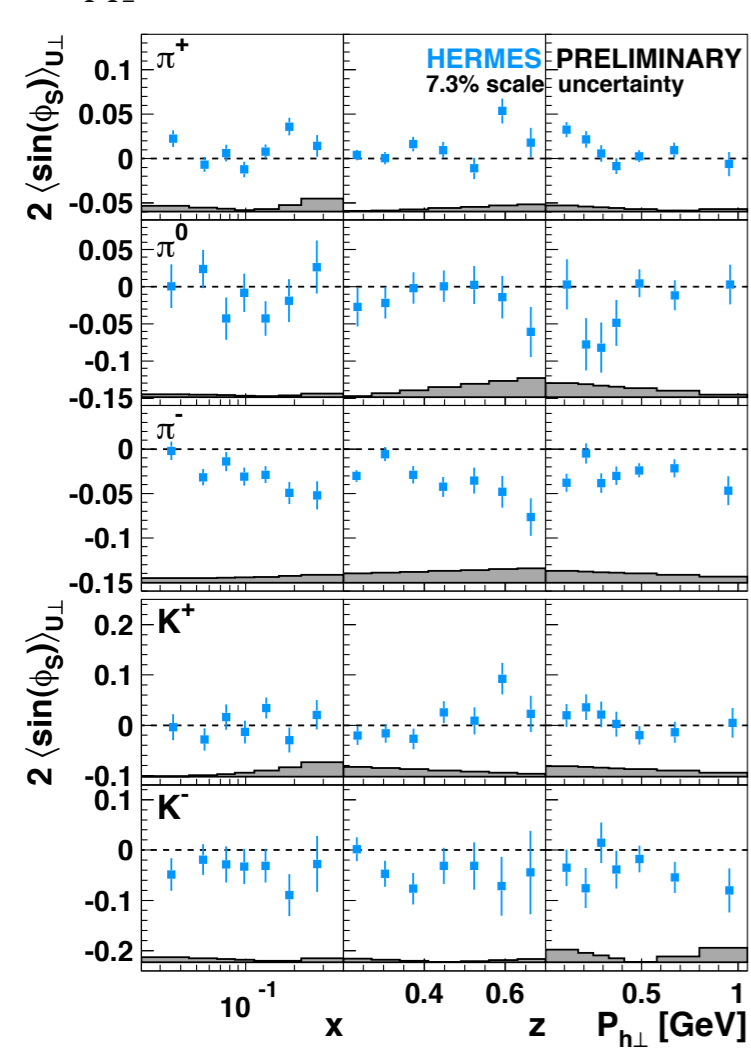
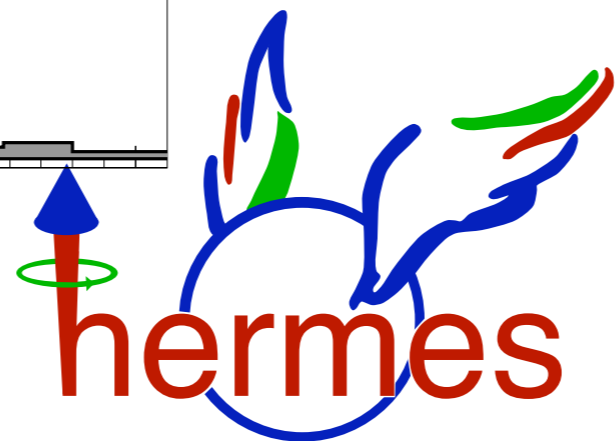
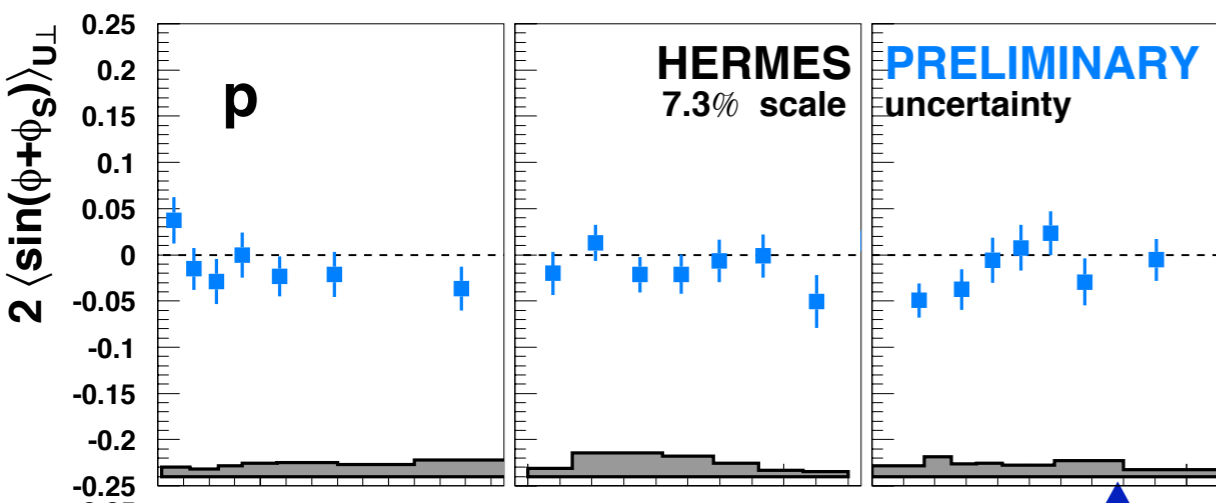
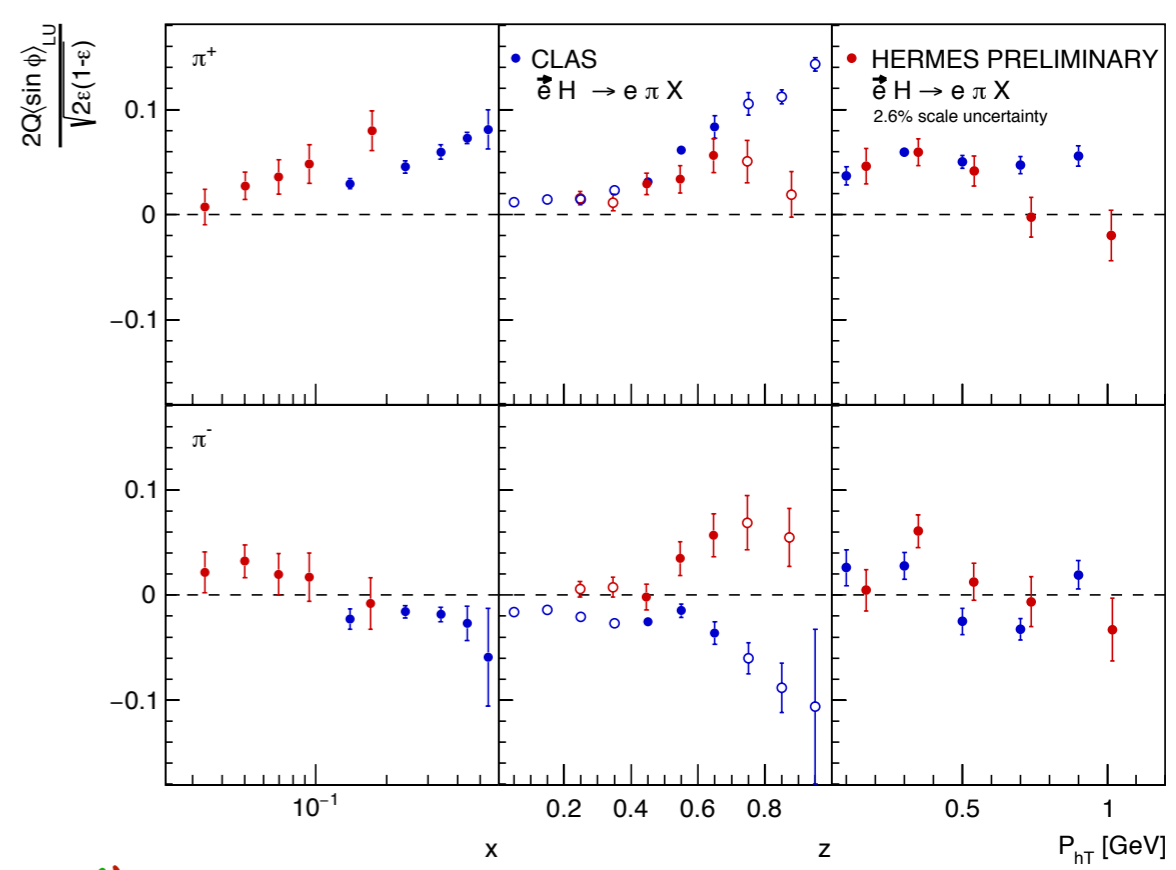
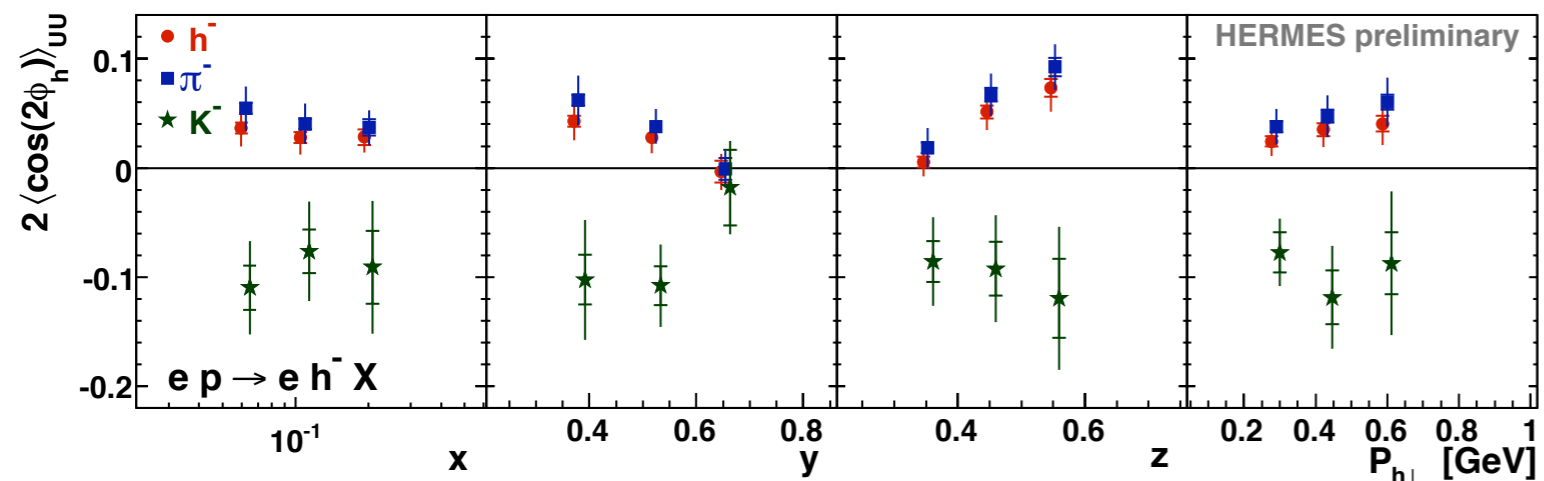
- opposite behavior at HERMES/CLAS of negative pions in z projection due to different x -range probed
- CLAS more sensitive to $e(x)$ Collins term due to higher x probed?

Subleading twist III - $\langle \sin(\phi) \rangle_{LU}$

$$\frac{M_h}{M_z} h_1^\perp E \oplus x g^\perp D_1 \oplus \frac{M_h}{M_z} f_1 G^\perp \oplus x e H_1^\perp$$

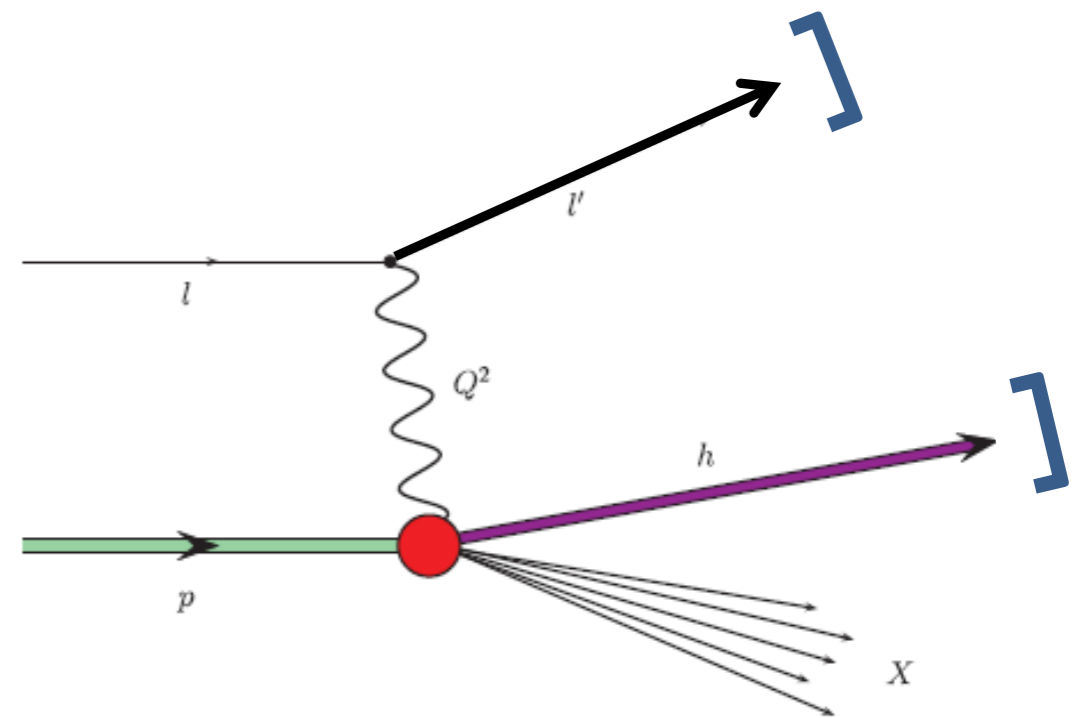


- consistent behavior for charged pions / hadrons at HERMES / COMPASS for isoscalar targets



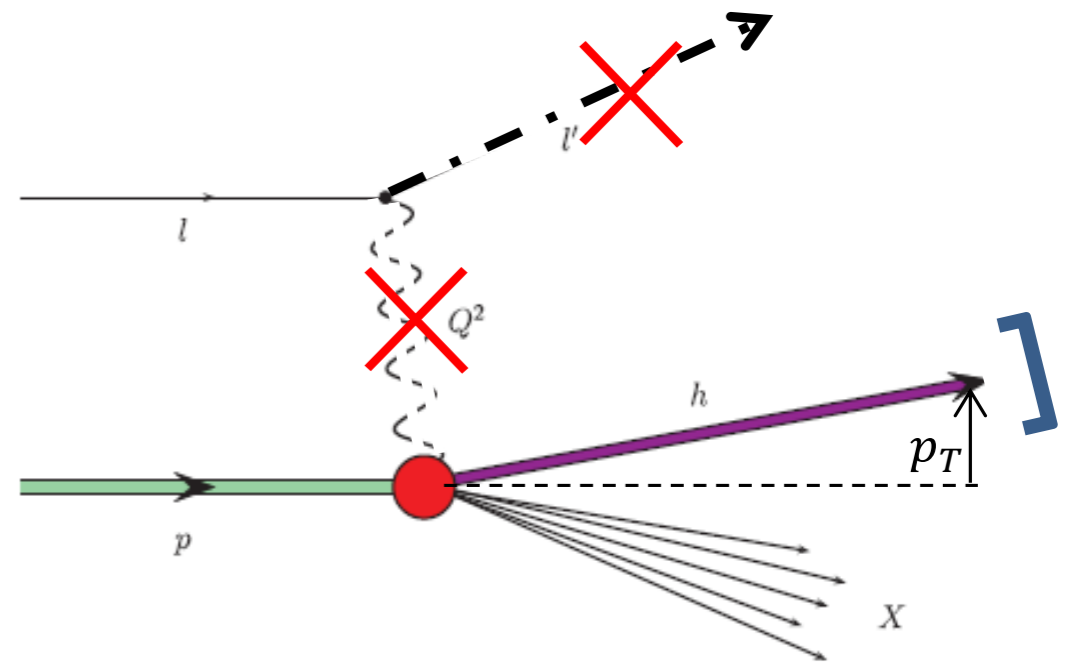
Backup slides

Semi-inclusive hadrons



[Airapetian et al., Phys. Lett. B 728, 183-190 (2014)]

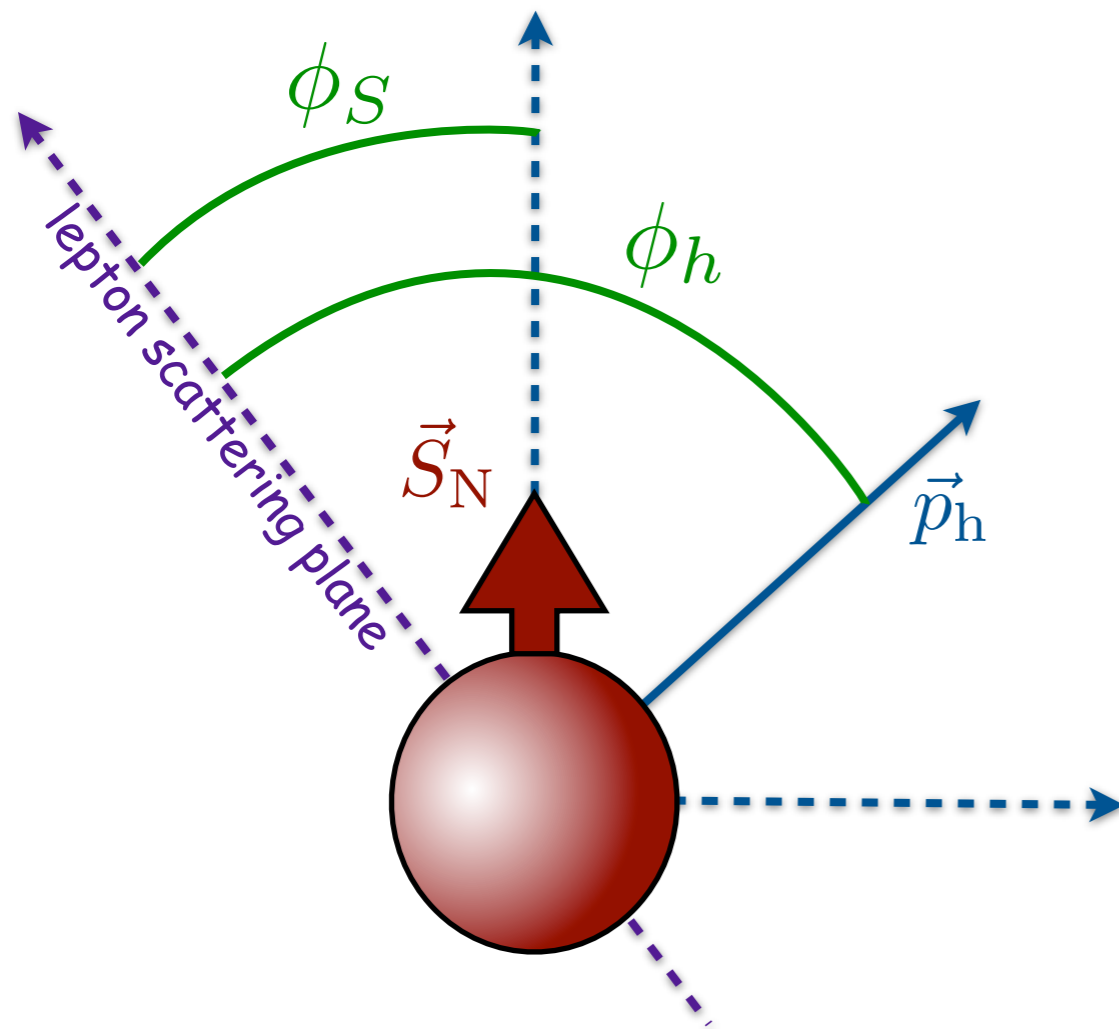
~~S~~emi-inclusive hadrons



[Airapetian et al., Phys. Lett. B 728, 183-190 (2014)]

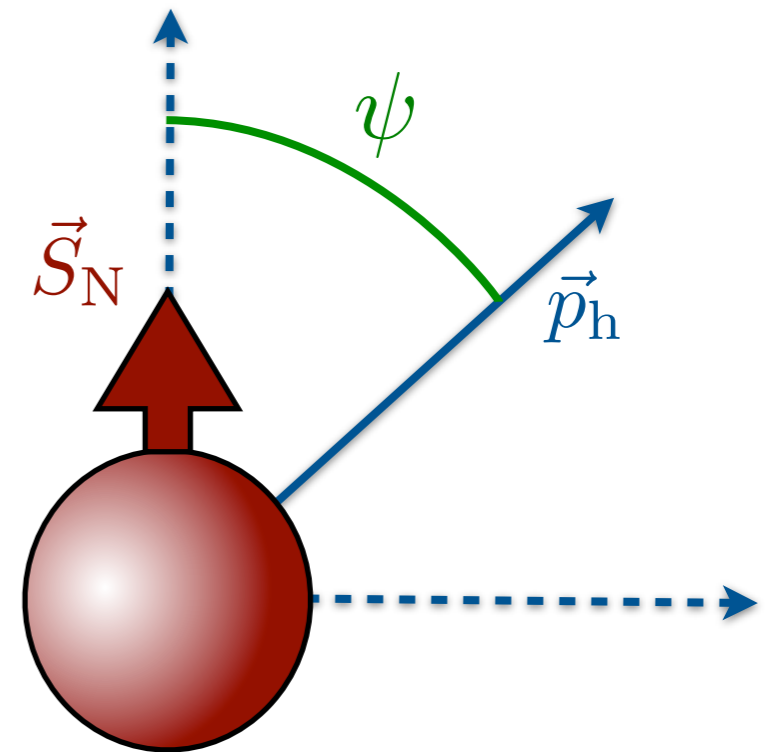
Inclusive hadron electro-production

$$ep^{\uparrow} \rightarrow ehX$$



virtual photon going
into the page

$$ep^{\uparrow} \rightarrow hX$$



lepton beam going
into the page

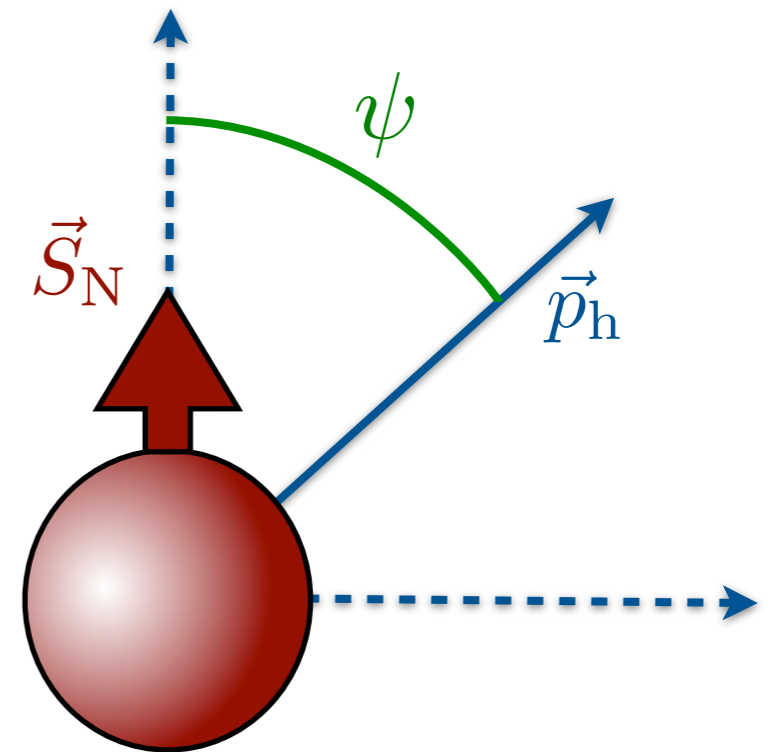
$$\psi \simeq \phi_h - \phi_S$$

→ "Sivers angle"

Inclusive hadron electro-production

- scattered lepton undetected
↳ lepton kinematics unknown

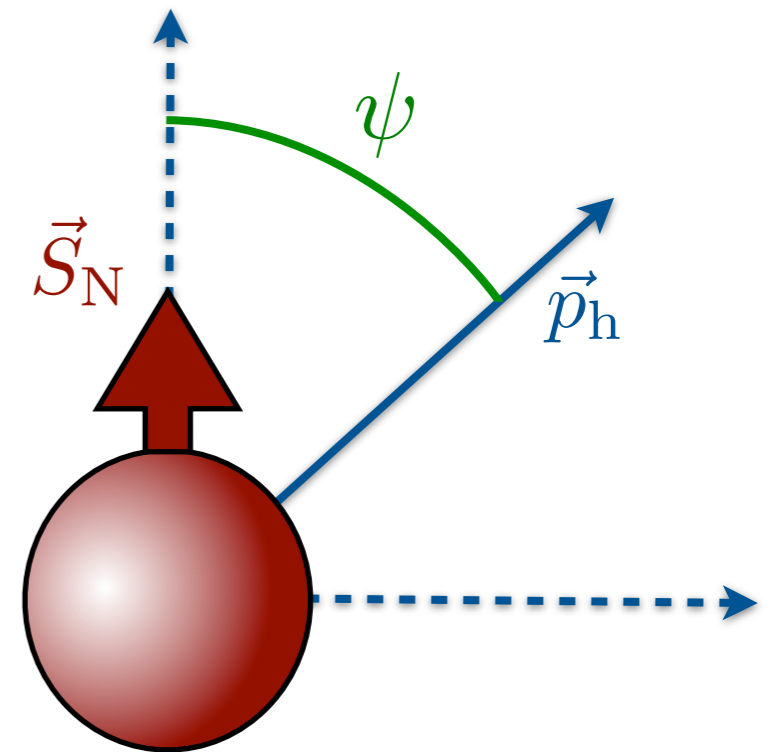
$$ep^{\uparrow} \rightarrow hX$$



Inclusive hadron electro-production

- scattered lepton undetected
↳ lepton kinematics unknown
- dominated by quasi-real photo-production (low Q^2)
↳ hadronic component of photon relevant?

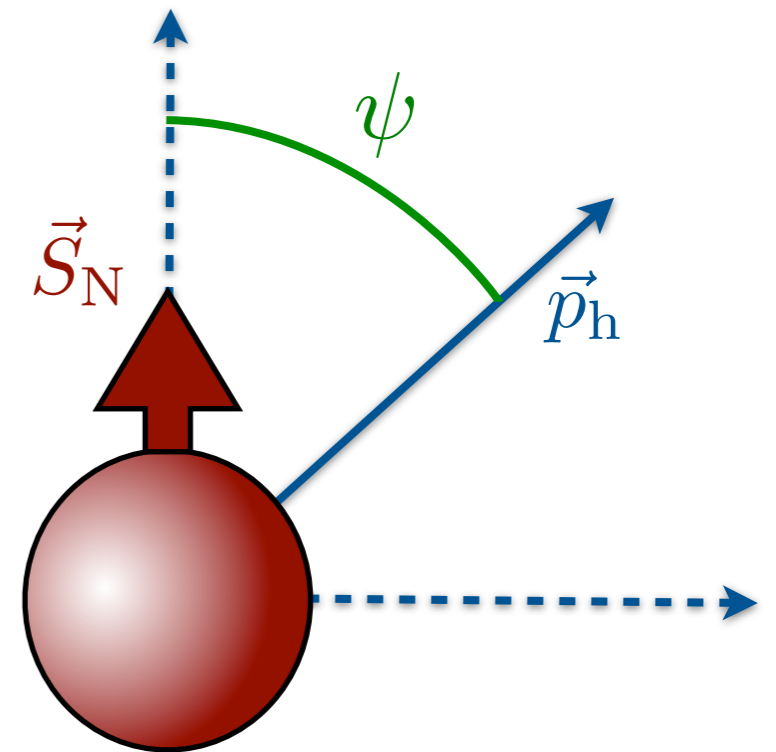
$$ep^{\uparrow} \rightarrow hX$$



Inclusive hadron electro-production

- scattered lepton undetected
 → lepton kinematics unknown
- dominated by quasi-real photo-production (low Q^2)
 → hadronic component of photon relevant?
- cross section proportional to $S_N (\mathbf{k} \times \mathbf{p}_h) \sim \sin\psi$

$$ep^{\uparrow} \rightarrow hX$$



$$A_{UT}(P_T, x_F, \psi) = A_{UT}^{\sin\psi}(P_T, x_F) \sin\psi$$

$$A_N \equiv \frac{\int_{\pi}^{2\pi} d\psi \sigma_{UT} \sin\psi - \int_0^{\pi} d\psi \sigma_{UT} \sin\psi}{\int_0^{2\pi} d\psi \sigma_{UU}}$$

$$= -\frac{2}{\pi} A_{UT}^{\sin\psi}$$

1D dependences of $A_{UT} \sin\psi$ amplitude

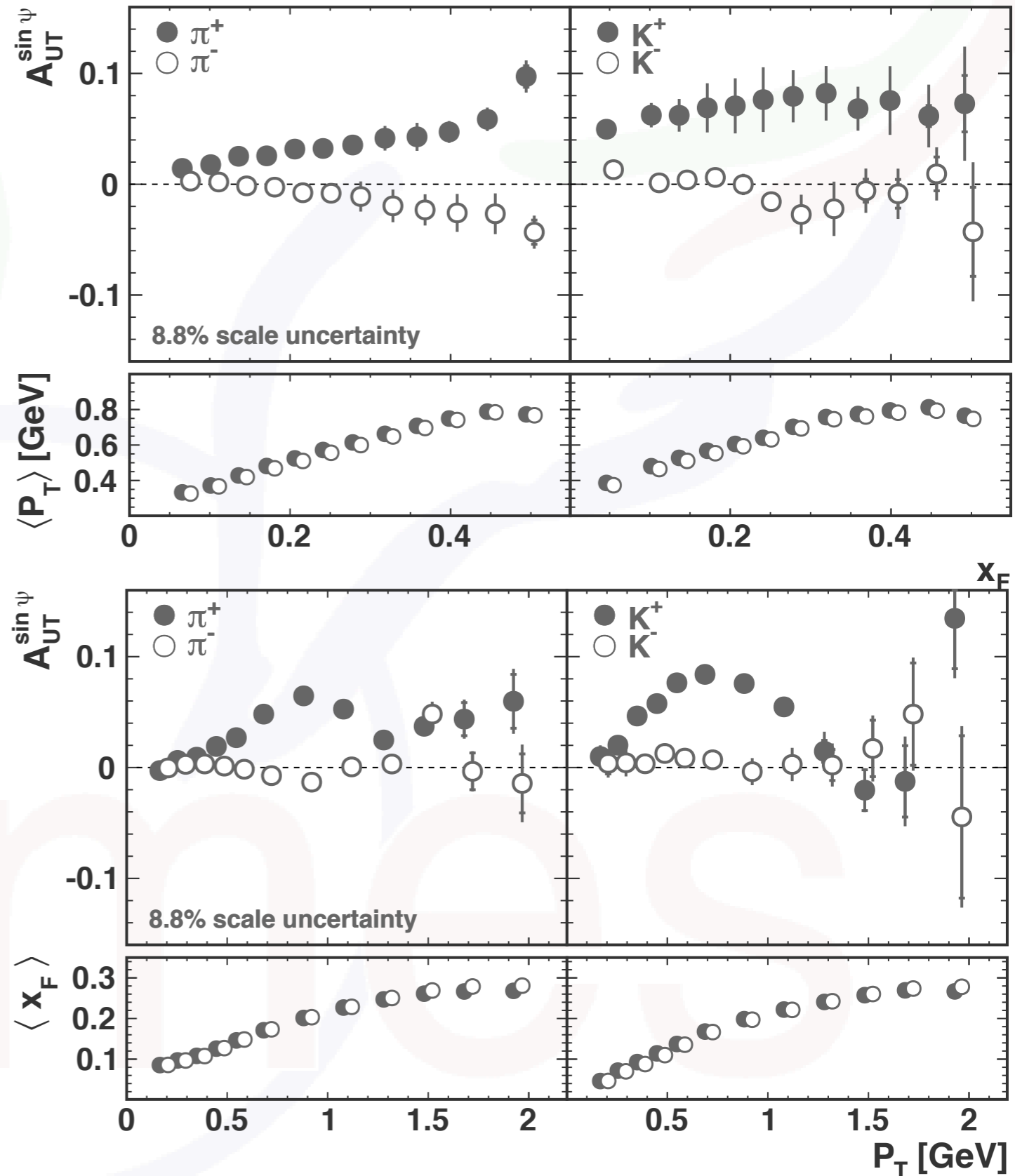
- clear left-right asymmetries for pions and positive kaons

- increasing with x_F (as in pp)

- initially increasing with P_T with a fall-off at larger P_T

- x_F and P_T correlated
 ➔ look at 2D dependences

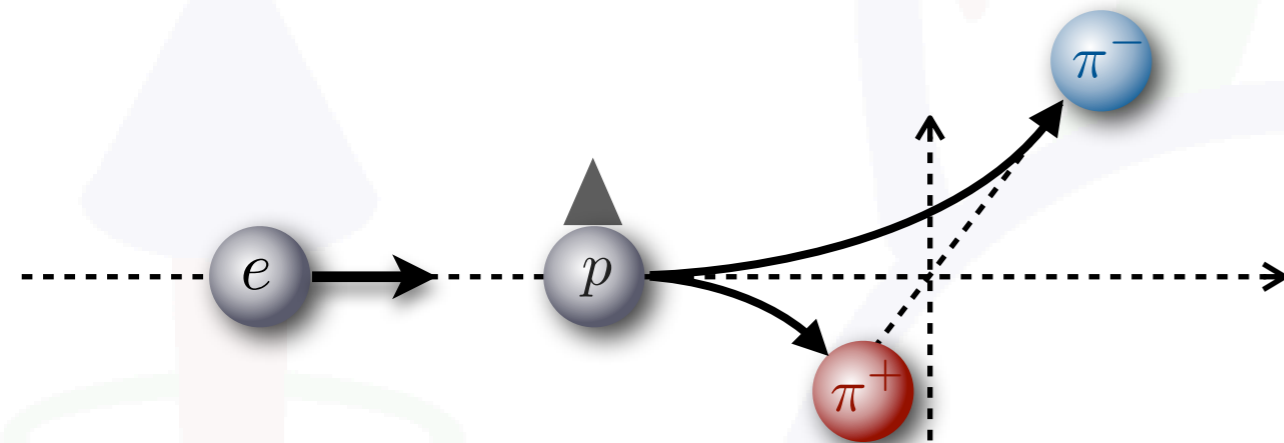
[Airapetian et al., Phys. Lett. B 728, 183-190 (2014)]



1D dependences of $A_{UT} \sin\psi$ amplitude

- clear left-right asymmetries for pions and positive kaons

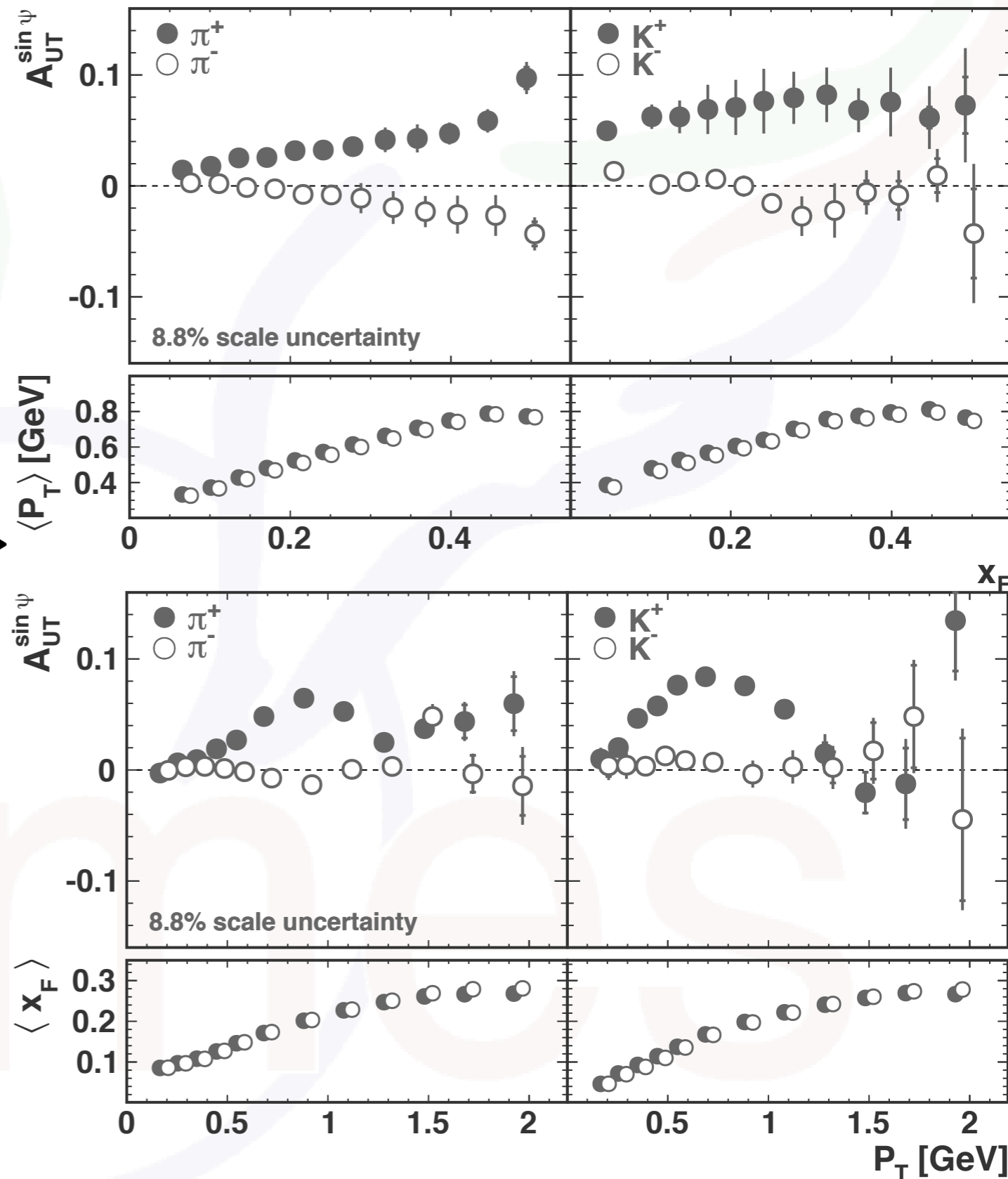
- increasing with x_F (as in pp)



- initially increasing with P_T with a fall-off at larger P_T

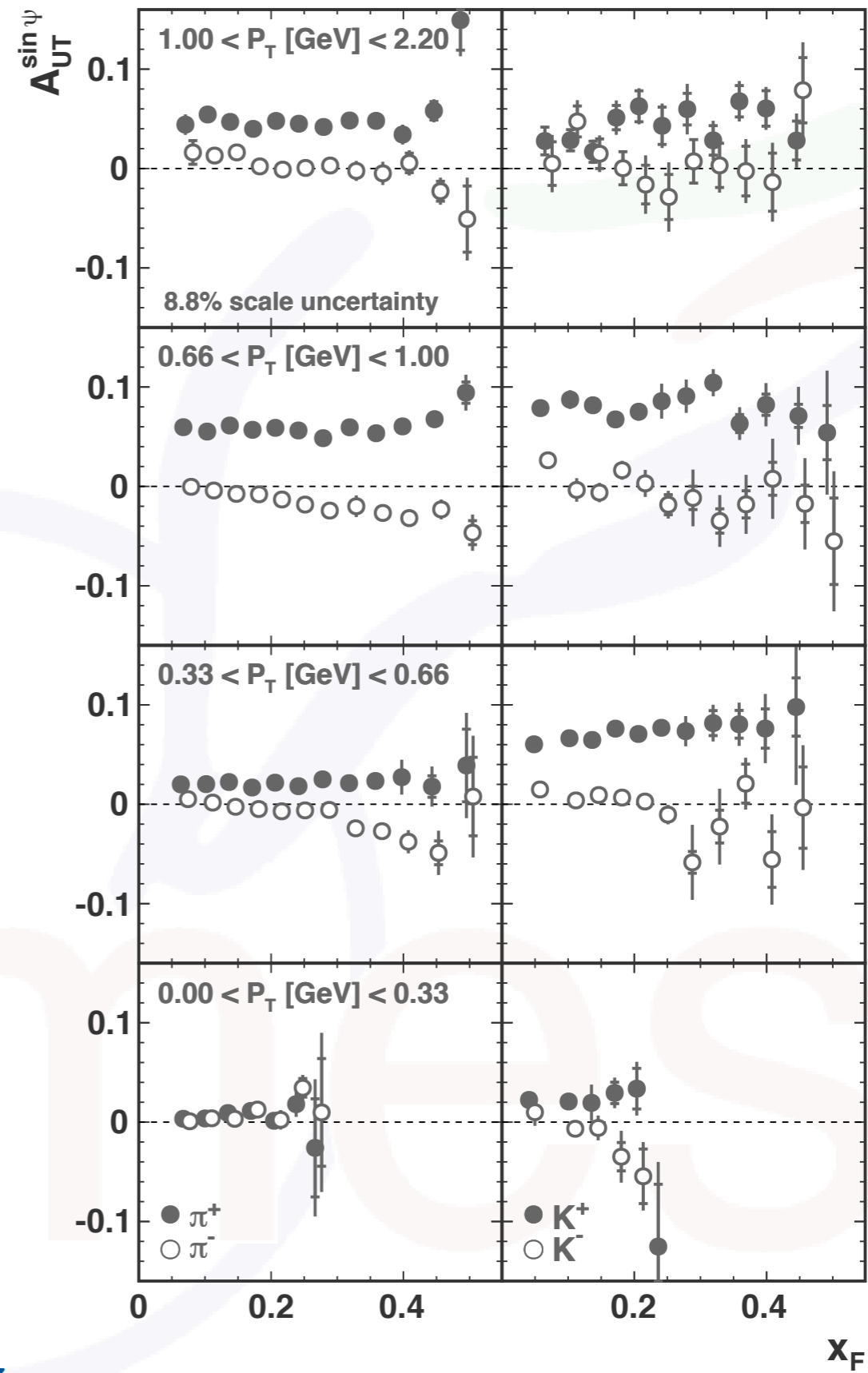
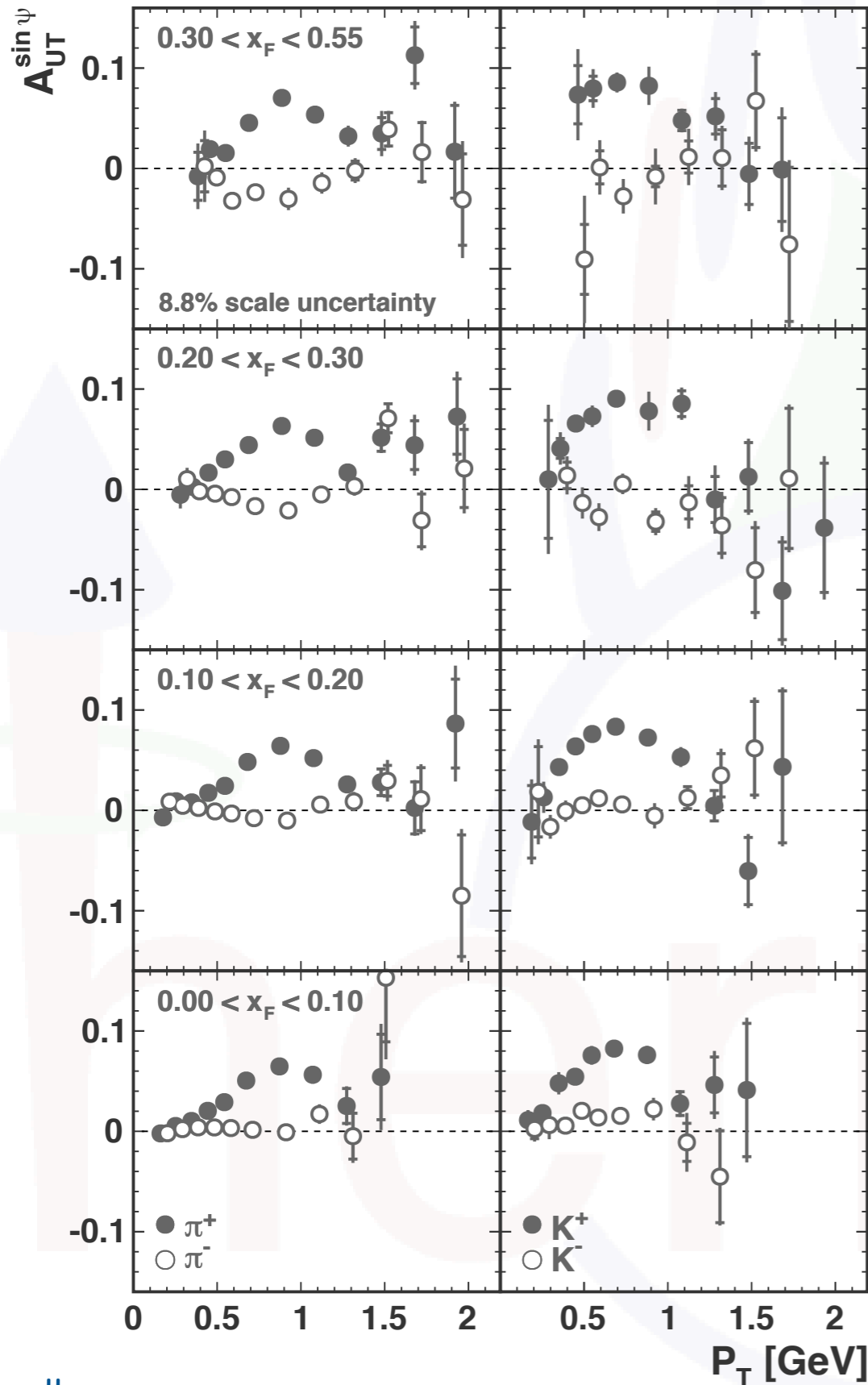
- x_F and P_T correlated
 ➔ look at 2D dependences

[Airapetian et al., Phys. Lett. B 728, 183-190 (2014)]



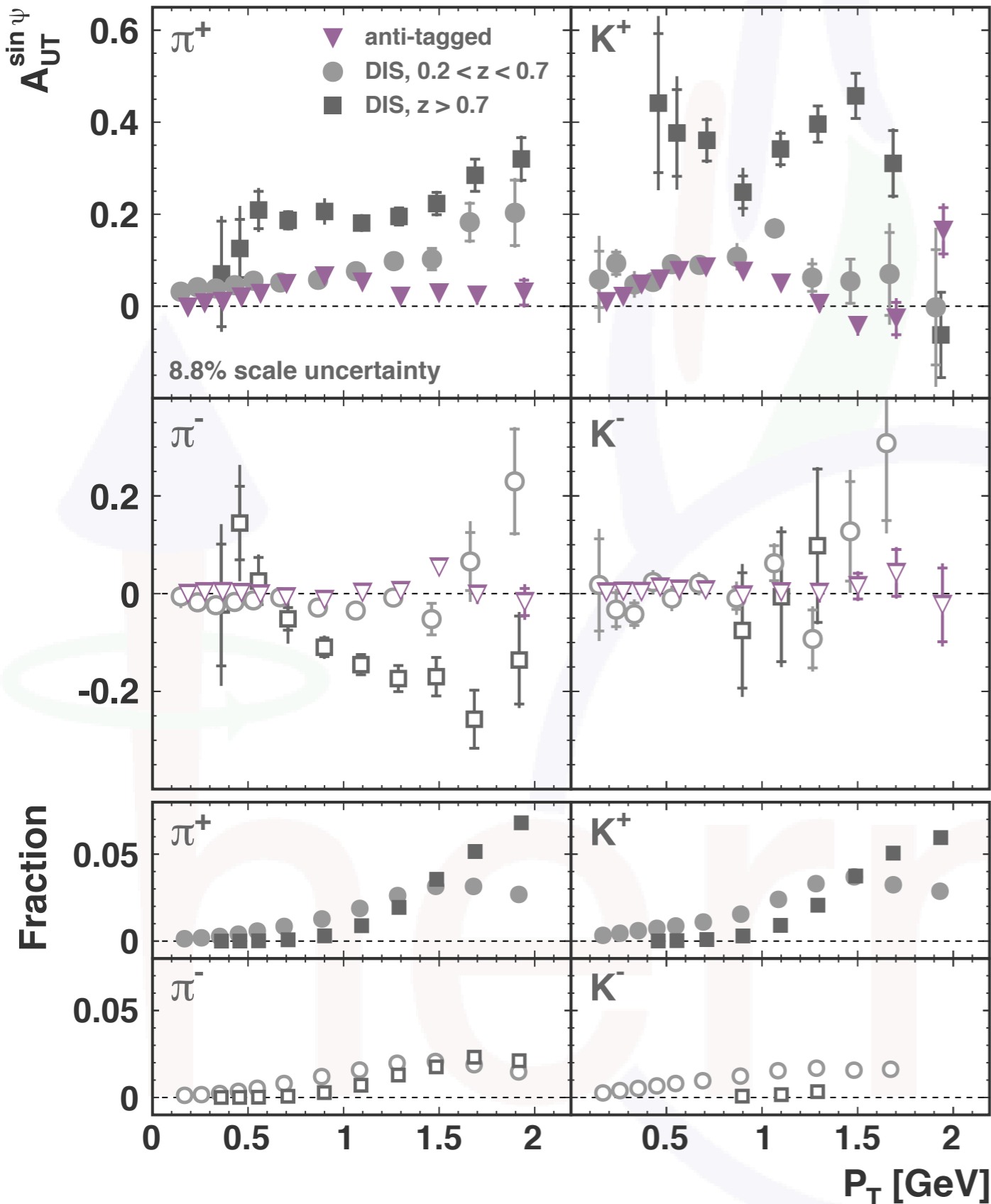
Inclusive hadrons: 2D dependences

[Airapetian et al., Phys. Lett. B 728, 183-190 (2014)]



Asymmetries of subprocesses

[Airapetian et al., Phys. Lett. B 728, 183-190 (2014)]



“anti-tagged”
no lepton in
acceptance

DIS
 $0.2 < z < 0.7$

DIS $z > 0.7$

- at large P_T significant contribution from DIS events ($Q^2 > 1$)
- asymmetries increase with larger z
- large asymmetries also for π^- in case of $z > 0.7$