

# Search for a two-photon exchange contribution to inclusive deep-inelastic scattering

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- Motivation: 2- $\gamma$  exchange in elastic ep scattering
- 2- $\gamma$  exchange in inelastic ep scattering:  $e + p^{\uparrow} \rightarrow e' + X$   
with transversely polarised target
- HERMES
- Results

# 2- $\gamma$ exchange in elastic ep scattering

● 2- $\gamma$  exchange best candidate to explain discrepancy in measurements of  $\mu_p G_E^p/G_M^p$  [Rosenbluth  $\leftrightarrow$  polarisation transfer (PT)]

A.J.R. Puckett et al., Phys. Rev. Lett. 104 (2010) 242301

● 2- $\gamma$  exchange has sizeable effect on parity-violating asymmetry

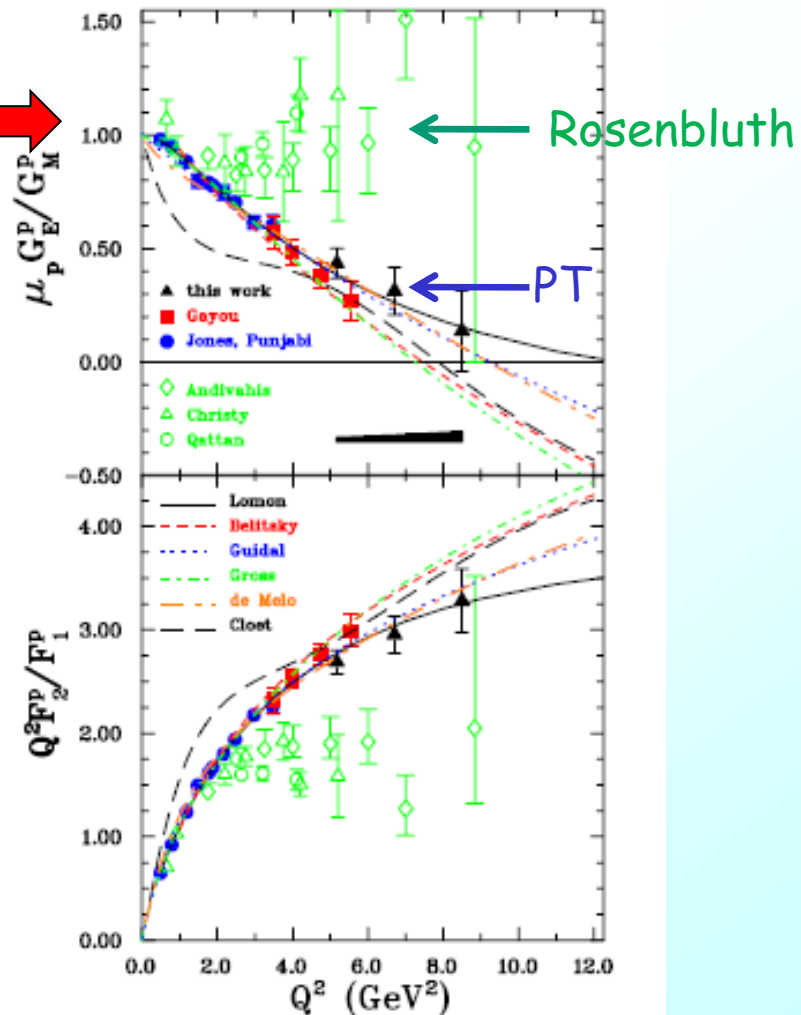
(A.V. Afanasev and C.E. Carlson, Phys. Rev. Lett. 94 (2005) 212301)

● 2- $\gamma$  exchange produces transverse-beam single-spin asymmetries (SSA) of order  $10^{-5}$  -  $10^{-6}$  in  $e^\uparrow p \rightarrow e'p'$

**SAMPLE**: S.P. Wells et al., Phys. Rev. C 63 (2001) 064001

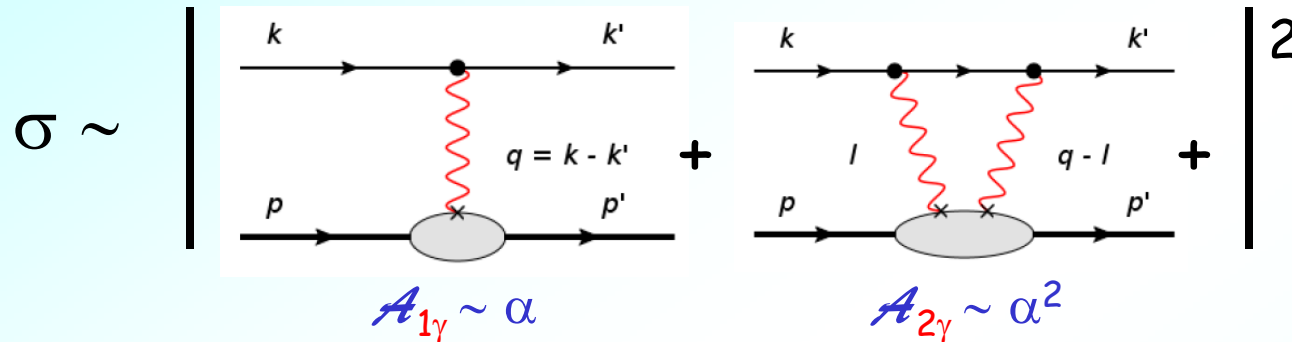
**A4**: F. Maas et al., Phys. Rev. Lett. 94 (2005) 082001

**GO**: D.S. Armstrong et al., Phys. Rev. Lett. 99 (2007) 092301



# 2- $\gamma$ exchange in elastic ep scattering

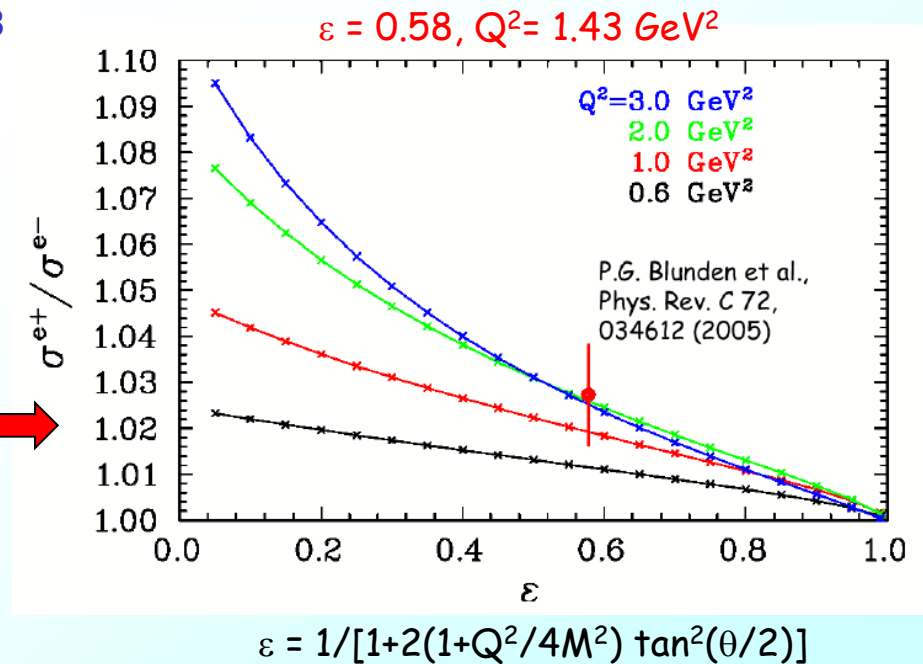
Precise determination of 2- $\gamma$ -exchange contribution:  $\sigma(e^+p)/\sigma(e^-p)$   
 CLAS, VEPP-3, OLYMPUS



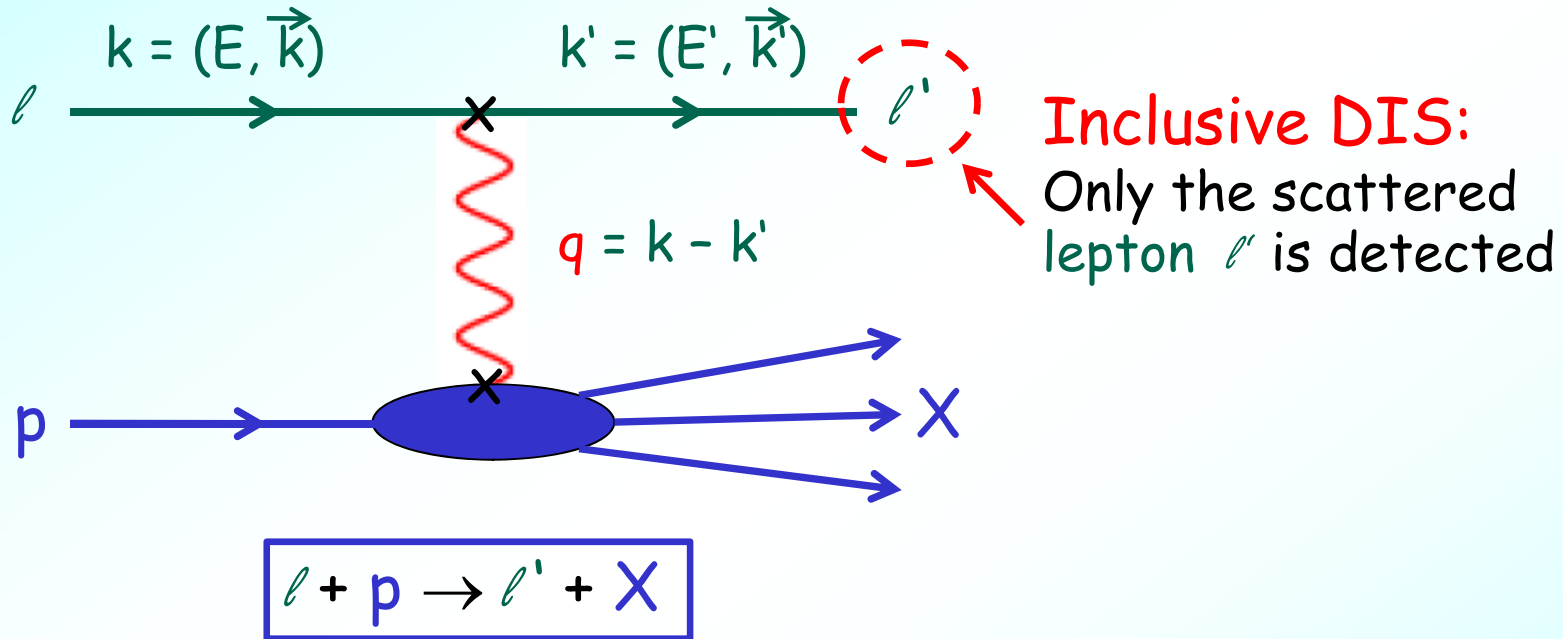
Interference term  $\sim A_{1\gamma} \times A_{2\gamma} \sim e_f \alpha^3$

$$\sigma(e^+p)/\sigma(e^-p) \approx 1 + 2 A_{2\gamma}/A_{1\gamma}$$

VEPP-3, preliminary data  
 Presented by D. Nikolenko at ICHEP 2010



# Inelastic ep scattering



■  $Q^2 = -q^2 = 2EE'(1 - \cos\theta)$   
negative squared four-momentum transfer

■  $x = Q^2 / 2M(E - E')$   
Bjorken scaling variable

# 2- $\gamma$ exchange in inelastic ep scattering

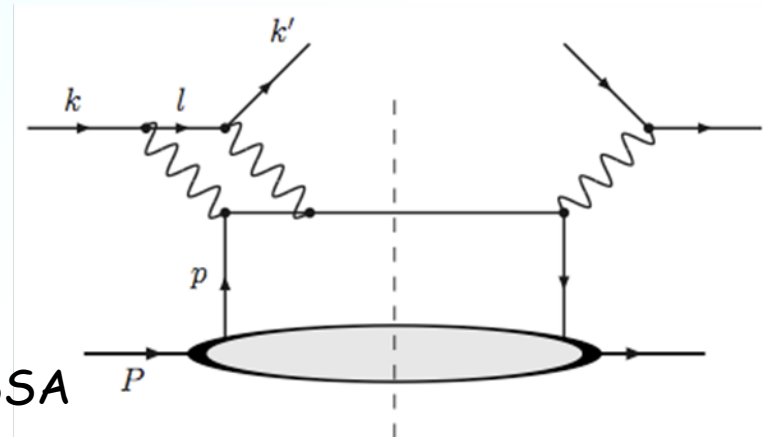
- **1- $\gamma$  exchange approximation:**

Forbids any SSA in inclusive DIS due to parity and time reversal invariance

- **2- $\gamma$  exchange contribution:**

Leads to a transverse-target and/or beam SSA arising from the interference of 1- $\gamma$  and 2- $\gamma$  exchange

(A. Metz, M. Schlegel and K. Goetze, Phys. Lett. B643 (2006) 319,  
A. Afanasev, M. Strikman, and C. Weiss, Phys. Rev. D77 (2008) 014028)



$$\sigma_{UT} \propto e_f \alpha \frac{m_{pol}}{Q} \epsilon_{\mu\nu\rho\sigma} S^\mu p^\nu k^\rho k'^\sigma C_T$$

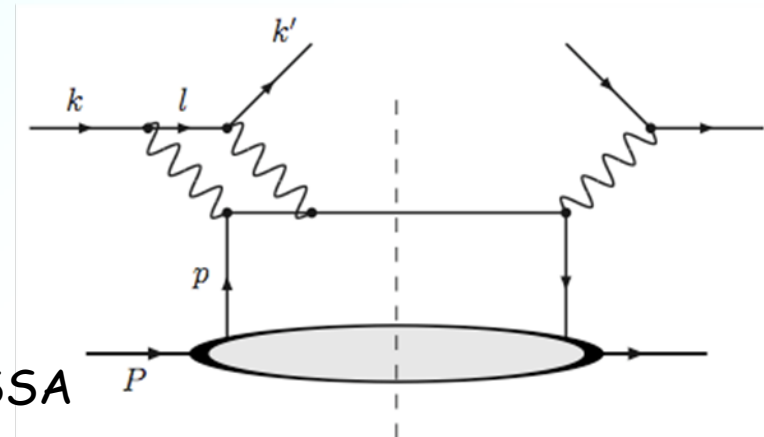
Spin dependent part of cross section for Unpolarised beam on a Transversely polarized target

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$$\sigma_{UT} \propto e_f \alpha \frac{m_{pol}}{Q} \epsilon_{\mu\nu\rho\sigma} S^\mu p^\nu k^\rho k'^\sigma C_T$$

Annotations in the diagram:  
 -  $e_f$ : beam charge  
 -  $m_{pol}$ : mass and spin of polarized particle  
 -  $S^\mu$ : mass and spin of polarized particle  
 -  $C_T$ : Higher-twist term

Spin dependent part of cross section for Unpolarised beam on a Transversely polarized target

Higher-twist term arising from q-q and q-g-q correlations, presently not yet calculable

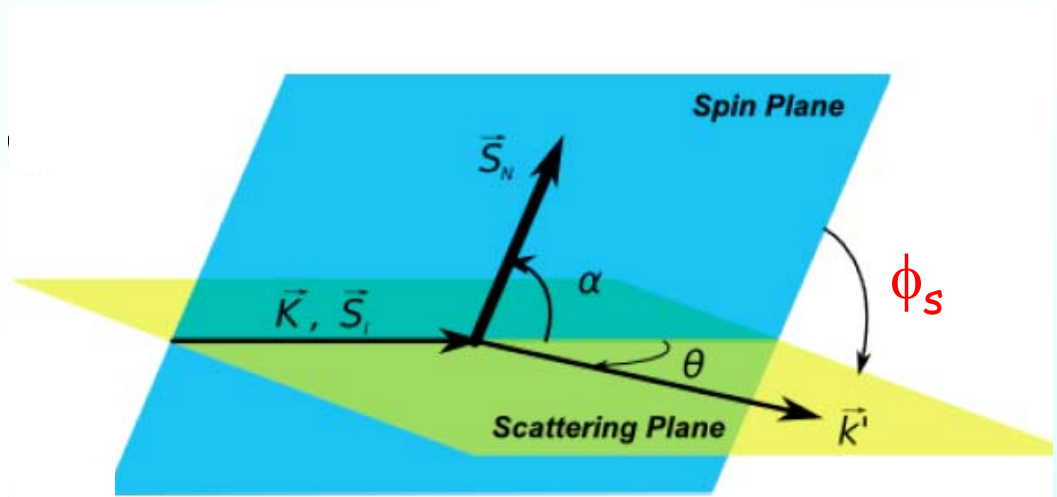
non-zero azimuthal SSA

# 2- $\gamma$ exchange in inelastic ep scattering

$$\sigma_{UT} \propto e_l \alpha (m_{\text{pol}}/Q) \left[ \varepsilon_{\mu\nu\rho\sigma} S^\mu p^\nu k^\rho k'^\sigma \right] C_T$$

$$\propto \vec{S} \cdot (\vec{k} \times \vec{k}')$$

Largest asymmetry when the spin vector  $\vec{S}$  is perpendicular to the lepton scattering plane



Azimuthal asymmetry:

Asymmetry amplitude

$$A(x, Q^2, \phi_s) = \frac{\sigma_{UT}(x, Q^2, \phi_s)}{\sigma_{UU}(x, Q^2)} = [A_{UT} \sin\phi_s(x, Q^2)] \sin\phi_s$$

$$A_N = \frac{\sigma_L - \sigma_R}{\sigma_L + \sigma_R} = \frac{2}{\pi} A_{UT} \sin\phi_s \text{ (left-right asymmetry)}$$



# 2- $\gamma$ exchange in inelastic ep scattering

## Early measurements:



$$\sigma(e^+p)/\sigma(e^-p)$$

DESY: H. Jostlein et al., Phys. Lett. B 52 (1974) 485  
DESY: S. Hartwig et al., Lett. Nuovo Cim. 15 (1976) 429  
SLAC: D.L. Fancher et al., Phys. Rev. Lett. 37 (1976) 1323  
SLAC: L.S. Rochester et al., Phys. Rev. Lett. 36 (1976) 1284  
DESY: S. Hartwig et al., Phys. Lett. B 82 (1979) 297



$$\sigma(\mu^+p)/\sigma(\mu^-p)$$

CERN-BCDMS: A. Argento et al., Phys. Lett. B 140 (1984) 142  
CERN-EMC: J.J. Aubert et al., Nucl. Phys. B 272 (1986) 158



$$A_N \text{ in } e^- + p^\uparrow \rightarrow e^- + X$$

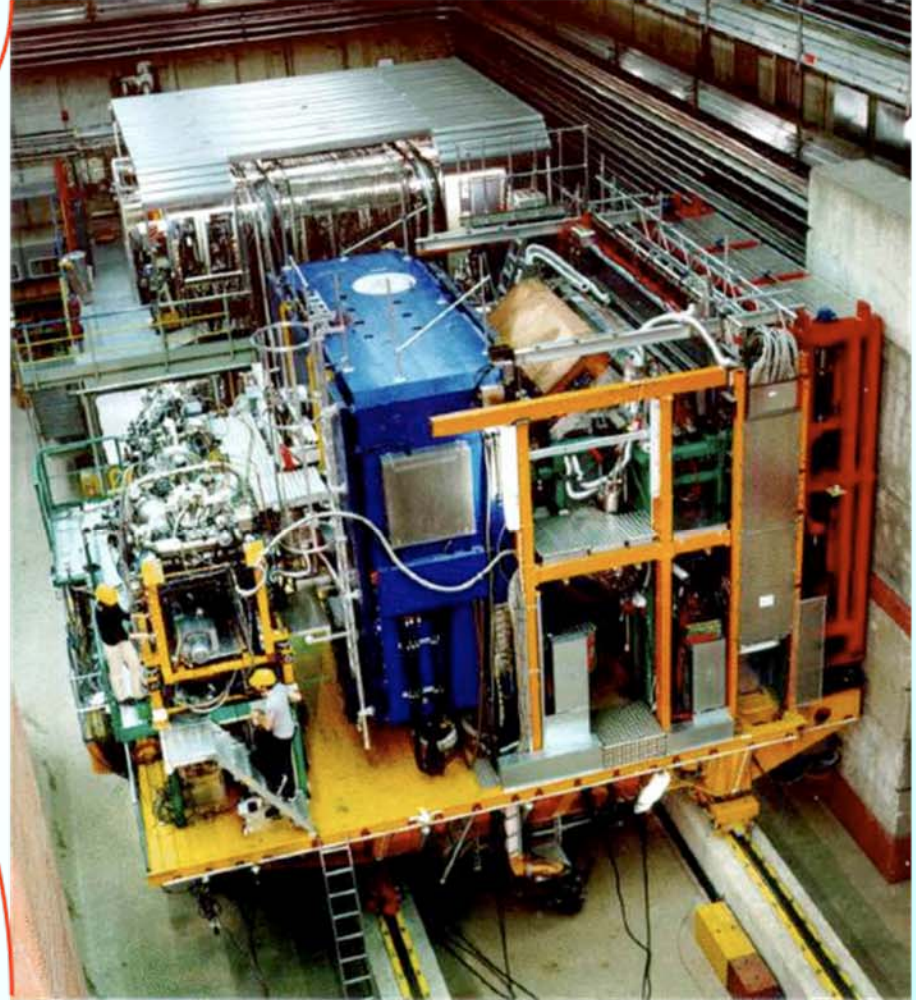
(region of nucleon resonances; search for violation of time-reversal invariance)

CEA: J.A. Appel et al., Phys. Rev. D 1 (1970) 1285  
CEA: J.R. Chen et al., Phys. Rev. Lett. 21 (1968) 1279  
SLAC: S. Rock et al., Phys. Rev. Lett. 24 (1970) 748

No effect within statistical accuracy of a few percent

Data taking: 1995-2007

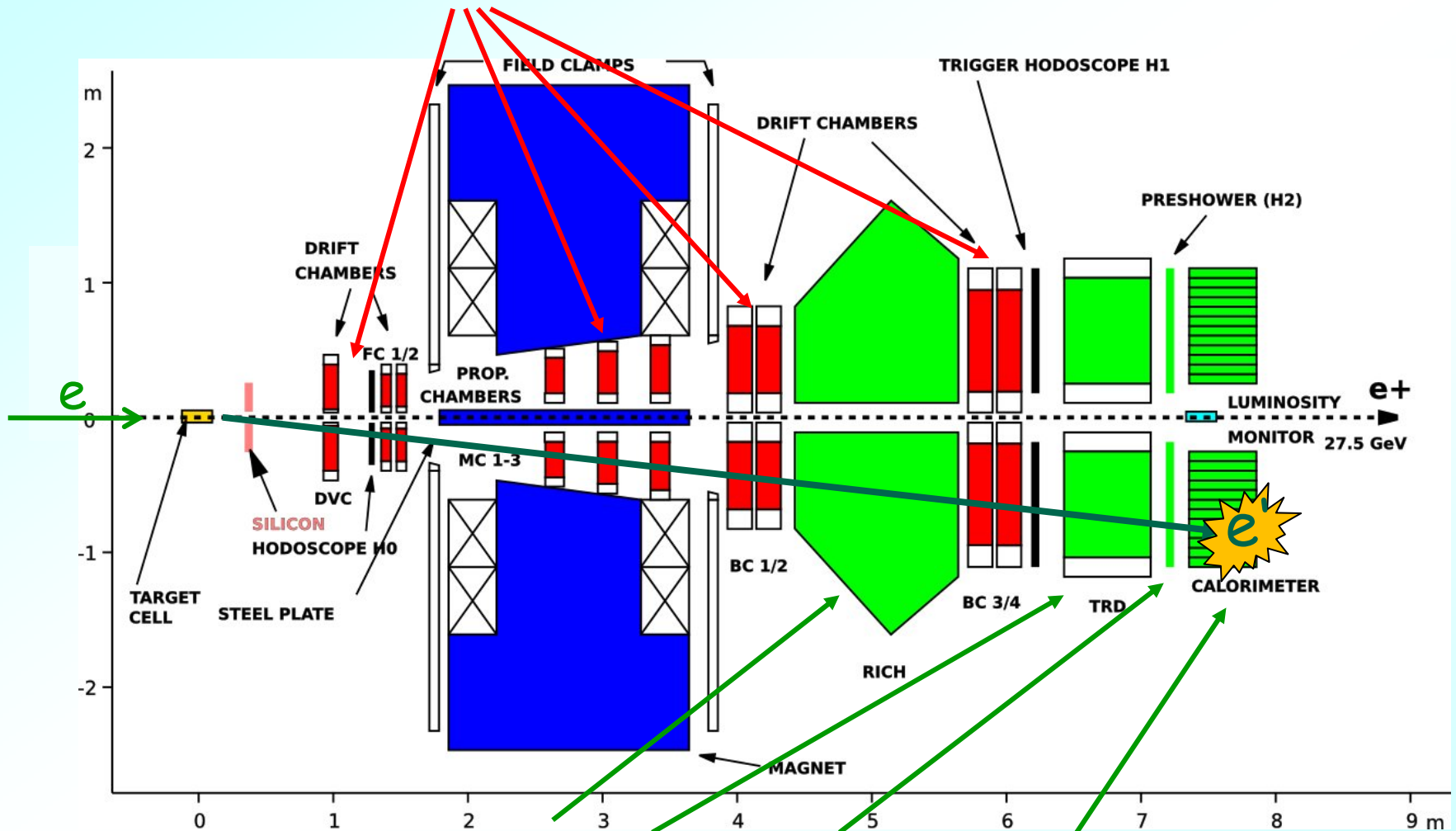
27.6 GeV  $e^+/e^-$  beam of HERA  
polarisation  $\leq 60\%$



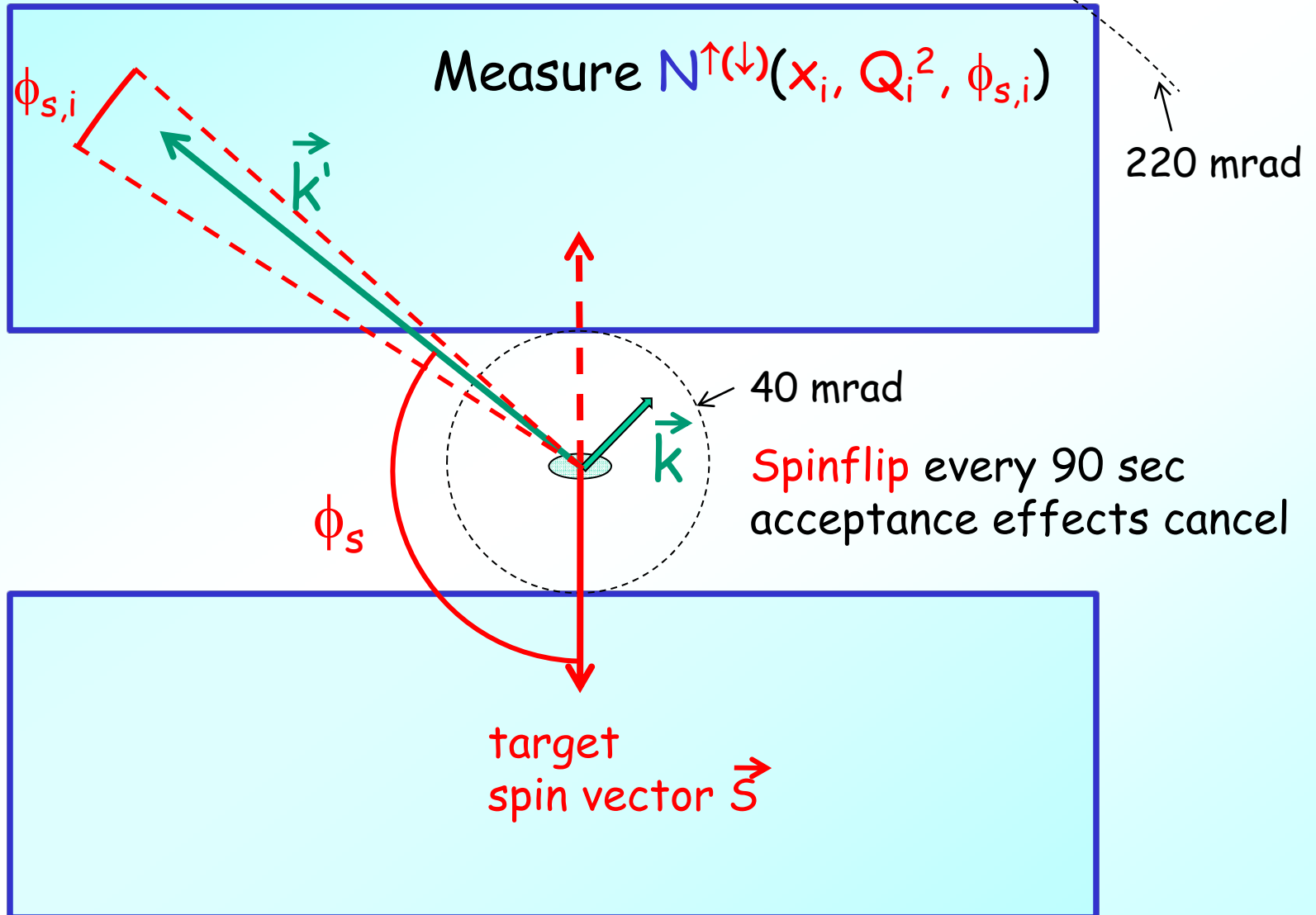
Internal gas targets

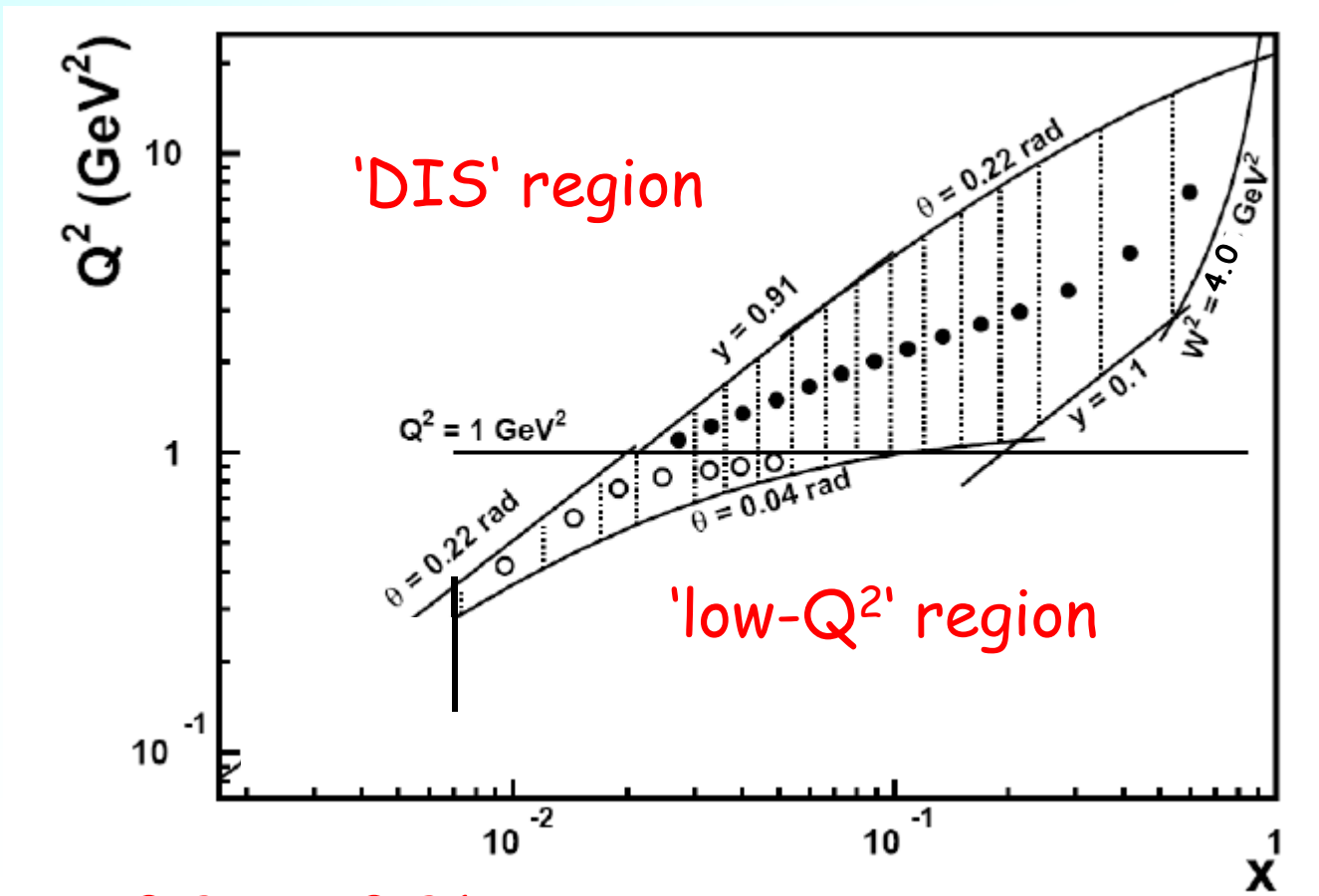
polarized :  $^1\text{H}$ ,  $^1\text{H}\uparrow$ ,  $^2\text{H}$ ,  $^3\text{He}$   
 unpolarized:  $^1\text{H}$ ,  $^2\text{H}$ ,  $^3\text{He}$ ,  $^4\text{He}$ ,  
 $\text{N}$ ,  $\text{Ne}$ ,  $\text{Kr}$ ,  $\text{Xe}$

- tracking:  $\delta p/p \sim 2\%$ ,  $\delta\Theta < 0.6$  mrad, 40-220 mrad



- PID: RICH, TRD, Preshower, Calorimeter  
hadron contamination in lepton sample for these data  $< 10^{-4}$ ,  
with lepton detection efficiency of 94%





$0.007 < x < 0.9, y < 0.91$   
 $0.25 \text{ GeV}^2 < Q^2 < 20 \text{ GeV}^2$   
 $W^2 > 4 \text{ GeV}^2$

# Single-spin azimuthal asymmetry

$$\begin{aligned}
 \frac{d^3 N^{\uparrow(\downarrow)}}{d\mathbf{x} dQ^2 d\phi_s} &= \left[ \overset{\text{luminosity}}{L^{\uparrow(\downarrow)}} d^3\sigma_{UU} + (-) \overset{\text{polarisation-weighted luminosity}}{L_P^{\uparrow(\downarrow)}} d^3\sigma_{UT} \right] \overset{\text{acceptance}}{\Omega(\mathbf{x}, Q^2, \phi_s)} \\
 &= d^3\sigma_{UU} \left[ L^{\uparrow(\downarrow)} + (-) L_P^{\uparrow(\downarrow)} A_{UT}^{\sin\phi_s}(\mathbf{x}, Q^2) \sin\phi_s \right] \Omega(\mathbf{x}, Q^2, \phi_s)
 \end{aligned}$$

$$A_{UT}(\mathbf{x}, Q^2, \phi_s) = \frac{N^{\uparrow}/L_P^{\uparrow} - N^{\downarrow}/L_P^{\downarrow}}{N^{\uparrow}/L^{\uparrow} + N^{\downarrow}/L^{\downarrow}} \cong A_{UT}^{\sin\phi_s}(\mathbf{x}, Q^2) \sin\phi_s$$

Acceptance effects cancel for  
small bin size or asymmetry



Possible apparative sources of false asymmetries:

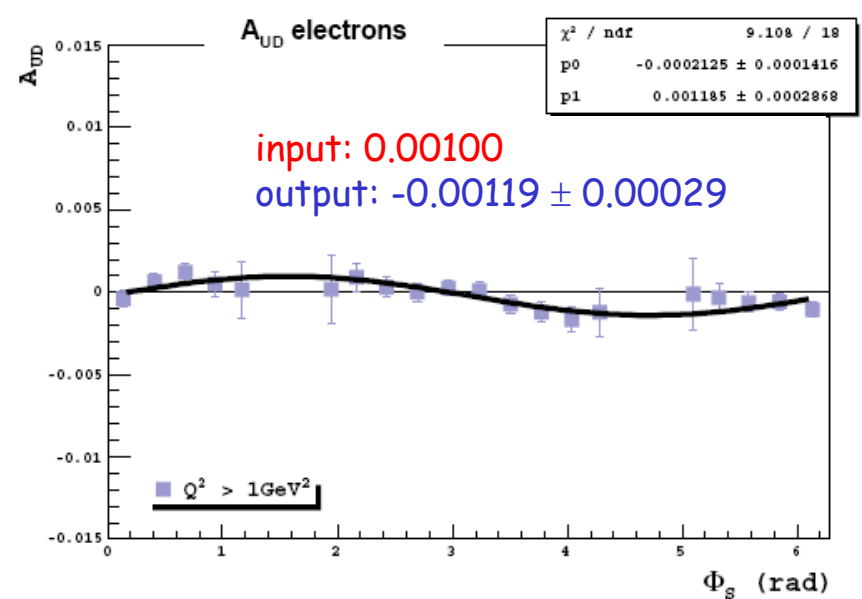
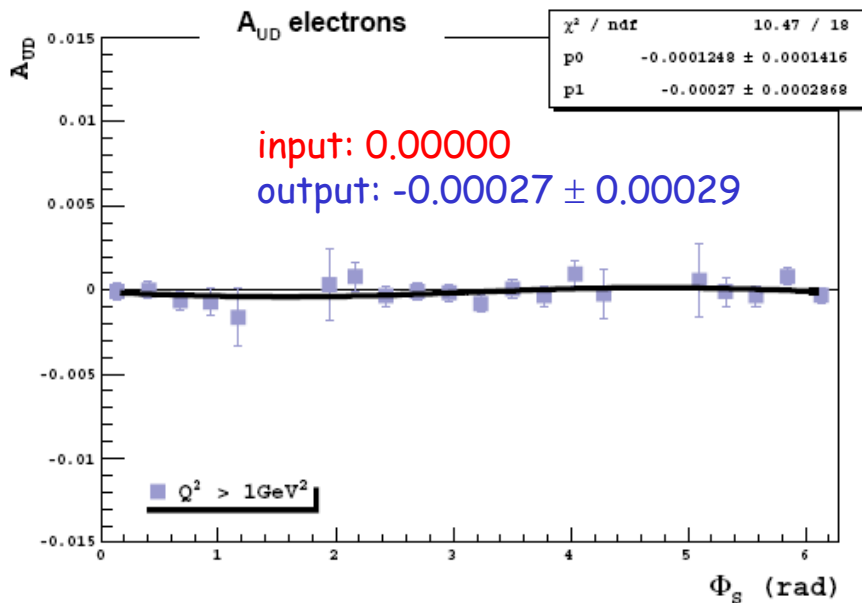
- misalignment of the detectors,
- beam position and slope at the interaction point,
- bending of  $e$  and  $e'$  in transverse holding field of target magnet



MC studies with artificial  $A_{UT} \sin\phi_S$  implemented



Examples:



Syst. uncertainty:  $\max[\delta^{MC}(\text{stat}), (\text{input}-\text{output})]$

	$\langle P \rangle$	Events
$e^+$	0.75	2.9 M
$e^-$	0.71	4.8 M

beam	$A_{UT}^{\sin\phi_S}$ $\times 10^{-3}$	$\delta A_{UT}^{\sin\phi_S}$ (stat.) $\times 10^{-3}$	$\delta A_{UT}^{\sin\phi_S}$ (syst.) $\times 10^{-3}$	$\langle x_B \rangle$	$\langle Q^2 \rangle$ [GeV <sup>2</sup> ]
$e^+$	-0.61	3.97	0.63	0.02	0.68
$e^-$	-6.55	3.40	0.63	0.02	0.68
$e^+$	-0.60	1.70	0.29	0.14	2.40
$e^-$	-0.85	1.50	0.29	0.14	2.40

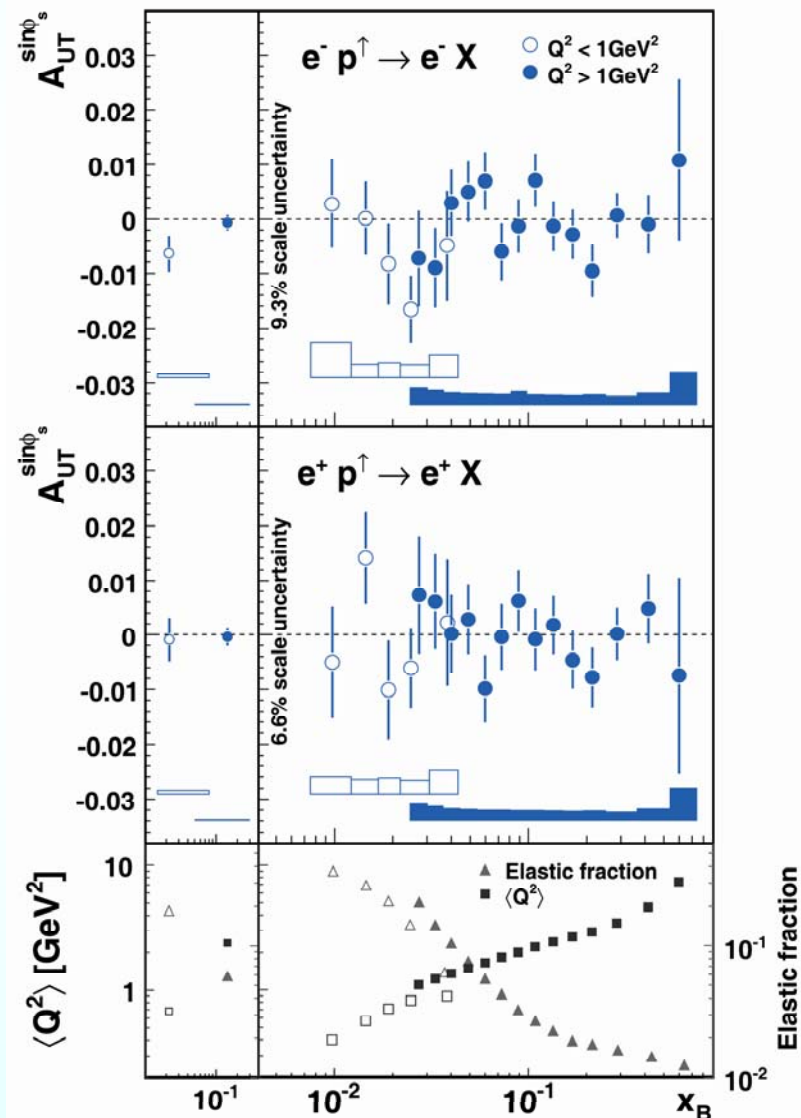
$$A_N = 2/\pi A_{UT}^{\sin\phi_S} = O(10^{-3})$$

same for  $e^+$  and  $e^-$



No sign for 2-photon exchange in DIS

PLB 682 (2010) 351





- HERMES has measured with high precision **single-spin asymmetries** in inclusive DIS  $e + p^{\uparrow} \rightarrow e' + X$  from a transversely polarised proton target
- No evidence for a **2-photon-exchange** contribution to the cross section has been observed within the experimental uncertainties of the order of  $10^{-3}$
- This sets the so far most precise limit of **2-photon exchange** in DIS