Recent Hermes results for SSAs and DSAs

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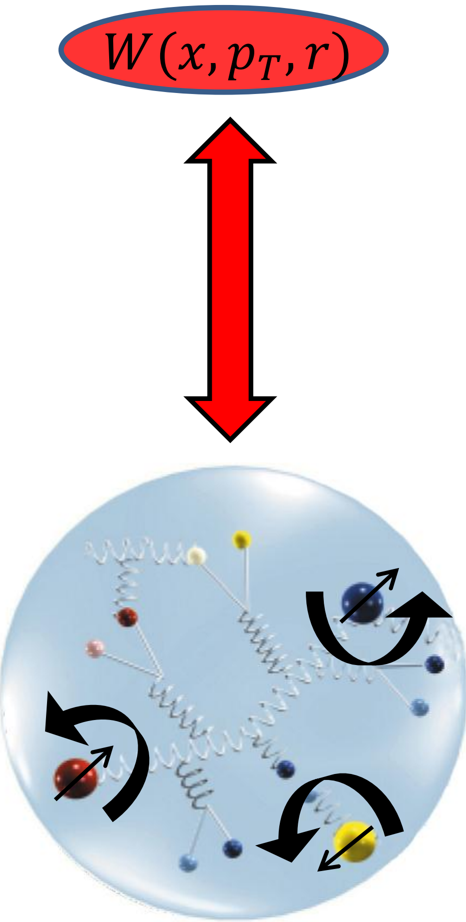
University of Ferrara

The phase-space distribution of partons

The **full phase-space distribution of the partons** encoded in the **Wigner function**



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The **full phase-space distribution of the partons** encoded in the **Wigner function**

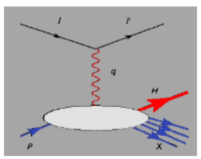
...but $\Delta x \Delta p \geq \frac{\hbar}{2}$ $\rightarrow$ no simultaneous knowledge of momentum and position

cannot be directly accessed experimentally $\rightarrow$ integrated quantities

$W(x, p_T, r)$

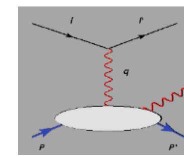
$\int d^3 r$ → $f(x, p_T)$

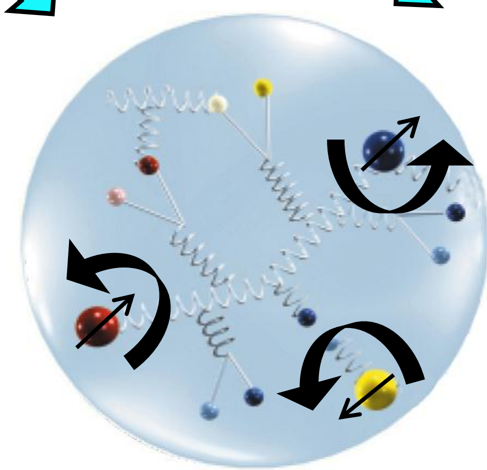
$\int d^2 p_T$ → $H(x, \xi, t)$



TMDs $f(x, p_T)$

$H(x, \xi, t)$ GPDs





The **non-collinear** structure of the nucleon

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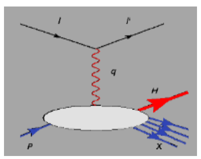
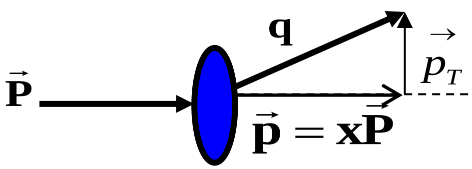
$\int d^3r$ $\qquad$ $\int d^2p_T$

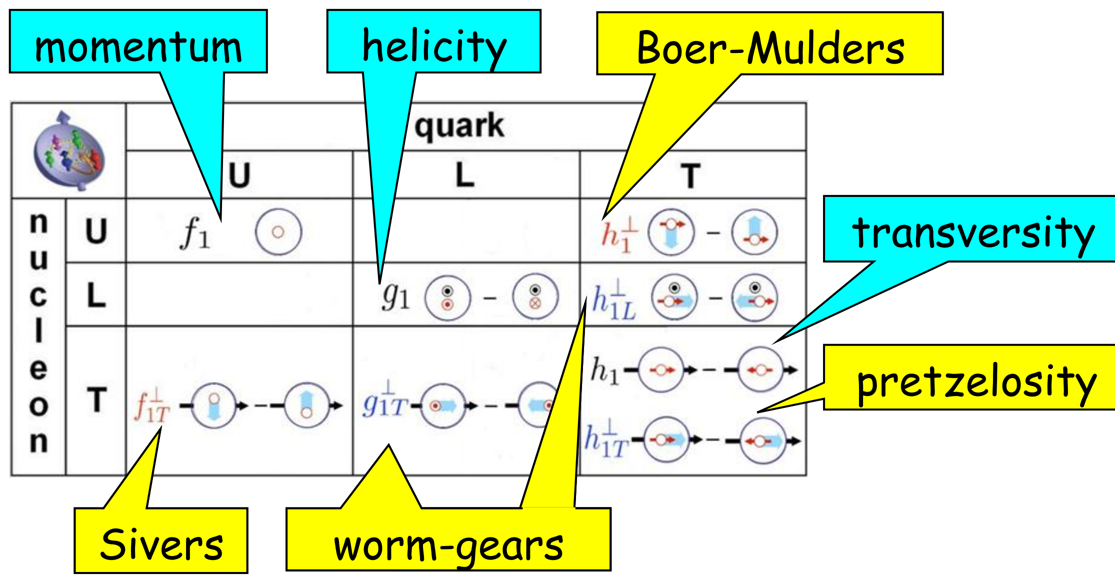


TMDs $\quad f(x, p_T)$

$\mathrm{H}(x, \xi, t)$ $\quad$ GPDs

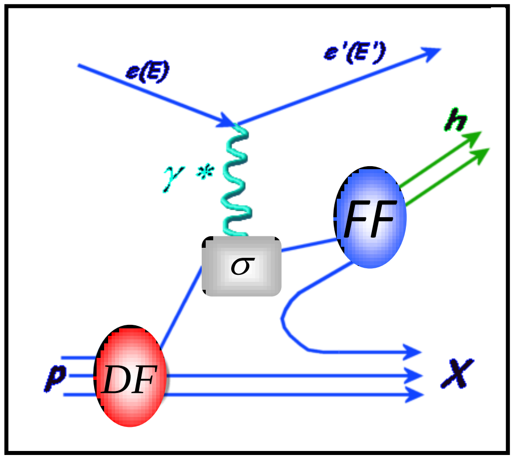


- TMDs depend on x and p_T
- Describe correlations between p_T and quark or nucleon spin (**spin-orbit correlations**)
- Provide **a 3-dim picture** of the nucleon in momentum space (**nucleon tomography**)

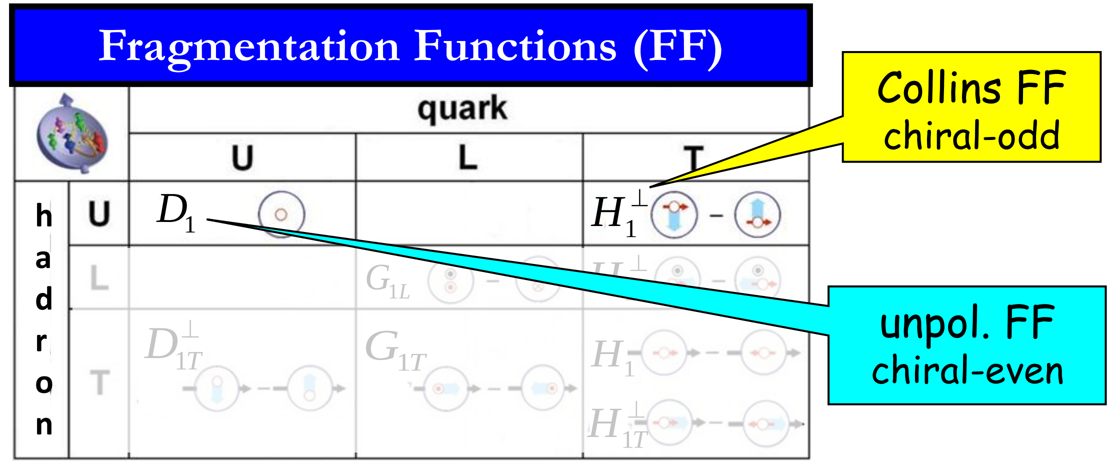


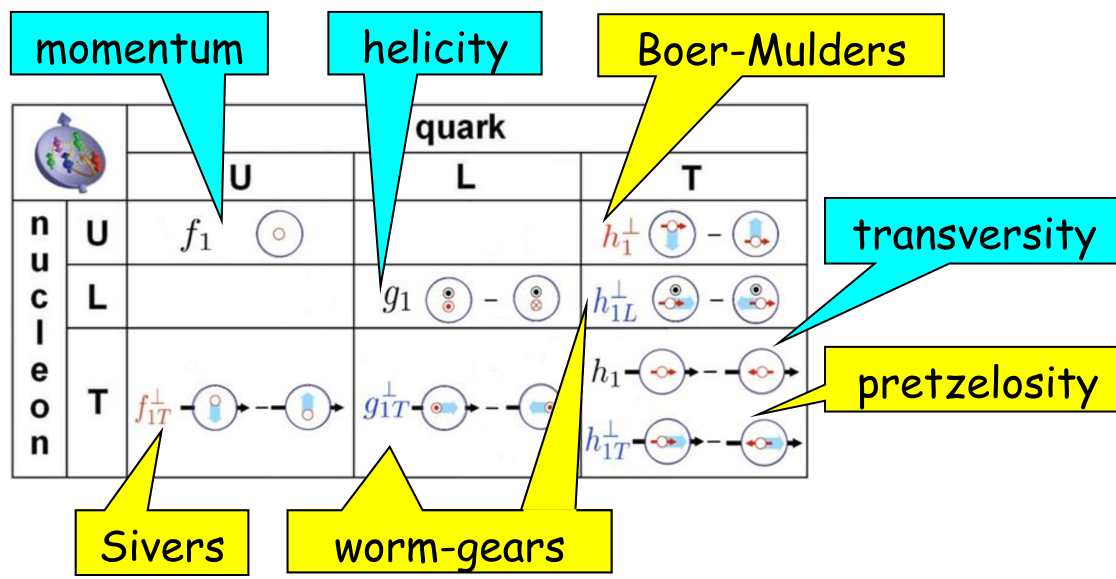
The non-collinear structure of the nucleon

Mostly investigated in **SIDIS**: detection of transverse momentum of produced hadrons gives access to p_T



$$\sigma^{ep\to ehX} = \sum_q DF \otimes \sigma^{eq\to eq} \otimes FF$$





- TMDs depend on x and p_T
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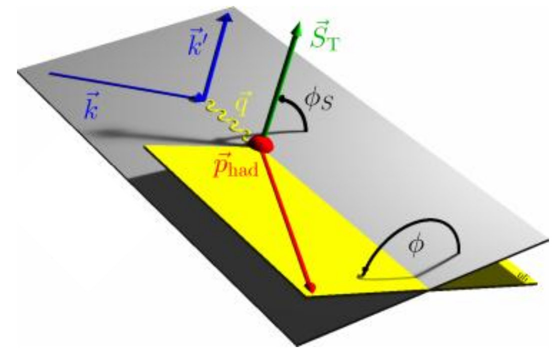
The SIDIS cross-section



$$\frac{d\sigma^h}{dx\,dy\,d\phi_S\,dz\,d\phi\,d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2}{xyQ^2}\frac{y^2}{2(1-\epsilon)}\left(1+\frac{\gamma^2}{2x}\right)$$

$$\Big\{ \Big[F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)}\cos(\phi)F_{UU}^{\cos(\phi)} + \epsilon\cos(2\phi)F_{UU}^{\cos(2\phi)}\Big]$$

$$+ \lambda_l \Big[\sqrt{2\epsilon(1-\epsilon)}\sin(\phi)F_{LU}^{\sin(\phi)}\Big]$$

$$+ S_L \Big[\sqrt{2\epsilon(1+\epsilon)}\sin(\phi)F_{UL}^{\sin(\phi)} + \epsilon\sin(2\phi)F_{UL}^{\sin(2\phi)}\Big]$$

$$+ S_L\,\lambda_l \Big[\sqrt{1-\epsilon^2}F_{LL} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi)F_{LL}^{\cos(\phi)}\Big]$$

$$+ S_T \Big[\sin(\phi-\phi_S)\left(F_{UT,T}^{\sin(\phi-\phi_S)} + \epsilon F_{UT,L}^{\sin(\phi-\phi_S)}\right) + \epsilon\sin(\phi+\phi_S)F_{UT}^{\sin(\phi+\phi_S)} + \epsilon\sin(3\phi-\phi_S)F_{UT}^{\sin(3\phi-\phi_S)} + \sqrt{2\epsilon(1+\epsilon)}\sin(\phi_S)F_{UT}^{\sin(\phi_S)} + \sqrt{2\epsilon(1+\epsilon)}\sin(2\phi-\phi_S)F_{UT}^{\sin(2\phi-\phi_S)}\Big]$$

$$+ S_T\,\lambda_l \Big[\sqrt{1-\epsilon^2}\cos(\phi-\phi_S)F_{LT}^{\cos(\phi-\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi_S)F_{LT}^{\cos(\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(2\phi-\phi_S)F_{LT}^{\cos(2\phi-\phi_S)}\Big]\Big\}$$

target polarization

$$F_{XY,Z} = F_{XY,Z}(x,y,z)$$

beam polarization

virtual photon polarization

The SIDIS cross-section

$$\frac{d\sigma^h}{dx\,dy\,d\phi_S\,dz\,d\phi\,d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2}{xyQ^2}\frac{y^2}{2(1-\epsilon)}\left(1+\frac{\gamma^2}{2x}\right)$$

$$\Big\{ \Big[F_{\mathrm{UU,T}} + \epsilon F_{\mathrm{UU,L}} + \sqrt{2\epsilon(1+\epsilon)}\cos(\phi) F_{\mathrm{UU}}^{\cos(\phi)} + \epsilon\cos(2\phi) F_{\mathrm{UU}}^{\cos(2\phi)}\Big]$$

unpolarized

$$+ \quad \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)}\sin(\phi) F_{\mathrm{LU}}^{\sin(\phi)}\right]$$

beam polarization

$$+ S_L \quad \left[\sqrt{2\epsilon(1+\epsilon)}\sin(\phi) F_{\mathrm{UL}}^{\sin(\phi)} + \epsilon\sin(2\phi) F_{\mathrm{UL}}^{\sin(2\phi)}\right]$$

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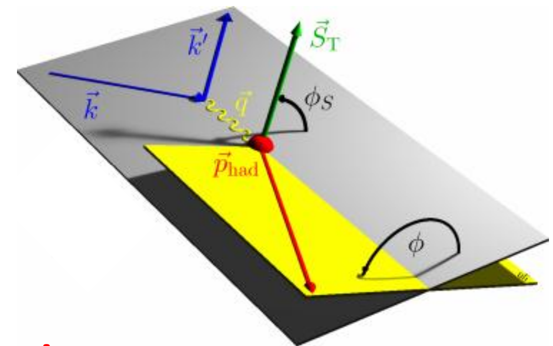
$$+ S_T\ \lambda_l \Big[\sqrt{1-\epsilon^2}\cos(\phi-\phi_S) F_{\mathrm{LT}}^{\cos(\phi-\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi_S) F_{\mathrm{LT}}^{\cos(\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(2\phi-\phi_S) F_{\mathrm{LT}}^{\cos(2\phi-\phi_S)}\Big]\Big\}$$

target polarization

beam and target polarization



target polarization

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beam polarization

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Selected **leading-twist** 1-hadron SIDIS results

Boer-Mulders function

$$\frac{d\sigma^h}{dx\,dy\,d\phi_S\,dz\,d\phi\,d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2}{xyQ^2}\frac{y^2}{2(1-\epsilon)}\left(1+\frac{\gamma^2}{2x}\right)$$

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$$+ \lambda_l \left[\sqrt{2\epsilon(1-\epsilon)}\sin(\phi) F_{\mathrm{LU}}^{\sin(\phi)}\right]$$

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$$+ S_T\,\lambda_l \left[\sqrt{1-\epsilon^2}\cos(\phi-\phi_S) F_{\mathrm{LT}}^{\cos(\phi-\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi_S) F_{\mathrm{LT}}^{\cos(\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(2\phi-\phi_S) F_{\mathrm{LT}}^{\cos(2\phi-\phi_S)}\right]\Big\}$$

$$F_{UU}^{\cos 2\phi_h} = \mathcal{C}\left[-\frac{2\,(\hat{\boldsymbol{h}}\cdot\boldsymbol{k}_T)\,(\hat{\boldsymbol{h}}\cdot\boldsymbol{p}_T) - \boldsymbol{k}_T\cdot\boldsymbol{p}_T}{M M_h} h_1^\perp H_1^\perp\right]$$

Describes correlation between quark transverse momentum and transverse spin in unpolarized nucleon

Distribution Functions

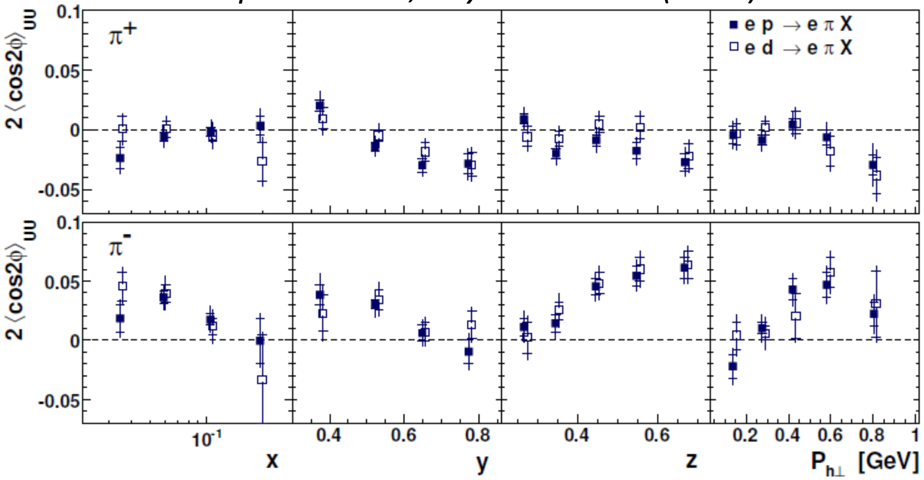
		quark U	quark L	quark T
nucleon	U	f_1		$h_1^\perp$
	L		g_1	$h_{1L}^\perp$
	T	$f_{1T}^\perp$	$g_{1T}^\perp$	h_1, $h_{1T}^\perp$

Fragmentation Functions

		quark U	quark L	quark T
h	U	D_1		$H_1^\perp$

The cos2ϕ amplitudes $\propto h_1^{\perp}(x, p_T^2) \otimes H_1^{\perp}(z, k_T^2)$

A. Airapetian et al, Phys. Rev. D 87 (2013) 012010



negative

positive

- Amplitudes are significant
→ clear evidence of BM effect

- similar results for H & D indicate $h_1^{\perp,u} \approx h_1^{\perp,d}$

- Opposite sign for π^+/π^- consistent with opposite signs of fav/unfav Collins

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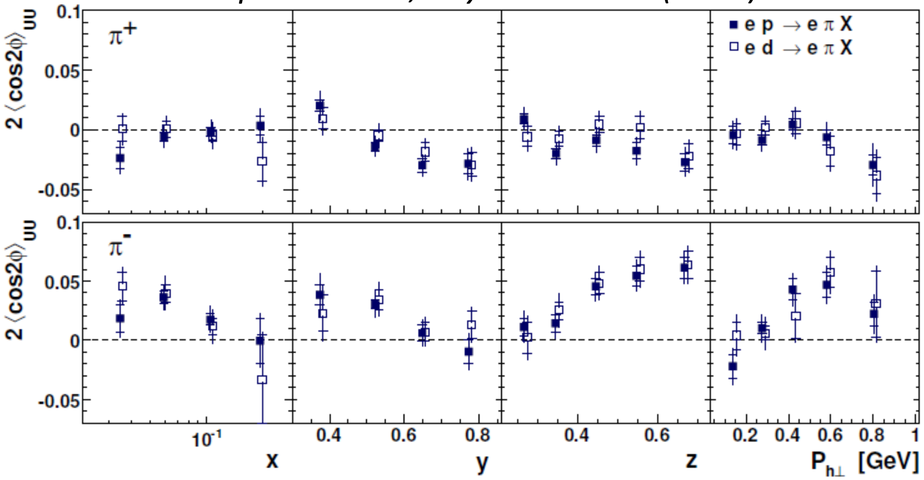

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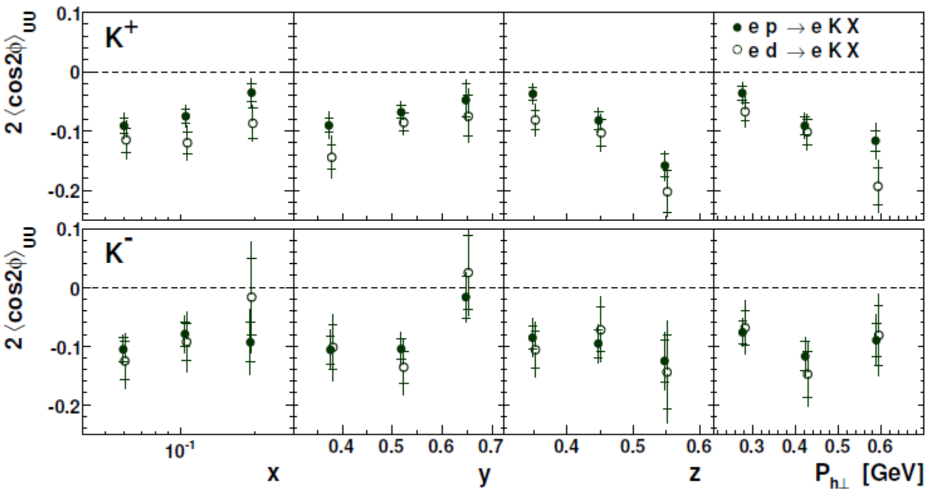

Large and negative

Large and negative

- K^+/K^- amplitudes are larger than for pions , have different kinematic dependencies than pions and have same sign
- different role of Collins FF for pions and kaons?
- Significant contribution from scattering off strange quarks?

Analysis multi-dimensional in x, y, z,and Pt

Create your own projections of results through: http://www-hermes.desy.de/cosnphi/

Transversity

$$F_{UT}^{\sin(\phi_h+\phi_S)} = \mathcal{C}\left[-\frac{\hat{\boldsymbol{h}}\cdot\boldsymbol{k}_T}{M_h} h_1 H_1^{\perp}\right]$$

$$\frac{d\sigma^h}{dx\,dy\,d\phi_S\,dz\,d\phi\,d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2}{xyQ^2}\frac{y^2}{2(1-\epsilon)}\left(1+\frac{\gamma^2}{2x}\right)$$

$$\Big\{ \Big[F_{\mathrm{UU,T}} + \epsilon F_{\mathrm{UU,L}} + \sqrt{2\epsilon(1+\epsilon)}\cos(\phi) F_{\mathrm{UU}}^{\cos(\phi)} + \epsilon\cos(2\phi) F_{\mathrm{UU}}^{\cos(2\phi)}\Big]$$

$$+ \lambda_l \Big[\sqrt{2\epsilon(1-\epsilon)}\sin(\phi) F_{\mathrm{LU}}^{\sin(\phi)}\Big]$$

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Describes probability to find transversely polarized quarks in a transversely polarized nucleon

Distribution Functions

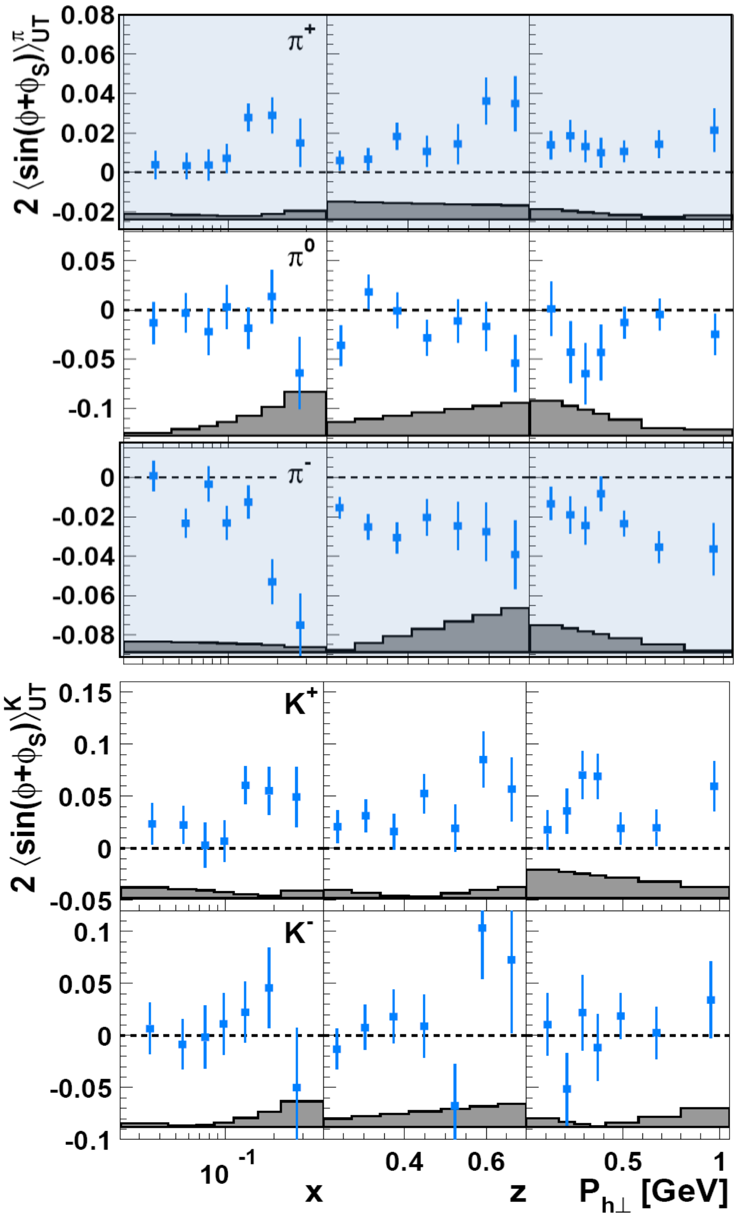
		quark U	quark L	quark T
nucleon	U	f_1		$h_1^{\perp}$
	L		g_1	$h_{1L}^{\perp}$
	T	$f_{1T}^{\perp}$	$g_{1T}^{\perp}$	h_1, $h_{1T}^{\perp}$

Fragmentation Functions

		quark U	quark L	quark T
h	U	D_1		$H_1^{\perp}$

Collins amplitudes $\propto h_1(x, p_T^2) \otimes H_1^{\perp}(z, k_T^2)$

[Airapetian *et al.*, Phys. Lett. B 693 (2010) 11-16]



- positive
- consistent with zero (isospin-symmetry)
- large and negative!

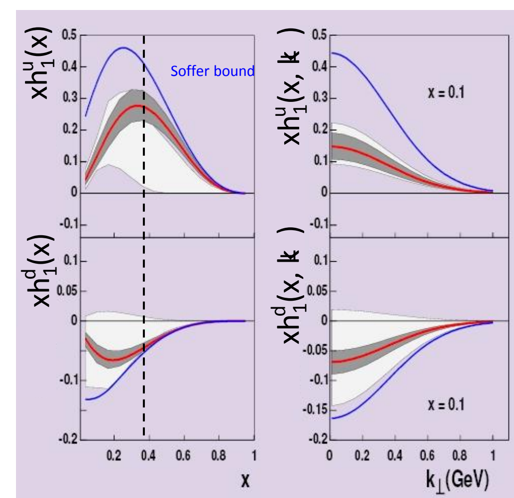
$$\begin{cases} u \to \pi^- \\ d \to \pi^+ \end{cases} \qquad \begin{cases} u \to \pi^+ \\ d \to \pi^- \end{cases}$$

$$H_1^{\perp,unfav}(z) \approx -H_1^{\perp,fav}(z)$$

Consistent with Belle/BaBar measurements in e^+e^-

$$e^+e^- \to \pi^+_{jet1}\, \pi^-_{jet2}\, X$$

- significantly positive
- consistent with zero



Anselmino et al. Phys. Rev. D 75 (2007)

hermes

BELLE

Sivers function

$$F_{UT,T}^{\sin(\phi_h-\phi_S)} = \mathcal{C}\left[-\frac{\hat{h}\cdot p_T}{M} f_{1T}^{\perp} D_1\right]$$

Describes correlation between quark transverse momentum and nucleon transverse polarization

$$\frac{d\sigma^h}{dx\,dy\,d\phi_S\,dz\,d\phi\,d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2}{xyQ^2}\frac{y^2}{2(1-\epsilon)}\left(1+\frac{\gamma^2}{2x}\right)$$
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$$+ \lambda_l \Big[\sqrt{2\epsilon(1-\epsilon)}\sin(\phi)F_{\mathrm{LU}}^{\sin(\phi)}\Big]$$
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$$+ S_L\,\lambda_l \Big[\sqrt{1-\epsilon^2}F_{\mathrm{LL}} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi)F_{\mathrm{LL}}^{\cos(\phi)}\Big]$$
$$+ S_T \Big[\sin(\phi-\phi_S)\left(F_{\mathrm{UT,T}}^{\sin(\phi-\phi_S)} + \epsilon F_{\mathrm{UT,L}}^{\sin(\phi-\phi_S)}\right) + \epsilon\sin(\phi+\phi_S)F_{\mathrm{UT}}^{\sin(\phi+\phi_S)} + \epsilon\sin(3\phi-\phi_S)F_{\mathrm{UT}}^{\sin(3\phi-\phi_S)}$$
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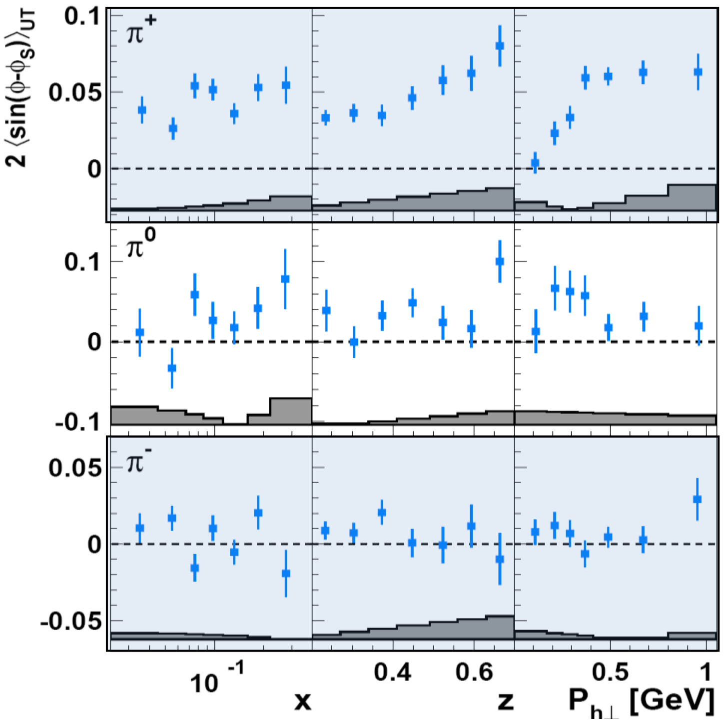
Fragmentation Functions

h \ quark	U	L	T
U	D_1		$H_1^{\perp}$

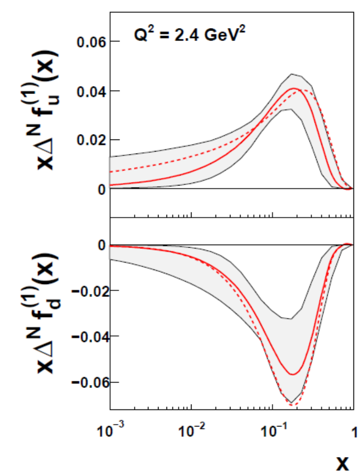
Sivers amplitudes $\propto f_{1T}^{\perp}(x, p_T^2) \otimes D_1(z, k_T^2)$

[Airapetian *et al.*, Phys. Rev. Lett. 103 (2009) 152002]



- **significantly positive**
- **slightly positive** (isospin-symmetry)
- **consistent with zero**

consistent with Sivers func. of opposite sign for u and d quarks



[Anselmino *et al.*,Eur.Phys.J.A3,2009]

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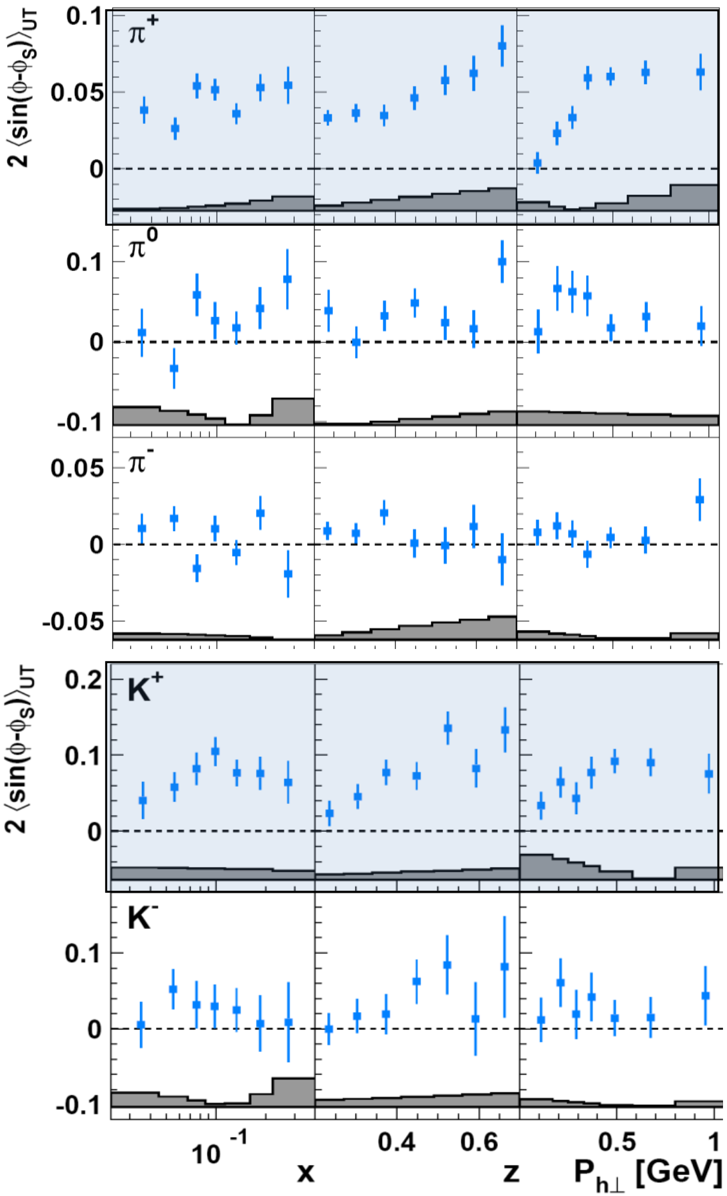

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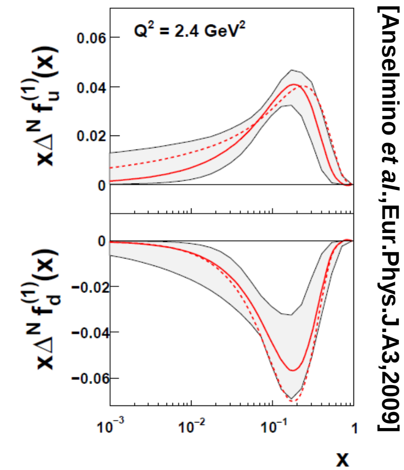
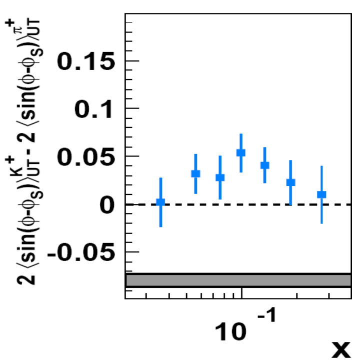
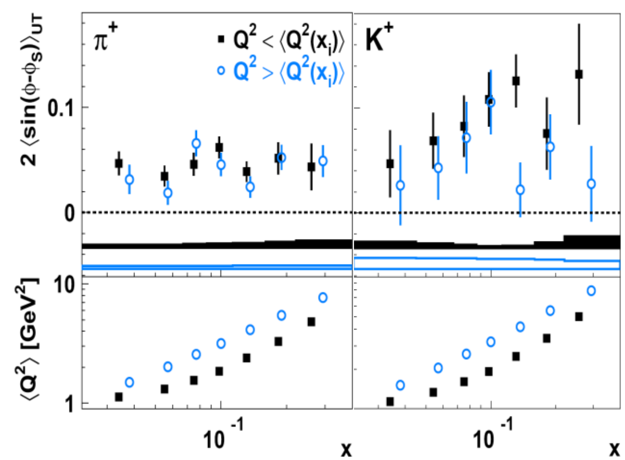


- **Larger than π^+!!**
  - role of sea quarks ?

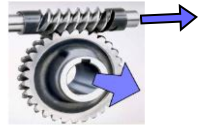


diff. comes from low Q^2



- Higher-twist contrib for K^+ ?

Worm-gear $g^{\perp}_{1T}$



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$$\Big\{ \Big[F_{UU,T}+\epsilon F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)}\cos(\phi)F_{UU}^{\cos(\phi)} + \epsilon\cos(2\phi)F_{UU}^{\cos(2\phi)}\Big]$$
$$+\ \lambda_l\Big[\sqrt{2\epsilon(1-\epsilon)}\sin(\phi)F_{LU}^{\sin(\phi)}\Big]$$
$$+\ S_L\Big[\sqrt{2\epsilon(1+\epsilon)}\sin(\phi)F_{UL}^{\sin(\phi)} + \epsilon\sin(2\phi)F_{UL}^{\sin(2\phi)}\Big]$$
$$+\ S_L\,\lambda_l\Big[\sqrt{1-\epsilon^2}F_{LL} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi)F_{LL}^{\cos(\phi)}\Big]$$
$$+\ S_T\Big[\sin(\phi-\phi_S)\Big(F_{UT,T}^{\sin(\phi-\phi_S)}+\epsilon F_{UT,L}^{\sin(\phi-\phi_S)}\Big) + \epsilon\sin(\phi+\phi_S)F_{UT}^{\sin(\phi+\phi_S)} + \epsilon\sin(3\phi-\phi_S)F_{UT}^{\sin(3\phi-\phi_S)} + \sqrt{2\epsilon(1+\epsilon)}\sin(\phi_S)F_{UT}^{\sin(\phi_S)} + \sqrt{2\epsilon(1+\epsilon)}\sin(2\phi-\phi_S)F_{UT}^{\sin(2\phi-\phi_S)}\Big]$$
$$+\ S_T\,\lambda_l\Big[\sqrt{1-\epsilon^2}\cos(\phi-\phi_S)F_{LT}^{\cos(\phi-\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi_S)F_{LT}^{\cos(\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(2\phi-\phi_S)F_{LT}^{\cos(2\phi-\phi_S)}\Big]\Big\}$$

Describes the probability to find longitudinally polarized quarks in a transversely polarized nucleon!

- requires interference between wave funct. components that differ by 1 unit of **OAM**
- Can be accessed in **LT DSAs**

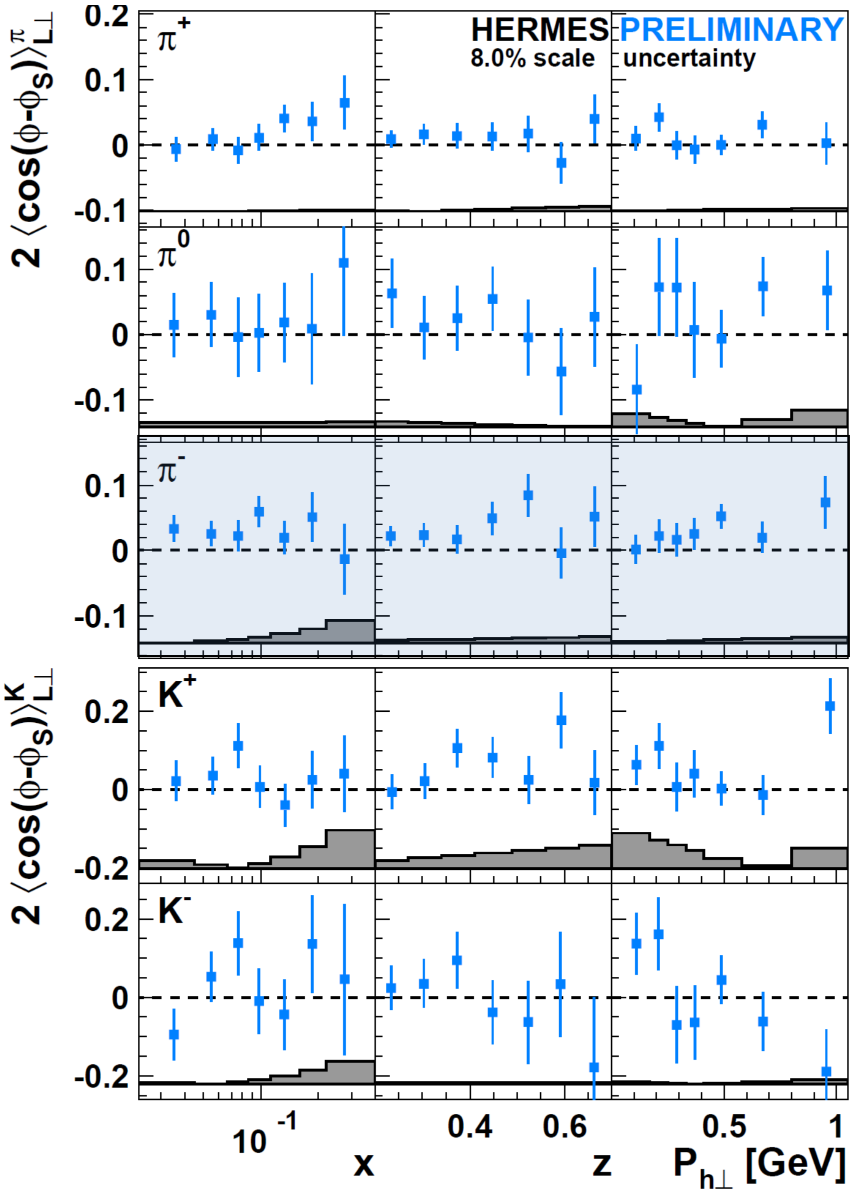
Distribution Functions

nucleon \ quark	U	L	T
U	f_1		$h_1^{\perp}$
L		g_1	$h_{1L}^{\perp}$
T	$f_{1T}^{\perp}$	$g_{1T}^{\perp}$	h_1, $h_{1T}^{\perp}$

Fragmentation Functions

h \ quark	U	L	T
U	D_1		$H_1^{\perp}$

The cos(ϕ-ϕ$_S$) amplitudes $\propto g_{1T}^{\perp}(x, p_T^2) \otimes D_1(z, k_T^2)$



- slightly positive ?
- consistent with zero
- positive!!
- slightly positive ?
- consistent with zero

similar observations from Hall-A and COMPASS

Selected **higher-twist** 1-hadron SIDIS results

Subleading twist

$$F_{UT}^{\sin\phi_S} = \frac{2M}{Q}\,\mathcal{C}\Big\{\Big(x f_T D_1 - \frac{M_h}{M}\,h_1\frac{\tilde{H}}{z}\Big) - \frac{\boldsymbol{k}_T\cdot\boldsymbol{p}_T}{2MM_h}\Big[\Big(x h_T H_1^\perp + \frac{M_h}{M}\,g_{1T}\frac{\tilde{G}^\perp}{z}\Big) - \Big(x h_T^\perp H_1^\perp - \frac{M_h}{M}\,f_{1T}^\perp\frac{\tilde{D}^\perp}{z}\Big)\Big]\Big\}$$

$$\frac{d\sigma^h}{dx\,dy\,d\phi_S\,dz\,d\phi\,d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2}{xyQ^2}\frac{y^2}{2(1-\epsilon)}\left(1+\frac{\gamma^2}{2x}\right)$$

$$\Big\{ \Big[F_{\mathrm{UU,T}} + \epsilon F_{\mathrm{UU,L}} + \sqrt{2\epsilon(1+\epsilon)}\cos(\phi)F_{\mathrm{UU}}^{\cos(\phi)} + \epsilon\cos(2\phi)F_{\mathrm{UU}}^{\cos(2\phi)}\Big]$$

$$+ \lambda_l\Big[\sqrt{2\epsilon(1-\epsilon)}\sin(\phi)F_{\mathrm{LU}}^{\sin(\phi)}\Big]$$

$$+ S_L\Big[\sqrt{2\epsilon(1+\epsilon)}\sin(\phi)F_{\mathrm{UL}}^{\sin(\phi)} + \epsilon\sin(2\phi)F_{\mathrm{UL}}^{\sin(2\phi)}\Big]$$

$$+ S_L\,\lambda_l\Big[\sqrt{1-\epsilon^2}F_{\mathrm{LL}} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi)F_{\mathrm{LL}}^{\cos(\phi)}\Big]$$

$$+ S_T\Big[\sin(\phi-\phi_S)\Big(F_{\mathrm{UT,T}}^{\sin(\phi-\phi_S)} + \epsilon F_{\mathrm{UT,L}}^{\sin(\phi-\phi_S)}\Big) + \epsilon\sin(\phi+\phi_S)F_{\mathrm{UT}}^{\sin(\phi+\phi_S)} + \epsilon\sin(3\phi-\phi_S)F_{\mathrm{UT}}^{\sin(3\phi-\phi_S)} + \sqrt{2\epsilon(1+\epsilon)}\sin(\phi_S)F_{\mathrm{UT}}^{\sin(\phi_S)} + \sqrt{2\epsilon(1+\epsilon)}\sin(2\phi-\phi_S)F_{\mathrm{UT}}^{\sin(2\phi-\phi_S)}\Big]$$

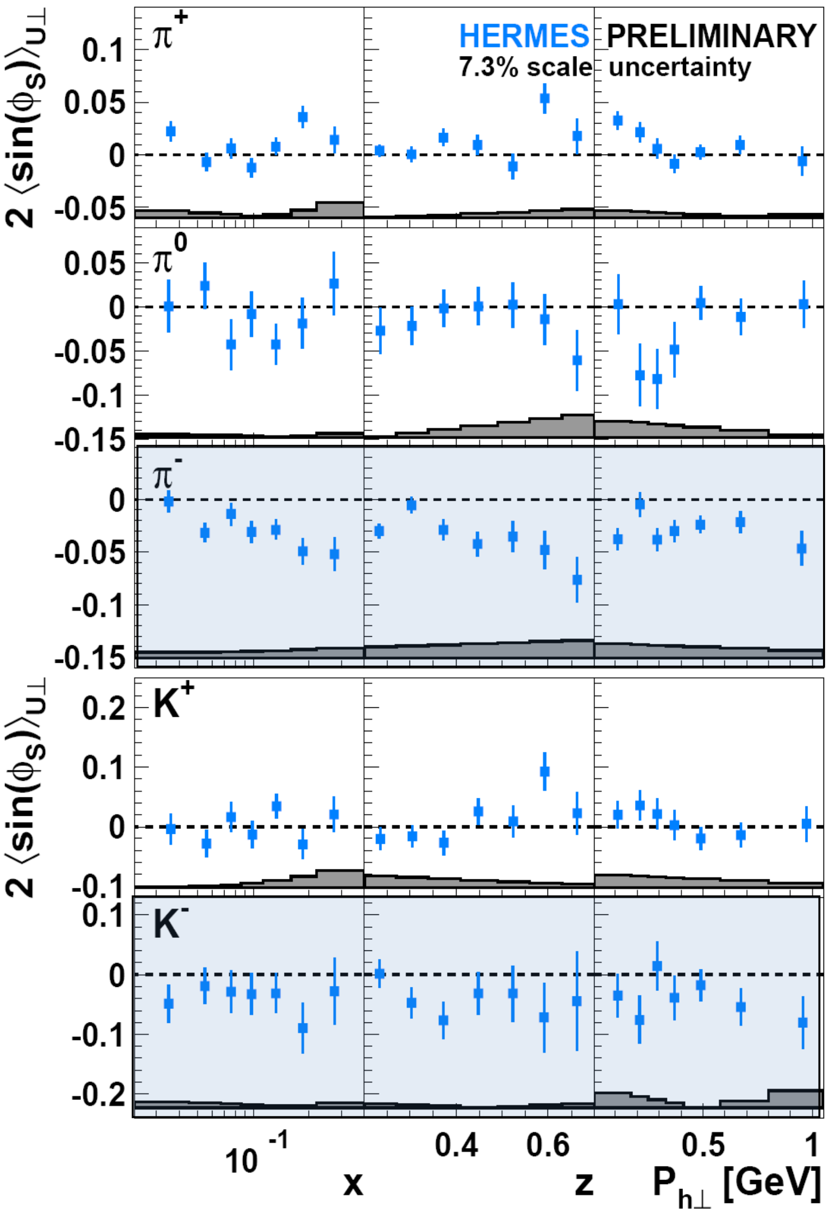
$$+ S_T\,\lambda_l\Big[\sqrt{1-\epsilon^2}\cos(\phi-\phi_S)F_{\mathrm{LT}}^{\cos(\phi-\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi_S)F_{\mathrm{LT}}^{\cos(\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(2\phi-\phi_S)F_{\mathrm{LT}}^{\cos(2\phi-\phi_S)}\Big]\Big\}$$

Sensitive to worm-gear $g_{1T}^\perp$, sivers, transversity + higher-twist DF and FF

Distribution Functions

nucleon \ quark	U	L	T
U	f_1		$h_1^\perp$
L		g_1	$h_{1L}^\perp$
T	$f_{1T}^\perp$	$g_{1T}^\perp$	h_1, $h_{1T}^\perp$

Subleading-twist $\sin(\phi_S)$ Fourier component

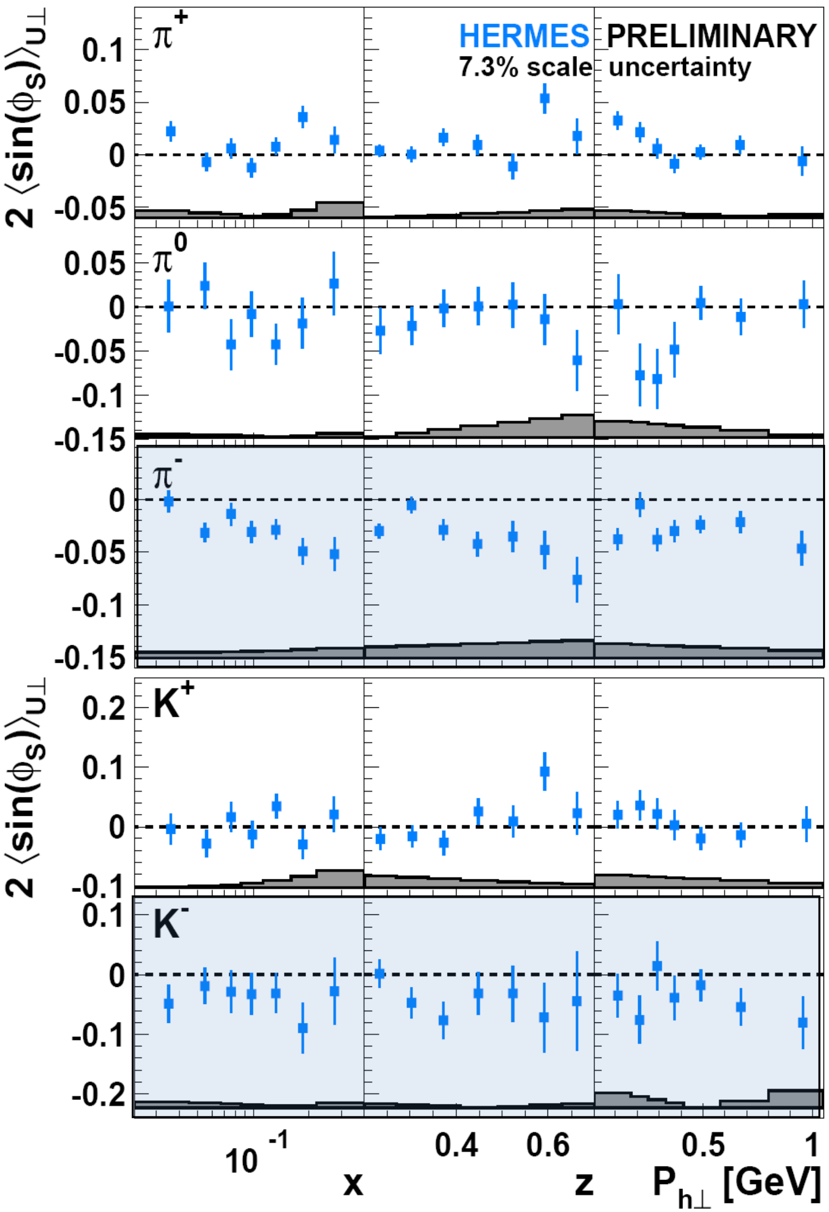


- sensitive to **worm-gear** $g_{1T}^{\perp}$, **Sivers function**, **Transversity**, etc
- **significant non-zero signal for π^- and K^- !**

☛ Large and negative

☛ negative

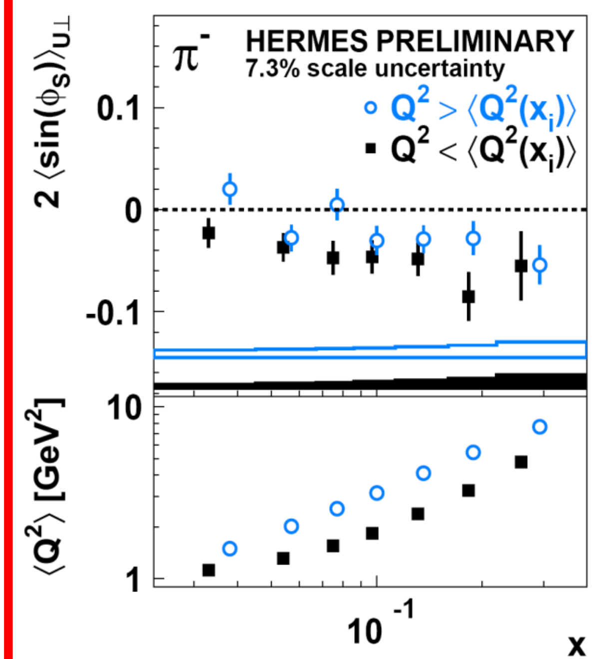
Subleading-twist $\sin(\phi_S)$ Fourier component



- sensitive to **worm-gear** $g_{1T}^{\perp}$, **Sivers function**, **Transversity**, etc
- **significant non-zero signal for π^- and K- !**

Large and negative

negative



- low-Q^2 amplitude larger
- hint of Q^2 dependence for π^-

$F_{LU}{}^{\sin\phi}$

$$F_{LU}^{\sin\phi_h} = \frac{2M}{Q}\,\mathcal{C}\left[-\frac{\hat{\boldsymbol{h}}\cdot\boldsymbol{k}_T}{M_h}\left(xe\,H_1^{\perp} + \frac{M_h}{M}\,f_1\frac{\tilde{G}^{\perp}}{z}\right) + \frac{\hat{\boldsymbol{h}}\cdot\boldsymbol{p}_T}{M}\left(xg^{\perp}D_1 + \frac{M_h}{M}\,h_1^{\perp}\frac{\tilde{E}}{z}\right)\right]$$

$$\frac{d\sigma^h}{dx\,dy\,d\phi_S\,dz\,d\phi\,d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2}{xyQ^2}\frac{y^2}{2(1-\epsilon)}\left(1+\frac{\gamma^2}{2x}\right)$$

$$\Big\{ \Big[F_{\rm UU,T} + \epsilon F_{\rm UU,L} + \sqrt{2\epsilon(1+\epsilon)}\cos(\phi)F_{\rm UU}^{\cos(\phi)} + \epsilon\cos(2\phi)F_{\rm UU}^{\cos(2\phi)}\Big]$$

$$+ \lambda_l\left[\sqrt{2\epsilon(1-\epsilon)}\sin(\phi)F_{\rm LU}^{\sin(\phi)}\right]$$

$$+ S_L\left[\sqrt{2\epsilon(1+\epsilon)}\sin(\phi)F_{\rm UL}^{\sin(\phi)} + \epsilon\sin(2\phi)F_{\rm UL}^{\sin(2\phi)}\right]$$

$$+ S_L\,\lambda_l\left[\sqrt{1-\epsilon^2}F_{\rm LL} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi)F_{\rm LL}^{\cos(\phi)}\right]$$

$$+ S_T\Big[\sin(\phi-\phi_S)\left(F_{\rm UT,T}^{\sin(\phi-\phi_S)} + \epsilon F_{\rm UT,L}^{\sin(\phi-\phi_S)}\right) + \epsilon\sin(\phi+\phi_S)F_{\rm UT}^{\sin(\phi+\phi_S)} + \epsilon\sin(3\phi-\phi_S)F_{\rm UT}^{\sin(3\phi-\phi_S)} + \sqrt{2\epsilon(1+\epsilon)}\sin(\phi_S)F_{\rm UT}^{\sin(\phi_S)} + \sqrt{2\epsilon(1+\epsilon)}\sin(2\phi-\phi_S)F_{\rm UT}^{\sin(2\phi-\phi_S)}\Big]$$

$$+ S_T\,\lambda_l\Big[\sqrt{1-\epsilon^2}\cos(\phi-\phi_S)F_{\rm LT}^{\cos(\phi-\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi_S)F_{\rm LT}^{\cos(\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(2\phi-\phi_S)F_{\rm LT}^{\cos(2\phi-\phi_S)}\Big]\Big\}$$

Sensitive to f_1, Boer-Mulders + higher-twist DF and FF

Distribution Functions

nucleon \ quark	U	L	T
U	f_1		$h_1^{\perp}$
L		g_1	$h_{1L}^{\perp}$
T	$f_{1T}^{\perp}$	$g_{1T}^{\perp}$	h_1, $h_{1T}^{\perp}$

$F_{LU}^{\sin\phi}$

$$F_{LU}^{\sin\phi_h} = \frac{2M}{Q}\,\mathcal{C}\Big[-\frac{\hat{h}\cdot k_T}{M_h}\Big(xe\,H_1^\perp + \frac{M_h}{M}\,f_1\frac{\tilde{G}^\perp}{z}\Big) + \frac{\hat{h}\cdot p_T}{M}\Big(xg^\perp D_1 + \frac{M_h}{M}\,h_1^\perp\frac{\tilde{E}}{z}\Big)\Big]$$

$$\frac{d\sigma^h}{dx\,dy\,d\phi_S\,dz\,d\phi\,d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2}{xyQ^2}\frac{y^2}{2(1-\epsilon)}\left(1+\frac{\gamma^2}{2x}\right)$$

$$\Big\{ \Big[F_{\rm UU,T} + \epsilon F_{\rm UU,L} + \sqrt{2\epsilon(1+\epsilon)}\cos(\phi)F_{\rm UU}^{\cos(\phi)} + \epsilon\cos(2\phi)F_{\rm UU}^{\cos(2\phi)}\Big]$$

$$+ \lambda_l\Big[\sqrt{2\epsilon(1-\epsilon)}\sin(\phi)F_{\rm LU}^{\sin(\phi)}\Big]$$

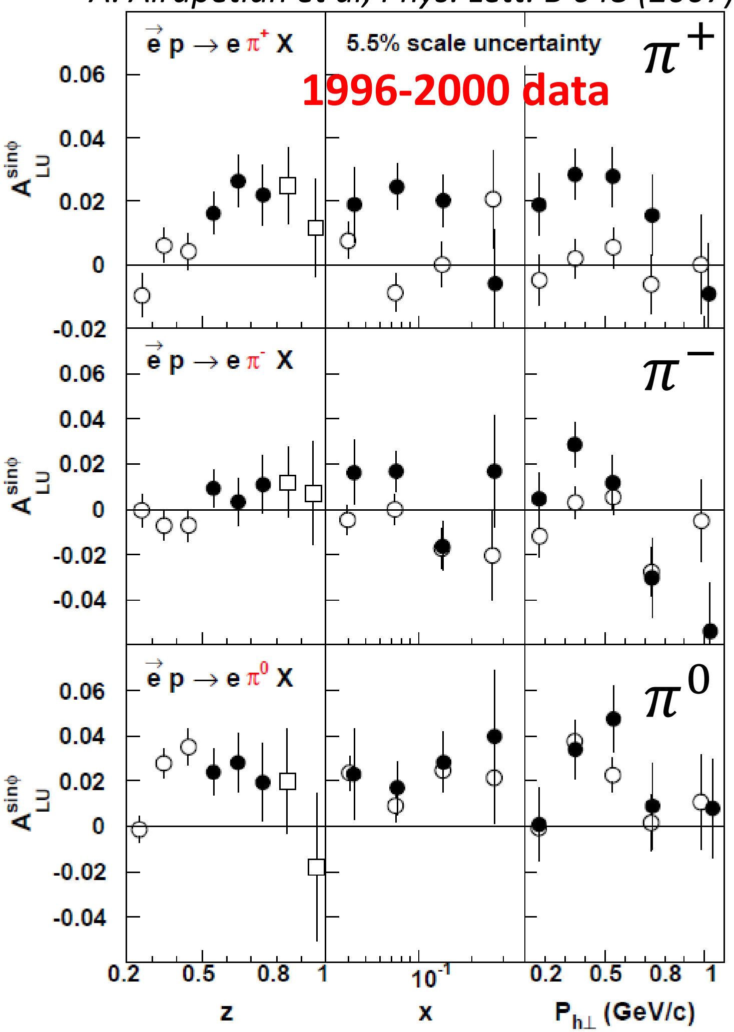
$$+ S_L\Big[\sqrt{2\epsilon(1+\epsilon)}\sin(\phi)F_{\rm UL}^{\sin(\phi)} + \epsilon\sin(2\phi)F_{\rm UL}^{\sin(2\phi)}\Big]$$

$$+ S_L\,\lambda_l\Big[\sqrt{1-\epsilon^2}F_{\rm LL} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi)F_{\rm LL}^{\cos(\phi)}\Big]$$

$$+ S_T\Big[\sin(\phi-\phi_S)\Big(F_{\rm UT,T}^{\sin(\phi-\phi_S)} + \epsilon F_{\rm UT,L}^{\sin(\phi-\phi_S)}\Big) + \epsilon\sin(\phi+\phi_S)F_{\rm UT}^{\sin(\phi+\phi_S)} + \epsilon\sin(3\phi-\phi_S)F_{\rm UT}^{\sin(3\phi-\phi_S)} + \sqrt{2\epsilon(1+\epsilon)}\sin(\phi_S)F_{\rm UT}^{\sin(\phi_S)} + \sqrt{2\epsilon(1+\epsilon)}\sin(2\phi-\phi_S)F_{\rm UT}^{\sin(2\phi-\phi_S)}\Big]$$

$$+ S_T\,\lambda_l\Big[\sqrt{1-\epsilon^2}\cos(\phi-\phi_S)F_{\rm LT}^{\cos(\phi-\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi_S)F_{\rm LT}^{\cos(\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(2\phi-\phi_S)F_{\rm LT}^{\cos(2\phi-\phi_S)}\Big]\Big\}$$

A. Airapetian et al, Phys. Lett. B 648 (2007)

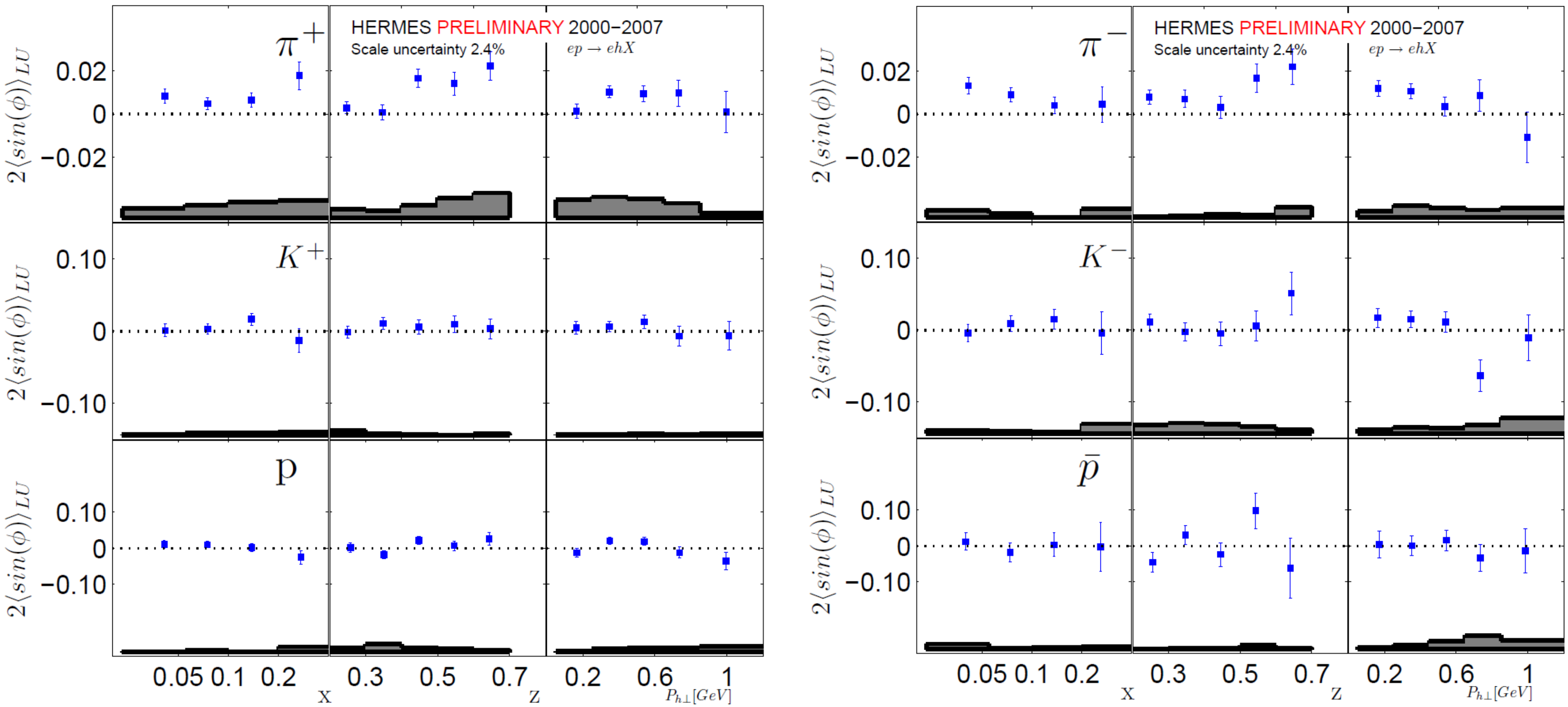


open circles 0.2<z<0.5
full circles 0.5<z<0.8
open squares: 0.8<z<1.0

$F_{LU}^{\sin\phi}$

$$F_{LU}^{\sin\phi_h} = \frac{2M}{Q}\,\mathcal{C}\left[-\frac{\hat{\boldsymbol{h}}\cdot\boldsymbol{k}_T}{M_h}\left(xe\,H_1^\perp + \frac{M_h}{M}\,f_1\frac{\tilde{G}^\perp}{z}\right) + \frac{\hat{\boldsymbol{h}}\cdot\boldsymbol{p}_T}{M}\left(xg^\perp D_1 + \frac{M_h}{M}\,h_1^\perp\frac{\tilde{E}}{z}\right)\right]$$

H target, 2000-2007 data 0.2<z<0.7

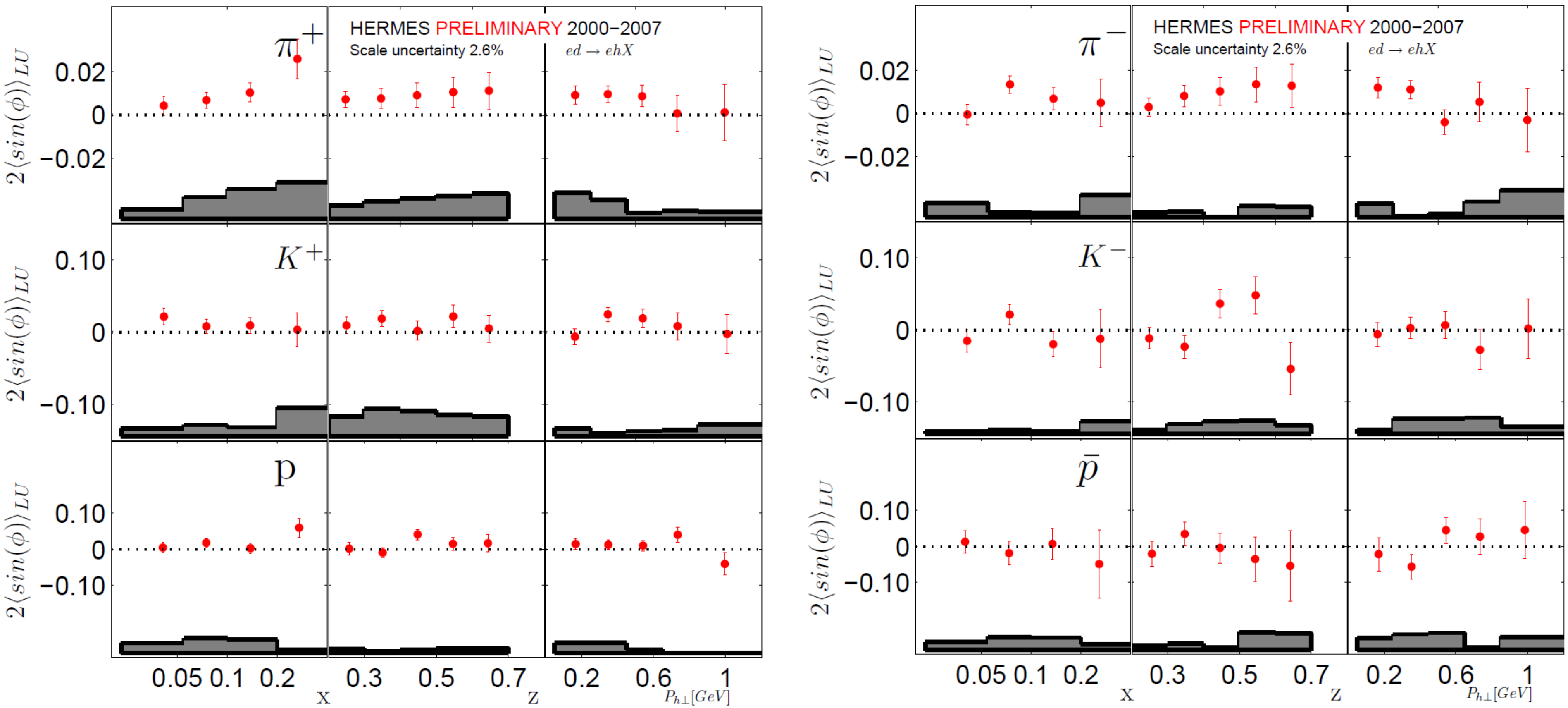


Released yesterday!!

$F_{LU}^{\sin\phi}$

$$F_{LU}^{\sin\phi_h} = \frac{2M}{Q}\,\mathcal{C}\Big[-\frac{\hat{\boldsymbol{h}}\cdot\boldsymbol{k}_T}{M_h}\Big(xe\,H_1^{\perp} + \frac{M_h}{M}\,f_1\frac{\tilde{G}^{\perp}}{z}\Big) + \frac{\hat{\boldsymbol{h}}\cdot\boldsymbol{p}_T}{M}\Big(xg^{\perp}D_1 + \frac{M_h}{M}\,h_1^{\perp}\frac{\tilde{E}}{z}\Big)\Big]$$

D target, 2000-2007 data 0.2<z<0.7



Released yesterday!!

2-hadron SIDIS results

Following formalism developed by **Steve Gliske**

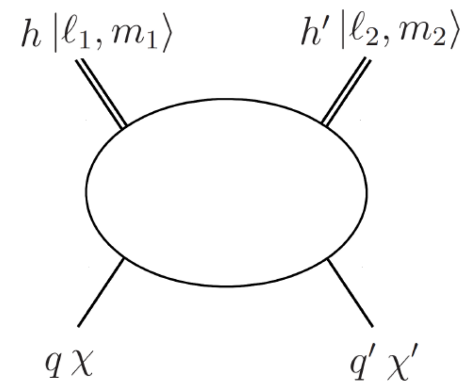
Find details in

Transverse Target Moments of Dihadron Production in Semi-inclusive Deep Inelastic Scattering at HERMES
S. Gliske, PhD thesis, University of Michigan, 2011

A short digression on di-hadron fragmentation functions

Standard definition of di-hadron FF assume no polarization of final state hadrons (pseudo-scalar mesons) or define mixtures of certain partial waves as new FFs

In the **new formalism there are only two di-hadron FFs**. Names and symbols are entirely associated with the quark spin states (D_1 for $\chi = \chi'$ and $H_1{}^{\perp}$ (*generalized Collins*) for $\chi \neq \chi'$), whereas the partial waves of the produced hadrons ($|l_1 m_1\rangle, |l_2 m_2\rangle$) are associated with partial waves of FFs.



$$D_1 = \sum_{\ell=1}^{\infty} \sum_{m=-\ell}^{\ell} P_{\ell,m}(\cos\vartheta) e^{im(\phi_R - \phi_k)} D_1^{|\ell,m\rangle}(z, M_h, |\boldsymbol{k}_T|)$$

$$H_1^{\perp} = \sum_{\ell=1}^{\infty} \sum_{m=-\ell}^{\ell} P_{\ell,m}(\cos\vartheta) e^{im(\phi_R - \phi_k)} H_1^{\perp|\ell,m\rangle}(z, M_h, |\boldsymbol{k}_T|)$$

The cross-section is identical to the ones in literature, the only difference is the interpretation of the FFs:

$$D_1^{|0,0\rangle} = D_{1,OO} = \left(\frac{1}{4} D_{1,OO}^{s} + \frac{3}{4} D_{1,OO}^{p}\right)$$

$$D_1^{|1,0\rangle} = D_{1,OL},$$

$$D_1^{|1,\pm1\rangle} = D_{1,OT} \mp \frac{|\boldsymbol{k}_T|\,|\boldsymbol{R}|}{M_h^2} G_{1,OT}^{\perp},$$

$$D_1^{|2,0\rangle} = \frac{1}{2} D_{1,LL},$$

$$D_1^{|2,\pm1\rangle} = \frac{1}{2}\left(D_{1,LT} \mp \frac{|\boldsymbol{k}_T|\,|\boldsymbol{R}|}{M_h^2} G_{1,LT}^{\perp}\right),$$

$$D_1^{|2,\pm2\rangle} = D_{1,TT} \mp \frac{1}{2}\frac{|\boldsymbol{k}_T|\,|\boldsymbol{R}|}{M_h^2} G_{1,TT}^{\perp},$$

$$H_1^{\perp|0,0\rangle} = H_{1,OO}^{\perp} = \frac{1}{4} H_{1,OO}^{\perp s} + \frac{3}{4} H_{1,OO}^{\perp p},$$

$$H_1^{\perp|1,1\rangle} = H_{1,OT}^{\perp} + \frac{|\boldsymbol{R}|}{|\boldsymbol{k}_T|} \bar{H}_{1,OT}^{\sphericalangle} = \frac{|\boldsymbol{R}|}{|\boldsymbol{k}_T|} H_{1,OT}^{\sphericalangle}$$

$$H_1^{\perp|1,0\rangle} = H_{1,OL}^{\perp}$$

$$H_1^{\perp|1,-1\rangle} = H_{1,OT}^{\perp}$$

$$H_1^{\perp|2,2\rangle} = H_{1,TT}^{\perp} + \frac{|\boldsymbol{R}|}{|\boldsymbol{k}_T|} \bar{H}_{1,TT}^{\sphericalangle} = \frac{|\boldsymbol{R}|}{|\boldsymbol{k}_T|} H_{1,TT}^{\sphericalangle},$$

$$H_1^{\perp|2,1\rangle} = \frac{1}{2} H_{1,LT}^{\perp} + \frac{1}{2}\frac{|\boldsymbol{R}|}{|\boldsymbol{k}_T|} \bar{H}_{1,LT}^{\sphericalangle} = \frac{1}{2}\frac{|\boldsymbol{R}|}{|\boldsymbol{k}_T|} H_{1,LT}^{\sphericalangle},$$

$$H_1^{\perp|2,0\rangle} = \frac{1}{2} H_{1,LL}^{\perp},$$

$$H_1^{\perp|2,-1\rangle} = \frac{1}{2} H_{1,LT}^{\perp},$$

$$H_1^{\perp|2,-2\rangle} = H_{1,TT}^{\perp}.$$

The di-hadron SIDIS cross-section

$$
\begin{aligned}
d\sigma_{UT} \;=\; & \frac{\alpha^2 M_h P_{h\perp}}{2\pi x y Q^2}\left(1+\frac{\gamma^2}{2x}\right)|\boldsymbol{S}_\perp| \\
& \times \sum_{\ell=0}^{2}\sum_{m=-\ell}^{\ell}\Bigg\{ A(x,y)\Big[P_{\ell,m}\sin((m+1)\phi_h - m\phi_R - \phi_S)) \\
& \qquad \times \left(F_{UT,T}^{P_{\ell,m}\sin((m+1)\phi_h - m\phi_R-\phi_S)} + \epsilon F_{UT,L}^{P_{\ell,m}\sin((m+1)\phi_h - m\phi_R-\phi_S)}\right)\Big] \\
& \quad + B(x,y)\Big[P_{\ell,m}\sin((1-m)\phi_h + m\phi_R + \phi_S) \boxed{F_{UT}^{P_{\ell,m}\sin((1-m)\phi_h + m\phi_R+\phi_S)}} \\
& \qquad + P_{\ell,m}\sin((3-m)\phi_h + m\phi_R - \phi_S) F_{UT}^{P_{\ell,m}\sin((3-m)\phi_h + m\phi_R-\phi_S)}\Big] \\
& \quad + V(x,y)\Big[P_{\ell,m}\sin(-m\phi_h + m\phi_R + \phi_S) F_{UT}^{P_{\ell,m}\sin(-m\phi_h + m\phi_R+\phi_S)} \\
& \qquad + P_{\ell,m}\sin((2-m)\phi_h + m\phi_R - \phi_S) F_{UT}^{P_{\ell,m}\sin((2-m)\phi_h + m\phi_R-\phi_S)}\Big]\Bigg\}.
\end{aligned}
$$

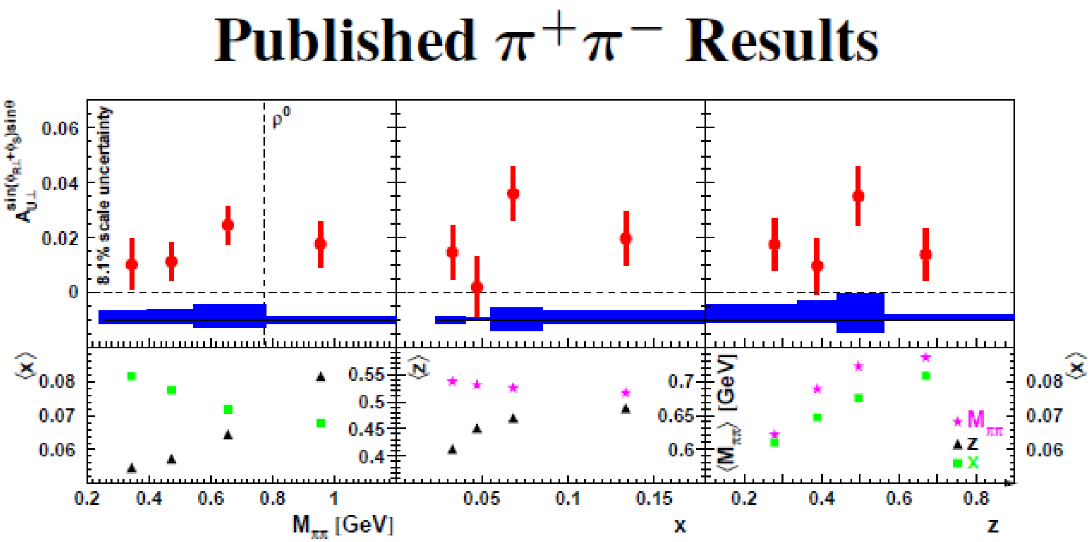
l and m correspond to the l and m in $|lm\rangle$ angular momentum state of the hadron

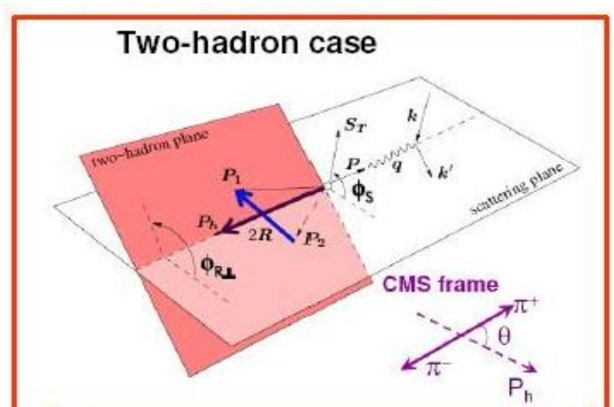
Considering all terms ($d\sigma_{UU}, d\sigma_{LU}, d\sigma_{UL}, d\sigma_{LL}, d\sigma_{UT}, d\sigma_{LT}$) there are **144 non-zero structure functions** at twist-3 level. The most known is

$$
F_{UT}^{P_{\ell,m}\sin((1-m)\phi_h+m\phi_R+\phi_S)} = -\mathcal{I}\left[\frac{|\boldsymbol{k}_T|}{M_h}\cos\left((m-1)\phi_h - \phi_p - m\phi_k\right)\; h_1 H_1^{\perp|\ell,m\rangle}\right]
$$

which for $l = 1$ and $m = 1$ reduces to the well known colliner $F_{UT}{}^{\sin\vartheta\,\sin(\phi_R+\phi_S)}$ related to transversity

The di-hadron SIDIS cross-section



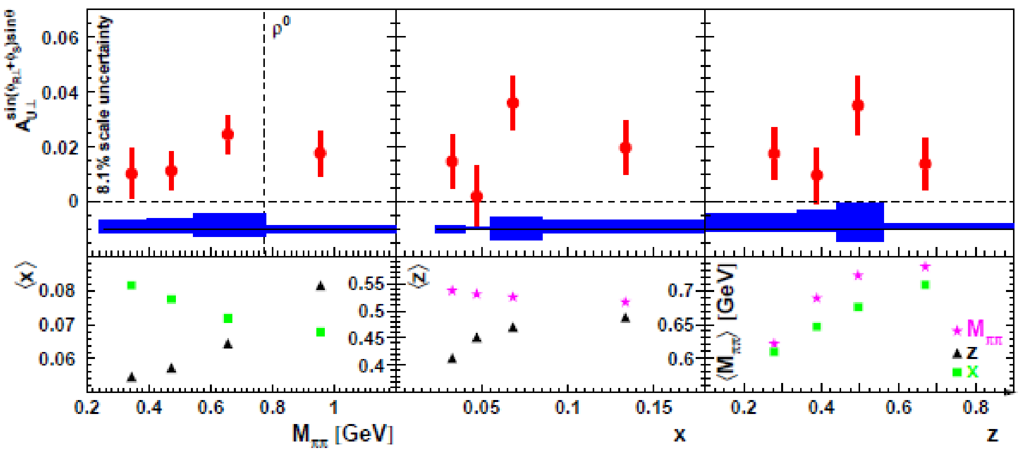


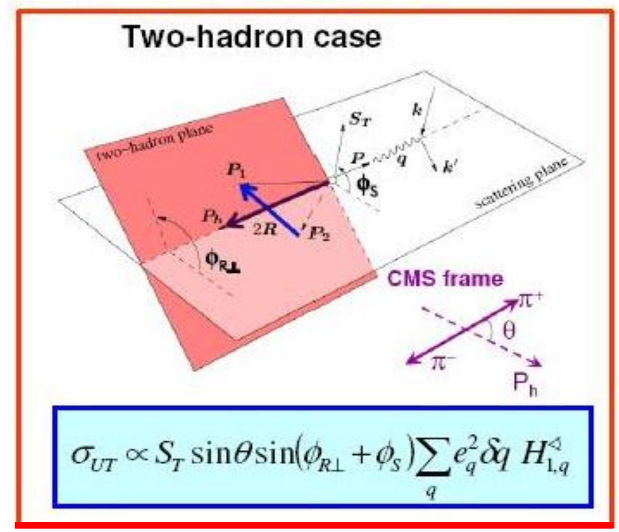
$$\sigma_{UT} \propto S_T \sin\theta \sin(\phi_{R\perp} + \phi_S) \sum_q e_q^2 \delta q \; H_{1,q}^{\sphericalangle}$$

- independent way to access transversity
- significantly positive amplitudes
- 1^{st} evidence of non zero dihadron FF
- no convolution integral involved
- limited statistical power (v.r.t. 1 hadron)

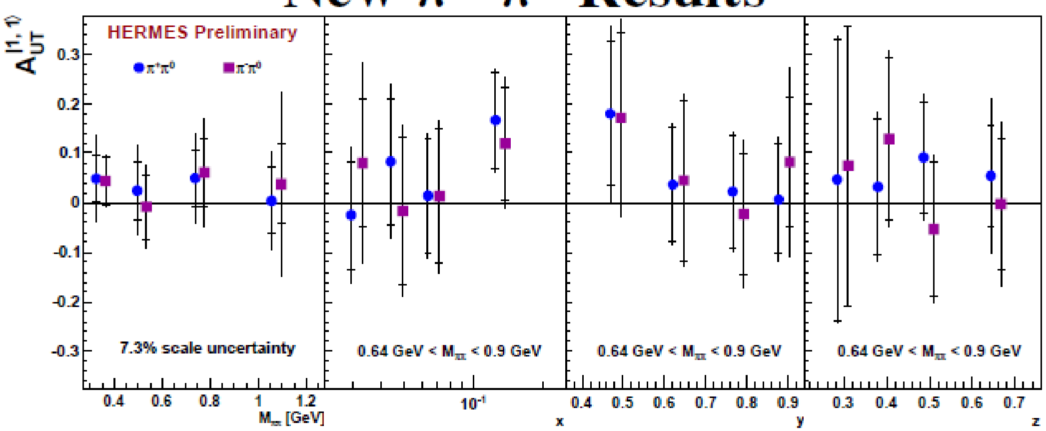
The di-hadron SIDIS cross-section

Published $\pi^+\pi^-$ Results





New $\pi^\pm\pi^0$ Results



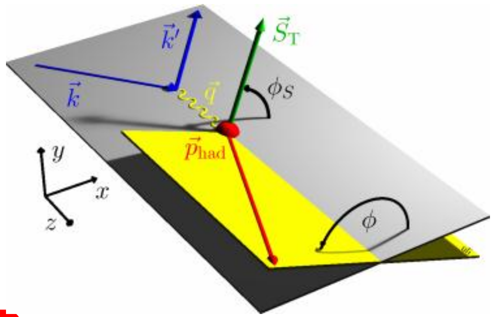
- New tracking, new PID, use of ϕ_R rather than $\phi_{R\perp}$
- Different fitting procedure and function
- Acceptance correction

- independent way to access transversity
- significantly positive amplitudes
- 1^{st} evidence of non zero dihadron FF
- no convolution integral involved
- limited statistical power (v.r.t. 1 hadron)
- signs are consistent for all $\pi\pi$ species
- statistics much more limited for $\pi^\pm\pi^0$
- despite uncertainties may still help to constrain global fits and may assist in $u-d$ flavor separation

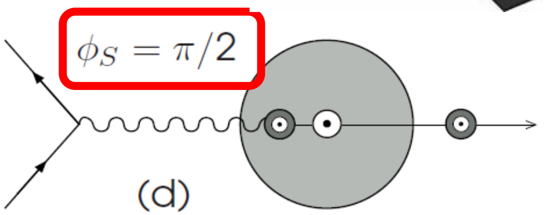
A short digression on the Lund/Artru string fragmentation model

(a phenomenological explanation of the Collins effect)

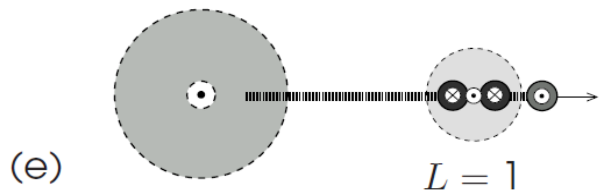
In the cross-section the Collins FF is always paired withy a distrib. function involving a transv. pol. quark.



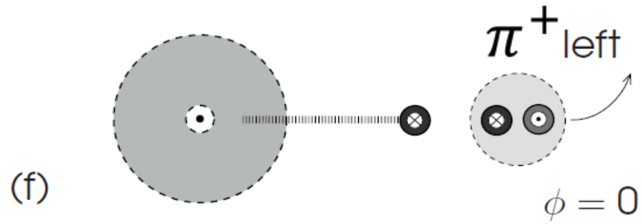
1. Assume u quark and proton have (transverse) spin alligned in the direction $\phi_S = \pi/2$. The model assumes that the struck quark is initially connected with the remnant via a gluon-flux tube (string)



2. When the string breaks, a $q\bar{q}$ pair is created with vacuum quantum numbers $J^P = 0^+$. The positive parity requires that the spins of q and $\bar{q}$ are aligned, thus an OAM $\boldsymbol{L} = \mathbf{1}$ has to compensate the spins



3. This OAM generates a transverse momentum of the produced pseudo-scalar meson (e.g. $\boldsymbol{\pi}^+$) and deflects the meson to the **left side** w.r.t. the struck quark direction, generating left-righ azimuthal asymmetries



A short digression on the Lund/Artru string fragmentation model

Relative to the proton transv. spin, the fragmenting quark can have spin parallel or antiparallel to $\left|\frac{1}{2}, \pm\frac{1}{2}\right\rangle$

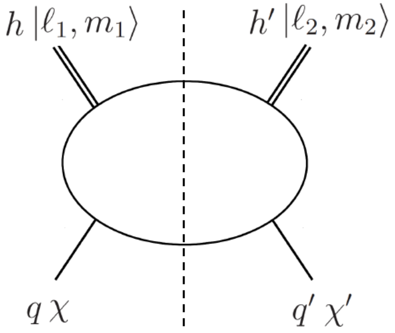
Then combining the spins of the formed di-quark systems one can get:

$$\frac{1}{2} \otimes \frac{1}{2} = 1 \oplus 0 \;\Rightarrow \begin{cases} 1\ spin\ 0\ state\ |0,0\rangle & \text{1 pseudo-scalar meson (PSM)} \\ 3\ spin\ 1\ states \begin{cases} |1,0\rangle \\ |1,\pm 1\rangle \end{cases} & \begin{matrix} \text{1 Longitudinal VM} \\ \text{2 transvrse VM} \end{matrix} \end{cases}$$

Lund/Artru prediction at the amplitude level: the asymmetry for PSM has opposite sign to that for transversely polarized VM (left vs. right side), and the amplitude for $|1,0\rangle$ is 0

Lund/Artru model makes predictions for the individual di-hadrons, but the Collins function includes pairs of di-hadrons

$\rightarrow$ to make predictions for the Collins function one needs to consider the cross-section level, i.e. the sum of contributing amplitudes times their complex conjugate



Using the Clebsch-Gordan algebra one obtains: $|1,\pm 1\rangle\,|1,\pm 1\rangle \equiv |2,\pm 2\rangle$

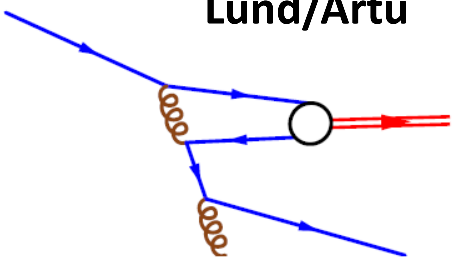
Lund/Artru prediction at the cross-section level: the $|2,\pm 2\rangle$ partial waves of the Collins func. for SIDIS VM production have the opposite sign as the respective PS Collins func.

"gluon radiaton model" vs. Lund/Artru model

The Lund/Artru model only accounts for favored Collins fragmentation. An extension of the model (the **gluon radiation model**), elaborated by **S. Gliske** accounts for the disfavored case

1. Struck quark emits a gluon in such a way that most of its momentum is transferred to the gluon
2. The struck quark then becomes part of the remnant
3. The radiated gluon produces a $q\bar{q}$ pair that eventually converts into a meson
4. For PSM the di-quark must interact further with the remnant to get the PSM quantum numbers. In case of VM the di-quark directly forms the meson

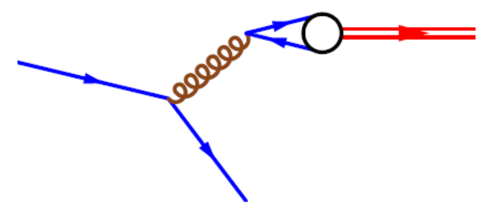
Lund/Artu



- Di-quark has q.n. of vacuum
- **Struck quark** joins the anti-quark in the final state $\rightarrow$ **favored fragment.**

Prediction: the $|2,\pm2\rangle$ partial wave of the Collins funct. for SIDIS VM production have the opposite sign as the respective PS Collins function

Gluon radiation



- Di-quark has q.n. of observed final state
- **Produced quark** joins the anti-quark in the final state $\rightarrow$ **disfavored fragment.**

Prediction: the disfavored $|2,\pm2\rangle$ Collins frag. also is expected to have opposite sign as the respective PS Collins function.

$\Rightarrow$ Models predict: fav = disfav for VM

Data say: fav $\cong -$ disfav for PSM (Collins π^+ $vs.$ π^-)

…and now let's look at the results

u dominance

Fragment. process	Fav/disfav	Deflection	Sign of amplitude
$\boldsymbol{u} \rightarrow \boldsymbol{\pi}^{+}$	fav PSM	left $(\phi_h \rightarrow 0)$	> 0 (Collins π^{+})
$\boldsymbol{u} \rightarrow \boldsymbol{\pi}^{-}$	disfav PSM	ight $(\phi_h \rightarrow \pi)$	< 0 (Collins π^{-})
$\boldsymbol{u} \rightarrow \boldsymbol{\rho}^{+} \rightarrow \pi^{+}\pi^{0}$	fav VM	right $(\phi_h \rightarrow \pi)$	< 0
$\boldsymbol{u} \rightarrow \boldsymbol{\rho}^{-} \rightarrow \pi^{-}\pi^{0}$	disfav VM	right $(\phi_h \rightarrow \pi)$	< 0
$\boldsymbol{u} \rightarrow \boldsymbol{\rho}^{\mathbf{0}} \rightarrow \pi^{+}\pi^{-}$	mixed VM	right $(\phi_h \rightarrow \pi)$	0 or < 0

from data (first two rows)

from models (last three rows)

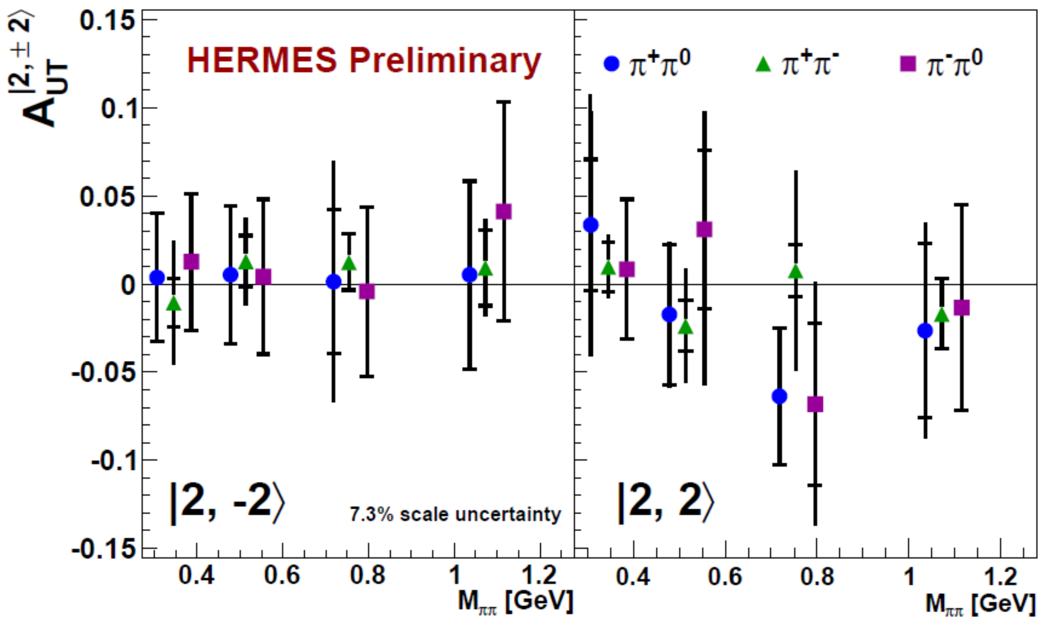
…and now let's look at the results

u dominance

Fragment. process	Fav/disfav	Deflection	Sign of amplitude
$\boldsymbol{u} \rightarrow \boldsymbol{\pi}^{+}$	fav PSM	left ($\phi_h \rightarrow 0$)	> 0 (Collins π^{+})
$\boldsymbol{u} \rightarrow \boldsymbol{\pi}^{-}$	disfav PSM	ight ($\phi_h \rightarrow \pi$)	< 0 (Collins π^{-})
$\boldsymbol{u} \rightarrow \boldsymbol{\rho}^{+} \rightarrow \pi^{+}\pi^{0}$	fav VM	right ($\phi_h \rightarrow \pi$)	< 0
$\boldsymbol{u} \rightarrow \boldsymbol{\rho}^{-} \rightarrow \pi^{-}\pi^{0}$	disfav VM	right ($\phi_h \rightarrow \pi$)	< 0
$\boldsymbol{u} \rightarrow \boldsymbol{\rho}^{0} \rightarrow \pi^{+}\pi^{-}$	mixed VM	right ($\phi_h \rightarrow \pi$)	0 or < 0

from data (rows 1–2)

from models (rows 3–5)



$|2,-2\rangle$ consistent with zero for all flavors

Not in contraddiction with models: if the transversity function causes the fragmenting quark to have positive polarization than Collins $|2,-2\rangle$ must be zero as this partial wave requires fragmenting quark with negative polarization

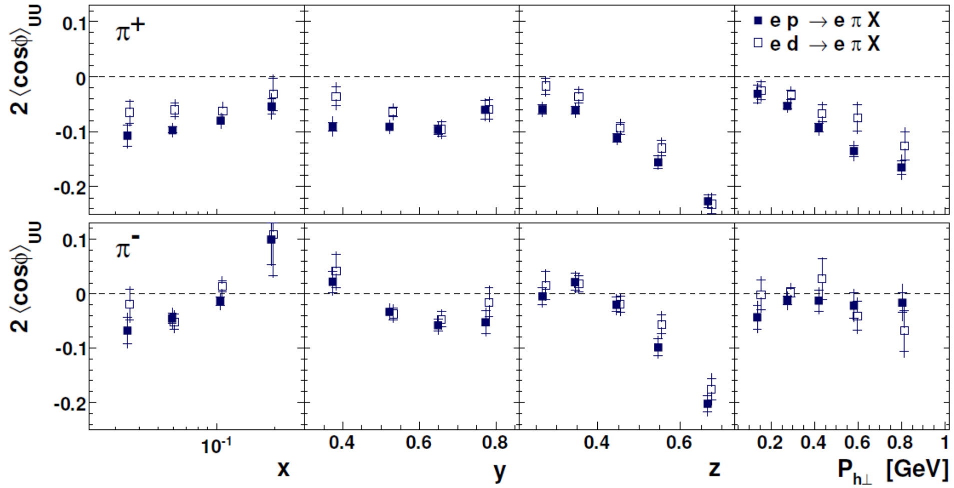
$|2,+2\rangle$ consistent with model expect:

- No signal outside ρ-mass bin
  $\rightarrow$ no non-resonant pion-pairs in p-wave
- Negative for $\rho^{\pm}$ (model predictions)
- very small for ρ^{0} (consistent with small Collins π^{0})

Back-up

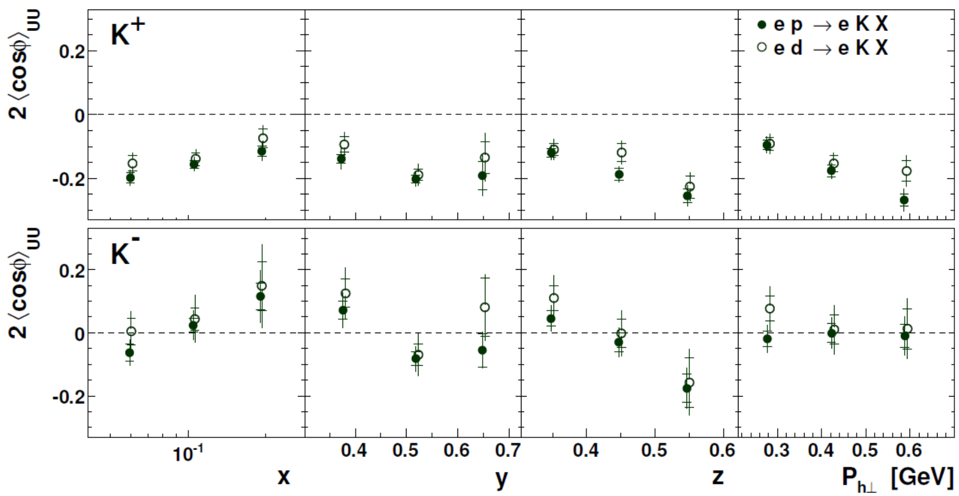
The cosϕ amplitudes

A. Airapetian et al, Phys. Rev. D 87 (2013) 012010



negative

negative



Large and negative

Consist. with 0

Analysis multi-dimensional in x, y, z,and Pt

Create your own projections of results through: http://www-hermes.desy.de/cosnphi/

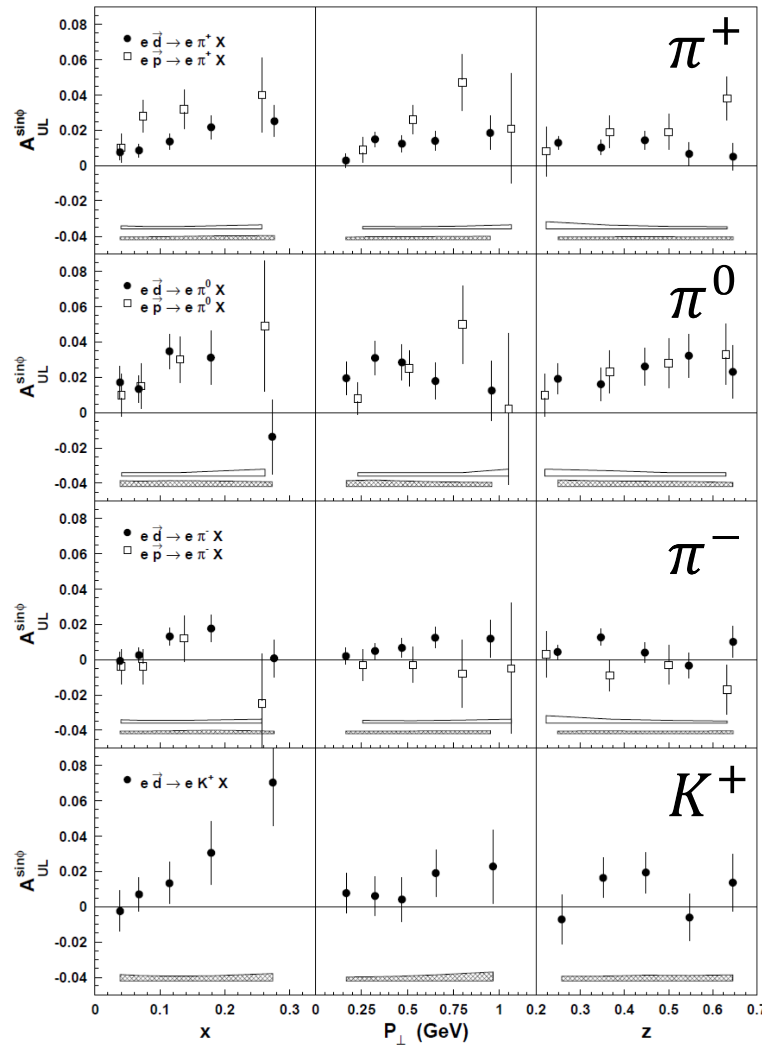
$F_{UL}{}^{\sin\phi}$

$$F_{UL}^{\sin\phi_h} = \frac{2M}{Q}\,\mathcal{C}\left[-\frac{\hat{\boldsymbol{h}}\cdot\boldsymbol{k}_T}{M_h}\left(x h_L H_1^{\perp} + \frac{M_h}{M}\, g_{1L}\frac{\tilde{G}^{\perp}}{z}\right) + \frac{\hat{\boldsymbol{h}}\cdot\boldsymbol{p}_T}{M}\left(x f_L^{\perp} D_1 - \frac{M_h}{M}\, h_{1L}^{\perp}\frac{\tilde{H}}{z}\right)\right]$$

$$\frac{d\sigma^h}{dx\,dy\,d\phi_S\,dz\,d\phi\,d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2}{xyQ^2}\frac{y^2}{2(1-\epsilon)}\left(1+\frac{\gamma^2}{2x}\right)$$

$$\begin{aligned}
\Big\{ \quad & \Big[F_{\mathrm{UU,T}} + \epsilon F_{\mathrm{UU,L}} \\
& + \sqrt{2\epsilon(1+\epsilon)}\cos(\phi) F_{\mathrm{UU}}^{\cos(\phi)} + \epsilon\cos(2\phi) F_{\mathrm{UU}}^{\cos(2\phi)}\Big] \\
+ \quad \lambda_l \quad & \Big[\sqrt{2\epsilon(1-\epsilon)}\sin(\phi) F_{\mathrm{LU}}^{\sin(\phi)}\Big] \\
+ S_L \quad & \Big[\sqrt{2\epsilon(1+\epsilon)}\sin(\phi) F_{\mathrm{UL}}^{\sin(\phi)} + \epsilon\sin(2\phi) F_{\mathrm{UL}}^{\sin(2\phi)}\Big] \\
+ S_L\ \lambda_l \quad & \Big[\sqrt{1-\epsilon^2} F_{\mathrm{LL}} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi) F_{\mathrm{LL}}^{\cos(\phi)}\Big] \\
+ S_T \quad & \Big[\sin(\phi-\phi_S)\left(F_{\mathrm{UT,T}}^{\sin(\phi-\phi_S)} + \epsilon F_{\mathrm{UT,L}}^{\sin(\phi-\phi_S)}\right) \\
& + \epsilon\sin(\phi+\phi_S) F_{\mathrm{UT}}^{\sin(\phi+\phi_S)} + \epsilon\sin(3\phi-\phi_S) F_{\mathrm{UT}}^{\sin(3\phi-\phi_S)} \\
& + \sqrt{2\epsilon(1+\epsilon)}\sin(\phi_S) F_{\mathrm{UT}}^{\sin(\phi_S)} \\
& + \sqrt{2\epsilon(1+\epsilon)}\sin(2\phi-\phi_S) F_{\mathrm{UT}}^{\sin(2\phi-\phi_S)}\Big] \\
+ S_T\ \lambda_l \quad & \Big[\sqrt{1-\epsilon^2}\cos(\phi-\phi_S) F_{\mathrm{LT}}^{\cos(\phi-\phi_S)} \\
& + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi_S) F_{\mathrm{LT}}^{\cos(\phi_S)} \\
& + \sqrt{2\epsilon(1-\epsilon)}\cos(2\phi-\phi_S) F_{\mathrm{LT}}^{\cos(2\phi-\phi_S)}\Big] \Big\}
\end{aligned}$$

A. Airapetian et al, Phys. Lett. B562 (2003)



Worm-gear $h^{\perp}{}_{1L}$

$$F_{UL}^{\sin 2\phi_h} = \mathcal{C}\left[-\frac{2\,(\hat{h}\cdot k_T)\,(\hat{h}\cdot p_T) - k_T\cdot p_T}{M M_h} h_{1L}^{\perp} H_1^{\perp} \right]$$

$$\frac{d\sigma^h}{dx\,dy\,d\phi_S\,dz\,d\phi\,d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2}{xyQ^2}\frac{y^2}{2(1-\epsilon)}\left(1+\frac{\gamma^2}{2x}\right)$$

$$\Big\{ \Big[F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)}\cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon\cos(2\phi) F_{UU}^{\cos(2\phi)}\Big]$$

$$+ \lambda_l \Big[\sqrt{2\epsilon(1-\epsilon)}\sin(\phi) F_{LU}^{\sin(\phi)}\Big]$$

$$+ S_L \Big[\sqrt{2\epsilon(1+\epsilon)}\sin(\phi) F_{UL}^{\sin(\phi)} + \epsilon\sin(2\phi) F_{UL}^{\sin(2\phi)}\Big]$$

$$+ S_L\,\lambda_l \Big[\sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi) F_{LL}^{\cos(\phi)}\Big]$$

$$+ S_T \Big[\sin(\phi-\phi_S)\Big(F_{UT,T}^{\sin(\phi-\phi_S)} + \epsilon F_{UT,L}^{\sin(\phi-\phi_S)}\Big) + \epsilon\sin(\phi+\phi_S) F_{UT}^{\sin(\phi+\phi_S)} + \epsilon\sin(3\phi-\phi_S) F_{UT}^{\sin(3\phi-\phi_S)} + \sqrt{2\epsilon(1+\epsilon)}\sin(\phi_S) F_{UT}^{\sin(\phi_S)} + \sqrt{2\epsilon(1+\epsilon)}\sin(2\phi-\phi_S) F_{UT}^{\sin(2\phi-\phi_S)}\Big]$$

$$+ S_T\,\lambda_l \Big[\sqrt{1-\epsilon^2}\cos(\phi-\phi_S) F_{LT}^{\cos(\phi-\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(\phi_S) F_{LT}^{\cos(\phi_S)} + \sqrt{2\epsilon(1-\epsilon)}\cos(2\phi-\phi_S) F_{LT}^{\cos(2\phi-\phi_S)}\Big]\Big\}$$

Describes the probability to find transversely polarized quarks in a longitudinally polarized nucleon

Distribution Functions

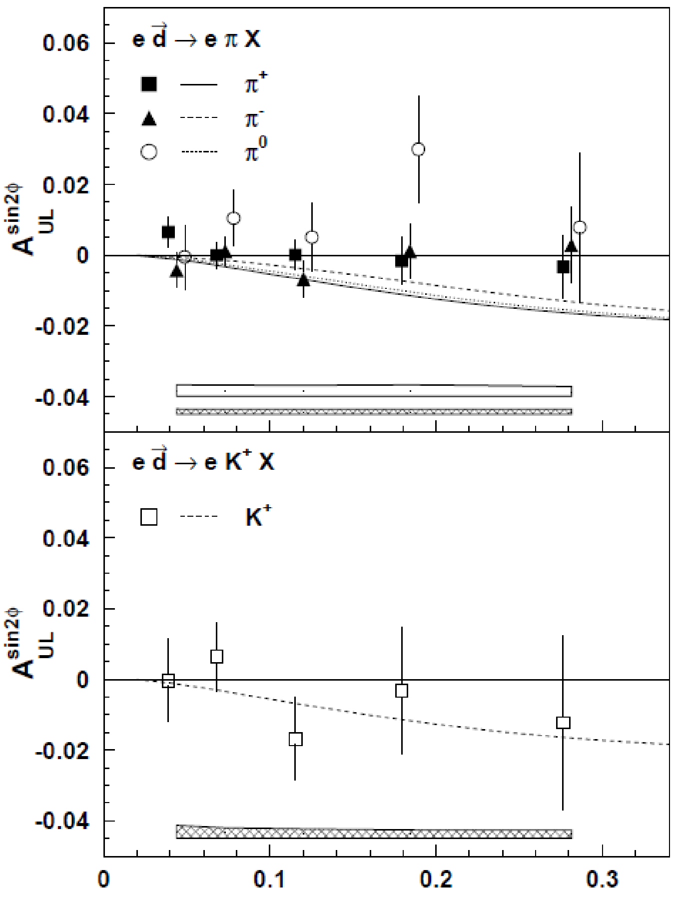
		quark U	quark L	quark T
nucleon	U	f_1		$h_1^{\perp}$
	L		g_1	$h_{1L}^{\perp}$
	T	$f_{1T}^{\perp}$	$g_{1T}^{\perp}$	h_1, $h_{1T}^{\perp}$

Fragmentation Functions

		quark U	quark L	quark T
h	U	D_1		$H_1^{\perp}$

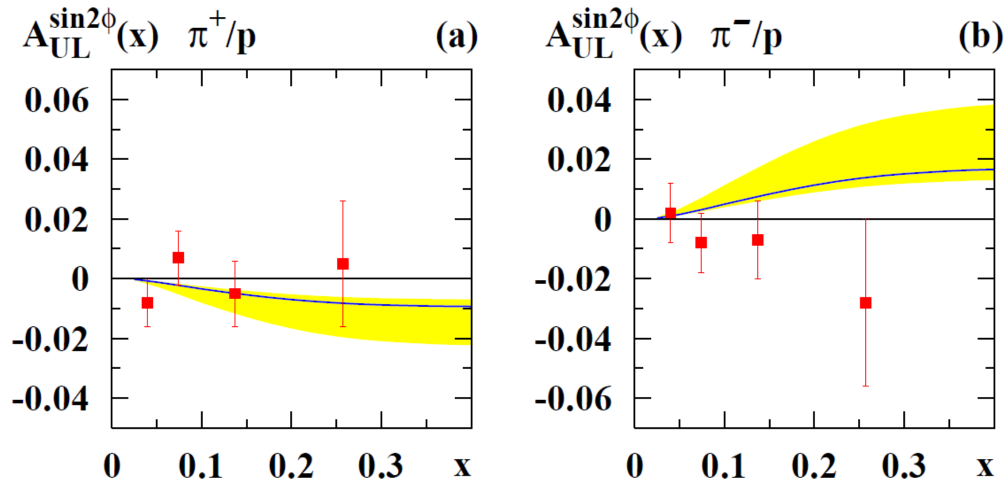
The sin(2ϕ) amplitude $\propto h_{1L}^{\perp}(x, p_T^2) \otimes H_1^{\perp}(z, k_T^2)$

Deuterium target



A. Airapetian et al, Phys. Lett. B562 (2003)

Hydrogen target



A. Airapetian et al, Phys. Rev. Lett. 84 (2000)

Amplitudes consistent with zero for all mesons and for both H and D targets

Pretzelosity

$$F_{UT}^{\sin(3\phi_h-\phi_S)} = \mathcal{C}\left[\frac{2\left(\hat{h}\cdot p_T\right)\left(p_T\cdot k_T\right)+p_T^2\left(\hat{h}\cdot k_T\right)-4\left(\hat{h}\cdot p_T\right)^2\left(\hat{h}\cdot k_T\right)}{2M^2M_h}\, h_{1T}^{\perp} H_1^{\perp}\right]$$

$$\frac{d\sigma^h}{dx\,dy\,d\phi_S\,dz\,d\phi\,d\mathbf{P}_{h\perp}^2} = \frac{\alpha^2}{xyQ^2}\frac{y^2}{2(1-\epsilon)}\left(1+\frac{\gamma^2}{2x}\right)$$

$$\Big\{ \left[F_{\mathrm{UU,T}}+\epsilon F_{\mathrm{UU,L}} + \sqrt{2\epsilon(1+\epsilon)}\cos(\phi)F_{\mathrm{UU}}^{\cos(\phi)}+\epsilon\cos(2\phi)F_{\mathrm{UU}}^{\cos(2\phi)}\right]$$

$$+\ \lambda_l\left[\sqrt{2\epsilon(1-\epsilon)}\sin(\phi)F_{\mathrm{LU}}^{\sin(\phi)}\right]$$

$$+\ S_L\left[\sqrt{2\epsilon(1+\epsilon)}\sin(\phi)F_{\mathrm{UL}}^{\sin(\phi)}+\epsilon\sin(2\phi)F_{\mathrm{UL}}^{\sin(2\phi)}\right]$$

$$+\ S_L\,\lambda_l\left[\sqrt{1-\epsilon^2}F_{\mathrm{LL}}+\sqrt{2\epsilon(1-\epsilon)}\cos(\phi)F_{\mathrm{LL}}^{\cos(\phi)}\right]$$

$$+\ S_T\left[\sin(\phi-\phi_S)\left(F_{\mathrm{UT,T}}^{\sin(\phi-\phi_S)}+\epsilon F_{\mathrm{UT,L}}^{\sin(\phi-\phi_S)}\right) +\epsilon\sin(\phi+\phi_S)F_{\mathrm{UT}}^{\sin(\phi+\phi_S)}+\epsilon\sin(3\phi-\phi_S)F_{\mathrm{UT}}^{\sin(3\phi-\phi_S)} +\sqrt{2\epsilon(1+\epsilon)}\sin(\phi_S)F_{\mathrm{UT}}^{\sin(\phi_S)} +\sqrt{2\epsilon(1+\epsilon)}\sin(2\phi-\phi_S)F_{\mathrm{UT}}^{\sin(2\phi-\phi_S)}\right]$$

$$+\ S_T\,\lambda_l\left[\sqrt{1-\epsilon^2}\cos(\phi-\phi_S)F_{\mathrm{LT}}^{\cos(\phi-\phi_S)} +\sqrt{2\epsilon(1-\epsilon)}\cos(\phi_S)F_{\mathrm{LT}}^{\cos(\phi_S)} +\sqrt{2\epsilon(1-\epsilon)}\cos(2\phi-\phi_S)F_{\mathrm{LT}}^{\cos(2\phi-\phi_S)}\right]\Big\}$$

Describes correlation between quark transverse momentum and transverse spin in a transversely pol. nucleon

- Sensitive to **non-spherical shape** of the nucleon

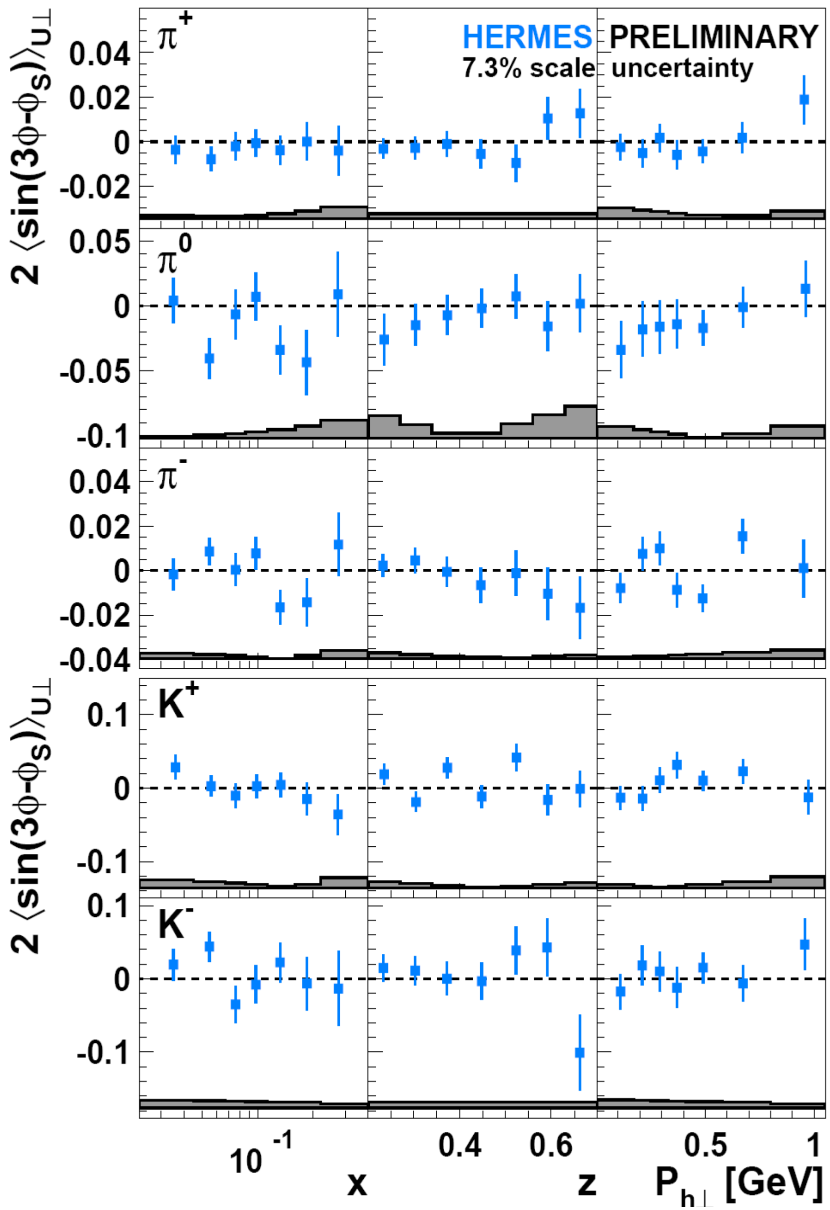
Distribution Functions

nucleon \ quark	U	L	T
U	f_1		$h_1^{\perp}$
L		g_1	$h_{1L}^{\perp}$
T	$f_{1T}^{\perp}$	$g_{1T}^{\perp}$	h_1, $h_{1T}^{\perp}$

Fragmentation Functions

h \ quark	U	L	T
U	D_1		$H_1^{\perp}$

The sin(3ϕ-ϕ_S) amplitude $\propto h_{1T}^{\perp}(x, p_T^2) \otimes H_1^{\perp}(z, k_T^2)$

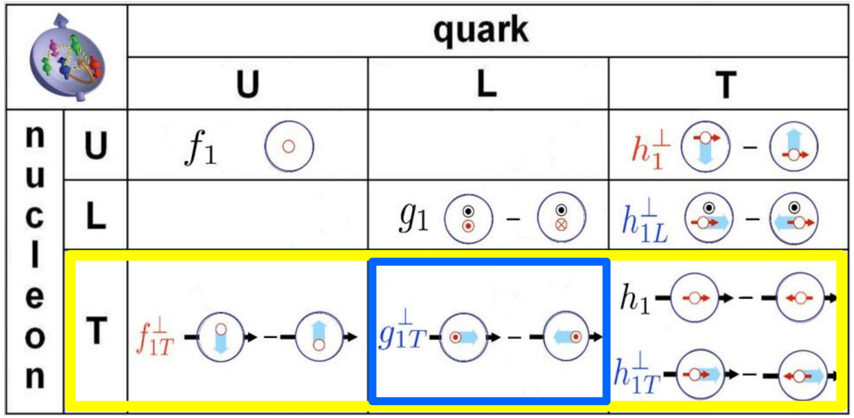


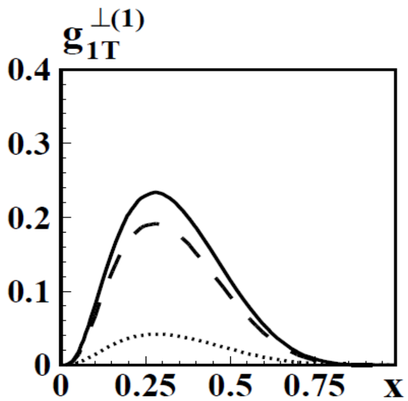
All amplitudes consistent with zero

…suppressed by two powers of $P_{h\perp}$ w.r.t. Collins and Sivers amplitudes

The worm-gear $g_{1T}{}^{\perp}$

- The only TMD that is both chiral-even and naïve-T-even
- requires interference between wave funct. components that differ by 1 unit of OAM





S. Boffi et al. (2009)
Phys. Rev. D 79 094012
Light cone constituent quark model
flavorless
dashed line: interf. L=0, L=1
dotted line: interf L=1, L=2

$\Rightarrow$ related to quark orbital motion inside nucleons

- Many models support simple relations among $g_{1T}^{\perp}$ and other TMDs:
  - $g_{1T}^{q} = -h_{1L}^{\perp q}$ (also supported by Lattice QCD and first data)
  - $g_{1T}^{q(1)}(x) \overset{WW-type}{\approx} x \int_x^1 \frac{dy}{y} g_1^q(y)$ (Wandzura-Wilczek appr.)

Probing $g_{1T}^{\perp}$ through Double Spin Asymmetries

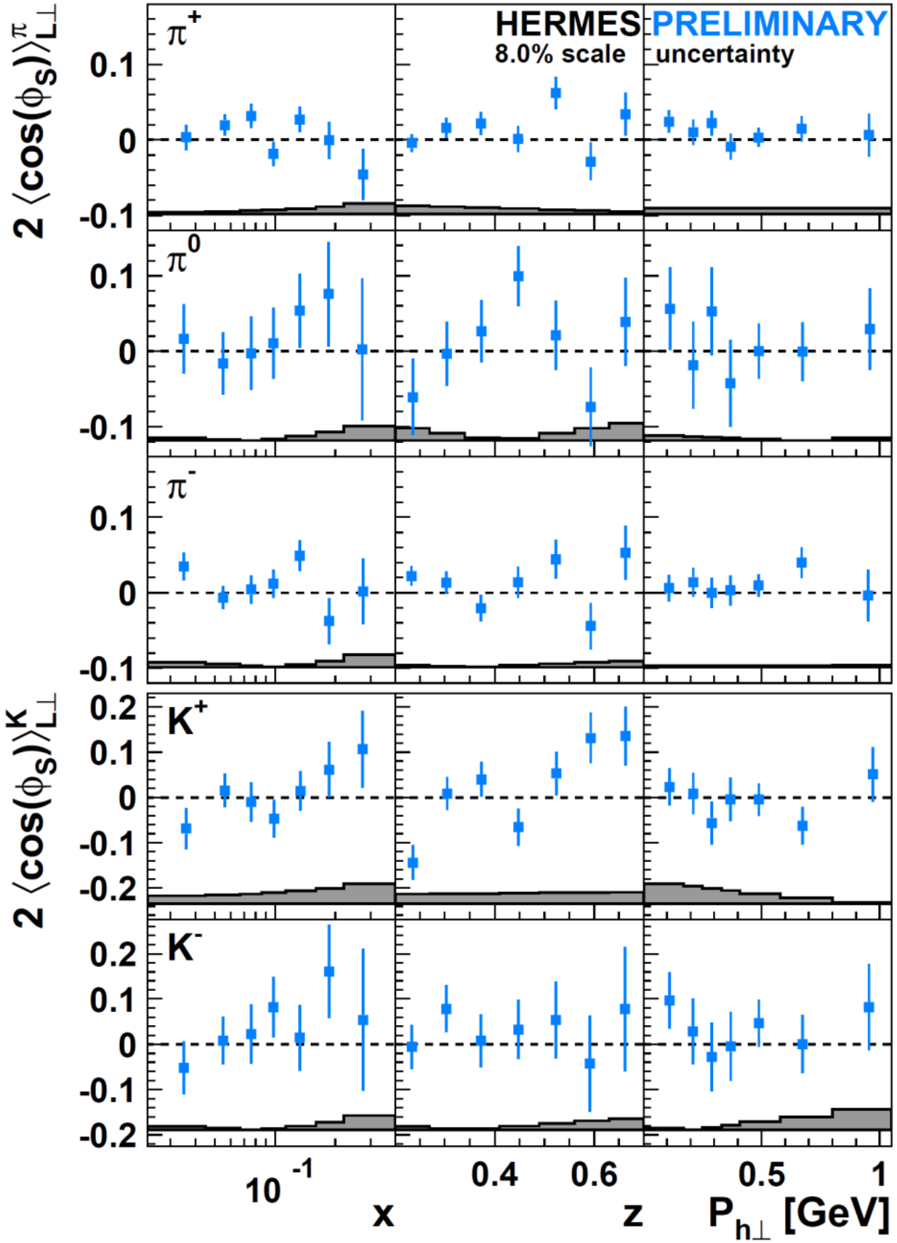
$$F_{LT}^{\cos(\phi_h-\phi_S)} = \mathcal{C}\left[\frac{\hat{\boldsymbol{h}}\cdot\boldsymbol{p}_T}{M}\, g_{1T}\, D_1\right]$$

$$\begin{aligned}F_{LT}^{\cos\phi_S} = \frac{2M}{Q}\,\mathcal{C}\Big\{-\Big(x g_T D_1 + \frac{M_h}{M}\, h_1 \frac{\tilde{E}}{z}\Big) \\ + \frac{\boldsymbol{k}_T\cdot\boldsymbol{p}_T}{2MM_h}\Big[\Big(x e_T H_1^{\perp} - \frac{M_h}{M}\, g_{1T} \frac{\tilde{D}^{\perp}}{z}\Big) + \Big(x e_T^{\perp} H_1^{\perp} + \frac{M_h}{M} f_{1T}^{\perp} \frac{\tilde{G}^{\perp}}{z}\Big)\Big]\Big\}\end{aligned}$$

$$\begin{aligned}F_{LT}^{\cos(2\phi_h-\phi_S)} = \frac{2M}{Q}\,\mathcal{C}\Big\{-\frac{2\,(\hat{\boldsymbol{h}}\cdot\boldsymbol{p}_T)^2 - \boldsymbol{p}_T^2}{2M^2}\Big(x g_T^{\perp} D_1 + \frac{M_h}{M}\, h_{1T}^{\perp} \frac{\tilde{E}}{z}\Big) \\ + \frac{2\,(\hat{\boldsymbol{h}}\cdot\boldsymbol{k}_T)\,(\hat{\boldsymbol{h}}\cdot\boldsymbol{p}_T) - \boldsymbol{k}_T\cdot\boldsymbol{p}_T}{2MM_h}\Big[\Big(x e_T H_1^{\perp} - \frac{M_h}{M}\, g_{1T} \frac{\tilde{D}^{\perp}}{z}\Big) \\ - \Big(x e_T^{\perp} H_1^{\perp} + \frac{M_h}{M} f_{1T}^{\perp} \frac{\tilde{G}^{\perp}}{z}\Big)\Big]\Big\}\end{aligned}$$

The simplest way to probe worm-gear $g_{1T}^{\perp}$ is through the $\cos(\phi - \phi_S)$ Fourier component

The cos(ϕ_S) Fourier component



≈ 0

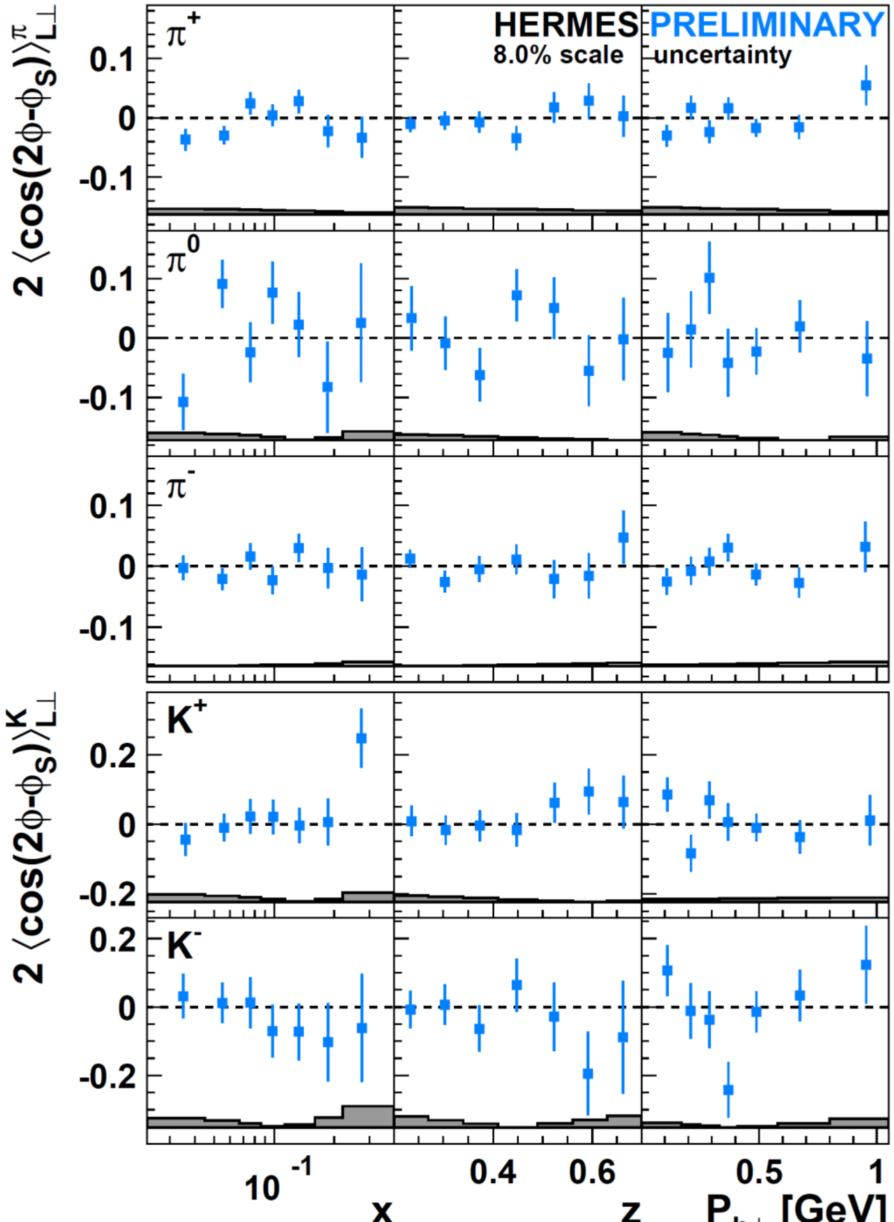
≈ 0

≈ 0

≈ 0
sign change?

≈ 0

The cos(2ϕ-ϕ$_S$) Fourier component



≈ 0

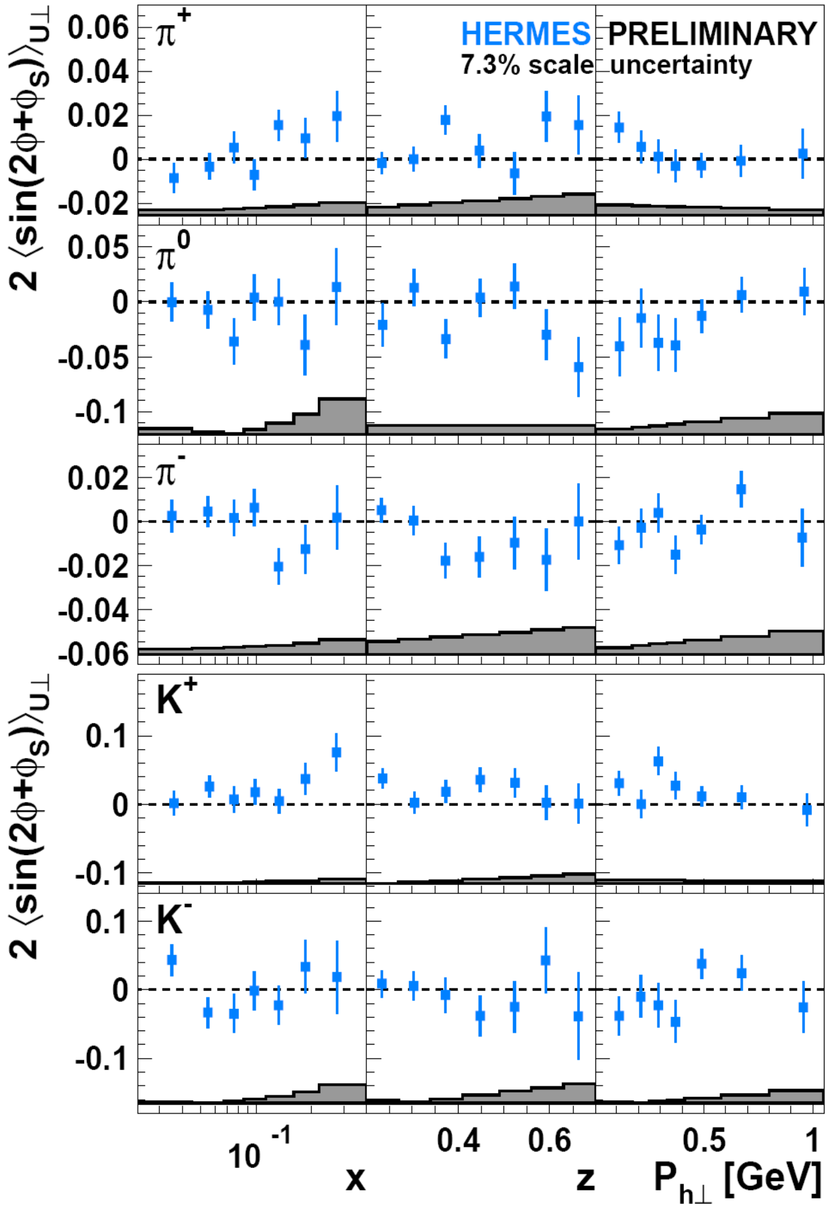
≈ 0

≈ 0

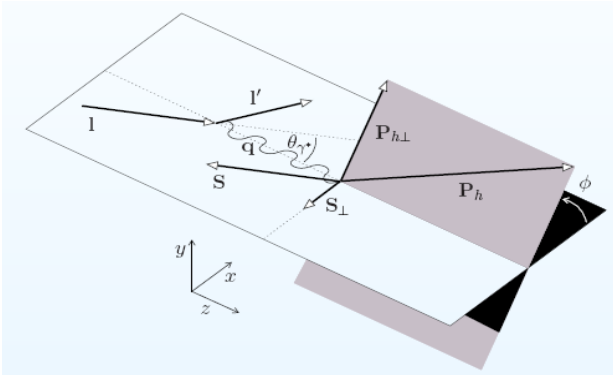
≈ 0

≈ 0

The sin(2ϕ+ϕ$_S$) Fourier component



- arises solely from longitudinal (w.r.t. virtual photon direction) component of the target spin

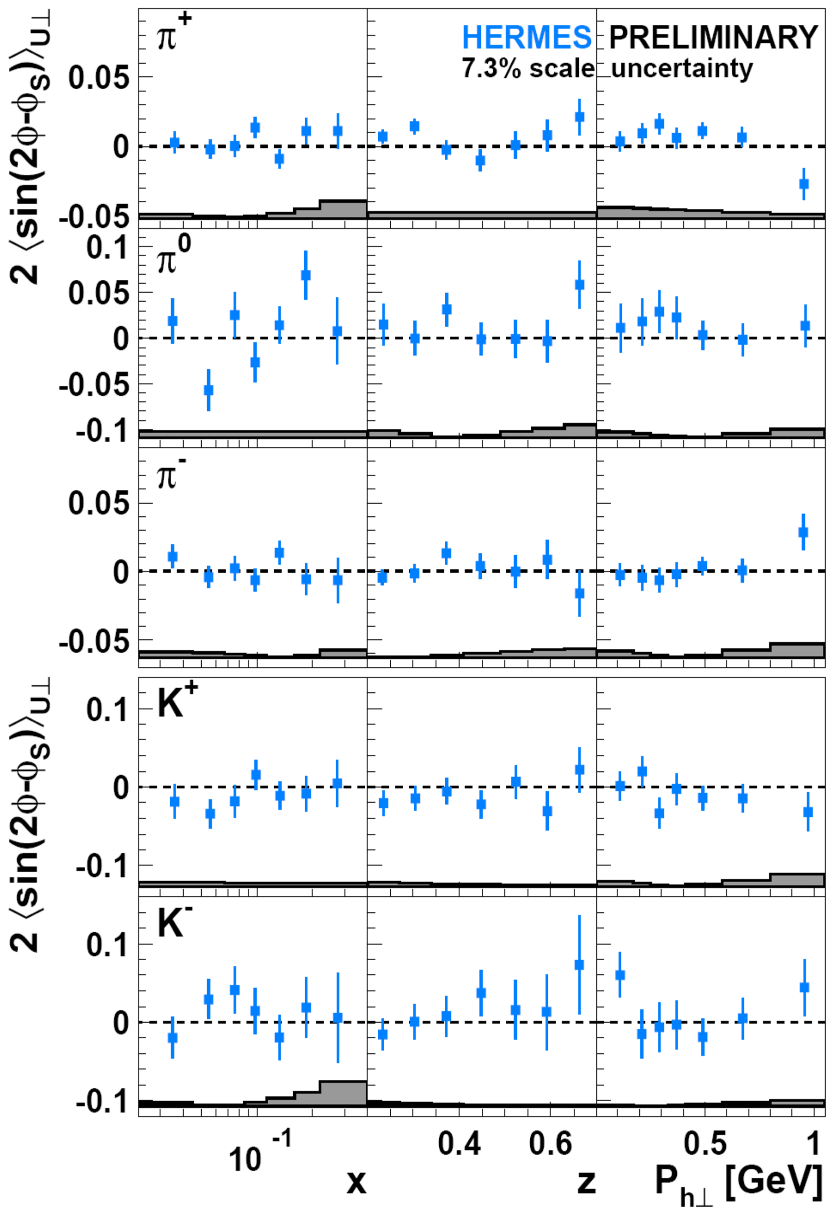


- related to <sin(2ϕ)>$_{UL}$ Fourier comp:

$$2\left\langle \sin(2\phi+\phi_S)\right\rangle_{UT}^{h} \propto \frac{1}{2}\sin(\vartheta_{l\gamma^*})2\left\langle \sin(2\phi)\right\rangle_{UL}^{h}$$

- sensitive to **worm-gear** $h_{1L}^{\perp}$
- suppressed by one power of $P_{h\perp}$ w.r.t. Collins and Sivers amplitudes
- **no significant signal observed (except maybe for K+)**

The subleading-twist sin(2ϕ-ϕ_S) Fourier component



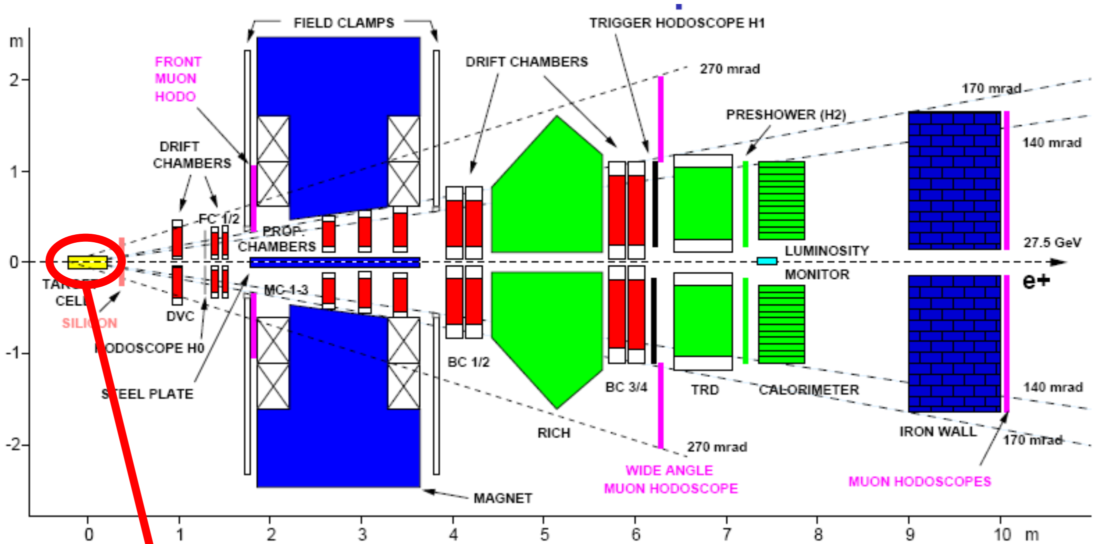
- sensitive to **worm-gear** $g_{1T}^{\perp}$, **Pretzelosity** and **Sivers function**:

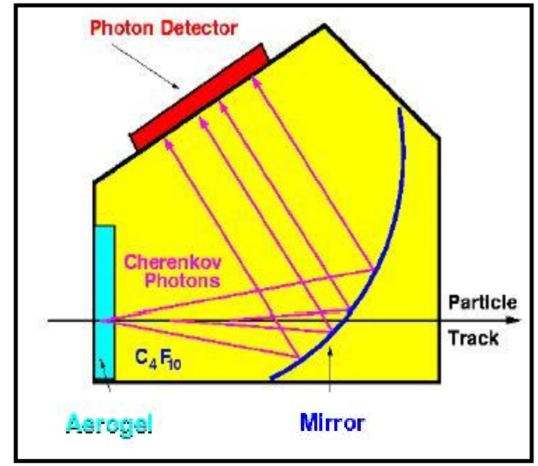
$$\propto \mathcal{W}_1(p_T,k_T,P_{h\perp})\left(x f_T^{\perp} D_1 - \frac{M_h}{M} h_{1T}^{\perp} \frac{\tilde{H}}{z}\right) - \mathcal{W}_2(p_T,k_T,P_{h\perp})\left[\left(x h_T H_1^{\perp} + \frac{M_h}{M} g_{1T} \frac{\tilde{G}^{\perp}}{z}\right) + \left(x h_T^{\perp} H_1^{\perp} - \frac{M_h}{M} f_{1T}^{\perp} \frac{\tilde{D}^{\perp}}{z}\right)\right]$$

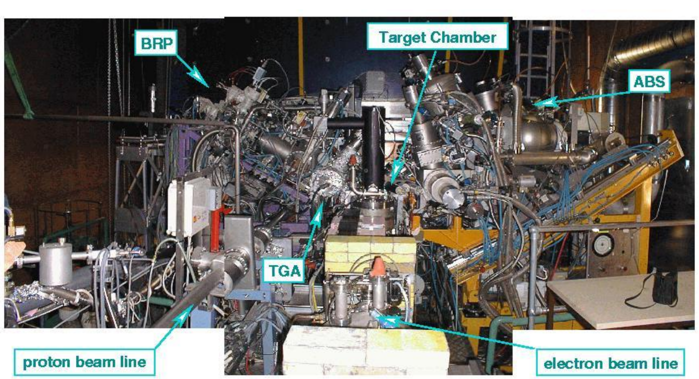
- suppressed by one power of $P_{h\perp}$ w.r.t. Collins and Sivers amplitudes

- **no significant non-zero signal observed**

The HERMES experiment at HERA

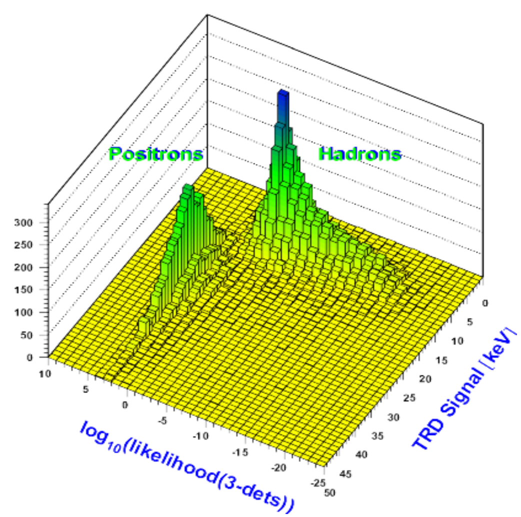




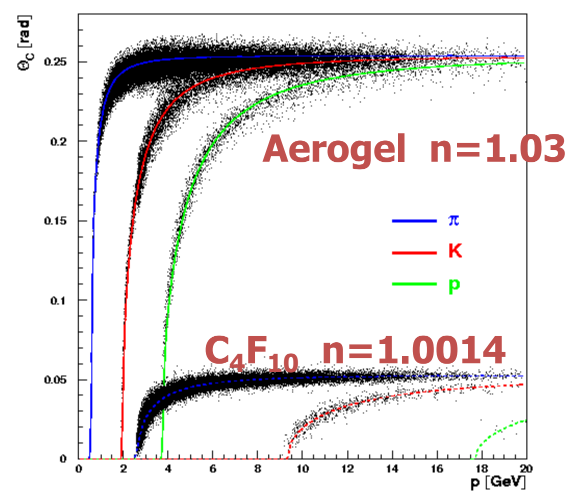


TRD, Calorimeter, preshower, RICH:

lepton-hadron > 98%

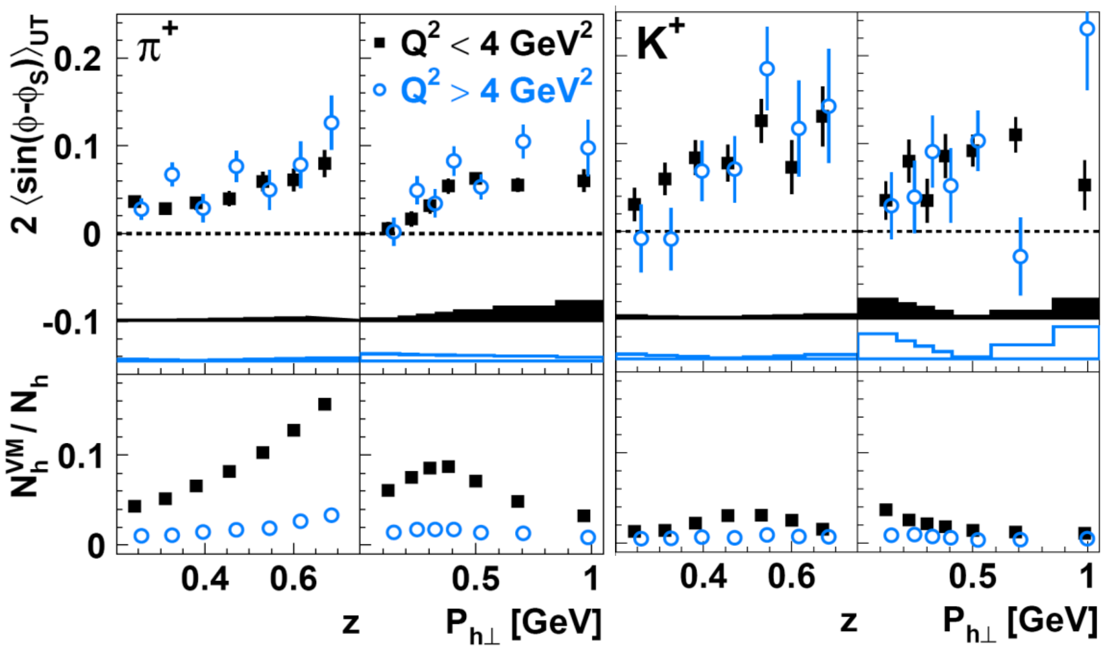


hadron separation



π ~ 98%, K ~ 88% , P ~ 85%

Siver amplitudes: additional studies

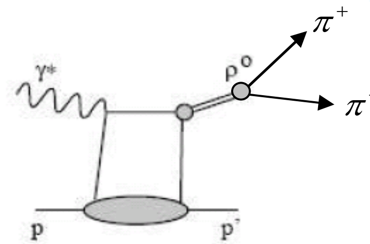


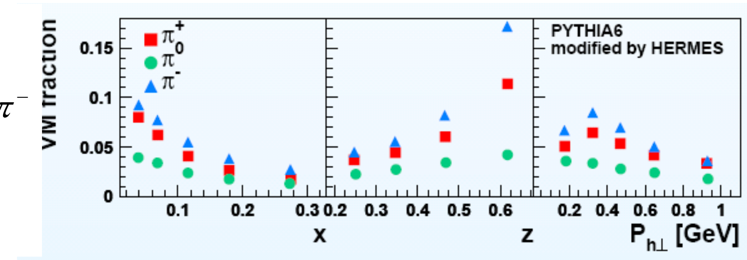
- No systematic shifts observed between high and low Q^2 amplitudes for both π^+ and K^+

  No indication of important contributions from exclusive VM

The pion-difference asymmetry

Contribution by decay of exclusively produced vector mesons (ρ^0,ω,ϕ) is not negligible (6-7% for pions and 2-3% for kaons), though substatially limited by the requirement $z<0.7$.

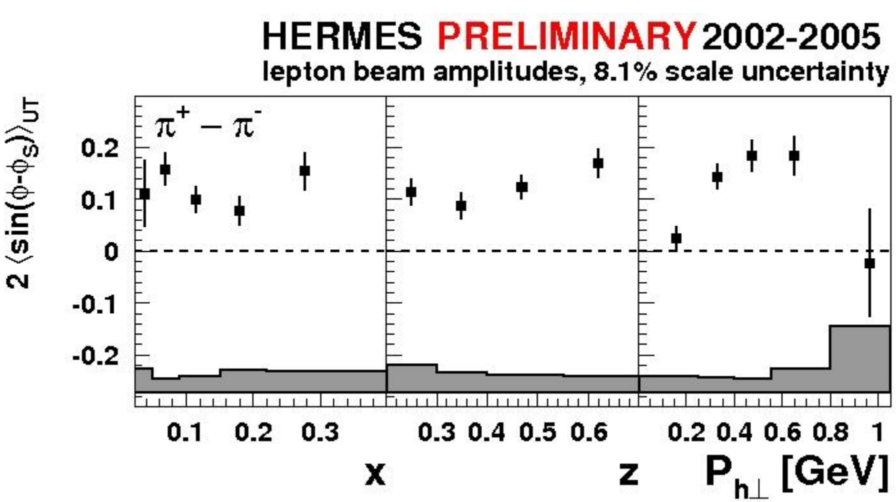


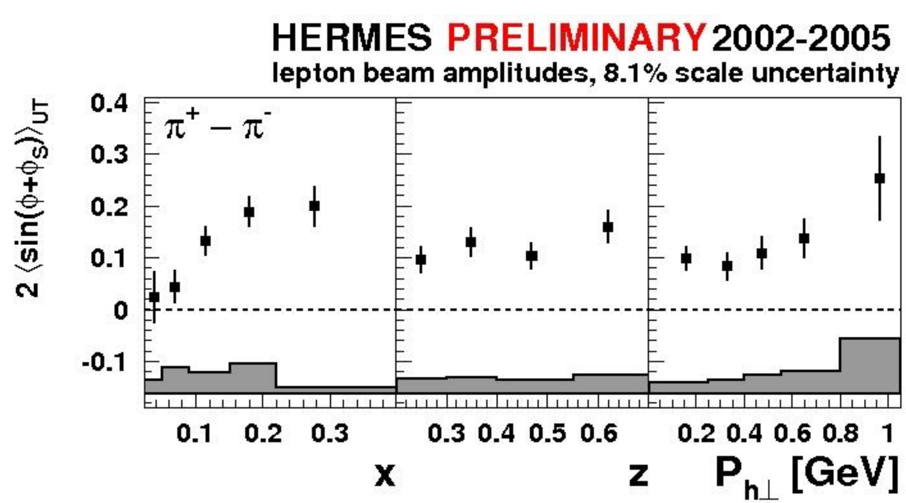


a new observable

$$A_{UT}^{\pi^+-\pi^-}(\phi,\phi_S) \equiv \frac{1}{P_T}\frac{(\sigma_{U\uparrow}^{\pi^+}-\sigma_{U\uparrow}^{\pi^-})-(\sigma_{U\downarrow}^{\pi^+}-\sigma_{U\downarrow}^{\pi^-})}{(\sigma_{U\uparrow}^{\pi^+}-\sigma_{U\uparrow}^{\pi^-})+(\sigma_{U\downarrow}^{\pi^+}-\sigma_{U\downarrow}^{\pi^-})}$$

Contribution from exclusive ρ^0 largely cancels out!





- significantly positive Sivers and Collins amplitudes are obtained
- measured amplitudes are not generated by exclusive VM contribution