

Pacific Spin 2019

The 11th Circum-Pan-Pacific Symposium on High Energy Spin Physics

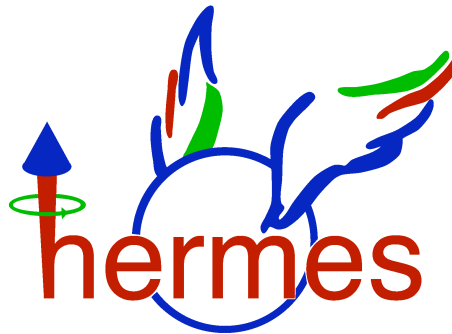
Date: **August 27-30, 2019**

Venue: **ANA Holiday Inn Resort Miyazaki**, Miyazaki in Kyusyu island of Japan

Longitudinal double-spin asymmetries in SIDIS studied at HERMES



Y. Miyachi, 山形大学
Yamagata University

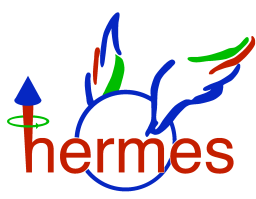


For the

collaboration

“Longitudinal double-spin asymmetries in semi-inclusive deep-inelastic scattering of electrons and positrons by protons and deuterons”, Phys. Rev. D 99 (2019) 112001

- Spin structure of the proton
- Semi-inclusive measurement of Deep Inelastic Scattering
- HERMES experiment
- Longitudinal double spin asymmetry
 - x , z , and $P_{h\perp}$ dependence
 - 3D results $(x, z, P_{h\perp})$
 - Charge difference asymmetry
 - Azimuthal modulation, $\cos\phi$
- Summary

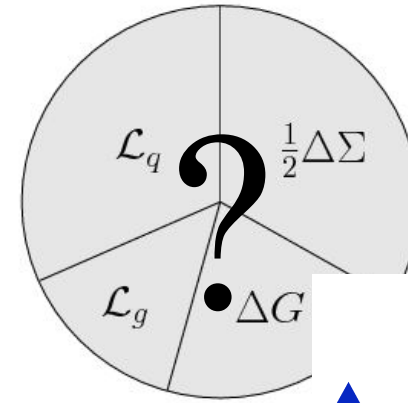


Spin structure of the proton studied with a fixed target



1980

$$\frac{1}{2} =$$



EMC, "Spin Crisis"

1st generation

1990

Inclusive DIS $\Rightarrow \Delta\Sigma \sim 0.3, \Delta G$

2nd generation

Semi-Inclusive DIS \Rightarrow Spin-Flavor structure

Transverse polarized target \Rightarrow TMD

Open Charm, high pt $\Rightarrow \Delta G$

Exclusive process \Rightarrow GPD

3rd generation

Polarized Drell-Yan

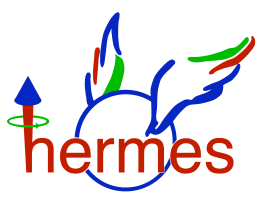
4th generation?

Pacific Spin

- 1996 Kobe
- 1999 Wako
- 2001 Beijing
- 2003 Washington
- 2005 Tokyo
- 2007 Vancouver
- 2009 Yamagata
- 2011 Cairns
- 2013 Ji'nan
- 2015 Taipei
- 2019 Miyazaki
- 2020



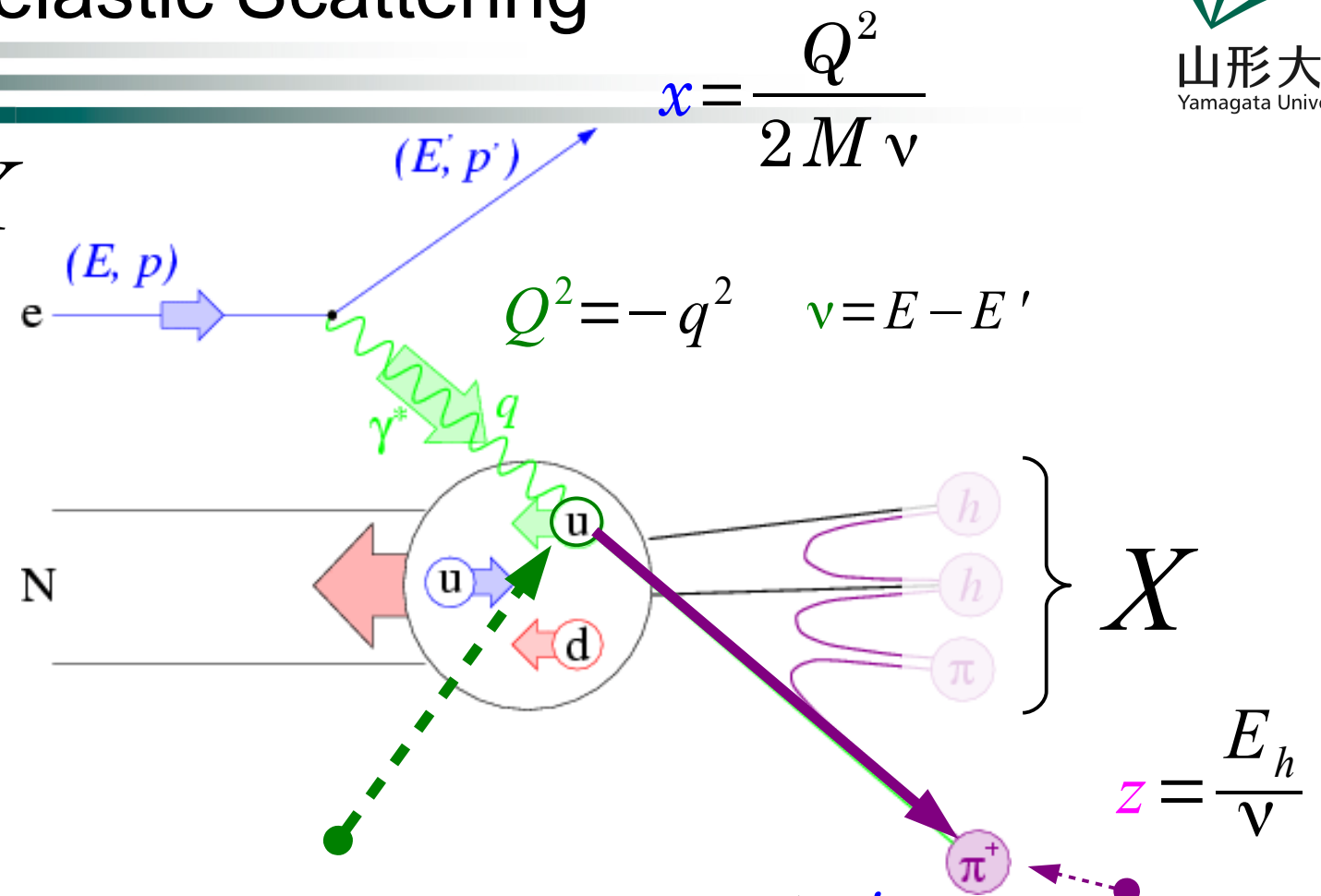
collider exp.



Semi-inclusive measurement of Deep Inelastic Scattering




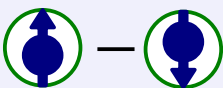

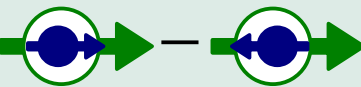
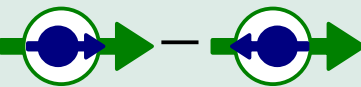


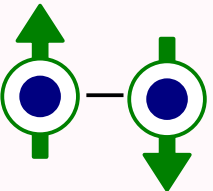
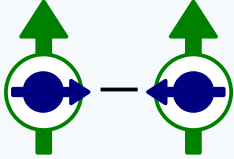

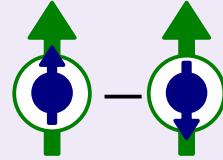
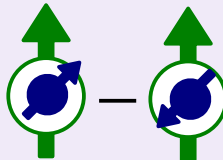



$$e N \rightarrow e' h X$$



$$\sigma^{eN \rightarrow e' h X} = \sum_q f^{N \rightarrow q} \otimes \hat{\sigma}^{eq \rightarrow e' q} \otimes F^{q \rightarrow h}$$

P arton D istribution F unction $f(x, Q^2)$	F ragmentation F unction $F(z, Q^2)$
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<p>Unpolarized Nucleon</p> 	<p>f_1 Parton Density</p> 		 <p>h_1^\perp Boer-Mulders</p>
<p>Longitudinally Polarized Nucleon</p> 	<p>$g_{1L} (= \Delta q)$ Helicity</p> 		 <p>h_{1L}^\perp Mulders-Kotzinian</p>
<p>Transversely Polarized Nucleon</p> 	<p>f_{1T}^\perp Sivers</p> 	<p>g_{1T} Worm-gear</p>  	<p>$h_{1T} (= \delta q)$ Transversity</p>  <hr style="border-top: 1px dashed black;"/> <p>h_{1T}^\perp Pretzelosity</p>  

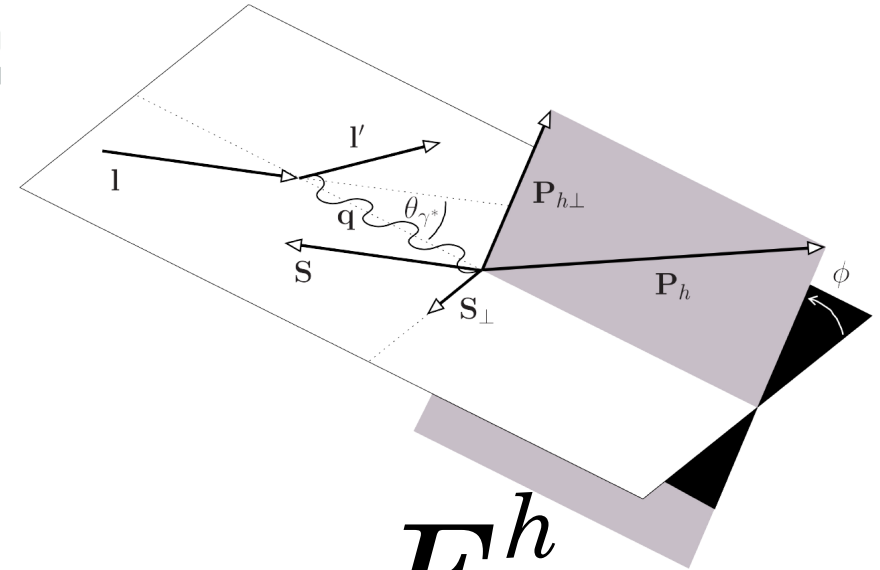


※ transverse nucleon polarization neglected

$$\frac{d\sigma^h}{dx dy dz dP_{h\perp}^2 d\phi} = \frac{2\pi\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x} \right)$$

$$\left\{ \begin{aligned} & F_{UU,T}^h + \epsilon F_{UU,L}^h + \lambda\Lambda\sqrt{1-\epsilon^2} F_{LL}^h \\ & + \sqrt{2}\epsilon \left[\lambda\sqrt{1-\epsilon} F_{LU}^{h,\sin\phi} + \Lambda\sqrt{1+\epsilon} F_{UL}^{h,\sin\phi} \right] \sin\phi \\ & + \sqrt{2}\epsilon \left[\lambda\Lambda\sqrt{1-\epsilon} F_{LL}^{h,\cos\phi} + \sqrt{1+\epsilon} F_{UU}^{h,\cos\phi} \right] \cos\phi \end{aligned} \right.$$

$$+ \Lambda\epsilon F_{UL}^{h,\sin\phi} \quad F_{LL}^h \propto \sum_q e_q^2 \left[g_{1L}^q(x, p_T^2) \otimes \mathcal{W}_1 D_1^{q \rightarrow h}(z, k_T^2) \right]$$



F^h
 beam **LL** target

Longitudinal double spin asymmetry

$$A_{LL}^h \equiv \frac{\sigma_{+-}^h - \sigma_{++}^h + \sigma_{-+}^h - \sigma_{--}^h}{\sigma_{+-}^h + \sigma_{++}^h + \sigma_{-+}^h + \sigma_{--}^h}$$

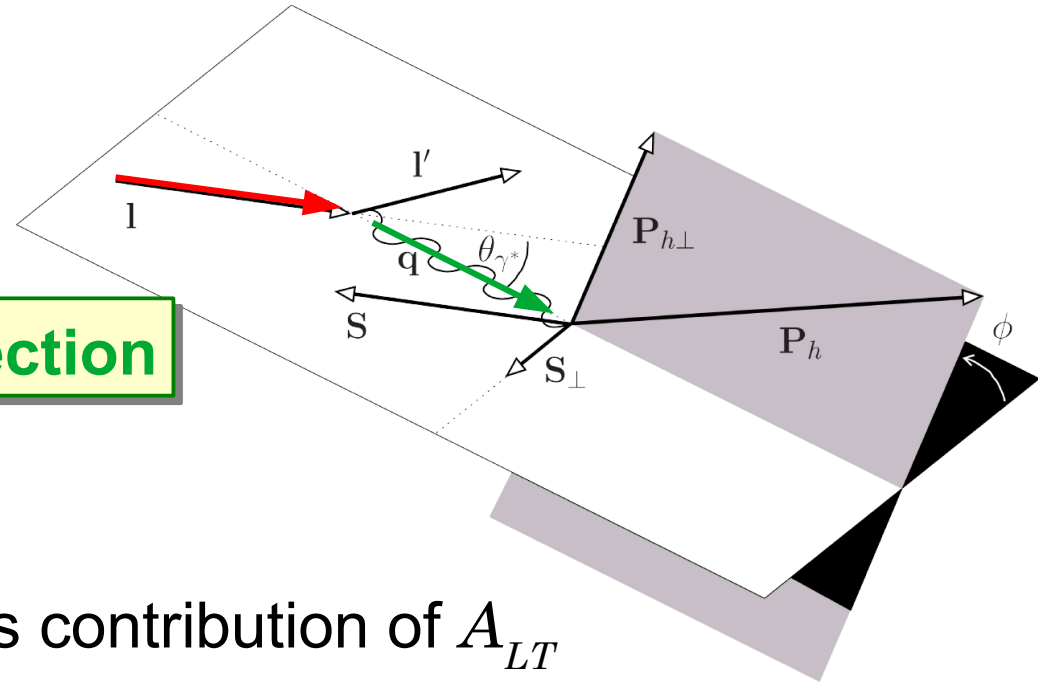
$$A_{LL}^h \equiv \frac{\sigma_{+-}^h - \sigma_{++}^h + \sigma_{-+}^h - \sigma_{--}^h}{\sigma_{+-}^h + \sigma_{++}^h + \sigma_{-+}^h + \sigma_{--}^h}$$

along **virtual photon direction**

along **beam direction**

$A_{||}$

includes contribution of A_{LT}



A_1^h

virtual-photon-nucleon asymmetry

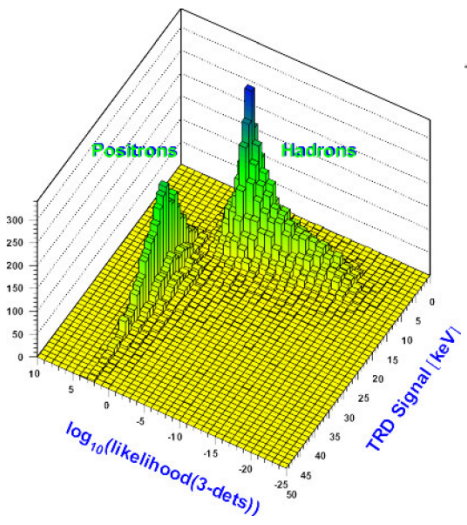
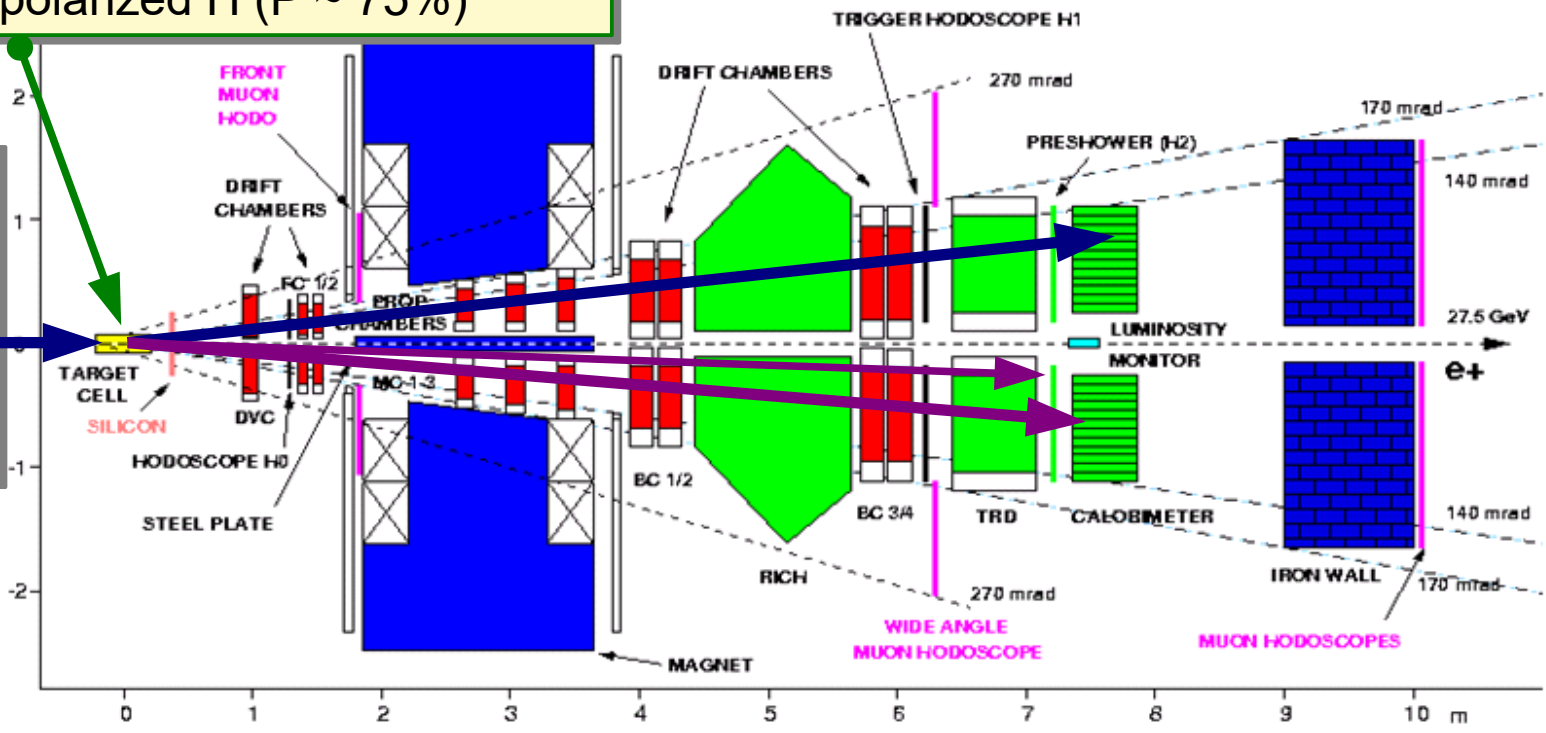
$$A_1^h = \frac{\sigma_{1/2}^h - \sigma_{3/2}^h}{\sigma_{1/2}^h + \sigma_{3/2}^h} = \frac{1}{D(1+\eta\gamma)} A_{||}^h$$

$$D = \frac{1-(1-y)\epsilon}{1+\epsilon R}$$

$$\eta = \frac{\epsilon\gamma y}{1-(1-y)\epsilon}$$

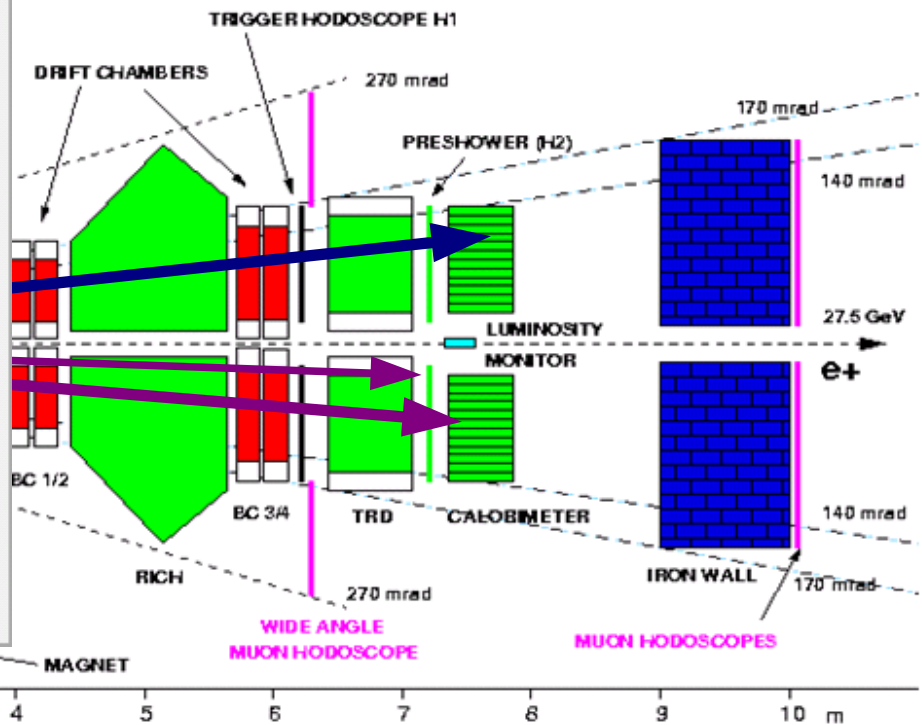
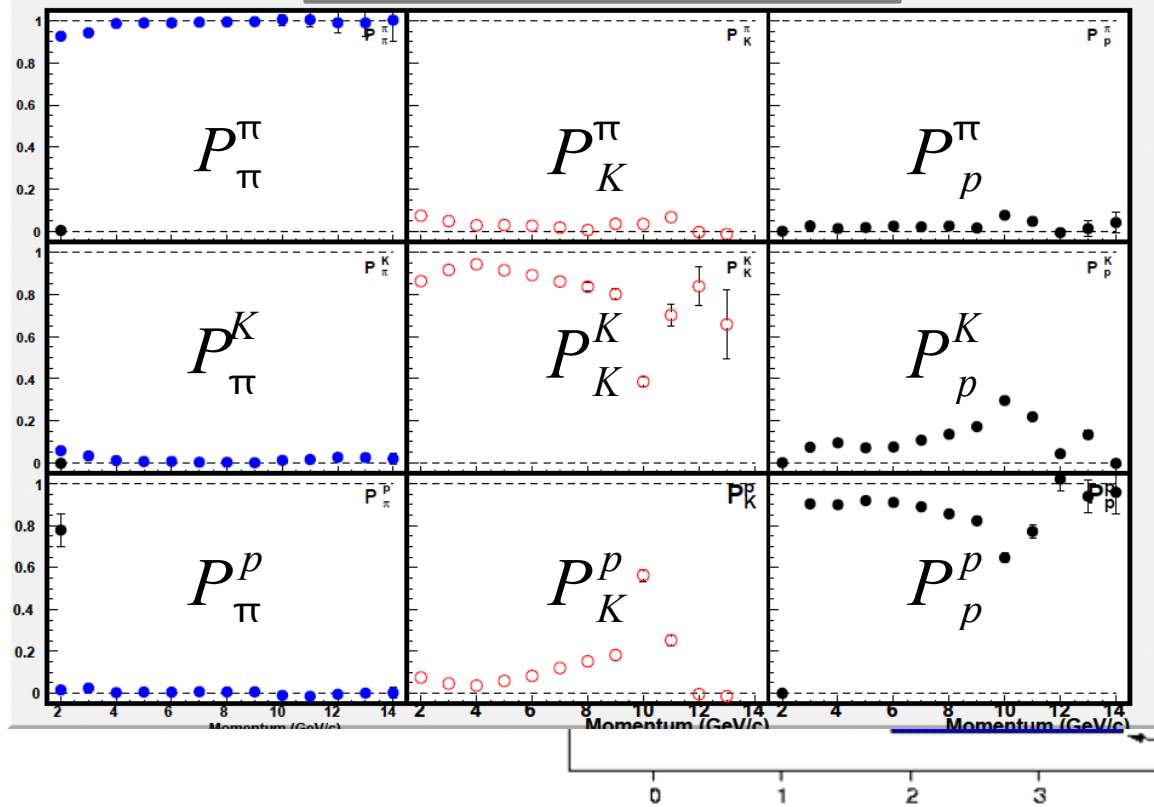
Targets: Unpolarized H, D, nuclei
Longitudinally polarized H, D ($P \sim 85\%$)
Transversely polarized H ($P \sim 75\%$)

HERA: 27.6 GeV
polarized
 electron/positron



Reconstruction: $\Delta p/p < 2\%$, $\Delta\theta < 0.6$ mrad
Lepton selection efficiency: > 99%
 with hadron contamination < 1%
 Hadron ID with RICH: π, K, p in $2 < p < 15$ GeV/c

RICH PID efficiency



Reconstruction: $\Delta p/p < 2\%$, $\Delta\theta < 0.6$ mrad

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Hadron ID with RICH: π, K, p in $2 < p < 15$ GeV/c

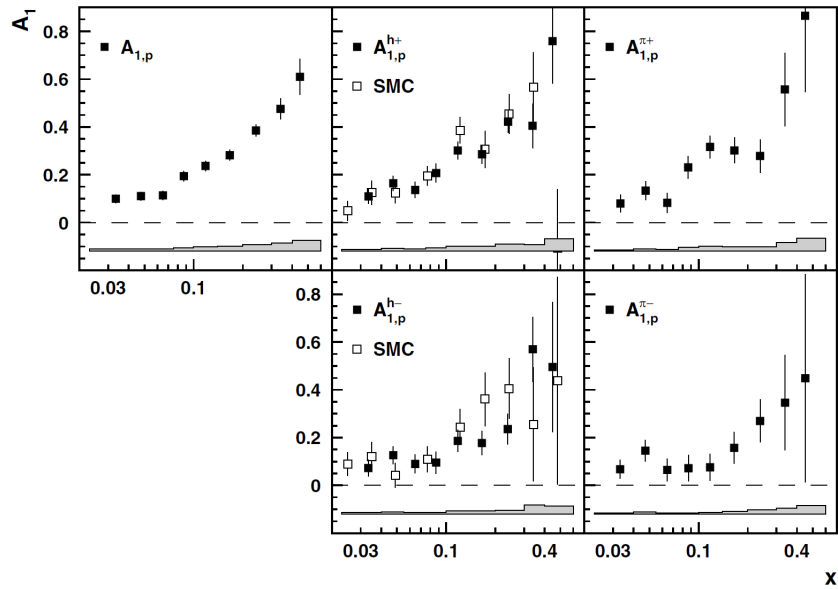
“Flavor Decomposition of the Polarized Quark Distributions in the Nucleon from Inclusive and Semi-Inclusive Deep-inelastic Scattering”, Phys. Lett. B464 (1999) 123-134

“Flavor Decomposition of the Sea Quark Helicity Distributions in the Nucleon from Semi-inclusive Deep-inelastic Scattering”, Phys. Rev. Lett. 92 (2004) 012005

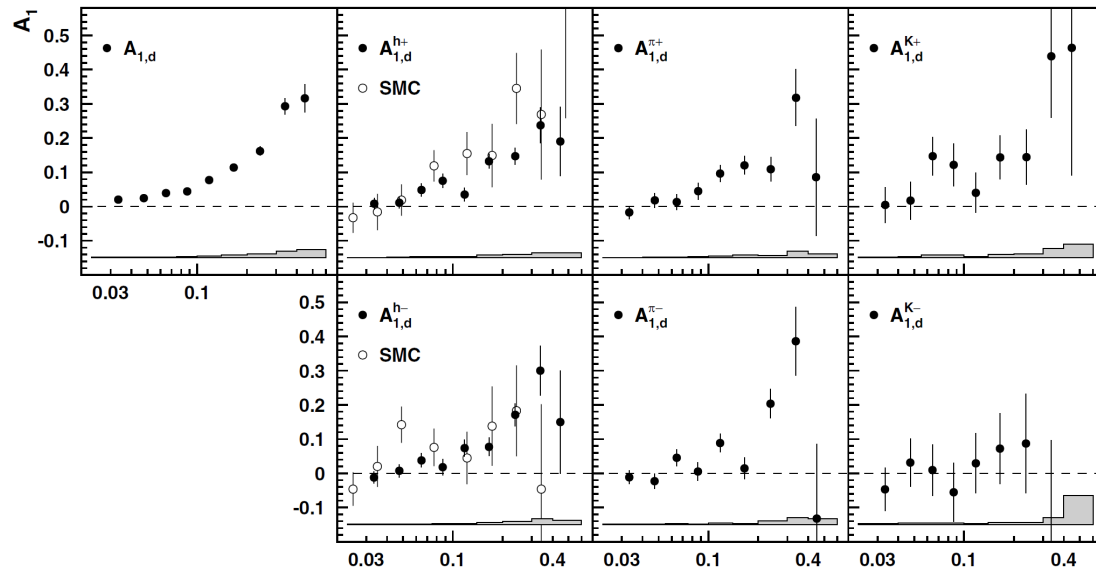
“Quark Helicity Distributions in the Nucleon for up, down, and strange Quarks from Semi-inclusive Deep-inelastic Scattering”, Phys. Rev D 71 (2005) 012003

“Measurement of Parton Distributions of Strange Quarks in the Nucleon from Charged-Kaon Production in Deep-Inelastic Scattering on the Deuteron”, Phys. Lett. B 666 (2008) 446

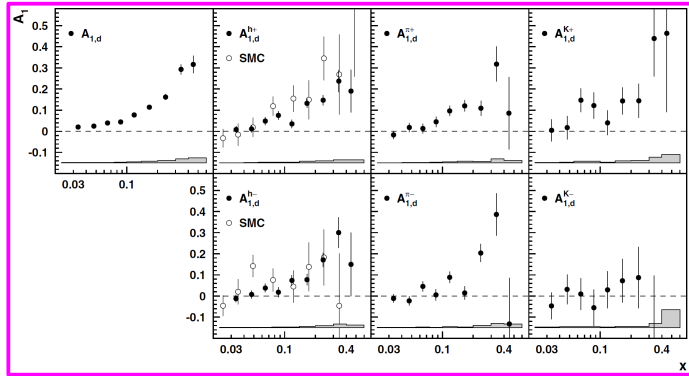
“Longitudinal double-spin asymmetries in semi-inclusive deep-inelastic scattering of electrons and positrons by protons and deuterons”, Phys. Rev. D 99 (2019) 112001



Year	Lepton beam charge	Average polarization	Fractional uncertainty
1996	e^+	52.8%	3.4%
1997	e^+	53.1%	3.4%
1998	e^-	52.1%	3.4%
1999	e^+	53.3%	1.8%
2000	e^+	53.3%	1.9%

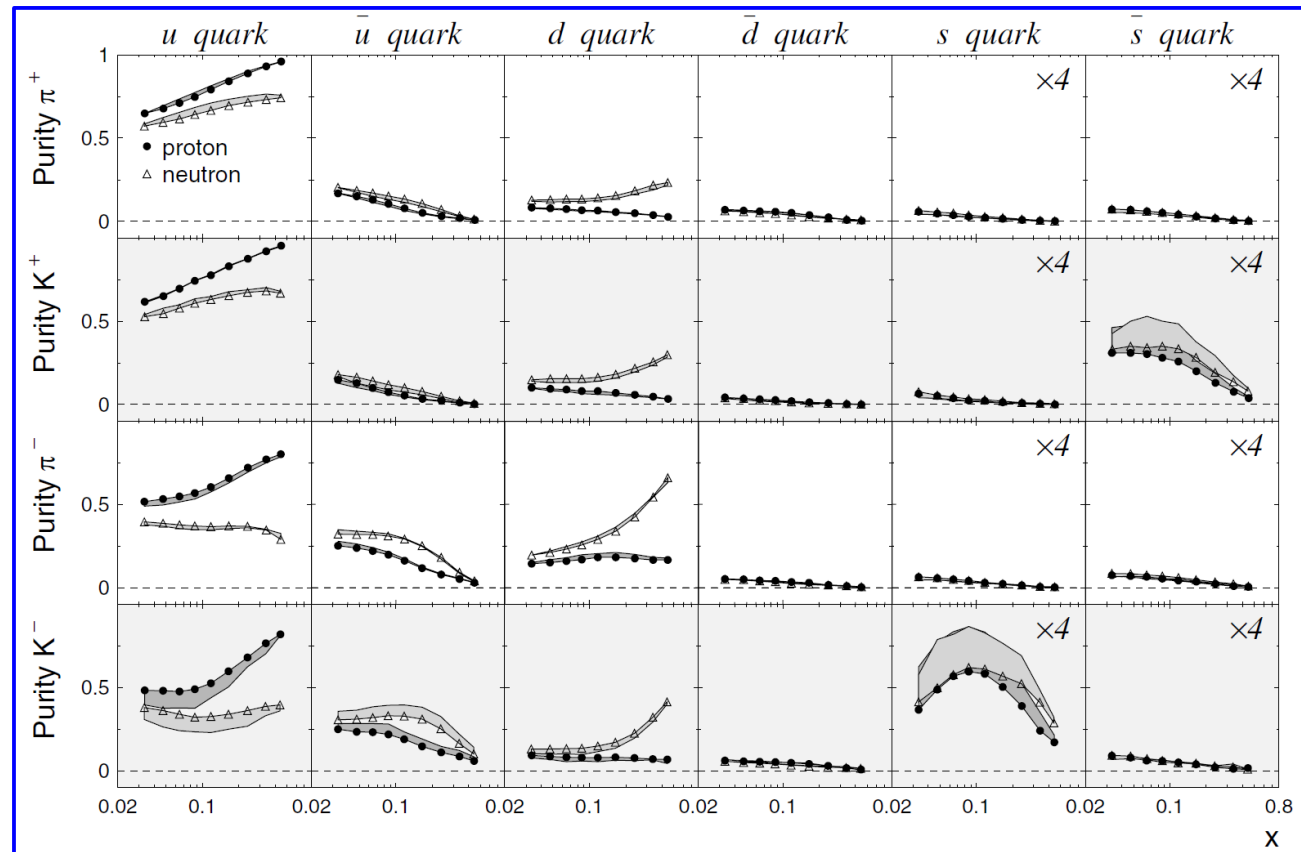


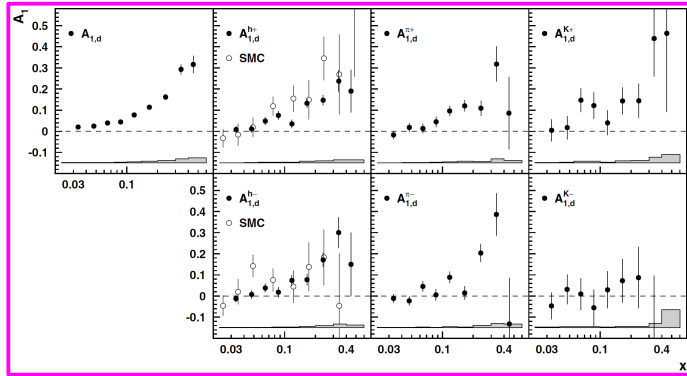
$$A_1^h = \frac{\sigma_{1/2}^h - \sigma_{3/2}^h}{\sigma_{1/2}^h + \sigma_{3/2}^h}$$



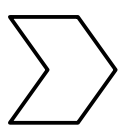
$$A_{1N}^h = P \cdot \frac{\Delta q}{q}$$

Estimated using MC tuned to the HERMES data

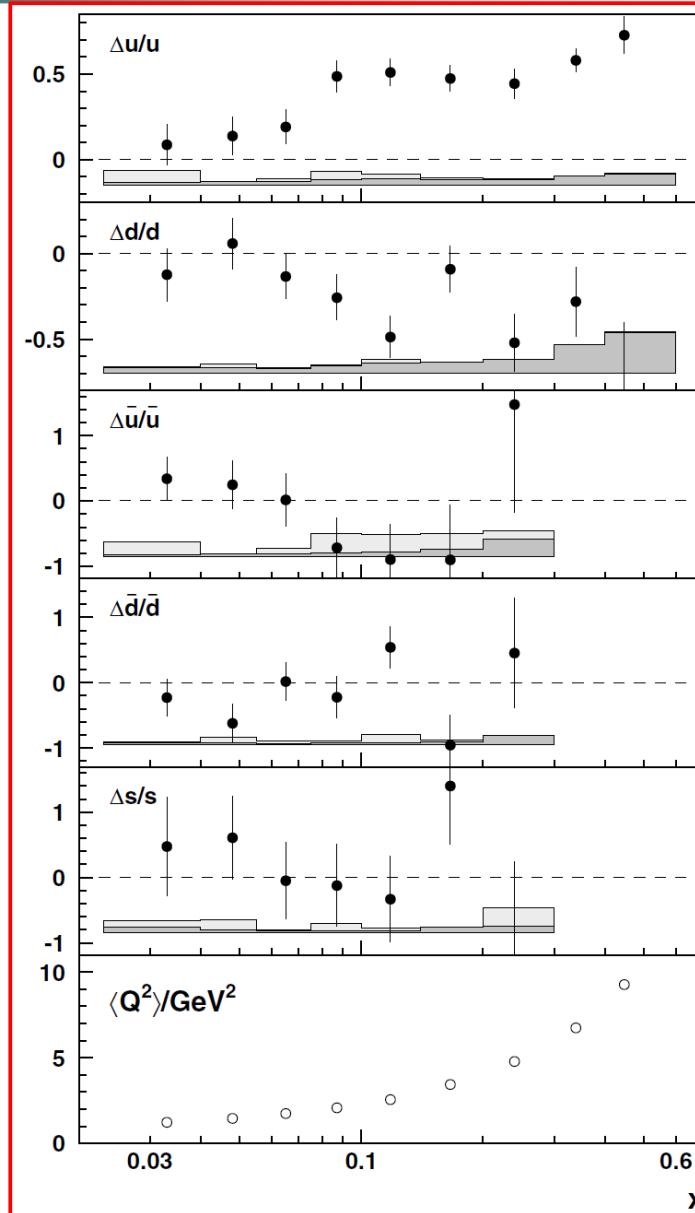




$$A_{1N}^h = P \cdot \frac{\Delta q}{q}$$



$$\frac{\Delta q}{q} = P^{-1} A_{1N}^h$$



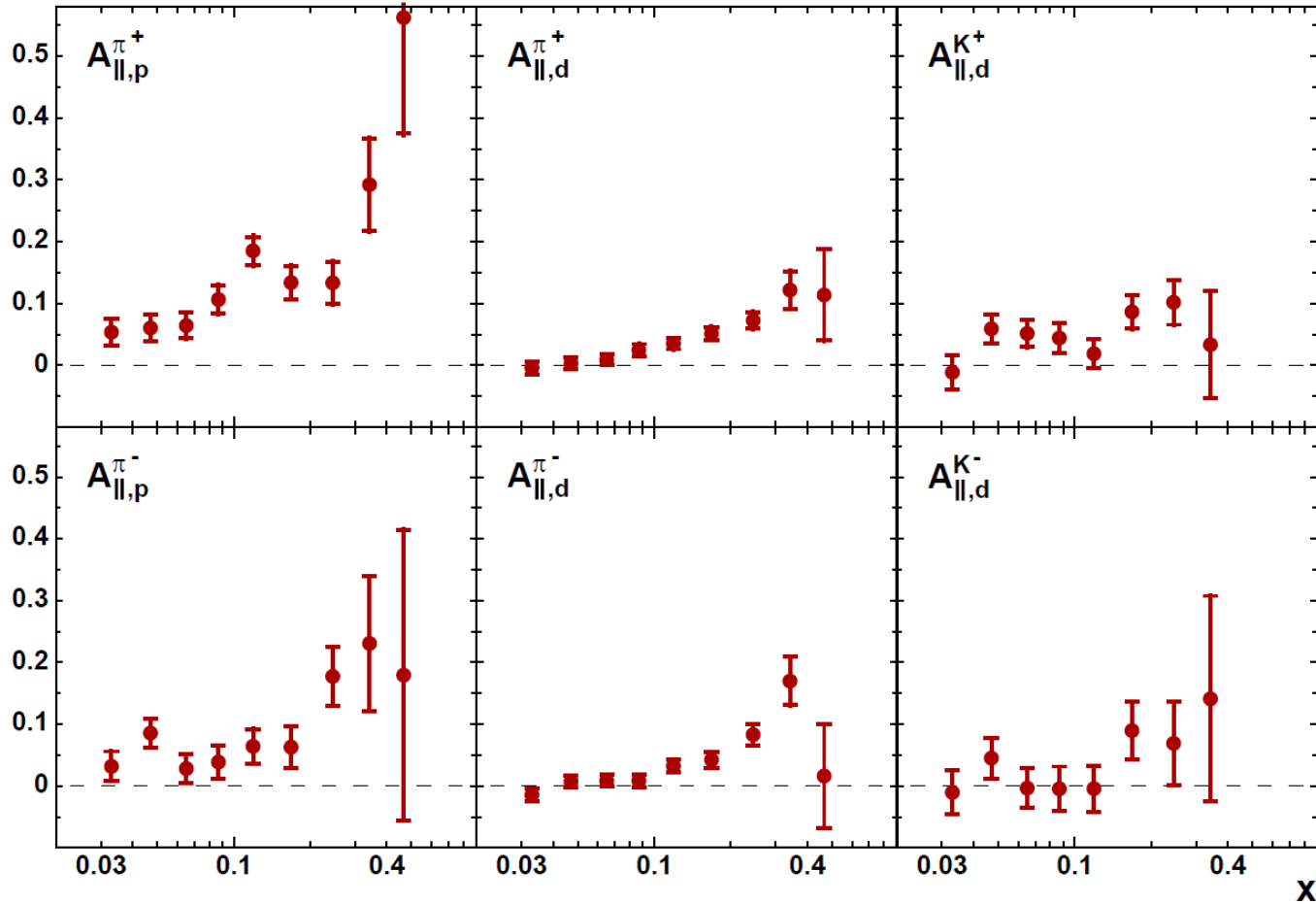
Moments in measured range

Δu	$0.601 \pm 0.039 \pm 0.049$
$\Delta \bar{u}$	$-0.002 \pm 0.036 \pm 0.023$
Δd	$-0.226 \pm 0.039 \pm 0.050$
$\Delta \bar{d}$	$-0.054 \pm 0.033 \pm 0.011$
Δs	$0.028 \pm 0.033 \pm 0.009$
$\Delta u + \Delta \bar{u}$	$0.599 \pm 0.022 \pm 0.065$
$\Delta d + \Delta \bar{d}$	$-0.280 \pm 0.026 \pm 0.057$
Δu_v	$0.603 \pm 0.071 \pm 0.040$
Δd_v	$-0.172 \pm 0.068 \pm 0.045$
$\Delta \bar{u} - \Delta \bar{d}$	$0.048 \pm 0.057 \pm 0.028$
$\Delta \Sigma$	$0.347 \pm 0.024 \pm 0.066$
Δq_3	$0.880 \pm 0.045 \pm 0.107$
Δq_8	$0.262 \pm 0.078 \pm 0.045$
$\Delta^{(2)}u$	$0.142 \pm 0.009 \pm 0.011$
$\Delta^{(2)}\bar{u}$	$-0.001 \pm 0.005 \pm 0.002$
$\Delta^{(2)}d$	$-0.049 \pm 0.010 \pm 0.013$
$\Delta^{(2)}\bar{d}$	$-0.003 \pm 0.004 \pm 0.001$
$\Delta^{(2)}s$	$0.001 \pm 0.003 \pm 0.001$
$\Delta^{(2)}u_v$	$0.144 \pm 0.013 \pm 0.011$
$\Delta^{(2)}d_v$	$-0.047 \pm 0.012 \pm 0.012$

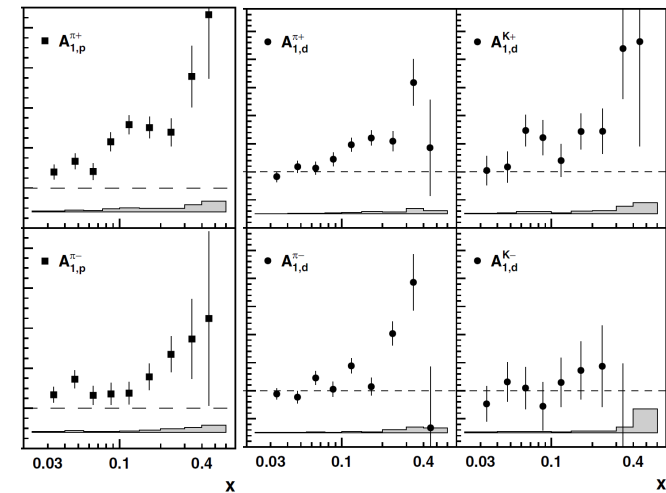
- Lower hadron momentum cut for the deuteron target
 - 4 GeV/c \Rightarrow 2 GeV/c, (\otimes 4 GeV/c for proton w/o RICH)
- Make A_{\parallel}^h available (w/ A_1^h)
 - R in SIDIS still to be measured
- D-state correction for deuteron on asymmetry level
- Better azimuthal asymmetry correction
 - Correction factor extracted using MC, in which data-driven model for azimuthal modulations (PRD87(2013)012010) are implemented
- Multi-dimensional $(x, z, P_{h\perp})$ dependences
- Charge difference asymmetry extraction
- Extract twist-3 cosine modulations

$$D = \frac{1 - (1 - y)\epsilon}{1 + \epsilon R}$$

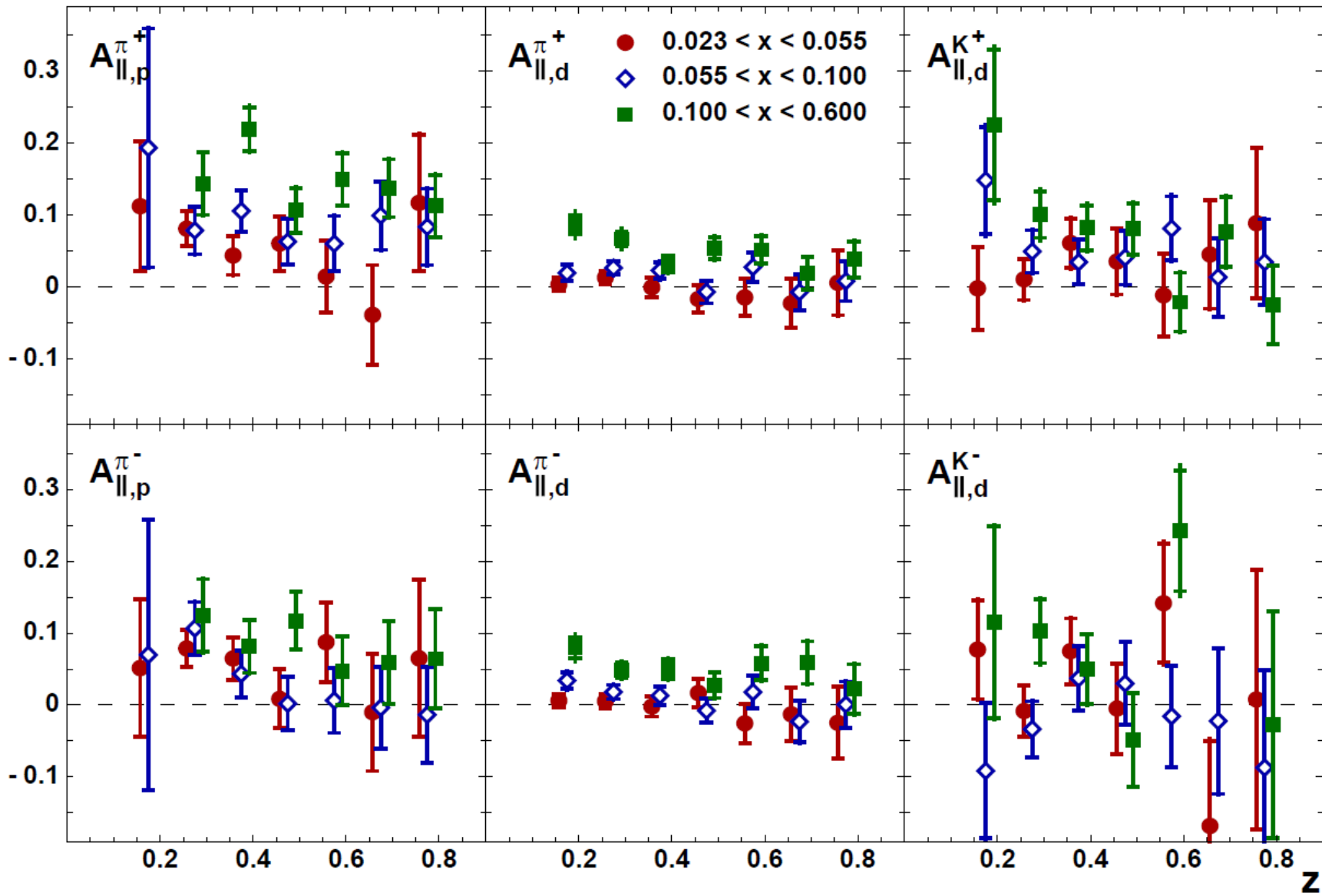
- Longitudinal double spin asymmetry
 - x -dependence
 - z -dependence
 - $P_{h\perp}$ dependence
 - $(x, z, P_{h\perp})$ dependence
 - Charge difference asymmetry
 - Azimuthal modulation: $\cos\phi$

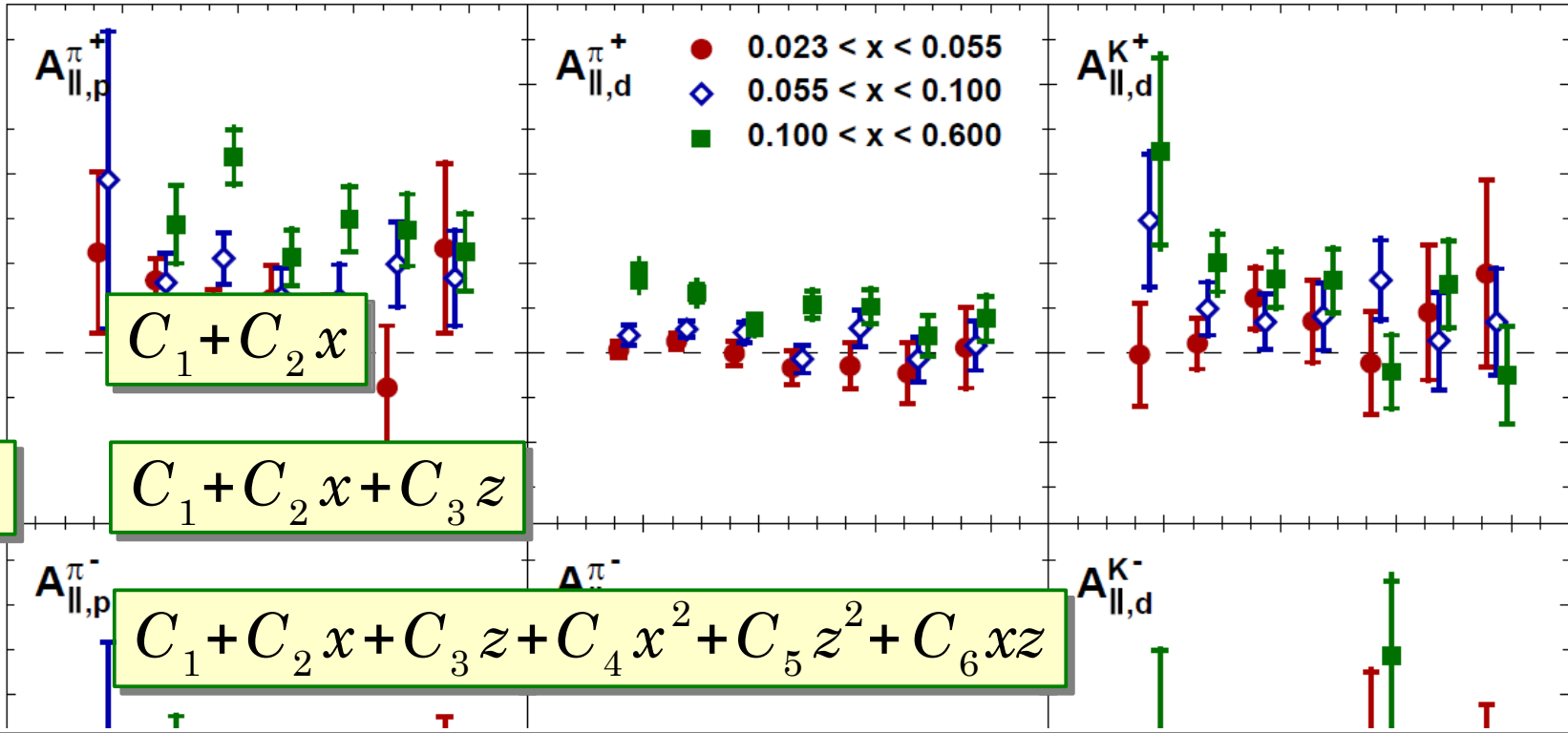


Consistent with
PRD71(2005)012003



No strong z dependence





Fit with

$$C_1 + C_2 x$$

$$C_1 + C_2 x + C_3 z$$

$$C_1 + C_2 x + C_3 z + C_4 x^2 + C_5 z^2 + C_6 xz$$

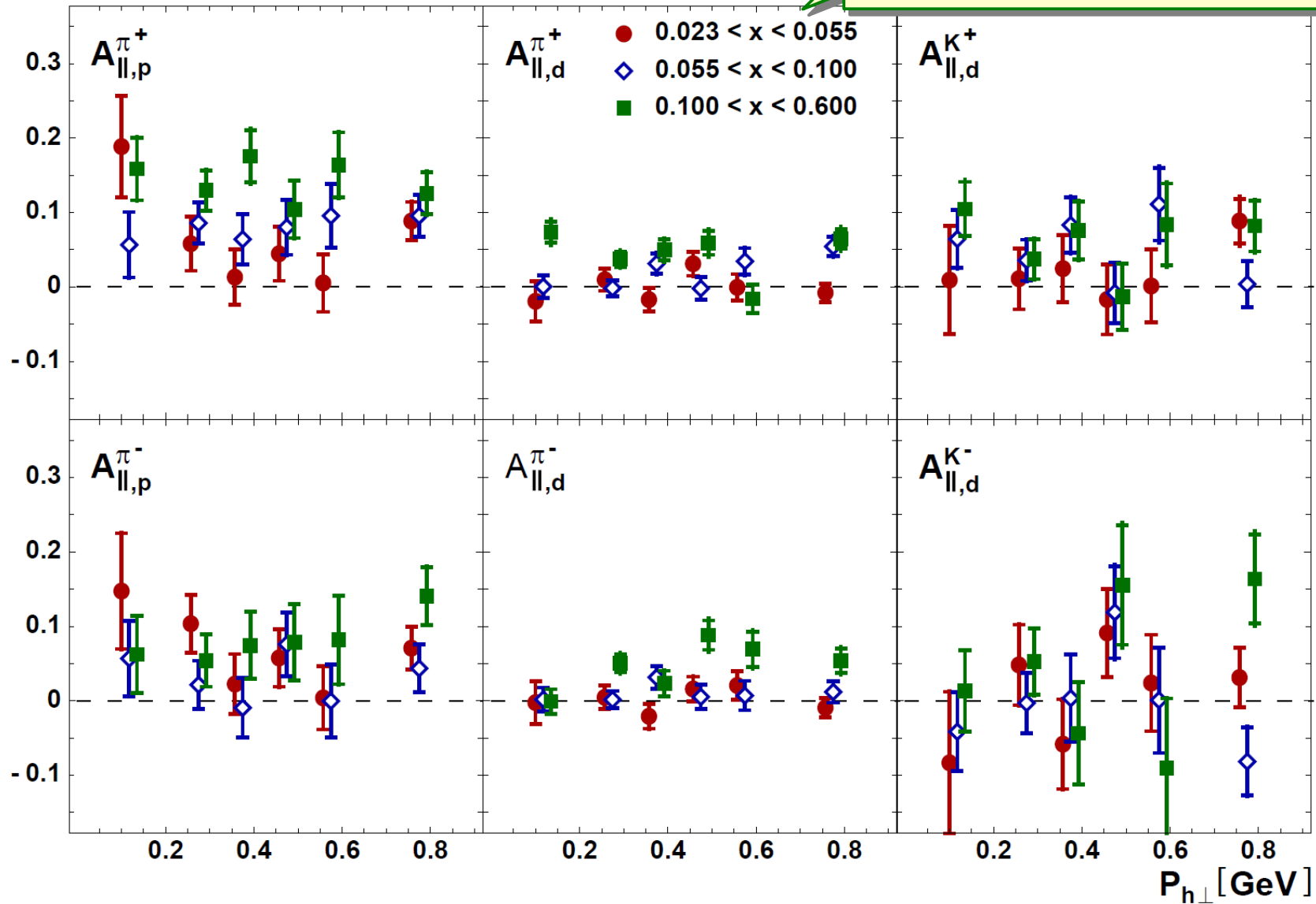
	$A_{1,p}^{\pi^+}$	$A_{1,p}^{\pi^-}$	$A_{1,d}^{\pi^+}$	$A_{1,d}^{\pi^-}$	$A_{1,d}^{K^+}$	$A_{1,d}^{K^-}$
χ^2 (NDF=16) $C_1^h + C_2^h x$	12.6	10.0	13.4	9.1	10.7	26.0
χ^2 (NDF=15) $C_1^h + C_2^h x + C_3^h z$	12.2	6.3	7.2	7.2	10.1	24.8
χ^2 (NDF=12) $C_1^h + C_2^h x + C_3^h z + C_4^h x^2 + C_5^h z^2 + C_6^h xz$	10.3	4.5	5.5	4.8	5.8	16.1

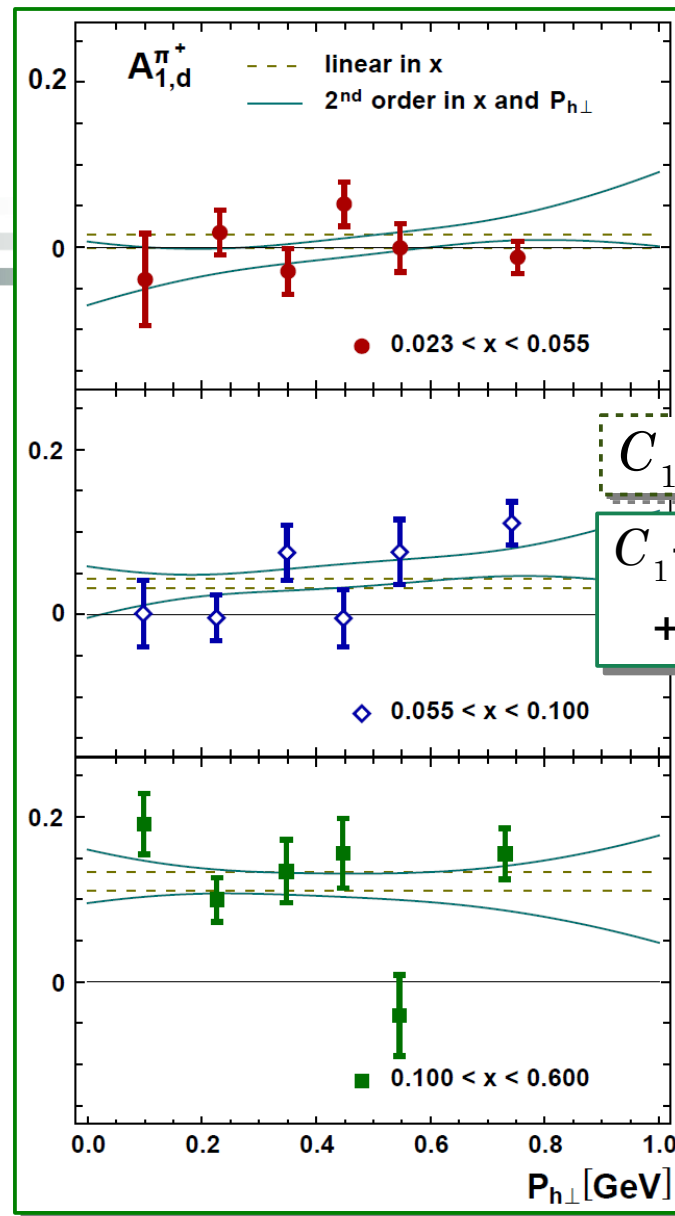
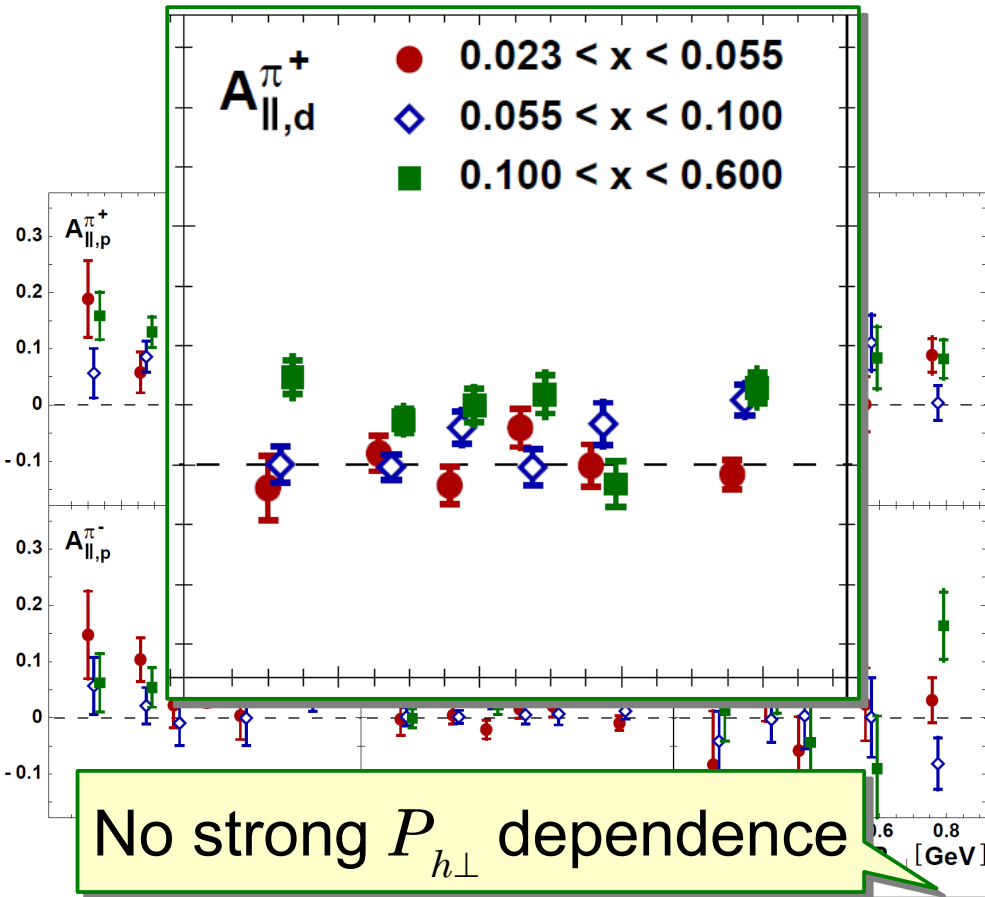
0.2 0.4 0.6 0.8 0.2 0.4 0.6 0.8 0.2 0.4 0.6 0.8

No strong z dependence

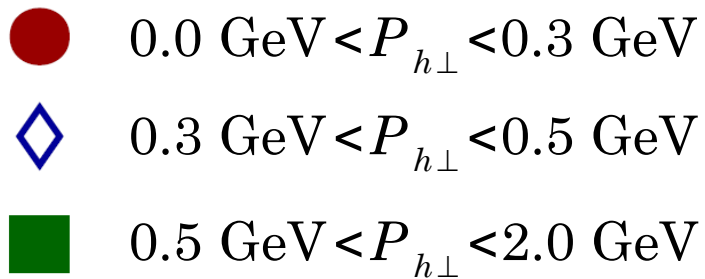
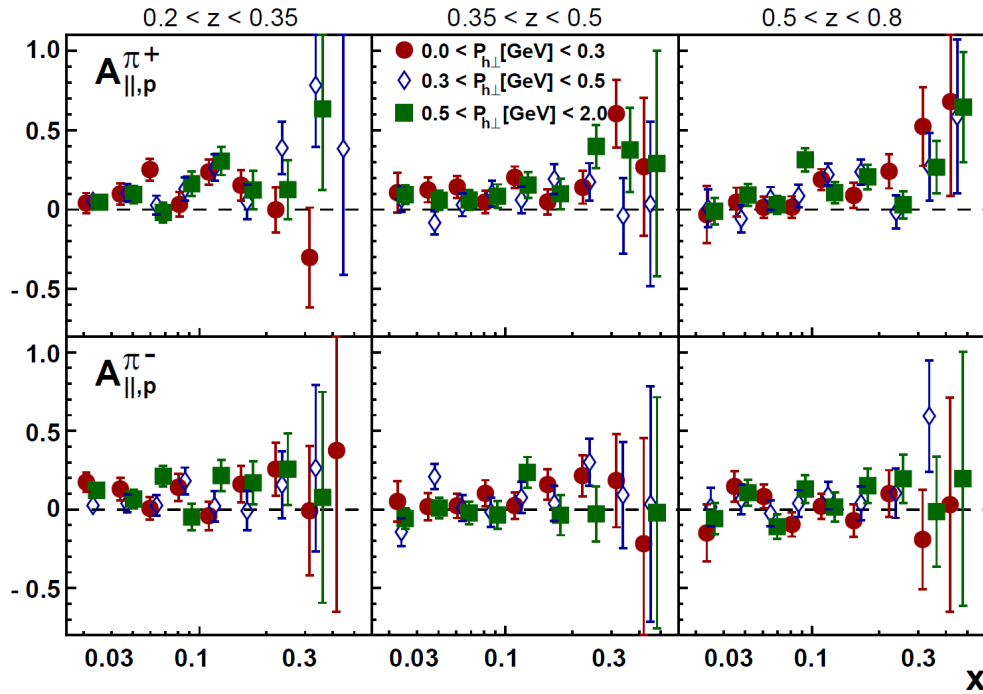


No strong $P_{h\perp}$ dependence

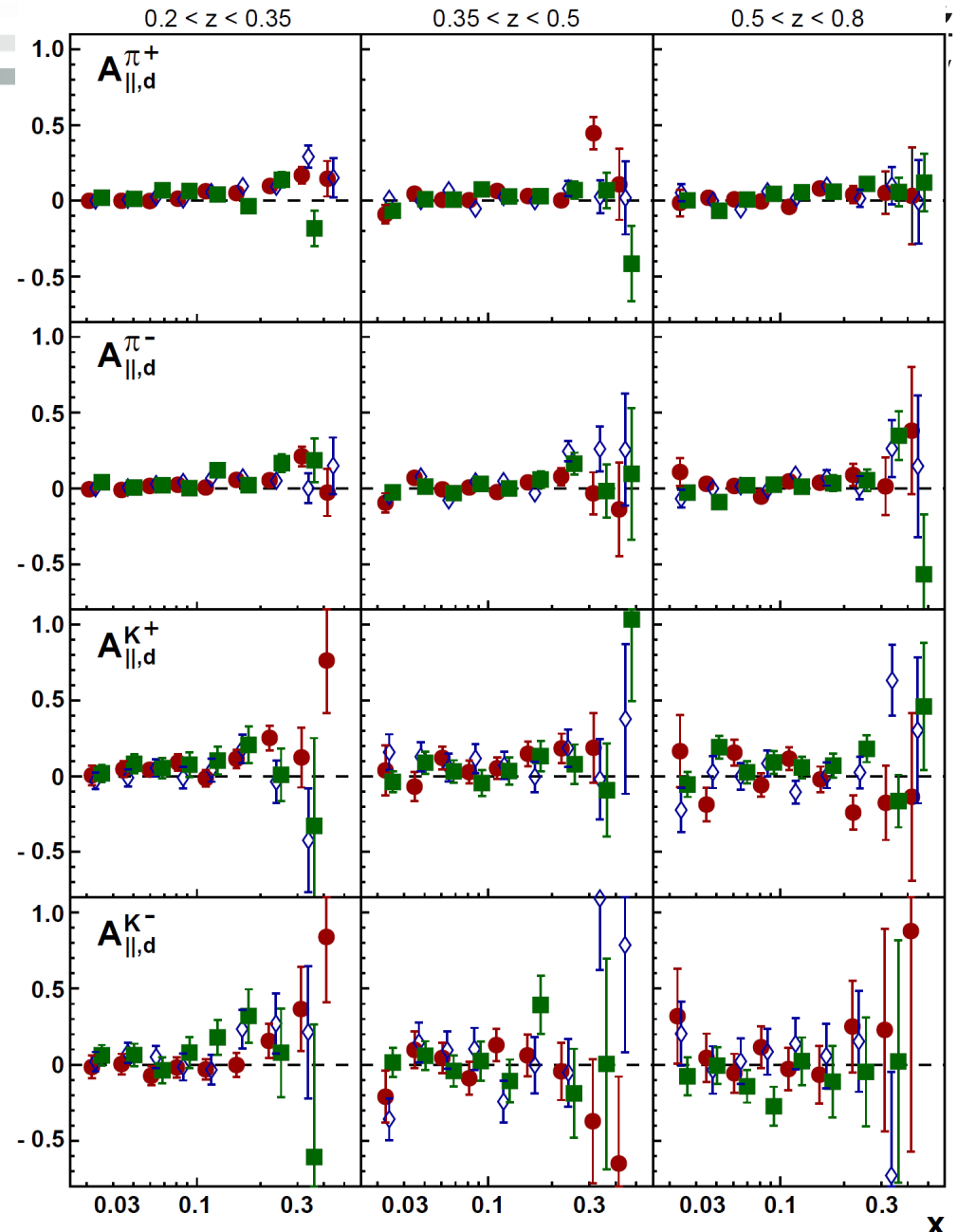


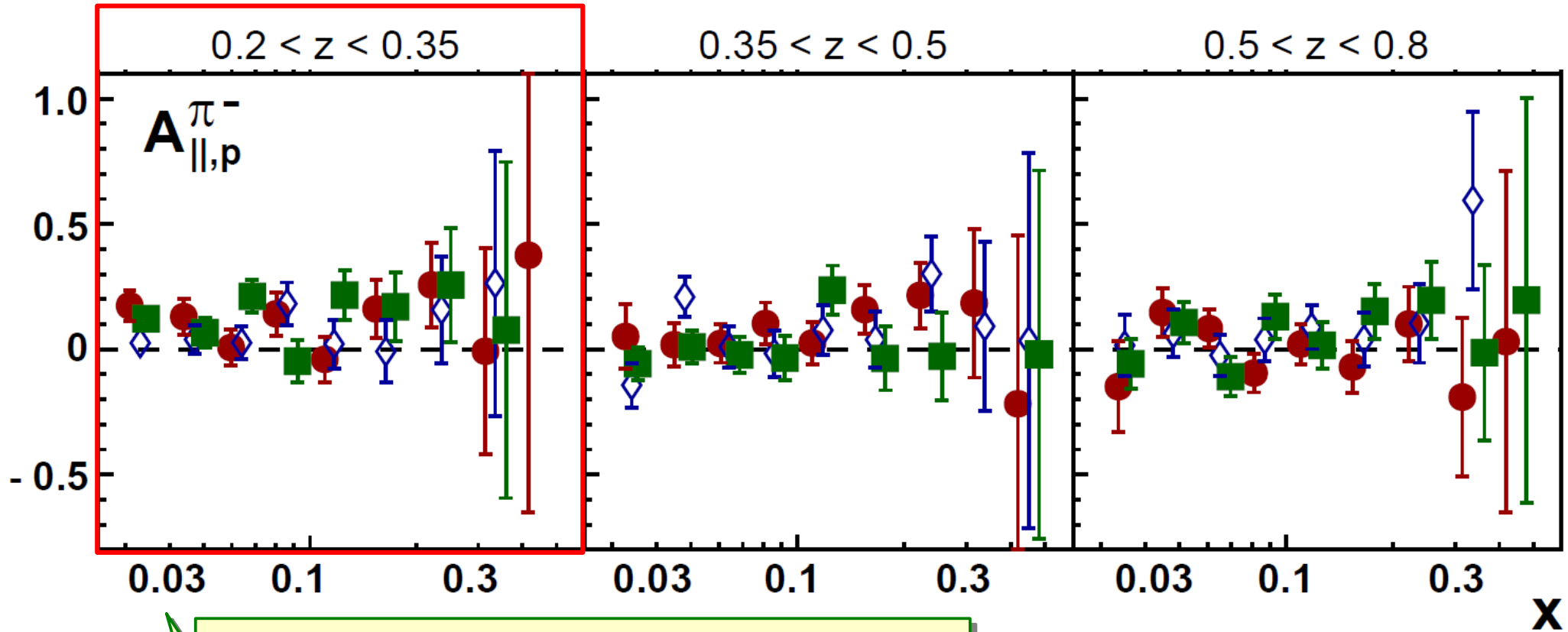


	$A_{1,p}^{\pi^+}$	$A_{1,p}^{\pi^-}$	$A_{1,d}^{\pi^+}$	$A_{1,d}^{\pi^-}$	$A_{1,d}^{K^+}$	$A_{1,d}^{K^-}$
$\chi^2 (NDF=16)$ $C_1^h + C_2^h x$	12.7	14.0	33.7	22.9	16.0	24.4
$\chi^2 (NDF=15)$ $C_1^h + C_2^h x + C_3^h P_{h\perp}$	12.7	13.9	31.9	20.6	16.0	23.6
$\chi^2 (NDF=12)$ $C_1^h + C_2^h x + C_3^h P_{h\perp} + C_4^h x^2 + C_5^h P_{h\perp}^2 + C_6^h x P_{h\perp}$	8.5	5.1	29.7	12.0	12.2	18.7



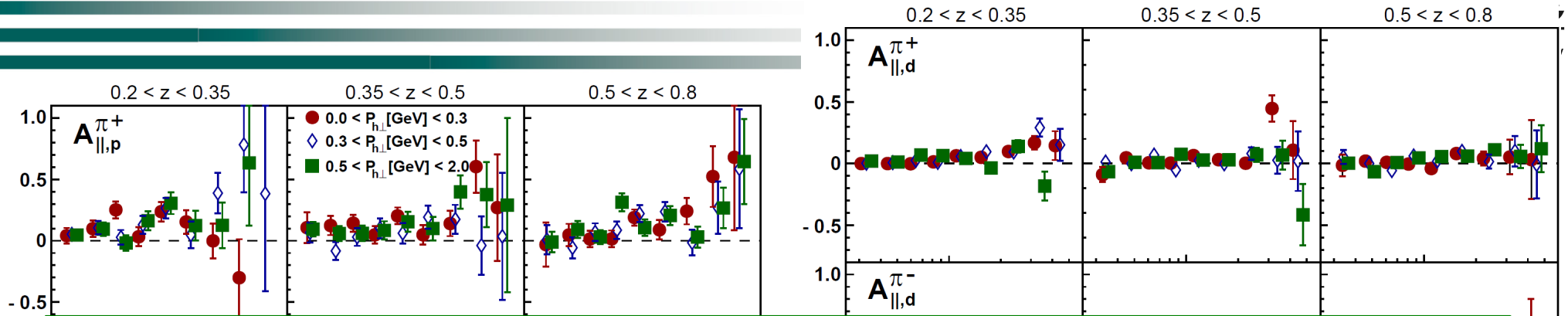
the most complete, unintegrated,
longitudinally polarized double-spin dataset





Possible contribution from $\Delta u > 0$ through dis-favored fragmentation process.

- $0.0 \text{ GeV} < P_{h\perp} < 0.3 \text{ GeV}$
- ◇ $0.3 \text{ GeV} < P_{h\perp} < 0.5 \text{ GeV}$
- $0.5 \text{ GeV} < P_{h\perp} < 2.0 \text{ GeV}$



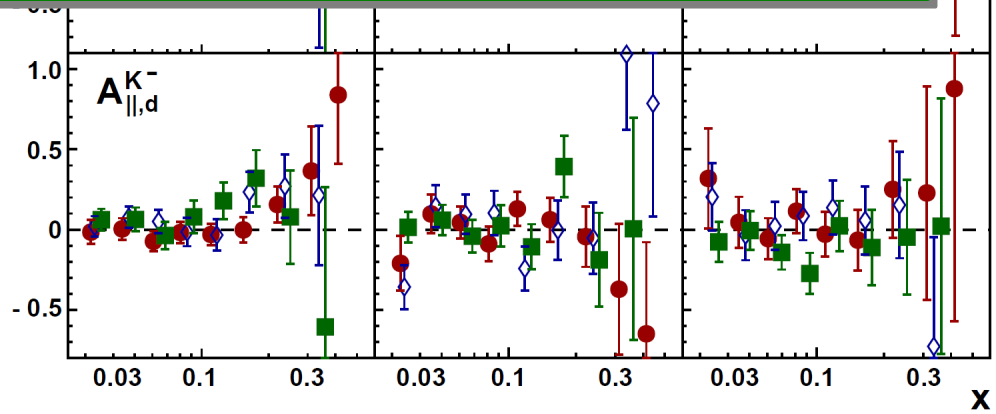
Statistical covariant matrix, other asymmetries are also available

mail-to: management@hermes.desy.de

<http://www-hermes.desy.de/notes/pub/publications.html>

- ◇ $0.3 \text{ GeV} < P_{h\perp} < 0.5 \text{ GeV}$
- $0.5 \text{ GeV} < P_{h\perp} < 2.0 \text{ GeV}$

the most complete, unintegrated,
longitudinally polarized double-spin dataset



$$A_1^{h^+ - h^-}(x) \equiv \frac{\left(\sigma_{1/2}^{h^+} - \sigma_{1/2}^{h^-}\right) - \left(\sigma_{3/2}^{h^+} - \sigma_{3/2}^{h^-}\right)}{\left(\sigma_{1/2}^{h^+} - \sigma_{1/2}^{h^-}\right) + \left(\sigma_{3/2}^{h^+} - \sigma_{3/2}^{h^-}\right)}$$

At **Leading Order** and **Leading Twist**

proton

FFs disappeared

deuteron

$$A_{1,p}^{h^+ - h^-} \stackrel{\text{LO LT}}{=} \frac{4g_1^{u_v} - g_1^{d_v}}{4f_1^{u_v} - f_1^{d_v}}$$

$$A_{1,d}^{h^+ - h^-} \stackrel{\text{LO LT}}{=} \frac{g_1^{u_v} + g_1^{d_v}}{f_1^{u_v} + f_1^{d_v}}$$

Independent on h

Charge conjugation symmetry

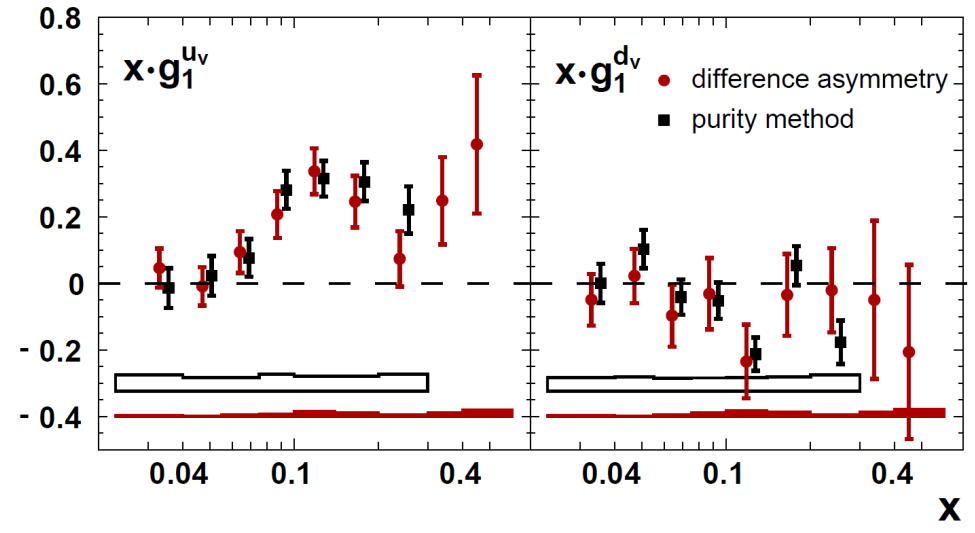
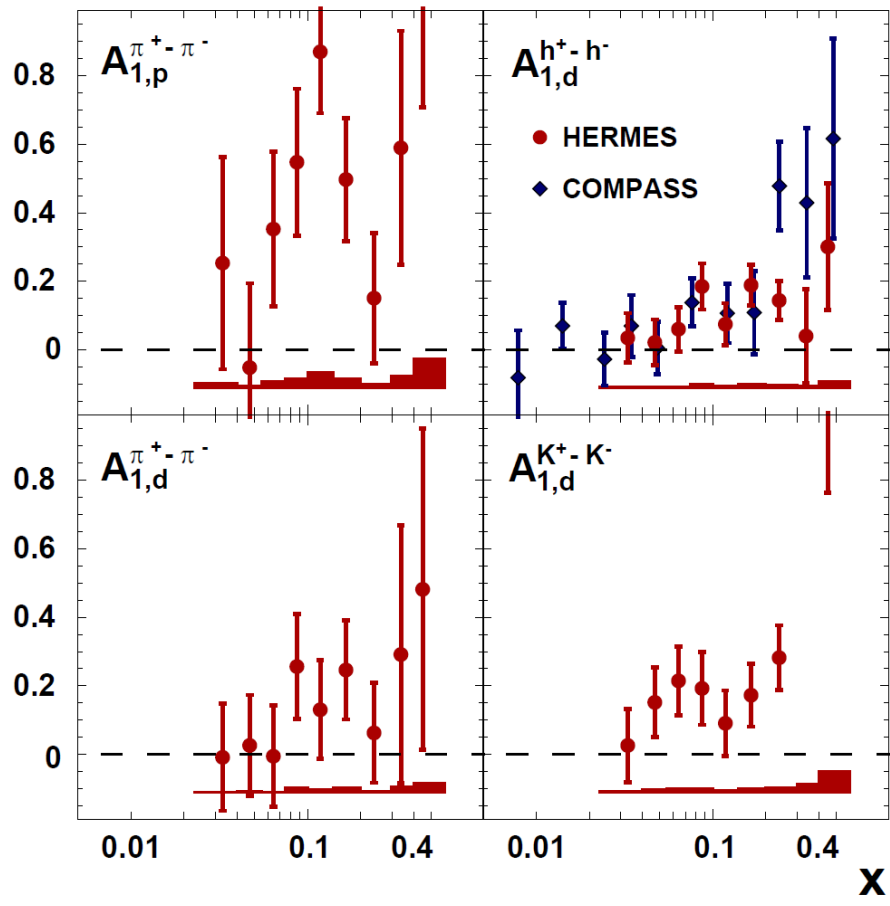
$$q \rightarrow h^+ = \bar{q} \rightarrow h^- \quad q \rightarrow h^+ = \bar{q} \rightarrow h^-$$

$$u \rightarrow \pi^\pm = d \rightarrow \pi^\mp$$

isospin symmetry

$$A_1^{h^+ - h^-}(x) \equiv \frac{(\sigma_{1/2}^{h^+} - \sigma_{1/2}^{h^-}) - (\sigma_{3/2}^{h^+} - \sigma_{3/2}^{h^-})}{(\sigma_{1/2}^{h^+} - \sigma_{1/2}^{h^-}) + (\sigma_{3/2}^{h^+} - \sigma_{3/2}^{h^-})}$$

$$\left\{ \begin{aligned} A_{1,p}^{h^+ - h^-} &\stackrel{\text{LO}}{=} \frac{4g_1^{u_v} - g_1^{d_v}}{4f_1^{u_v} - f_1^{d_v}} \\ A_{1,d}^{h^+ - h^-} &\stackrel{\text{LO}}{=} \frac{g_1^{u_v} + g_1^{d_v}}{f_1^{u_v} + f_1^{d_v}} \end{aligned} \right.$$

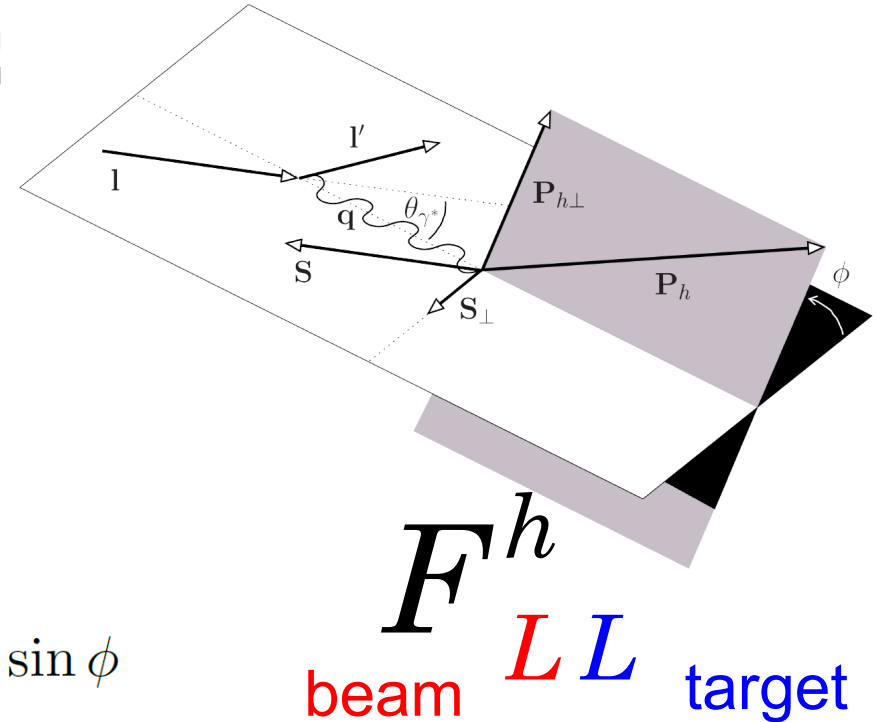




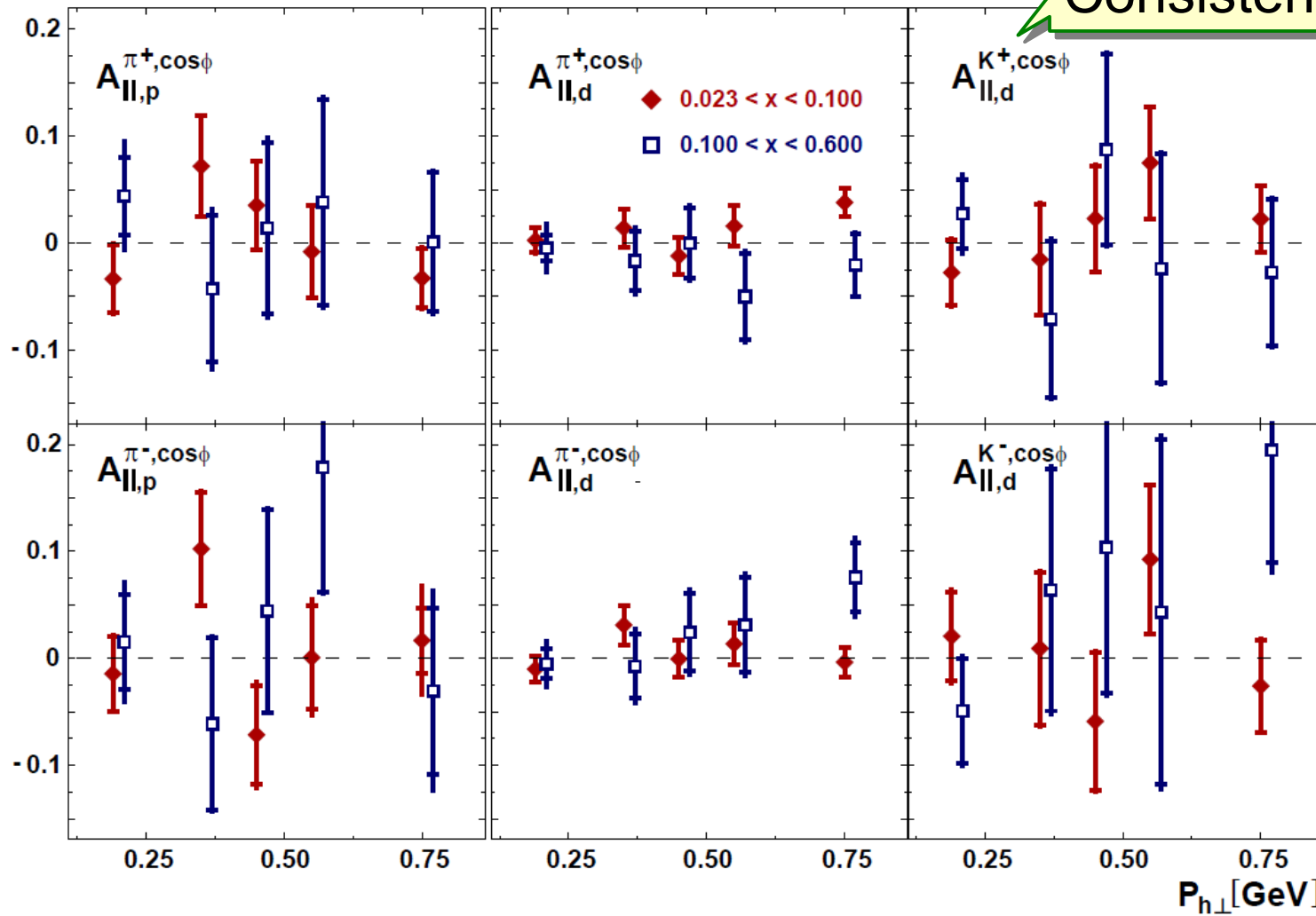
※ transverse nucleon polarization neglected

$$\frac{d\sigma^h}{dx dy dz dP_{h\perp}^2 d\phi} = \frac{2\pi\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left(1 + \frac{\gamma^2}{2x} \right)$$

$$\left\{ \begin{aligned} & F_{UU,T}^h + \epsilon F_{UU,L}^h + \lambda\Lambda\sqrt{1-\epsilon^2} F_{LL}^h \\ & + \sqrt{2\epsilon} \left[\lambda\sqrt{1-\epsilon} F_{LU}^{h,\sin\phi} + \Lambda\sqrt{1+\epsilon} F_{UL}^{h,\sin\phi} \right] \sin\phi \\ & + \sqrt{2\epsilon} \left[\lambda\Lambda\sqrt{1-\epsilon} F_{LL}^{h,\cos\phi} + \sqrt{1+\epsilon} F_{UU}^{h,\cos\phi} \right] \cos\phi \\ & + \Lambda\epsilon F_{UL}^{h,\sin 2\phi} \sin 2\phi + \epsilon F_{UU}^{h,\cos 2\phi} \cos 2\phi \end{aligned} \right\}$$



These cosine modulations were extracted.



$\cos 2\phi$ modulations were also found to be consistent with zero.

- Extracted longitudinal double-spin asymmetries in SIDIS using the full HERMES data set with updated analysis method
 - the most complete, unintegrated, longitudinally polarized double-spin dataset available
 - provide A_{\parallel} in addition to virtual-photon-nucleon asymmetry A_1
 - no significant dependence on z and $P_{h\perp}$ in A_1 within precision of the measurements
 - hadron-charge difference asymmetries
 - in agreement with COMPASS
 - used for LO, leading-twist extraction of valence helicity PDFs
 - Zero consistent $\cos\phi$ moments of semi-inclusive double-spin asymmetry obtained