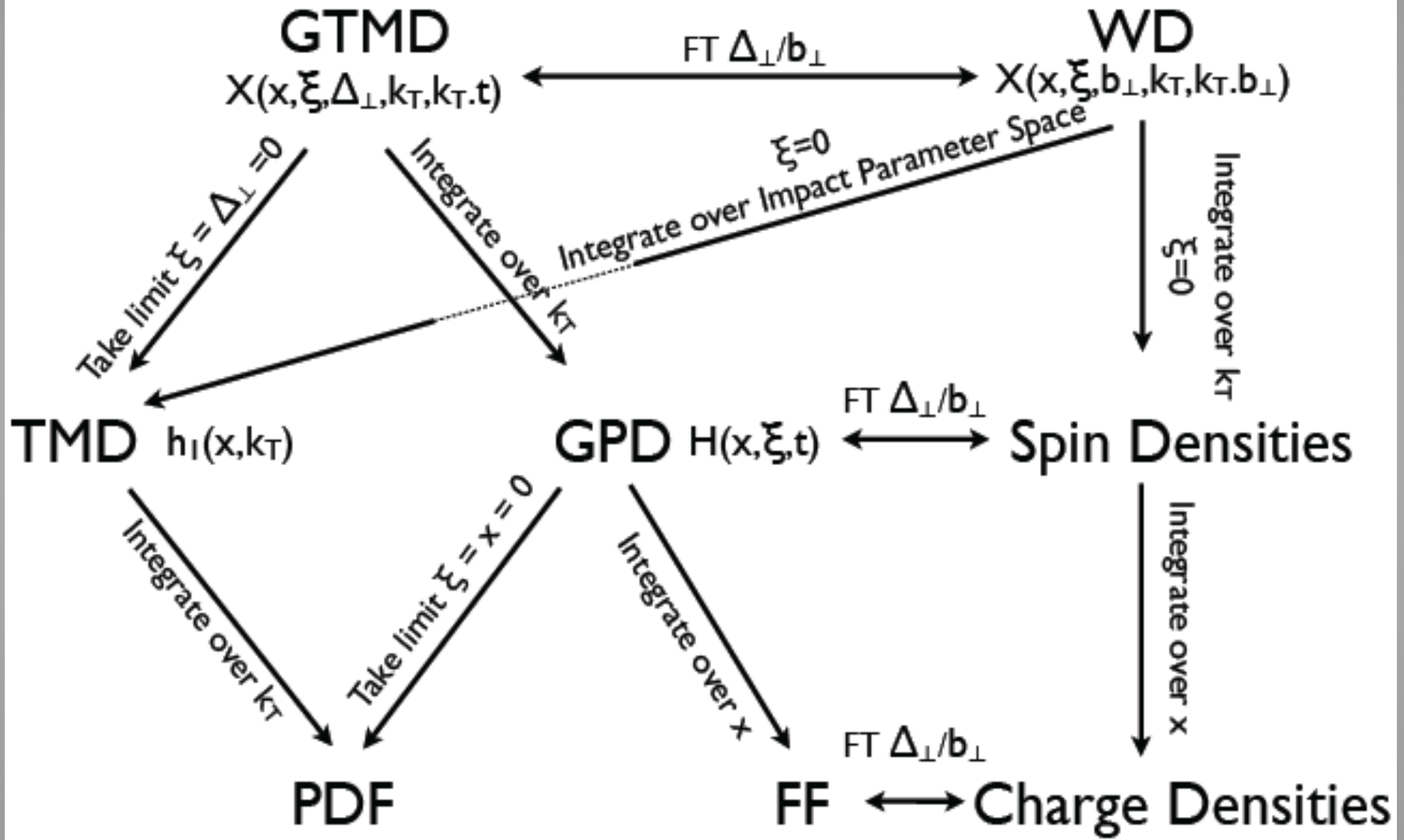


# Spin Structure Results from HERMES

Wolf-Dieter Nowak  
DESY Zeuthen

on behalf of the  hermes collaboration

# Contemporary Hierarchy of Partonic Distributions

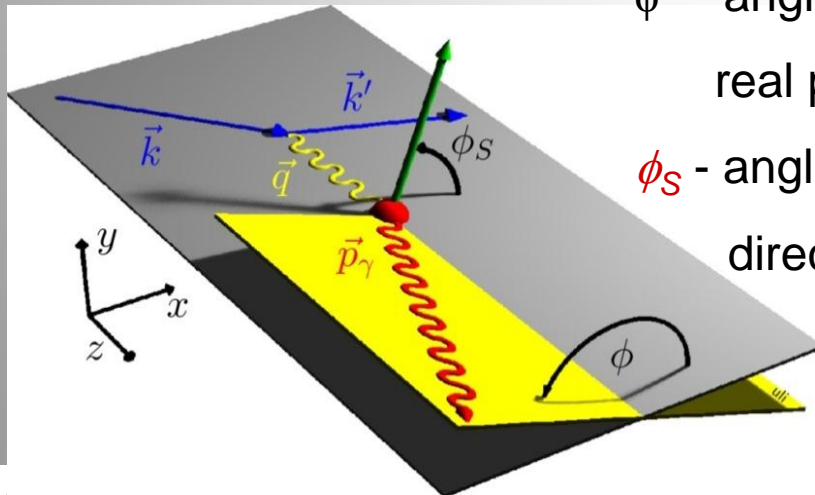


Courtesy M. Murray, Glasgow



# Unique Experimental Conditions at HERMES

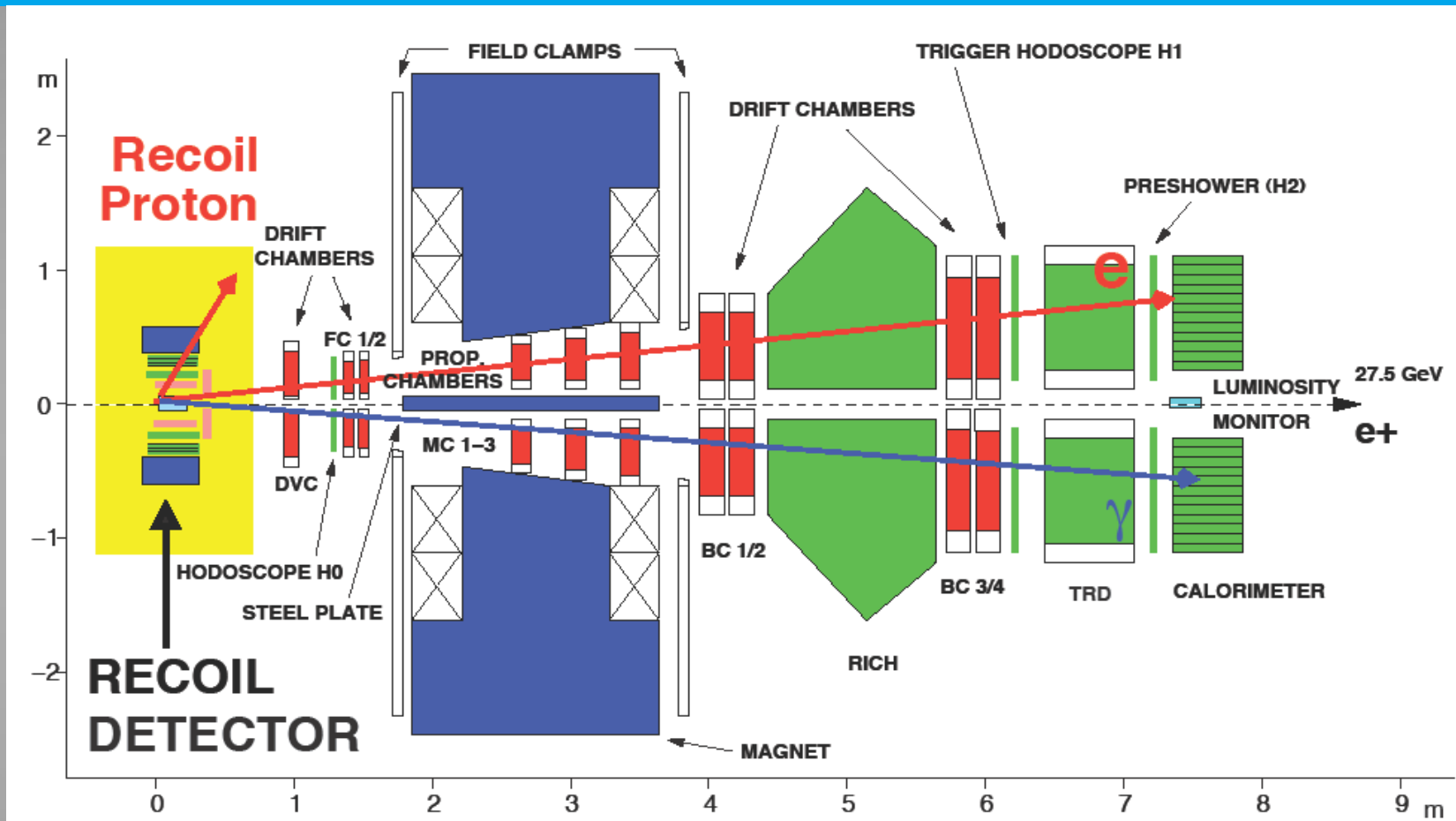
- Running 1996-2007 in 27.5 GeV HERA lepton beam line
- Both beam charges available (positrons, electrons)
- Longitudinal beam polarization (both helicities)
- Various internal gas targets were installed (no dilution)
- Forward spectrometer: tracking and PID (RICH since 1998)
- Recoil detector around target region: 2006-07



$\phi$  - angle between lepton scattering and real photon (hadron) production planes

$\phi_S$  - angle between (transverse) target spin direction and lepton scattering plane

# HERMES Spectrometer & Recoil Detector



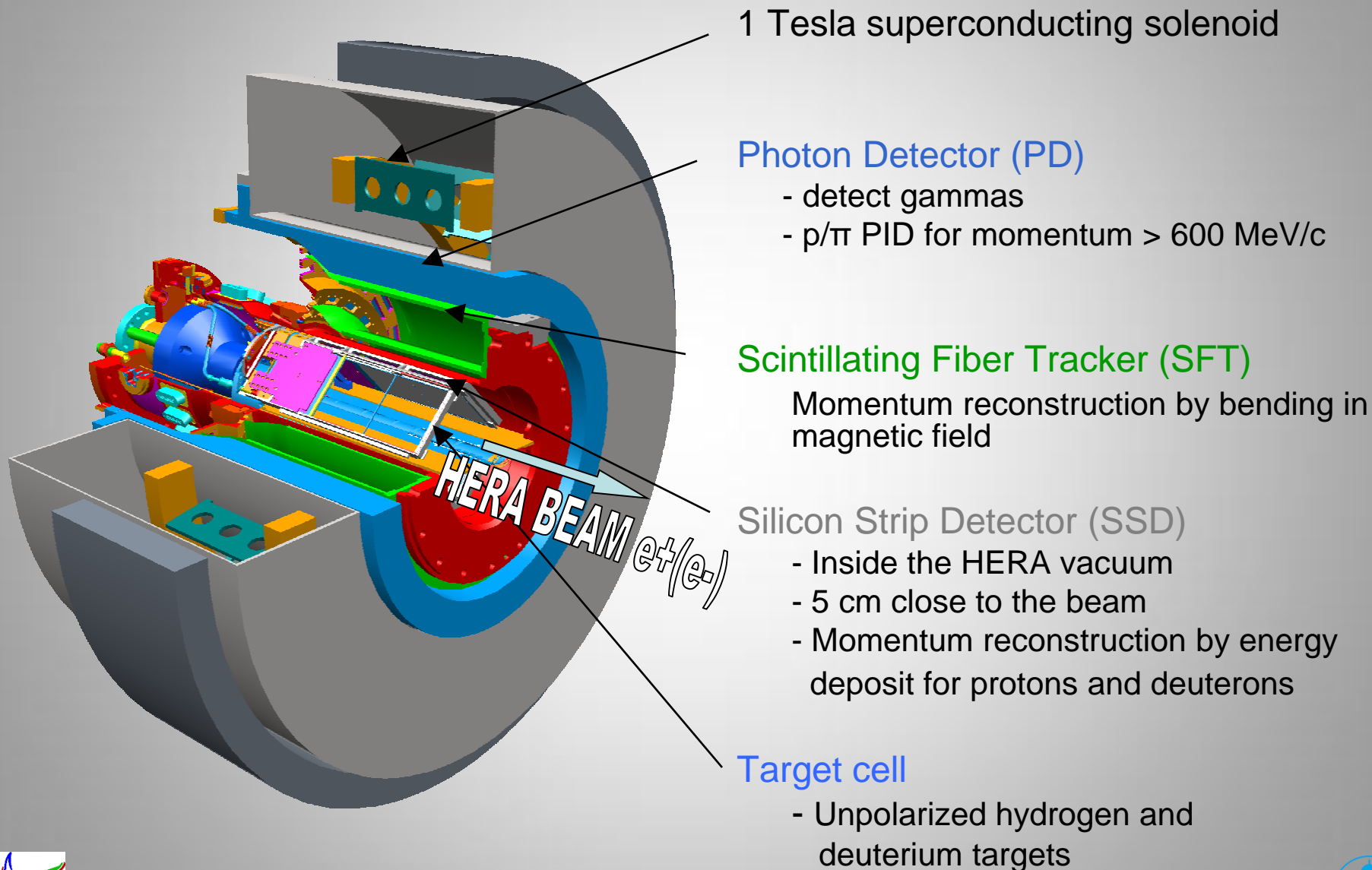
## Internal gas targets:

- Longitudinally polarized  $H, D$
- Transversely polarized  $H$
- Unpolarized  $H, D, {}^4\text{He}, N, Ne, Kr, Xe$

## Forward magnetic spectrometer

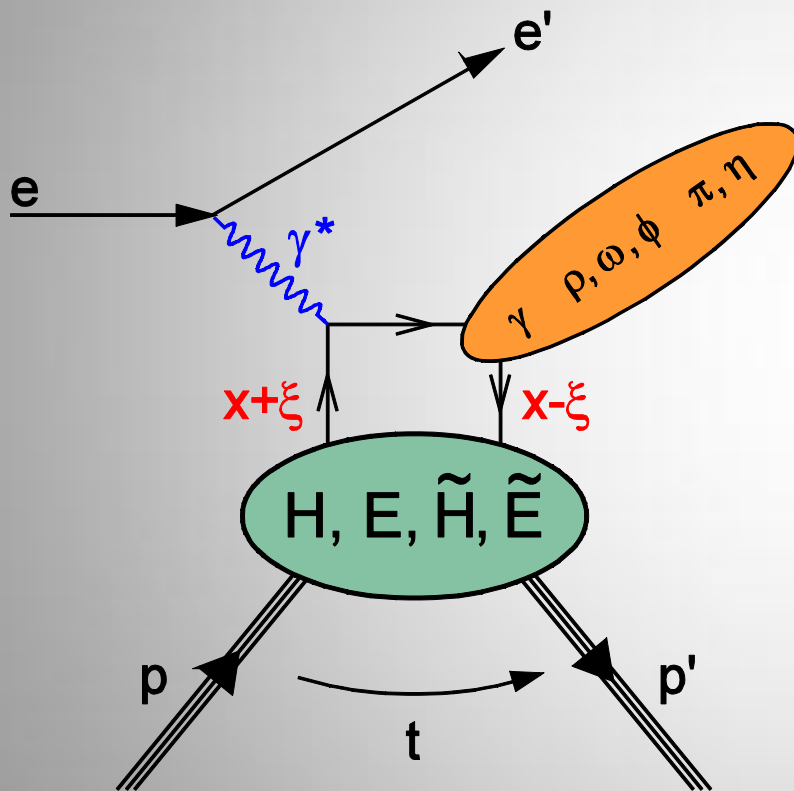
- Momentum resolution  $1-2\%$
- Particle identification:  
 $RICH, TRD, H(od)2, calorimeter$

# HERMES Recoil Detector



# Exclusive Processes: Access to GPDs

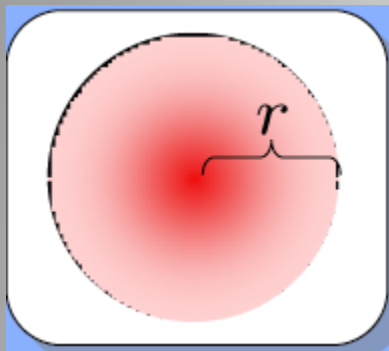
## Generalized Parton Distributions (GPDs):



- For spin-1/2 target 4 chiral-even leading-twist quark GPDs:  $H, E, \tilde{H}, \tilde{E}$
- $H, \tilde{H}$  conserve nucleon helicity,  $E, \tilde{E}$  involve nucleon helicity flip
- Different final states are sensitive to different (combinations of) GPDs:
- DVCS ( $\gamma$ )  $\rightarrow H, E, \tilde{H}, \tilde{E}$
- Vector mesons ( $\rho, \omega, \phi$ )  $\rightarrow H, E$
- Pseudoscalar mesons ( $\pi, \eta$ )  $\rightarrow \tilde{H}, \tilde{E}$

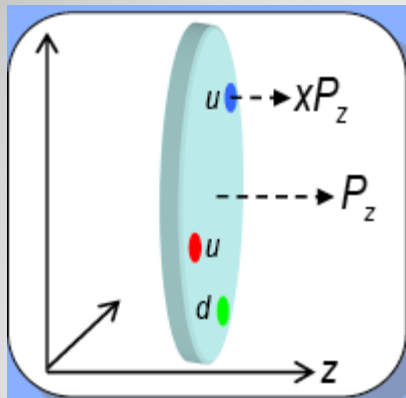
# Interpretation of GPDs

Elastic Form Factors

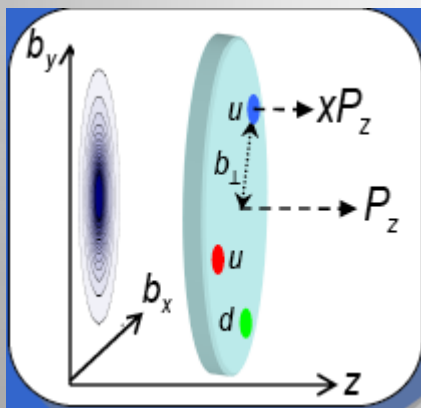


transverse position of partons

Parton Distribution Functions (PDFs)



longitudinal momentum of partons



Correlation between longitudinal momentum and transverse position

- GPDs include Form Factors and Parton Distribution Functions as moments and forward limits, resp.

- GPDs yield a multidimensional description of nucleon structure (longitudinal momentum vs. transverse position)

→ **NUCLEON TOMOGRAPHY**

- GPDs offer access to quark total angular momentum through the Ji relation (in principle also for gluons):

$$J_q = \lim_{t \rightarrow 0} \int_{-1}^1 dx x [H_q(x, \xi, \xi) + E_q(x, \xi, \xi)]$$

[X. Ji, *Phys. Rev. Lett.* 78 (1997) 610]

# Principle of Extraction of Asymmetry Amplitudes

- Distribution in expectation value of measured yield ( $A_{UL}, A_{LL}$  missing):

$$\langle N(e_l, P_l, S_t, \phi, \phi_S) \rangle \propto$$

$$\sigma_{UU}(\phi) \left[ 1 + e_l A_C + P_l A_{LU}^{DVCS} + e_l P_l A'_{LU} + S_t A_{UT}^{DVCS} + e_l S_t A'_{UT} + P_l S_t A_{LT}^{BH+DVCS} + e_l P_l S_t A'_{LT} \right]$$

- Examples of Fourier expansion of measured (X-section) asymmetries:

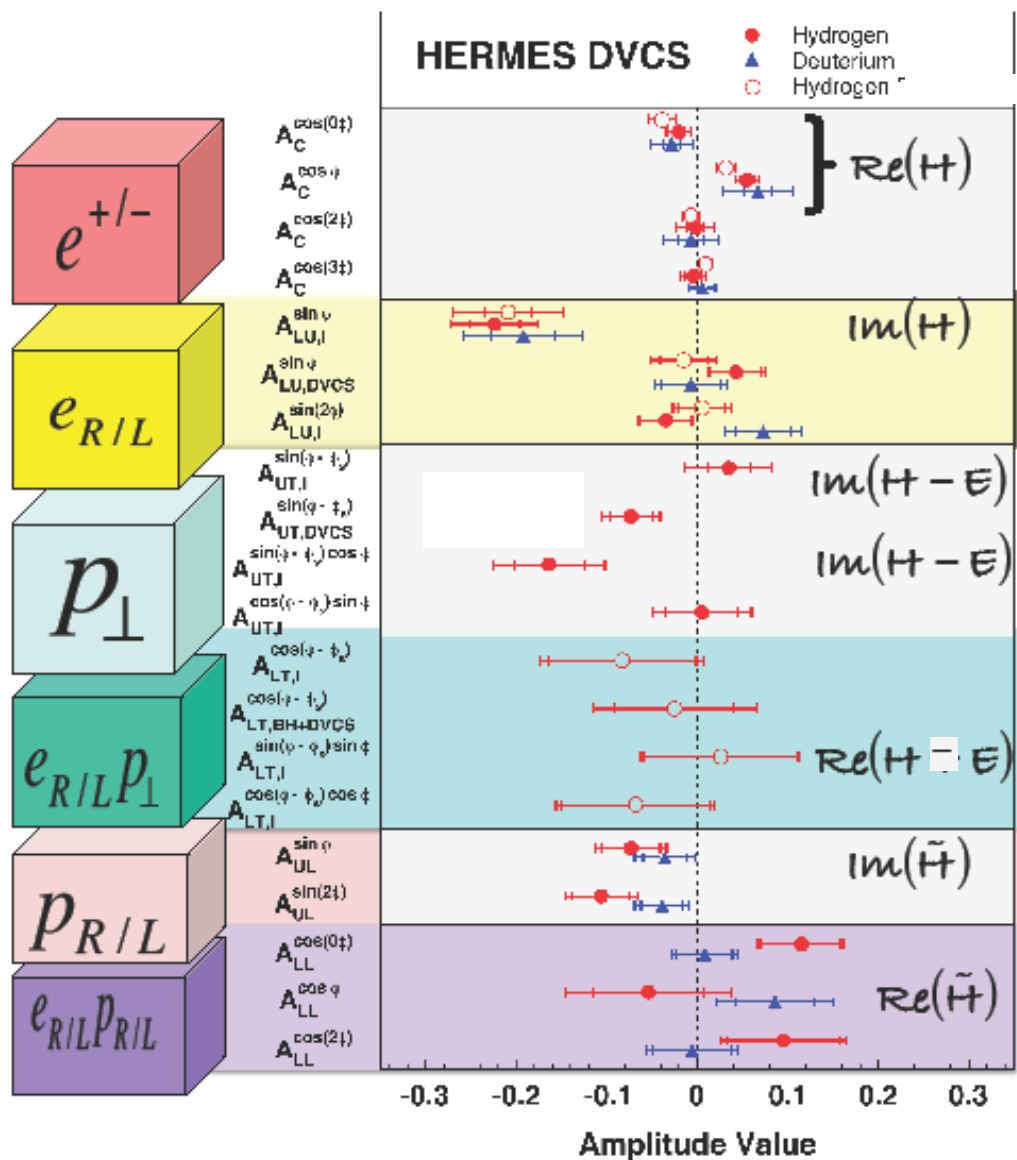
$$A_C \approx \sum_{n=0}^3 A_C^{\cos(n\phi)} \cos(n\phi) \quad \left[ A_{LU} \approx \sum_{n=1}^2 A_{LU}^{\sin(n\phi)} \sin(n\phi) \right]$$

- Simultaneous extraction of asymmetry amplitudes with Maximum Likelihood Method
- Asymmetry amplitudes provide information about Compton Form Factors (CFFs): convolution of GPDs with hard scattering amplitudes

$$\mathcal{F}(\xi, t) = \sum_q \int_{-1}^1 dx C_q^{\mp}(\xi, x) F^q(x, \xi, t)$$



# HERMES DVCS results on proton & deuteron targets



## Access to GPD $H, \tilde{H}, E$

- JHEP 11 (2009) 083
- Nucl. Phys. B829

sensitive to  $J_u$

- JHEP 06 (2008) 066

• Phys. Lett. B 704  
Oct. 5, 2011

- JHEP 06 (2010) 019
- Nucl. Phys. B 842

15

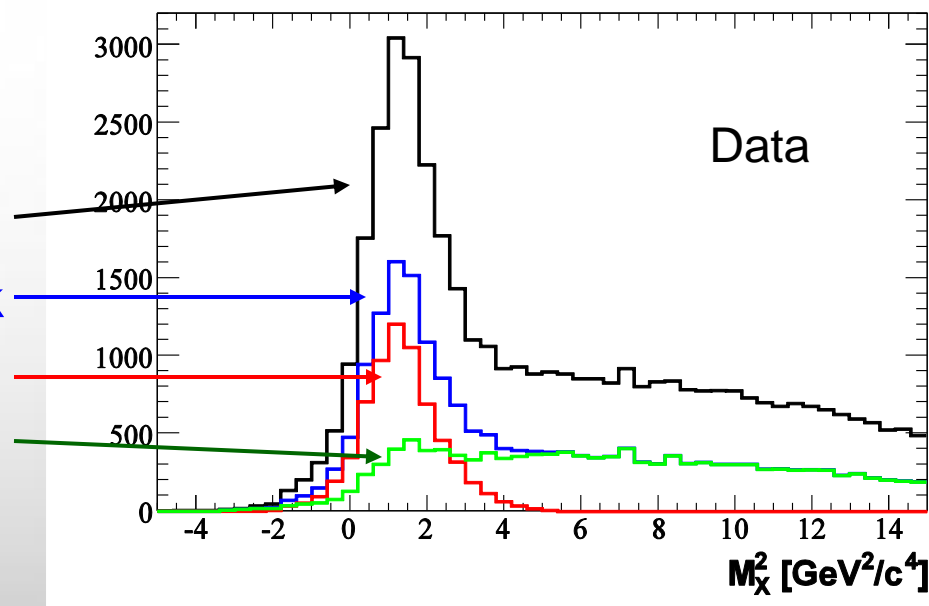


# Exclusivity by Recoil Detector & kinematic event fitting

see poster I. Brodsky

- Kinematic event fitting technique used for DVCS events:
  - All 3 particles in final state detected → 4 constraints from energy-momentum conservation
  - Selection of 'pure' BH/DVCS ( $ep \rightarrow epy$ ) with high efficiency ( $\sim 84\%$ )
  - Allows to suppress background from associated and semi-inclusive processes to a negligible level ( $\sim 0.1\%$ )
  - Makes even PID (Particle Identification) unnecessary

- DVCS missing-mass distribution:
  - No requirement for Recoil
  - Positively charged Recoil track
  - Kinematic fit probability  $> 1\%$
  - Kinematic fit probability  $< 1\%$

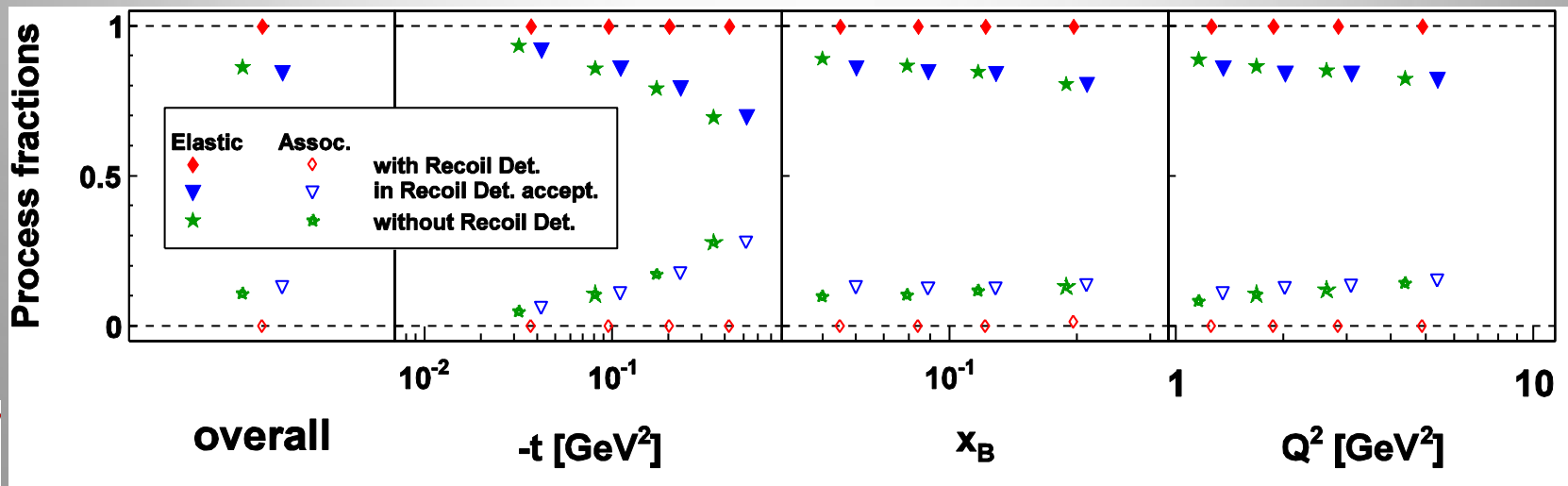
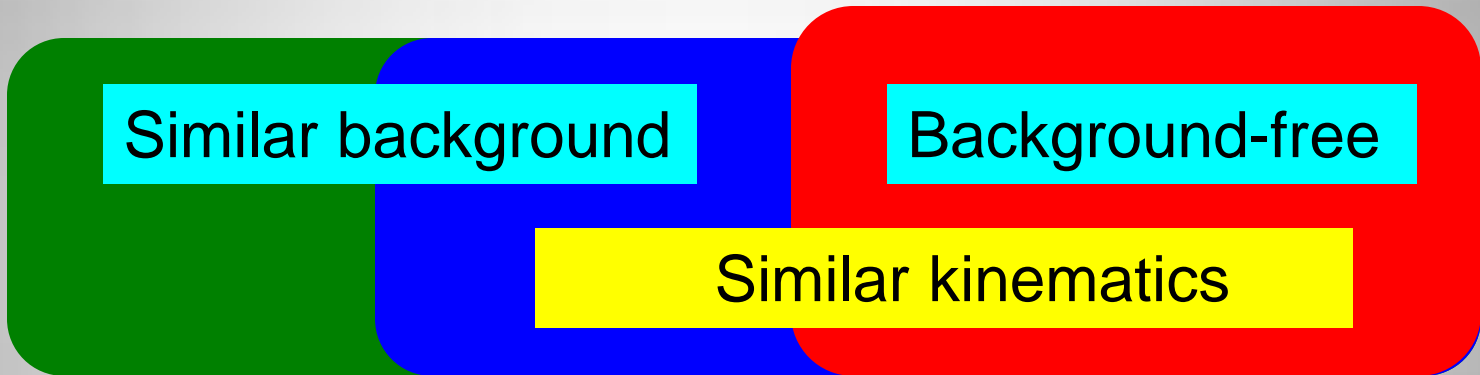


# 'Traditional' and 'Recoil-detector' DVCS samples ( $\mathcal{A}_{LU}$ )

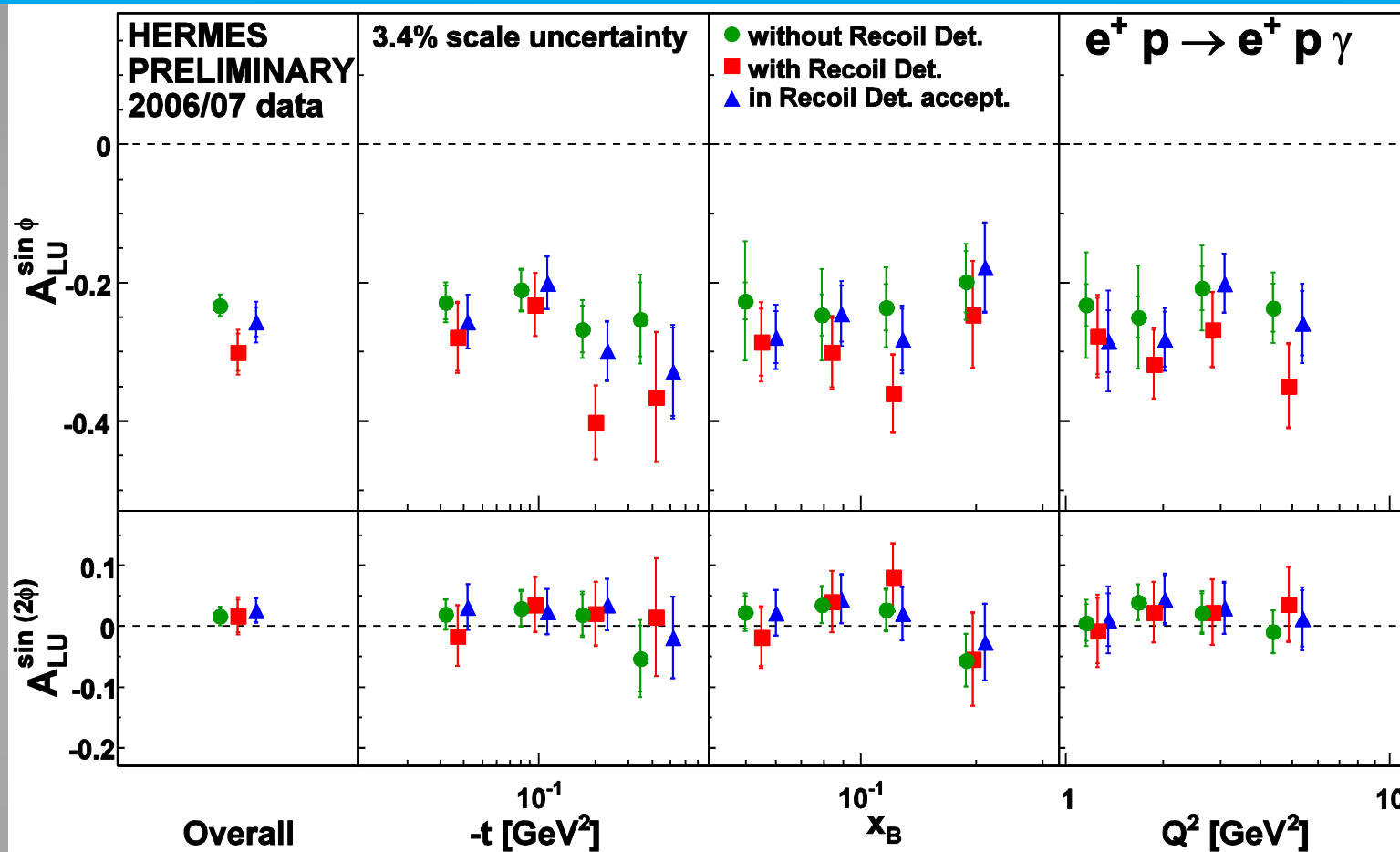
Without Recoil Detector ('traditional' or 'unresolved')

In Recoil Detector acceptance ('unresolved reference')

With Recoil Detector ('pure' sample)



# Comparison of traditional & recoil-detector ( $A_{LU}$ ) samples



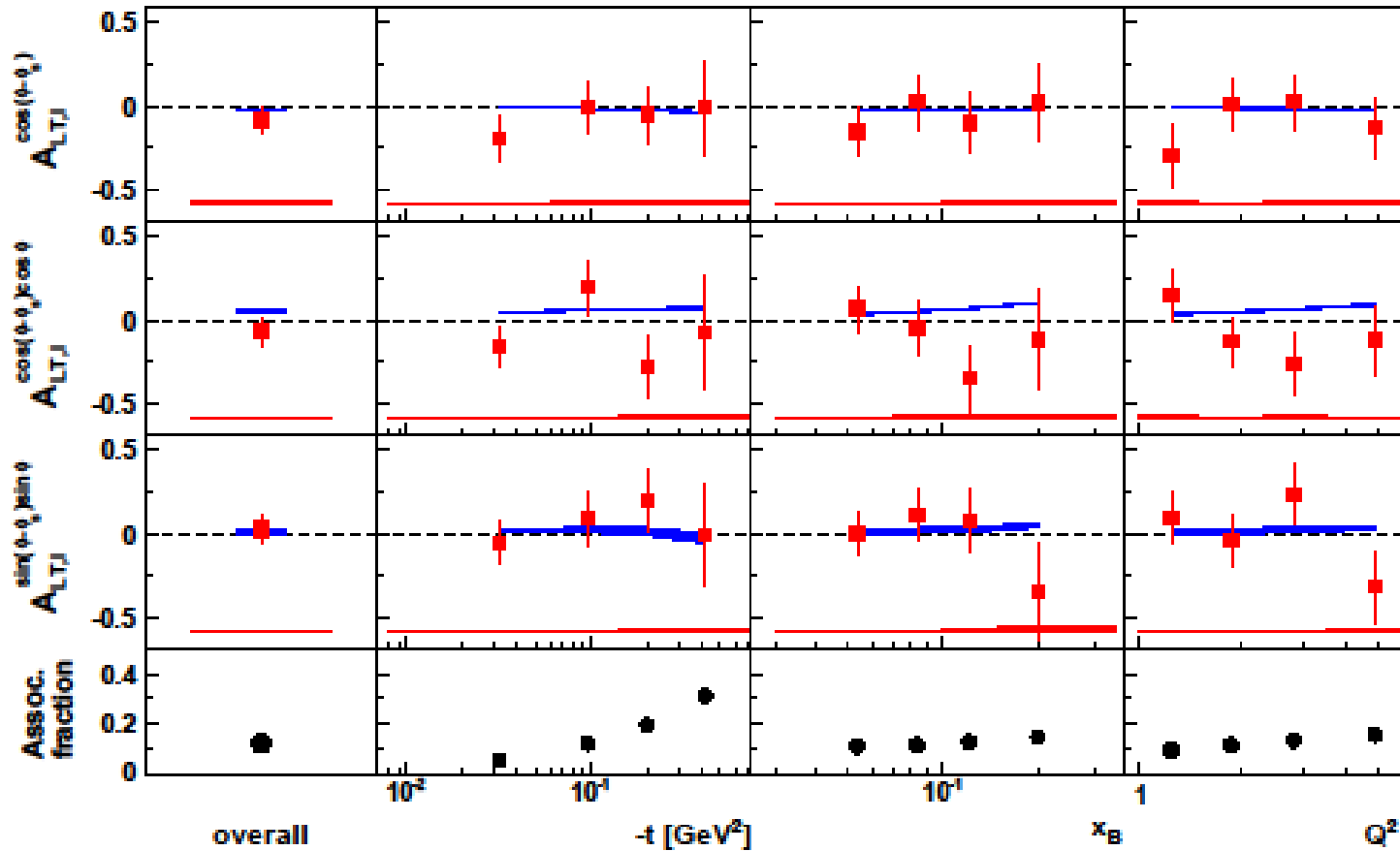
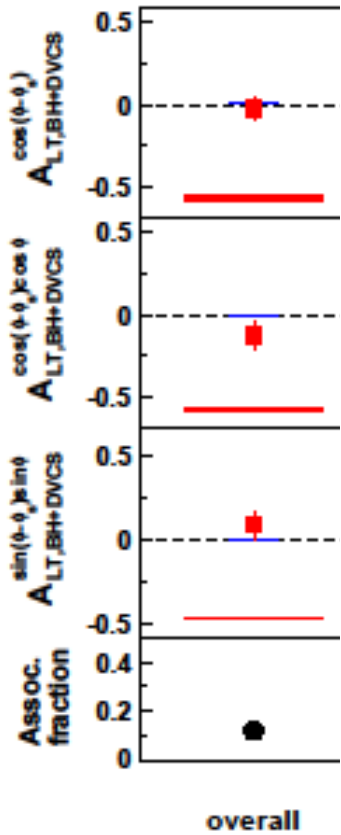
- Indication that the leading amplitude for **pure BH/DVCS** (background < 0.1%) is slightly larger in magnitude than the one in **Recoil Detector acceptance**
- Extraction of asymmetry amplitudes for associated processes is a subject of ongoing dedicated analysis



# Charge-difference double Spin LT Asymmetry

see poster A. Movsisyan

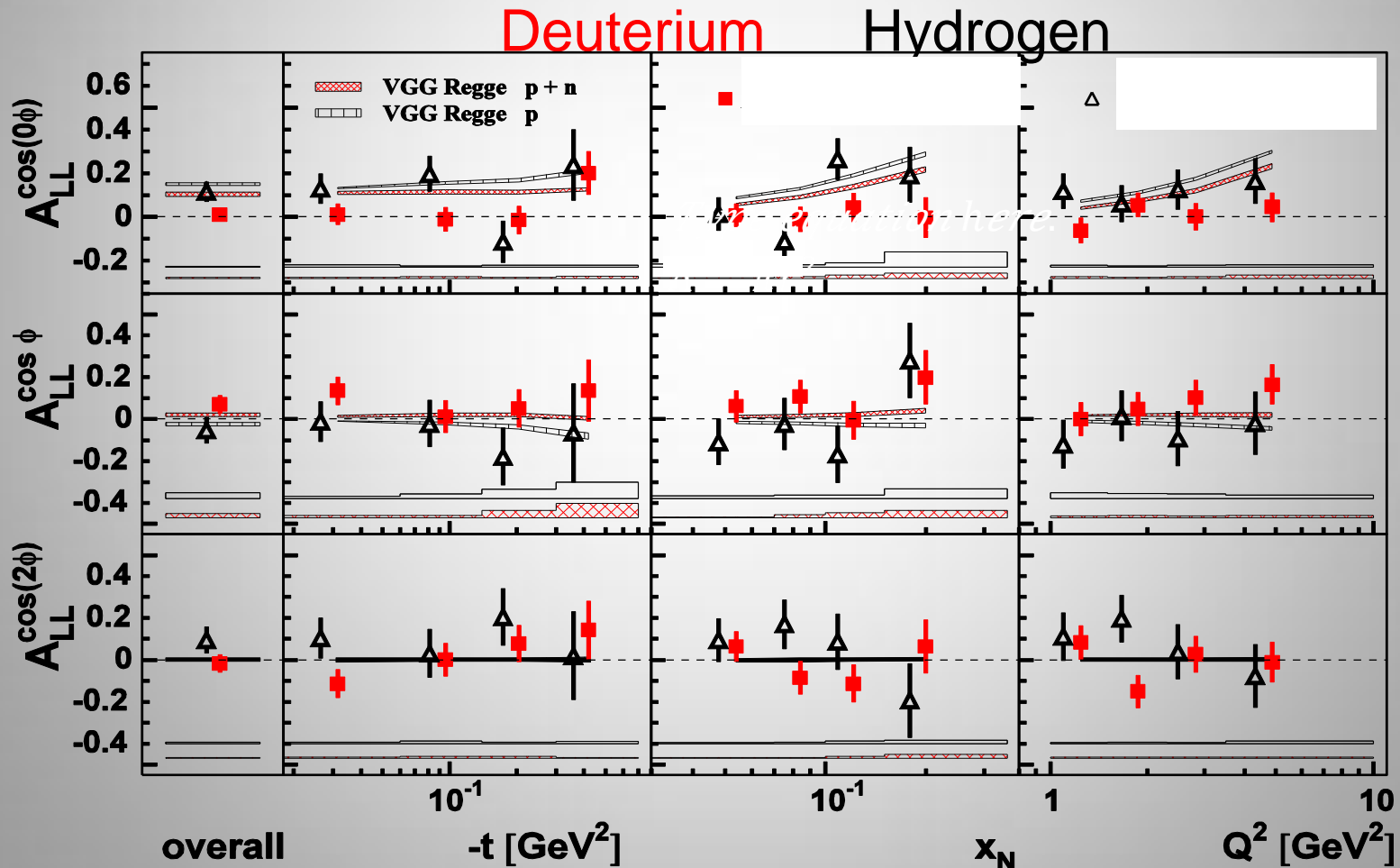
- Sensitive to Re (H+E)
- All consistent with zero. Still: **useful input for global GPD fits!**





# Single-charge double Spin LL Asymmetry

- Sensitive to  $(\tilde{H})$ , non-zero Hydrogen asymmetry amplitudes
- Useful input for global fits. (Deuterium data are rare!)



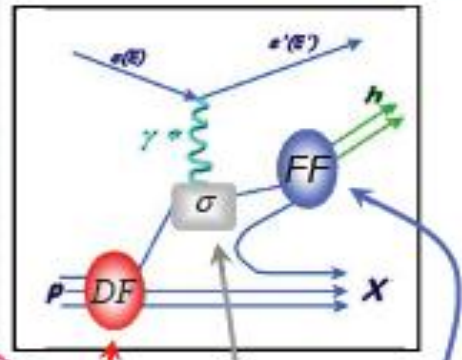
# Semi-inclusive deep inelastic scattering

quark polarisation

N/q	U	L	T
U	$f_1$ number density		
L		$g_1$ helicity	
T			$h_T$ transversity

nucleon polarisation

## Leading Twist PDFs



see poster  
E. Avetisyan

$$\sigma^{ep \rightarrow ehX} = \sum_{q} DF \otimes \sigma^{eq \rightarrow eq} \otimes FF^h$$

## Leading Twist TMDs

Intrinsic  $k_T$  quark distribution

### Distribution Functions (DF)

quark polarisation

N/q	U	L	T
U	$f_1$ number density		$h_1^+$ Boer-Mulders
L		$g_1$ helicity	$h_{1L}^+$ Boer-Mulders
T	$f_{1T}^+$ Sivers	$g_{1T}$ Boer-Mulders	$h_T^+$ transversity

nucleon polarisation

### Fragmentation Functions (FF)

N/q	U	L	T
U	$D_1$ unpolarized		$H_1^+$ Collins

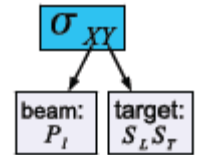
Off-diagonal elements are important objects:

Interference between wave functions with different angular momenta: contains infos about parton orbital angular momenta

Testing QCD at the amplitude level

# 1-hadron production cross section

$$\begin{aligned}
 d\sigma = & d\sigma_{UU}^0 + \cos(2\phi)d\sigma_{UU}^1 + \frac{1}{Q}\cos(\phi)d\sigma_{UU}^2 + P_t\frac{1}{Q}\sin(\phi)d\sigma_{LU}^3 \\
 & + S_L\left[\sin(2\phi)d\sigma_{UL}^4 + \frac{1}{Q}\sin(\phi)d\sigma_{UL}^5 + P_t\left(d\sigma_{LL}^6 + \frac{1}{Q}\cos(\phi)d\sigma_{LL}^7\right)\right] \\
 & + S_T\left[\sin(\phi - \phi_s)d\sigma_{UT}^8 + \sin(\phi + \phi_s)d\sigma_{UT}^9 + \sin(3\phi - \phi_s)d\sigma_{UT}^{10} + \frac{1}{Q}\sin(2\phi - \phi_s)d\sigma_{UT}^{11} + \frac{1}{Q}\sin(\phi_s)d\sigma_{UT}^{12}\right. \\
 & \left. P_t\left(\cos(\phi - \phi_s)d\sigma_{LT}^{13} + \frac{1}{Q}\cos(\phi_s)d\sigma_{LT}^{14} + \frac{1}{Q}\cos(2\phi - \phi_s)d\sigma_{LT}^{15}\right)\right]
 \end{aligned}$$



☛ disentangling the contributions:

☛ experiments with beam and target polarization states (U, L, T)

☛ extract the relevant Fourier amplitudes based on their azimuthal dependences

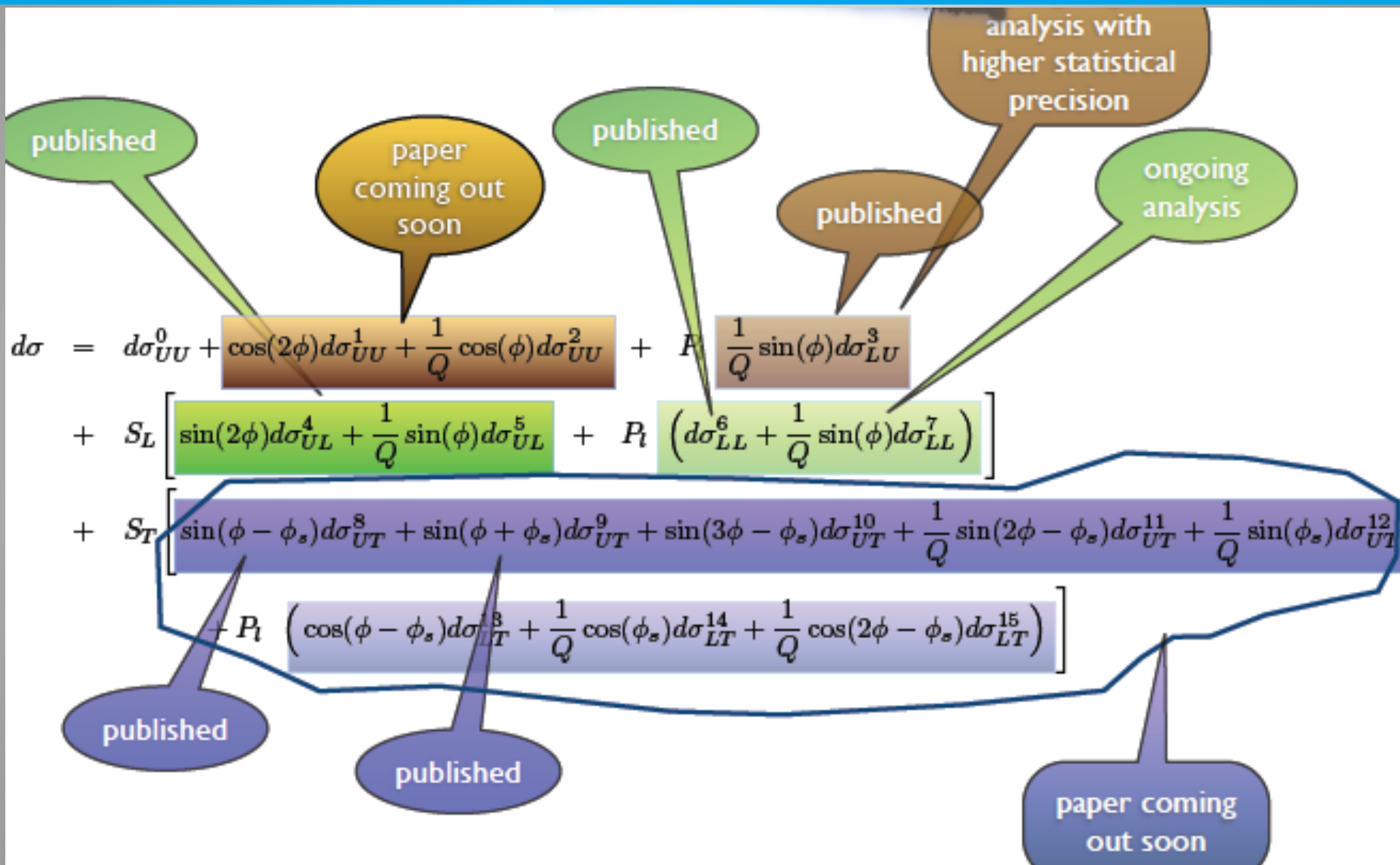
☛ if no perfect detection efficiency:

$$\begin{aligned}
 N(\phi, \phi_s) = & \epsilon(\phi, \phi_s) \sigma_{UU}^0 \left\{ 1 + 2\langle \cos \phi \rangle_{UU} \cos \phi + 2\langle \cos 2\phi \rangle_{UU} \cos 2\phi \right. \\
 & + S_T \left( 2\langle \sin(\phi - \phi_s) \rangle_{UT} \sin(\phi - \phi_s) + 2\langle \sin(\phi + \phi_s) \rangle_{UT} \sin(\phi + \phi_s) + \right. \\
 & \left. \left. 2\langle \sin(3\phi - \phi_s) \rangle_{UT} \sin(\phi + \phi_s) + \dots \right) \right. \\
 & + S_T P_t \left( 2\langle \cos(\phi - \phi_s) \rangle_{UT} \cos(\phi - \phi_s) + 2\langle \cos \phi_s \rangle_{UT} \cos \phi_s + \right. \\
 & \left. \left. 2\langle \cos(2\phi - \phi_s) \rangle_{UT} \cos(\phi - \phi_s) \right) \right\}
 \end{aligned}$$

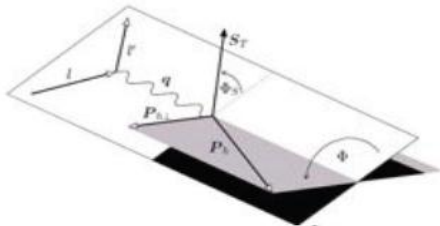
☛ fit the cross section asymmetry for opposite spin states



# Overview on HERMES activities in SIDIS sector



# Cahn and Boer-Mulders effect



$$\frac{d\sigma^h}{dx dy d\phi_S dz d\phi dP_{h\perp}^2} = \frac{\alpha^2 y^2}{xyQ^2 2(1-\epsilon)} \left( 1 + \frac{\gamma^2}{2x} \right)$$

$$\left\{ \begin{aligned} & [F_{UU}^{\perp T} + \epsilon F_{UU}^{\perp L} \\ & + \sqrt{2\epsilon(1+\epsilon)} \cos(\phi) F_{UU}^{\cos(\phi)} + \epsilon \cos(2\phi) F_{UU}^{\cos(2\phi)}] \end{aligned} \right.$$

+  $\lambda$

+  $S_L$

+  $S_L \lambda$

+  $S_T$

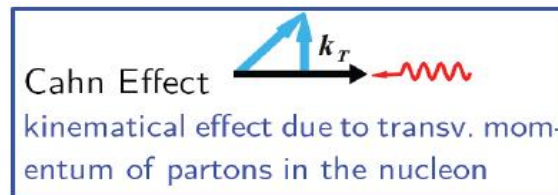
+  $S_T \lambda$

		quark		
		U	L	T
nucleon	U	$f_1$		$h_1^\perp$
	L		$g_1$	$h_{1L}^\perp$
	T	$f_{1T}^\perp$	$g_{1T}^\perp$	$h_{1T}^\perp$

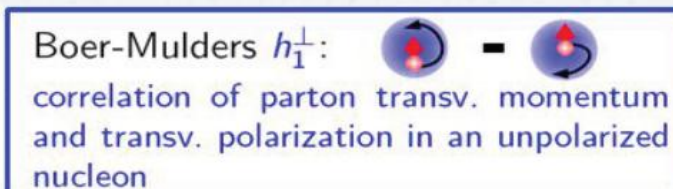
$$\sigma_{UU}^{\cos(\phi)} \propto [f_1 \otimes D_1 + h_1^\perp \otimes H_1^\perp + \dots] / Q$$

$$\sigma_{UU}^{\cos(2\phi)} \propto h_1^\perp \otimes H_1^\perp + [f_1 \otimes D_1 + \dots] / Q^2$$

- Cahn effect:



- Boer-Mulders effect: Boer-Mulders TMD



# Phi modulations in unpolarized SIDIS pion X-section

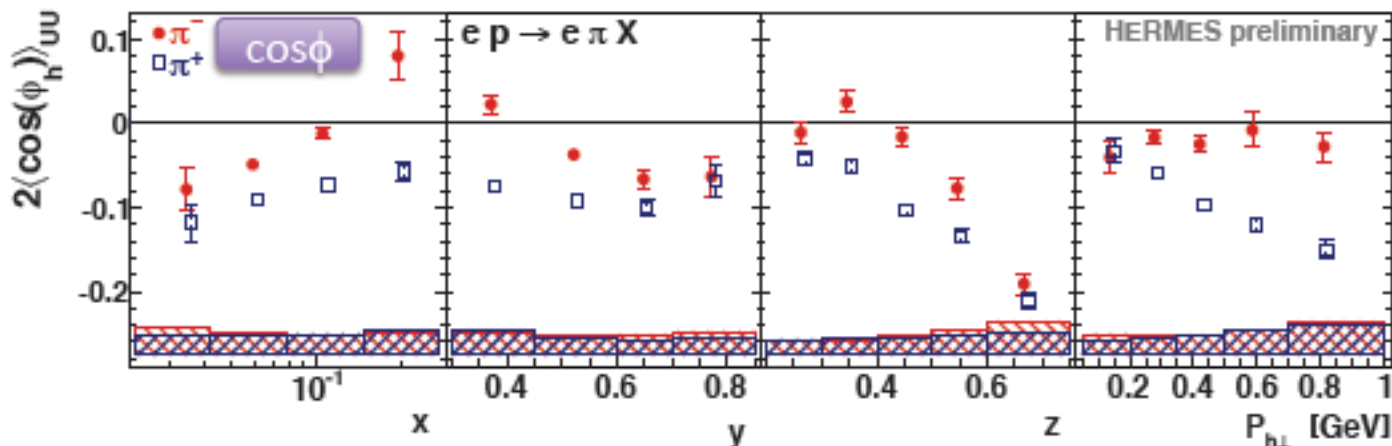
$\cos\phi$  large and negative !

$$\sigma_{UU}^{\cos(\phi)} \propto [f_1 \otimes D_1 + h_1^\perp \otimes H_1^\perp + \dots] / Q$$

Increasing with  $z$  and  $P_h$

Large difference in hadron charge !

Larger in magnitude for  $\pi^+$



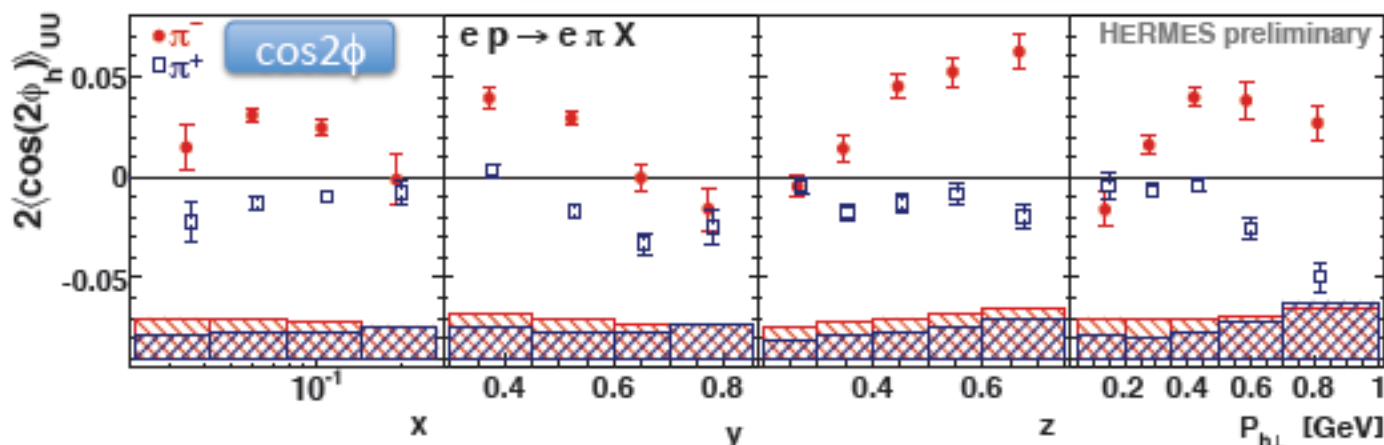
$\cos 2\phi$  non-zero !

$$\sigma_{UU}^{\cos(2\phi)} \propto h_1^\perp \otimes H_1^\perp + [f_1 \otimes D_1 + \dots] / Q^2$$

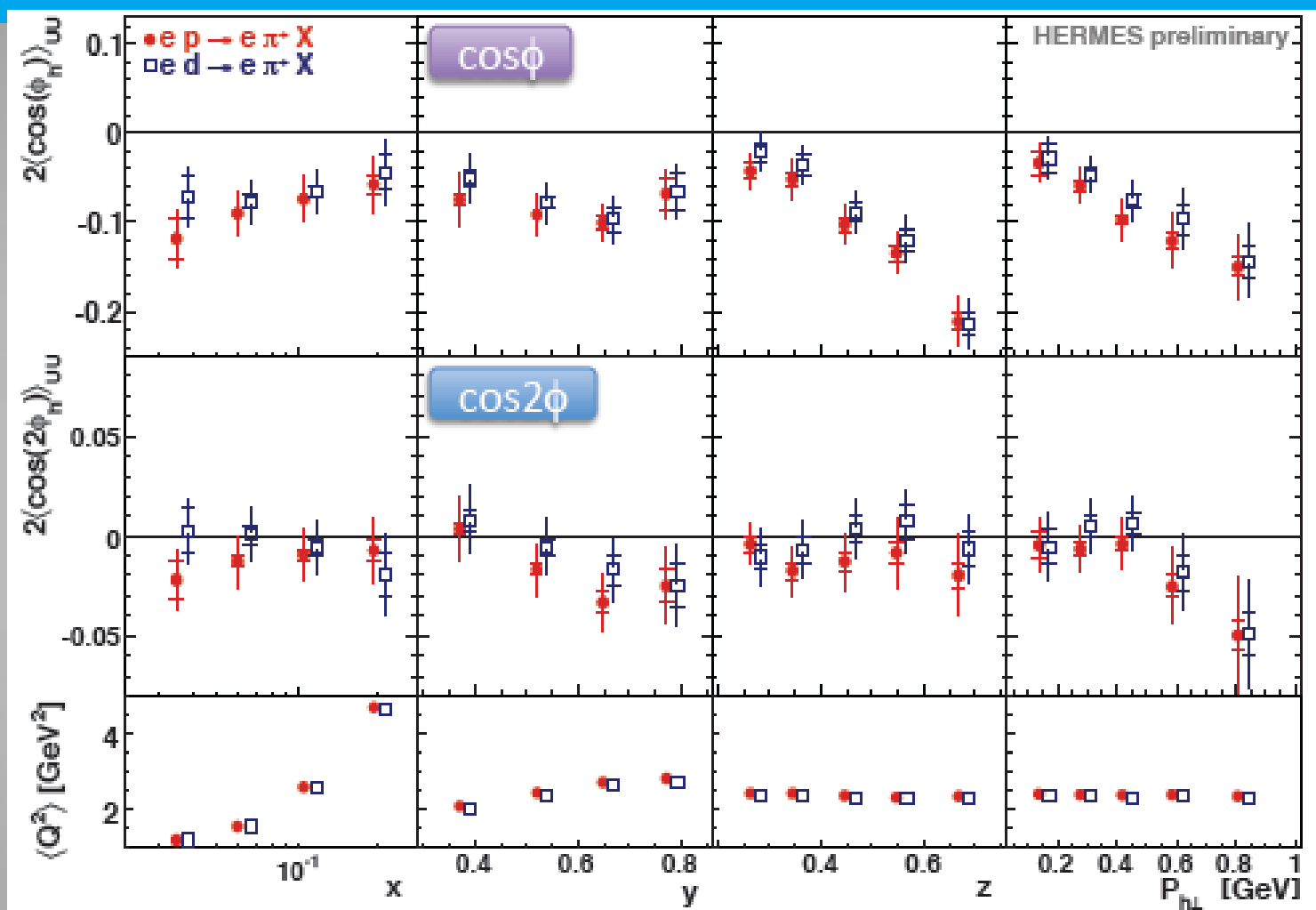
Difference in hadron charge !

Positive for  $\pi^-$

Negative for  $\pi^+$



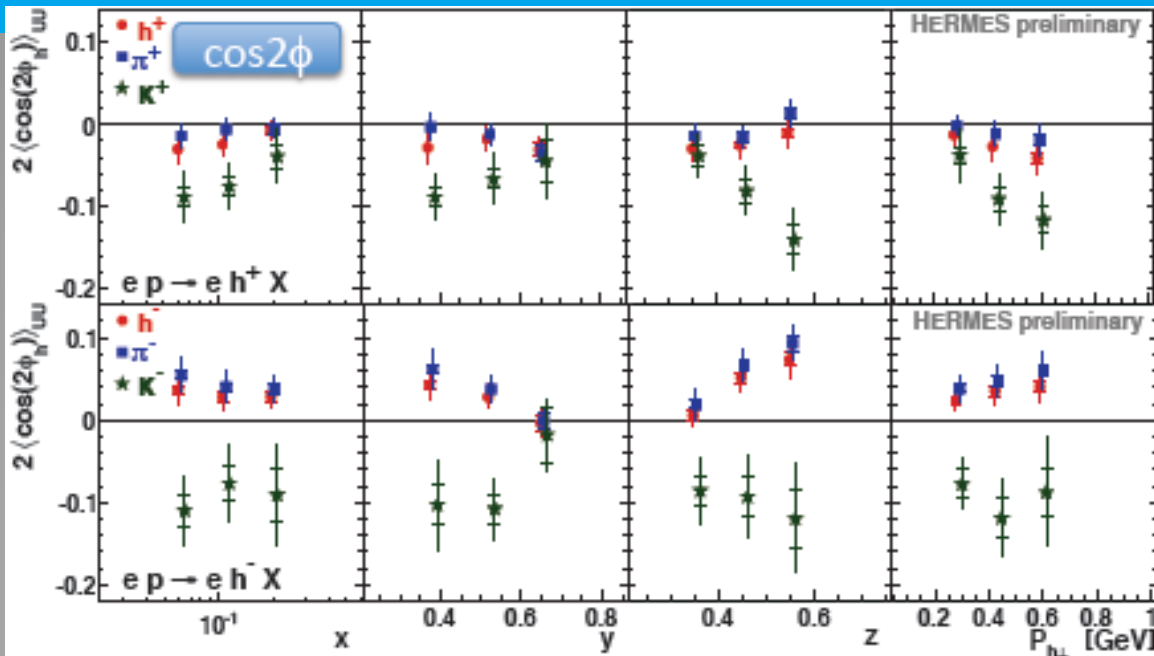
# Proton vs. deuteron data: u- vs d-quark Boer-Mulders fct



Quark d vs u contribution ?  
 DATA support Boer-Mulders of same sign for u and d



# Boer-Mulders: kaons vs pions



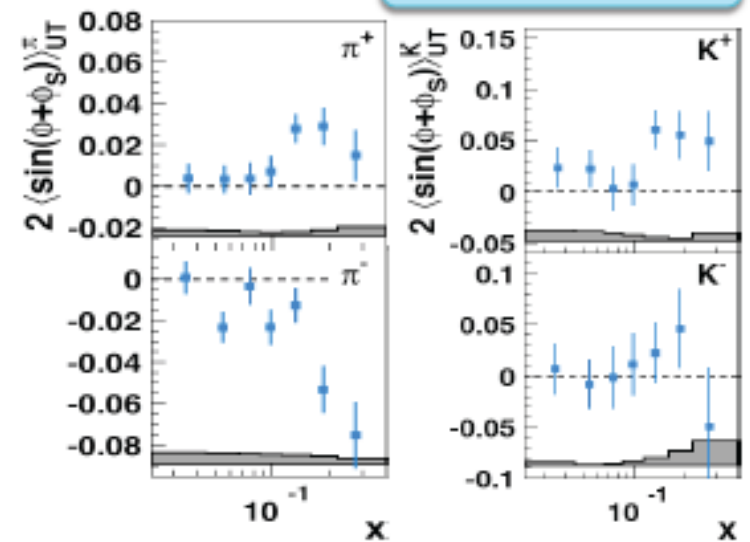
$$\sigma_{UU}^{\cos(2\phi)} \propto h_1^\perp \otimes H_1^\perp + [f_1 \otimes D_1 + \dots] / Q^2$$



Striking difference versus pions !

Phys. Lett. B 693 (2010) 11-16

$$\sigma_{UT}^{\sin(\phi+\phi_S)} \propto h_1 H_1^\perp$$



The kaon puzzle

Already found in  $A_{UT}$ : Collins+Transversity

Role of the sea in distribution and fragmentation functions



# 'Pretzelosity' DF from transversely polarized proton

$$\begin{aligned}
 d\sigma &= d\sigma_{UU}^0 + \cos(2\phi)d\sigma_{UU}^1 + \frac{1}{Q}\cos(\phi)d\sigma_{UU}^2 + P_t\frac{1}{Q}\sin(\phi)d\sigma_{LU}^3 \\
 &+ S_L\left[\sin(2\phi)d\sigma_{UL}^4 + \frac{1}{Q}\sin(\phi)d\sigma_{UL}^5 + P_t\left(d\sigma_{LL}^6 + \frac{1}{Q}\cos(\phi)d\sigma_{LL}^7\right)\right] \\
 &+ S_T\left[\sin(\phi - \phi_s)d\sigma_{UT}^8 + \sin(\phi + \phi_s)d\sigma_{UT}^9 + \sin(3\phi - \phi_s)d\sigma_{UT}^{10} + \frac{1}{Q}\sin(2\phi - \phi_s)d\sigma_{UT}^{11} + \frac{1}{Q}\sin(\phi_s)d\sigma_{UT}^{12}\right. \\
 &\quad \left. P_t\left(\cos(\phi - \phi_s)d\sigma_{LT}^{13} + \frac{1}{Q}\cos(\phi_s)d\sigma_{LT}^{14} + \frac{1}{Q}\cos(2\phi - \phi_s)d\sigma_{LT}^{15}\right)\right]
 \end{aligned}$$

☛ “pretzelosity” DF  $h_{1T}^{\perp,q}(x, p_T^2)$  gives a measure of the deviation of the nucleon shape from a sphere

☛ correlation between parton transverse momentum and parton transverse polarization in a transversely polarized nucleon

☛ it is expected to be suppressed at small and large  $x$  w.r.t.  $f_1^q$ ,  $g_1^q$ ,  $h_1^q$

☛ involve quark and nucleon helicity flips;  
is related to chiral-odd GPD

$$h_{1T}^{(0)\perp,q}(x, p_T^2) = \frac{3}{(1-x)^2} \tilde{H}_T^q(x, 0, 0)$$

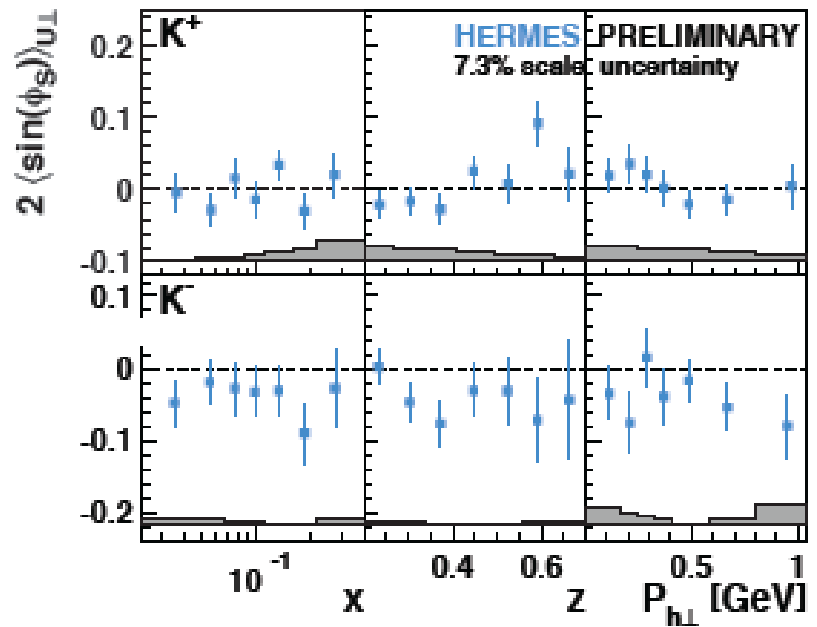
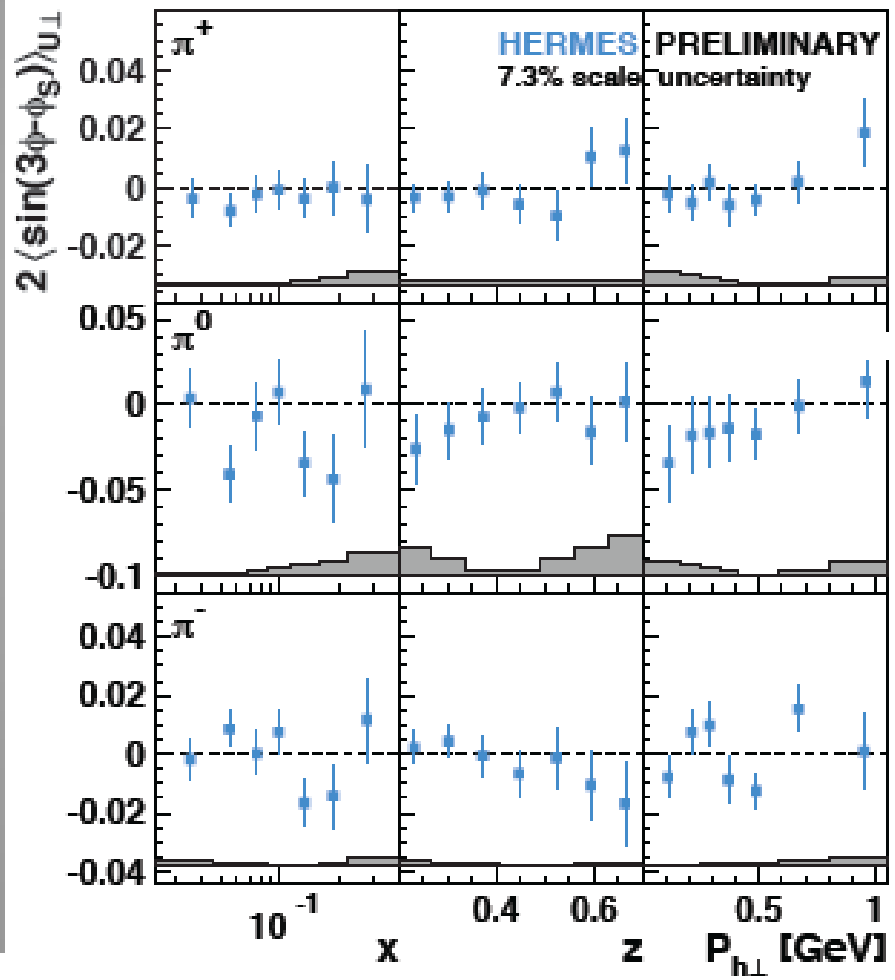
☛ gives the measure of ‘relativistic effects’ in the nucleon:  $\frac{p_T^2}{2M^2} h_{1T}^{\perp,q}(x, p_T^2) = g_1^q(x, p_T^2) - h_1^q(x, p_T^2)$





# Result on pretzelosity

$$2 \langle \sin(3\phi - \phi_s) \rangle_{UT} \propto \frac{\sum_q e_q^2 x h_{1T}^{\perp(1),q}(x) \otimes_w H_1^{\perp(1/2)q}(z)}{\sum_q e_q^2 f_1^q(x) \otimes D_1^q(z)}$$




- ➡ suppressed by two powers of  $P_{h\perp}$  compared to Collins and Sivers amplitudes
- ➡ compatible with zero within uncertainties
- ➡ pretzelosity might be non-zero at higher  $P_{h\perp}$




# 'Worm-gear' DFs from transversely polarized proton

$$\begin{aligned}
 d\sigma &= d\sigma_{UU}^0 + \cos(2\phi)d\sigma_{UU}^1 + \frac{1}{Q}\cos(\phi)d\sigma_{UU}^2 + P_1\frac{1}{Q}\sin(\phi)d\sigma_{LU}^3 \\
 &+ S_L\left[\sin(2\phi)d\sigma_{UL}^4 + \frac{1}{Q}\sin(\phi)d\sigma_{UL}^5 + P_1\left(d\sigma_{LL}^6 + \frac{1}{Q}\cos(\phi)d\sigma_{LL}^7\right)\right] \\
 &+ S_T\left[\sin(\phi - \phi_s)d\sigma_{UT}^8 + \sin(\phi + \phi_s)d\sigma_{UT}^9 + \sin(3\phi - \phi_s)d\sigma_{UT}^{10} + \frac{1}{Q}\sin(2\phi - \phi_s)d\sigma_{UT}^{11} + \frac{1}{Q}\sin(\phi_s)d\sigma_{UT}^{12}\right. \\
 &\quad \left. P_1\left[\cos(\phi - \phi_s)d\sigma_{LT}^{13} + \frac{1}{Q}\cos(\phi_s)d\sigma_{LT}^{14} + \frac{1}{Q}\cos(2\phi - \phi_s)d\sigma_{LT}^{15}\right]\right]
 \end{aligned}$$

 worm-gear DF  $g_{1T}^q(x, p_T^2)$  and  $h_{1L}^{\perp, q}(x, p_T^2)$  describes the probability to find a longitudinal/transverse polarized quark in a transversely/longitudinally polarized nucleon

 on a transversely target  $h_{1L}^{\perp, q}(x, p_T^2)$  accessible in the measurements through  $\sin(2\phi + \phi_s)$  Fourier component

 gives correlation between parton transverse momentum and parton longitudinal / transverse polarization in a longitudinal / transversely polarized nucleon

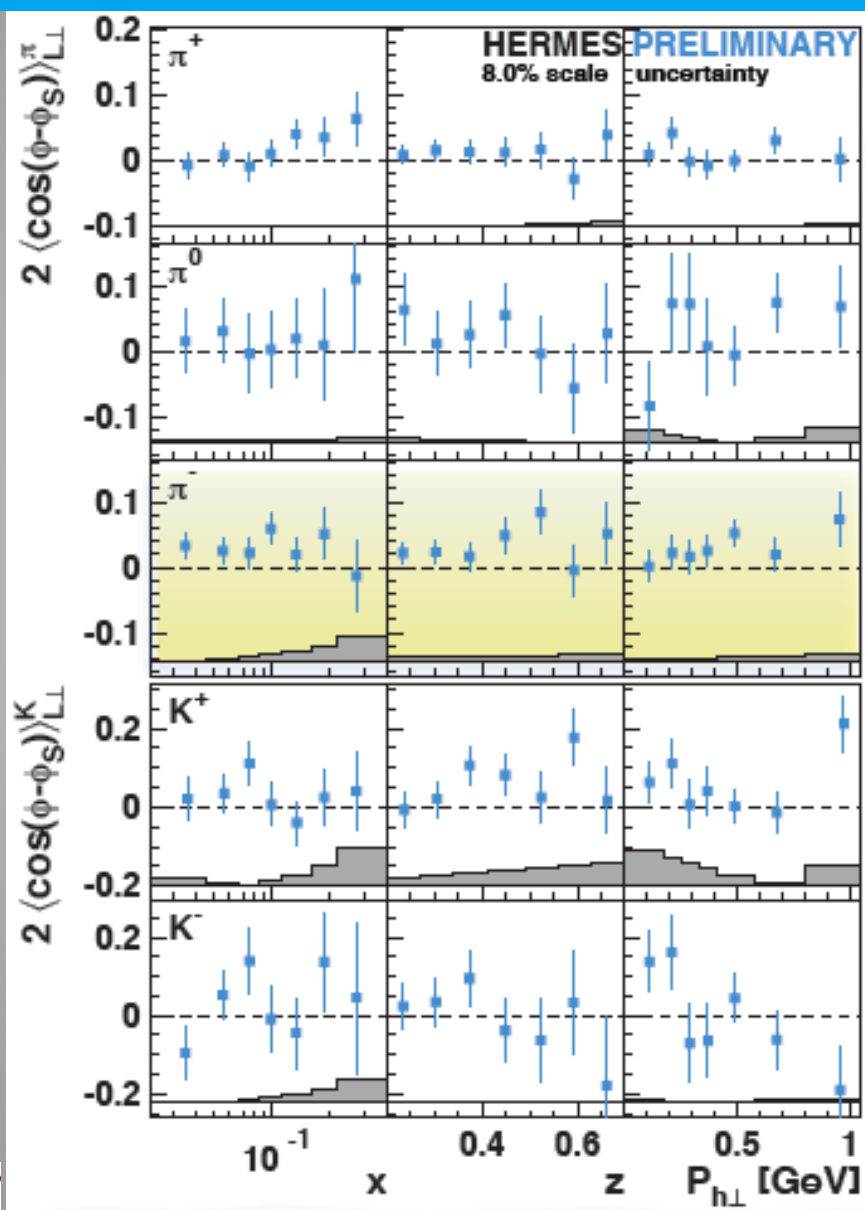
 model dependent relations:  $g_{1T}^{\perp, q}(x, p_T^2) \approx x \int_x^1 \frac{1}{y} g_1^q(y, p_T^2) dy$

$$h_{1L}^{\perp, q}(x, p_T^2) = -g_{1T}^{\perp, q}(x, p_T^2) \quad h_{1L}^{\perp, q}(x, p_T^2) \approx -x \int_x^1 \frac{1}{y} h_1^q(y, p_T^2) dy$$





# Worm-gear DF from long.pol. beam/transv.pol.target



*the  $\cos(\phi - \phi_s)$  amplitudes*

$$2\langle \cos(\phi - \phi_s) \rangle_{LT} \propto \frac{C \left[ -\frac{\hat{P}_{h\perp} \cdot \mathbf{p}_T}{M_h} g_{1T}^{\perp,q}(x, p_T^2) D_1^q(z, k_T^2) \right]}{C \left[ f_1^q(x, p_T^2) D_1^q(z, k_T^2) \right]}$$

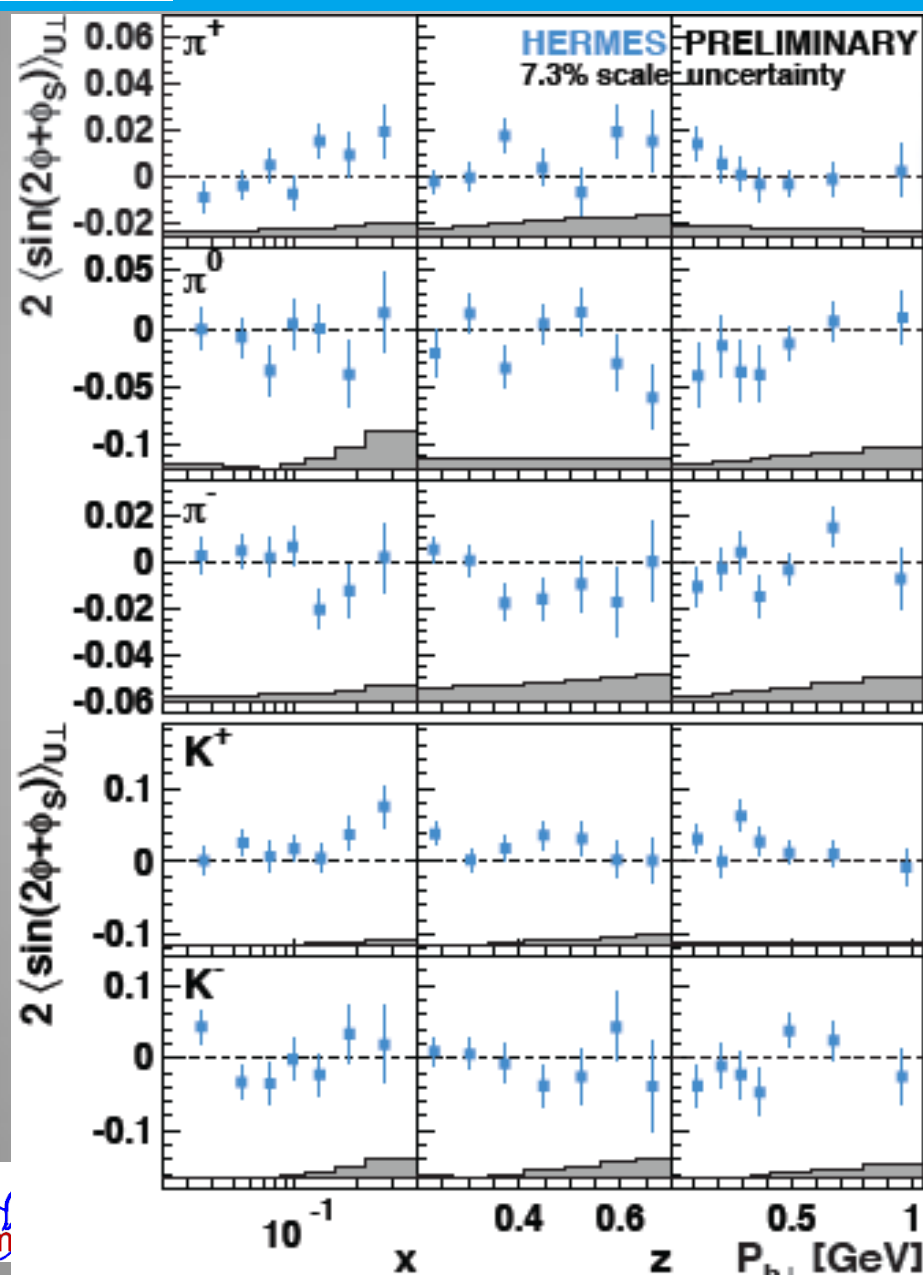
uncertainties are larger than in single-spin asymmetries scaled by the beam polarization value

- $\pi^+$  slightly positive
- $\pi^0$  compatible with zero
- $\pi^-$  positive  
 evidence for non-zero worm-gear distribution
- $K^+$  slightly positive
- $K^-$  compatible with zero



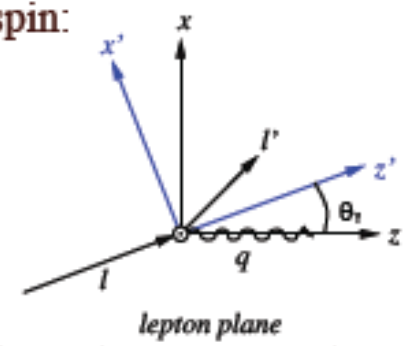


# Subleading-twist amplitudes for pions & kaons



*the subleading-twist  $\sin(2\phi + \phi_s)$  amplitudes*

arises solely from longitudinal component of the target spin:



$$P_T A_{U\perp}(\phi, \phi_s) = S_T A_{UT}(\phi, \phi_s) + S_L A_{UL}(\phi, \phi_s)$$

- longitudinal component of the target spin <15%
- expected to scale as  $\sin \theta_{\gamma} \cdot \langle \sin(2\phi)_{UL} \rangle$
- related to worm-gear DF  $h_{1L}^{\perp, q}$
- $\sin(2\phi + \phi_s)$  amplitude is suppressed by one power of  $P_{h\perp}$  compared to Collins and Sivers amplitudes
- compatible with zero within uncertainties except maybe  $K^+$

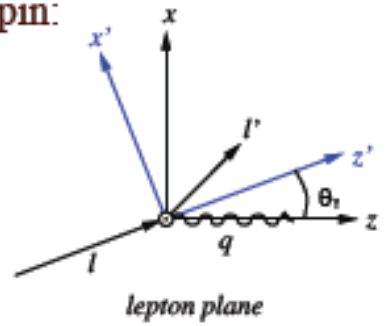




# Subleading-twist amplitudes for pions & kaons

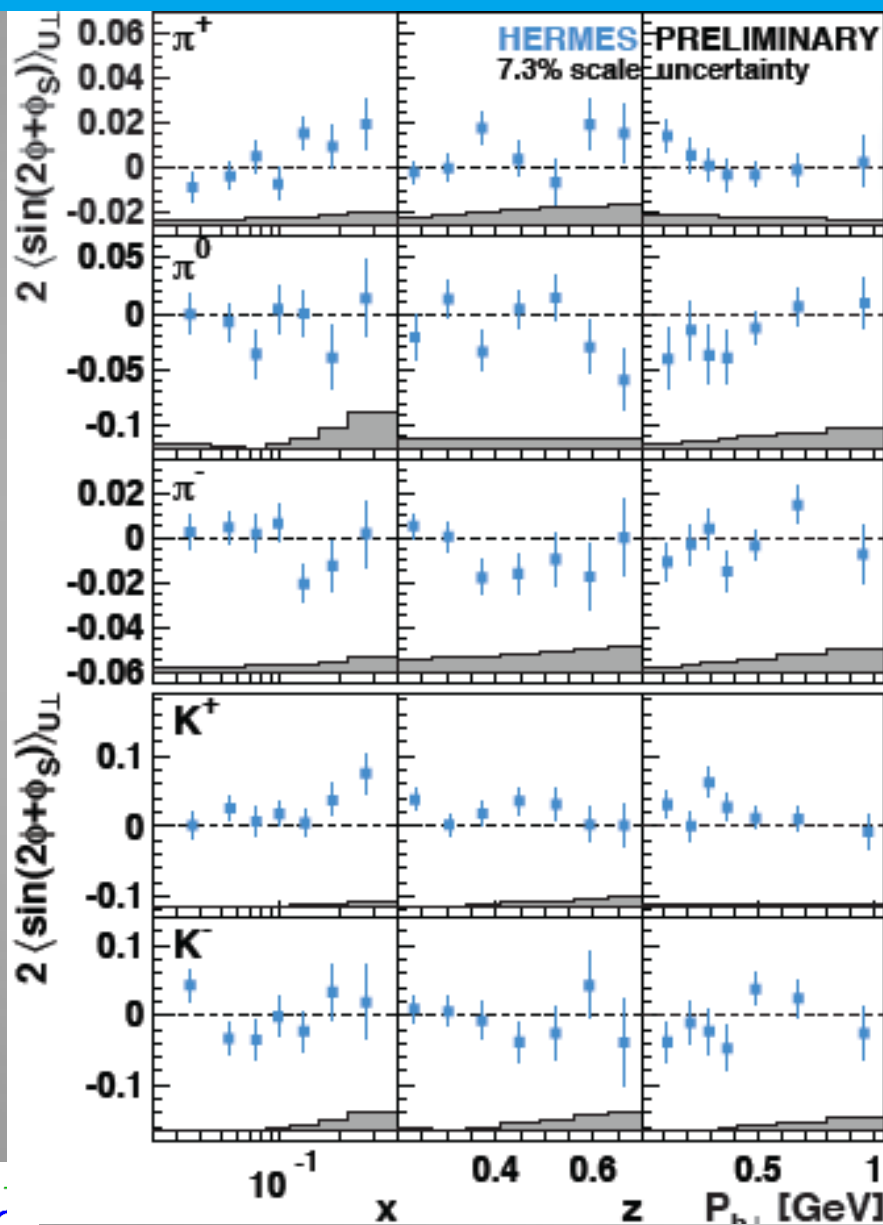
*the subleading-twist  $\sin(2\phi + \phi_s)$  amplitudes*

arises solely from longitudinal component of the target spin:



$$P_T A_{U\perp}(\phi, \phi_s) = S_T A_{UT}(\phi, \phi_s) + S_L A_{UL}$$

- longitudinal component of the target spin <15%
- expected to scale as  $\sin \theta_{\gamma^*} \langle \sin(2\phi)_{UL} \rangle$
- related to worm-gear DF  $h_{1L}^{\perp,q}$
- $\sin(2\phi + \phi_s)$  amplitude is suppressed by one power of  $P_{h\perp}$  compared to Collins and Sivers amplitudes
- compatible with zero within uncertainties except maybe  $K^+$



# Conclusions and Outlook

- HERMES has been pioneering the experimental study of the spin structure of the nucleon towards GPDs and TMDs
- HERMES will stay unique (for a while) offering exclusive and semi-inclusive data taken with both beam charges and undiluted Hydrogen and Deuterium targets
- HERMES has been publishing an almost full spectrum of **asymmetries in DVCS** and intends to publish more on **exclusive meson production** (Rho, Phi, Omega, Eta, Pi0)
- HERMES intends to soon have published results on a wide variety of **(flavor-separated) SIDIS asymmetries**

Spin results not addressed in this talk: Exclusive VM production **(see poster B. Marianski)**

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