

New results of exclusive ϕ^0 and ρ^0 vector meson production at HERMES

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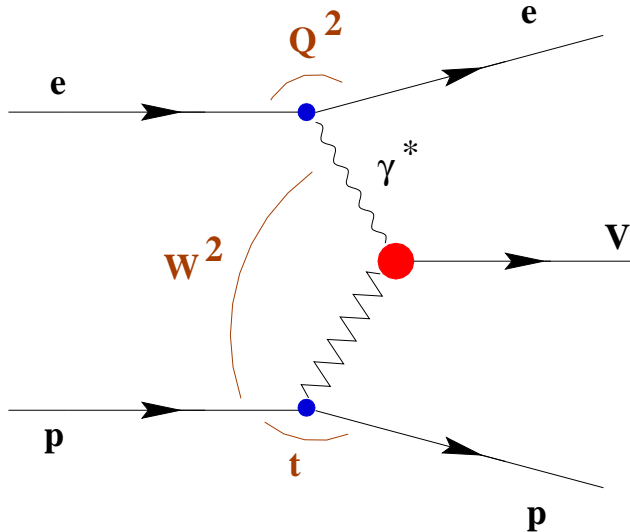
on behalf of HERMES Collaboration

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- Rudiments
- Spin Density Matrix Elements (SDME's) : definitions and their determination
- SDME's and Amplitudes for ρ^0 and ϕ vector mesons
- Dependencies of SDME's on Q^2 and t'
- The observables:
 - $R = \frac{\sigma_L}{\sigma_T}$
 - the signatures of the Natural or Unnatural Parity Exchange amplitudes
- Conclusions
- Outlook

$e + p \rightarrow e' + p' + V$: *Rudiments*



Kinematics:

- $\nu = 5 \div 24$ **GeV**, $\langle \nu \rangle = 13.3$ **GeV**,
- $Q^2 = 0.5 \div 7.0$ **GeV²**, $\langle Q^2 \rangle = 2.3$ **GeV²**
- $W = 3.0 \div 6.5$ **GeV**, $\langle W \rangle = 4.9$ **GeV**,
- $x_{Bj} = 0.01 \div 0.35$, $\langle x_{Bj} \rangle = 0.07$
- $t' = 0 \div 0.4$ **GeV²**, $\langle t' \rangle = 0.13$ **GeV²**

- In one photon approximation
 $\equiv \gamma^* + p \rightarrow p' + V$
- The amplitude of this process can be factorized:
 $A = \Phi_{\gamma^* \rightarrow q\bar{q}}^* \otimes A_{q\bar{q}+p \rightarrow q\bar{q}+p} \otimes \Phi_{q\bar{q} \rightarrow V}$.
In these three steps the interaction time ($q\bar{q}$) with target is shorter than fluctuation and formation of VM. (Collins, Frankfurt and Stirman Phys.Rev D56(1997)2982)
- $\gamma^* + N \rightarrow \rho^0(\phi) + N'$ it is good tool to study the helicity conservation:
 - helicity state of γ^* is easy to determine (QED)
 - $\rho^0 \rightarrow \pi^+ \pi^-$ decay determines the helicity of ρ^0

Exclusive ρ^0 and ϕ Meson Production

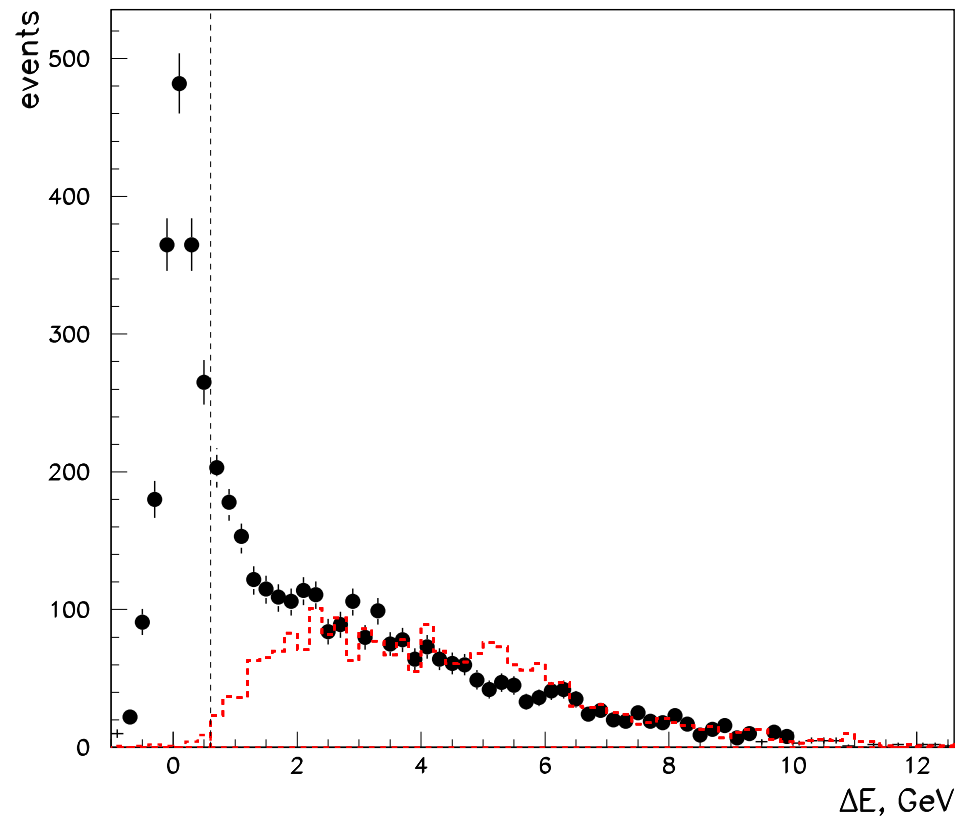
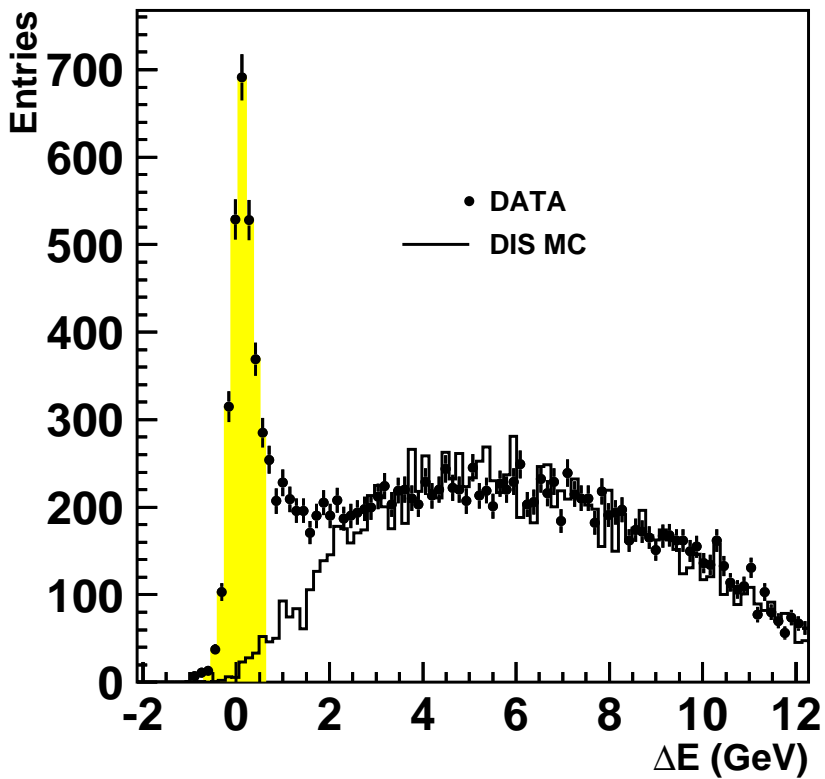
$$e + p \rightarrow e' + p' + \rho^0 \rightarrow \pi^+ \pi^-$$

$$e + p \rightarrow e' + p' + \phi \rightarrow K^+ K^-$$

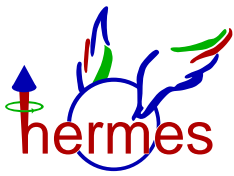
The exclusive events were selected from the missing energy spectra $\Delta E = \frac{M_X^2 - M_p^2}{2M_p}$ for:

ρ^0

ϕ .



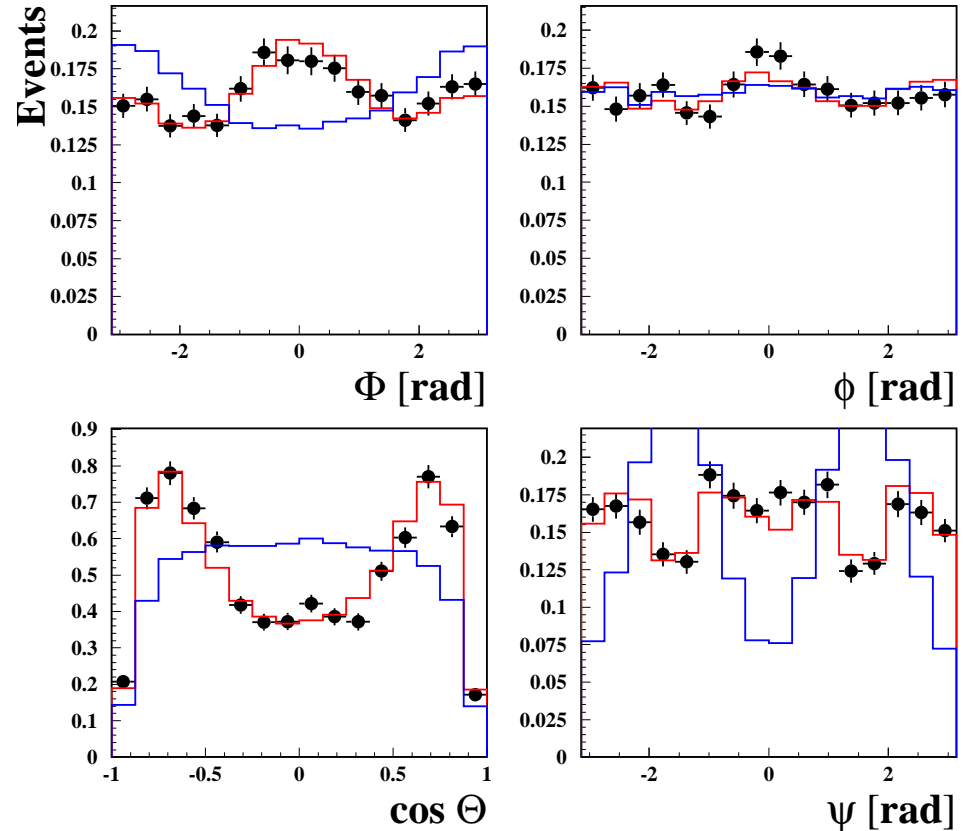
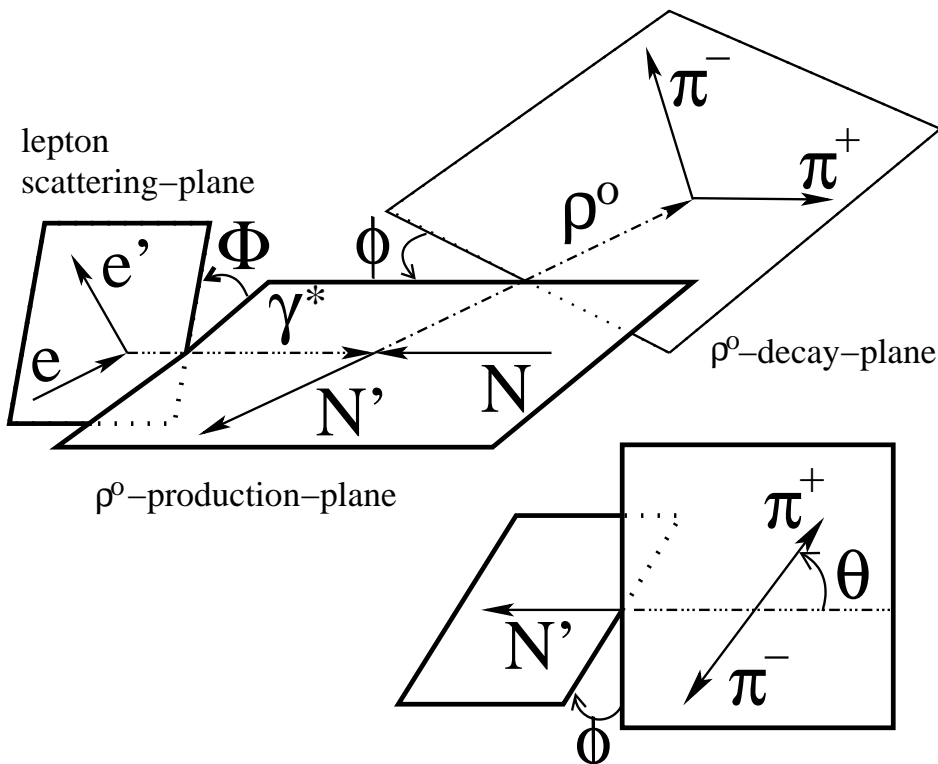
The background was simulated by code MC PYTHIA.



ρ^0 & ϕ -meson Spin Density Matrix Elements (SDMEs)

- SDMEs: $r_{\lambda_\rho \lambda'_\rho}^\alpha \sim \rho(V) = \frac{1}{2} T_{\lambda_V \lambda_\gamma} \rho(\gamma) T_{\lambda_V \lambda_\gamma}^+$
spin-density matrix of the vector meson $\rho(V)$ in terms of the photon matrix $\rho(\gamma)$ and helicity amplitude $T_{\lambda_V \lambda_\gamma}$
- presented according to K.Schilling and G.Wolf (Nucl. Phys. B61 (1973) 381)
 $\alpha = 04, 1 - 3, 5 - 8$ long. or trans. photon, $\lambda_\rho = -1, 0, 1$ - polarization of $\rho^0(\phi)$
- measured experimentally at $5 < W < 75$ GeV (HERMES, COMPASS, H1, ZEUS)
- provide access to *helicity amplitudes* $T_{\lambda_V \lambda_\gamma}$, which are:
 - extracted experimentally from SDMEs
 - calculated from GPDs: S.V.Goloskokov, P.Kroll arXiv:0708.3569 [hep-ph]27.08.07;
Eur.Phys.J. C 50,829 (2007) hep-ph/0601290; Eur.Phys.J. C 42,281 (2005)
hep-ph/0501242

Fit of Angular Distributions Using Max. Likelihood Method in MINUIT



- Simulated Events: matrix of fully reconstructed MC events from initial uniform angular distribution
- Binned Maximum Likelihood Method: $8 \times 8 \times 8$ bins of $\cos(\Theta)$, ϕ , Φ . Simultaneous fit of 23 SDMEs $r_{ij}^\alpha = W(\Phi, \phi, \cos \Theta)$ for data with negative and positive beam helicity ($\langle |P_b| \rangle = 53.5\%$, $\Psi = \Phi - \phi$)

$$W(\cos \Theta, \phi, \Phi) = W^{unpol} + W^{long.pol},$$

$$W^{unpol}(\cos \Theta, \phi, \Phi) =$$

$$\begin{aligned} & \frac{3}{8\pi^2} \left[\frac{1}{2}(1 - r_{00}^{04}) + \frac{1}{2}(3r_{00}^{04} - 1) \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^{04}\} \sin 2\Theta \cos \phi - r_{1-1}^{04} \sin^2 \Theta \cos 2\phi \right. \\ & - \epsilon \cos 2\Phi \left(r_{11}^1 \sin^2 \Theta + r_{00}^1 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^1\} \sin 2\Theta \cos \phi - r_{1-1}^1 \sin^2 \Theta \cos 2\phi \right) \\ & - \epsilon \sin 2\Phi \left(\sqrt{2}\text{Im}\{r_{10}^2\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^2\} \sin^2 \Theta \sin 2\phi \right) \\ & + \sqrt{2\epsilon(1 + \epsilon)} \cos \Phi \left(r_{11}^5 \sin^2 \Theta + r_{00}^5 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^5\} \sin 2\Theta \cos \phi - r_{1-1}^5 \sin^2 \Theta \cos 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1 + \epsilon)} \sin \Phi \left(\sqrt{2}\text{Im}\{r_{10}^6\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^6\} \sin^2 \Theta \sin 2\phi \right) \right], \end{aligned}$$

$$\begin{aligned} W^{long.pol.}(\cos \Theta, \phi, \Phi) = & \frac{3}{8\pi^2} P_{beam} \left[\sqrt{1 - \epsilon^2} \left(\sqrt{2}\text{Im}\{r_{10}^3\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^3\} \sin^2 \Theta \sin 2\phi \right) \right. \\ & + \sqrt{2\epsilon(1 - \epsilon)} \cos \Phi \left(\sqrt{2}\text{Im}\{r_{10}^7\} \sin 2\Theta \sin \phi + \text{Im}\{r_{1-1}^7\} \sin^2 \Theta \sin 2\phi \right) \\ & \left. + \sqrt{2\epsilon(1 - \epsilon)} \sin \Phi \left(r_{11}^8 \sin^2 \Theta + r_{00}^8 \cos^2 \Theta - \sqrt{2}\text{Re}\{r_{10}^8\} \sin 2\Theta \cos \phi - r_{1-1}^8 \sin^2 \Theta \cos 2\phi \right) \right] \end{aligned}$$

SDMEs and Amplitudes for: ρ^0, ϕ

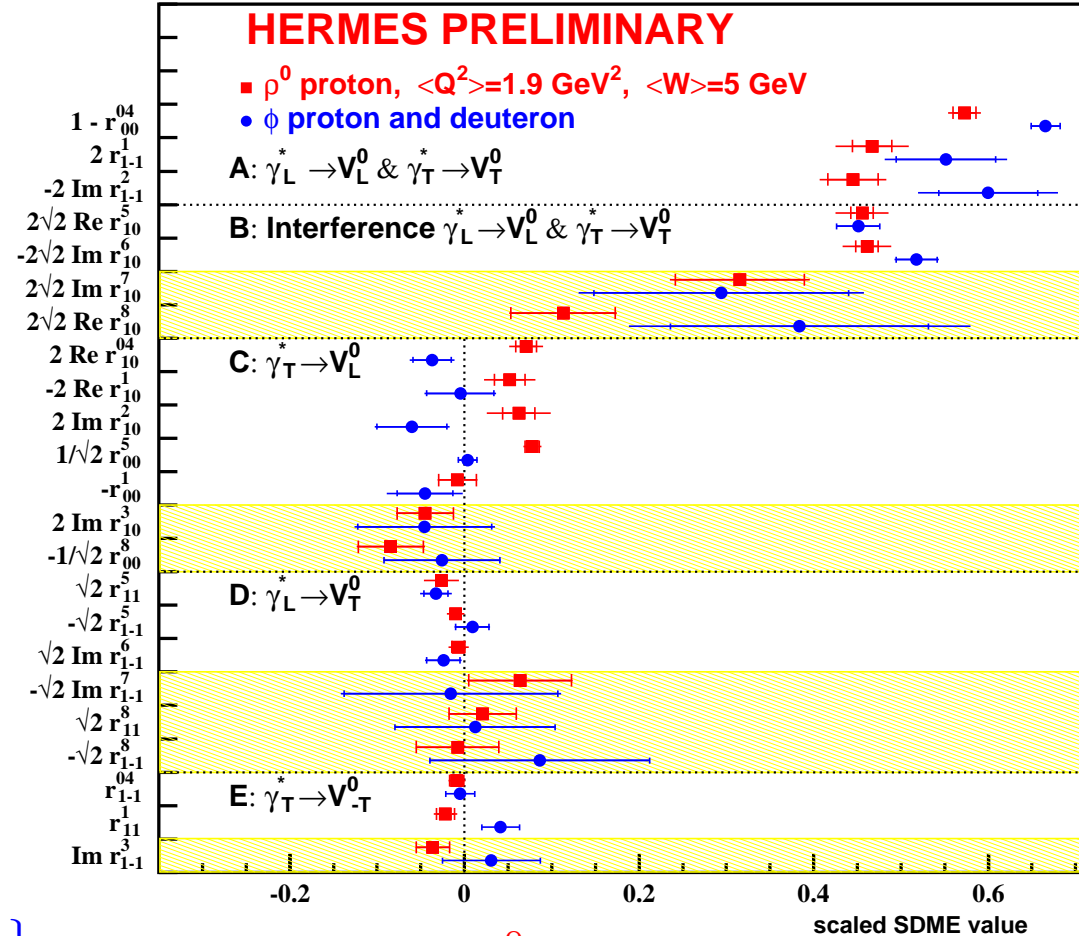
A- SCHC $\gamma_L^* \rightarrow \rho_L^0$ and $\gamma_T^* \rightarrow \rho_T^0$
 $|T_{11}|^2 \propto 1 - r_{00}^4 \propto r_{1-1}^1 \propto -\text{Im}\{r_{1-1}^2\}$

B- Interference: γ_L^*, ρ_T^0
 $\text{Re}\{T_{00}T_{11}^*\} \propto \text{Re}\{r_{10}^5\} \propto -\text{Im}\{r_{10}^6\}$
 $\text{Im}\{T_{11}T_{00}^*\} \propto \text{Im}\{r_{10}^7\} \propto \text{Re}\{r_{10}^8\}$

C- Spin Flip: $\gamma_T^* \rightarrow \rho_L^0$
 $\text{Re}\{T_{11}T_{01}^*\} \propto \text{Re}\{r_{10}^{04}\}$
 $\propto \text{Re}\{r_{10}^1\} \propto \text{Im}\{r_{10}^2\}$
 $\text{Re}\{T_{01}T_{00}^*\} \propto r_{00}^5$
 $|T_{01}|^2 \propto r_{00}^1$
 $\text{Im}\{T_{01}T_{11}^*\} \propto \text{Im}\{r_{10}^3\}$
 $\text{Im}\{T_{01}T_{00}^*\} \propto r_{00}^8$

D-Spin Flip: $\gamma_L^* \rightarrow \rho_T^0$
 $\text{Re}\{T_{10}T_{11}^*\} \propto r_{11}^5 \propto r_{1-1}^5 \propto \text{Im}\{r_{1-1}^6\}$
 $\text{Im}\{T_{10}T_{11}^*\} \propto \text{Im}\{r_{1-1}^7\} \propto r_{11}^8 \propto r_{1-1}^8$

E- Double Spin Flip: $\gamma_T^* \rightarrow \rho_{-T}^0$
 $\text{Re}\{T_{1-1}T_{11}^*\} \propto r_{1-1}^{04} \propto r_{11}^1$
 $\text{Im}\{T_{1-1}T_{11}^*\} \propto \text{Im}\{r_{1-1}^3\}$

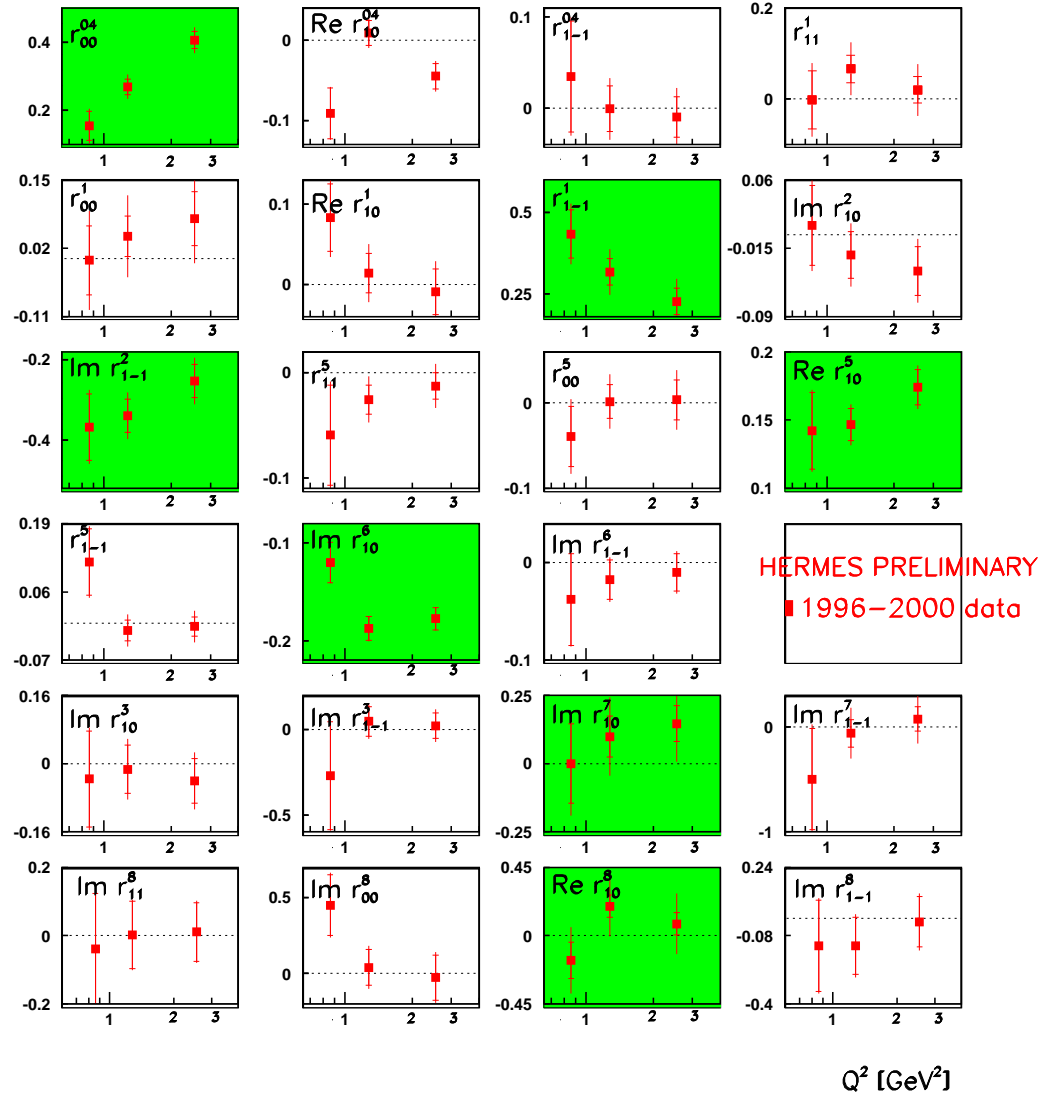


\Rightarrow **Hierarchy of ρ^0 amplitudes:**

$$|T_{00}| \sim |T_{11}| \gg |T_{01}| > |T_{10}| \gtrsim |T_{1-1}|, (0 \rightarrow L, 1 \rightarrow T)$$

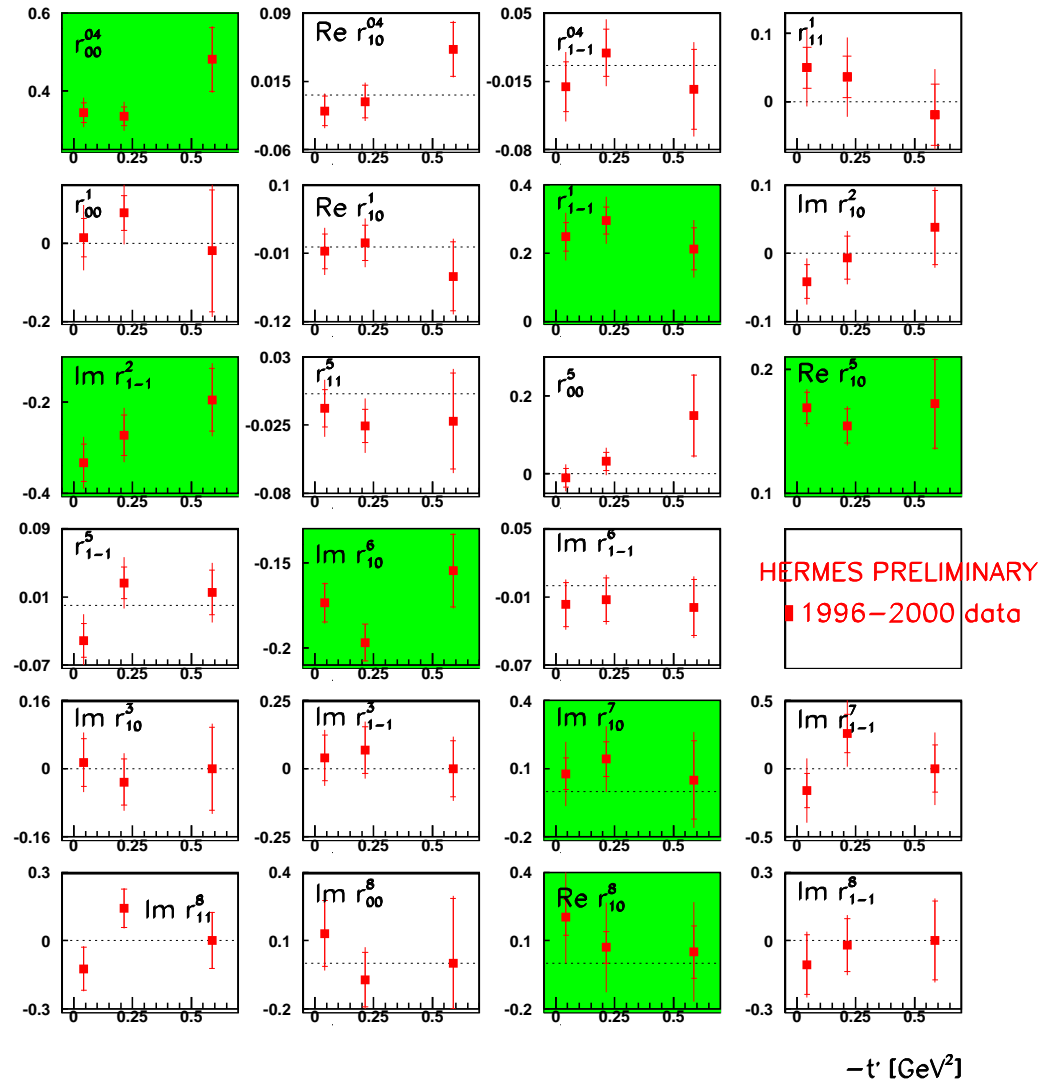
\Rightarrow ϕ meson SDMEs are consistent with SCHC, $|T_{00}| \sim |T_{11}|$

Dependencies of ϕ meson SDME's on Q^2

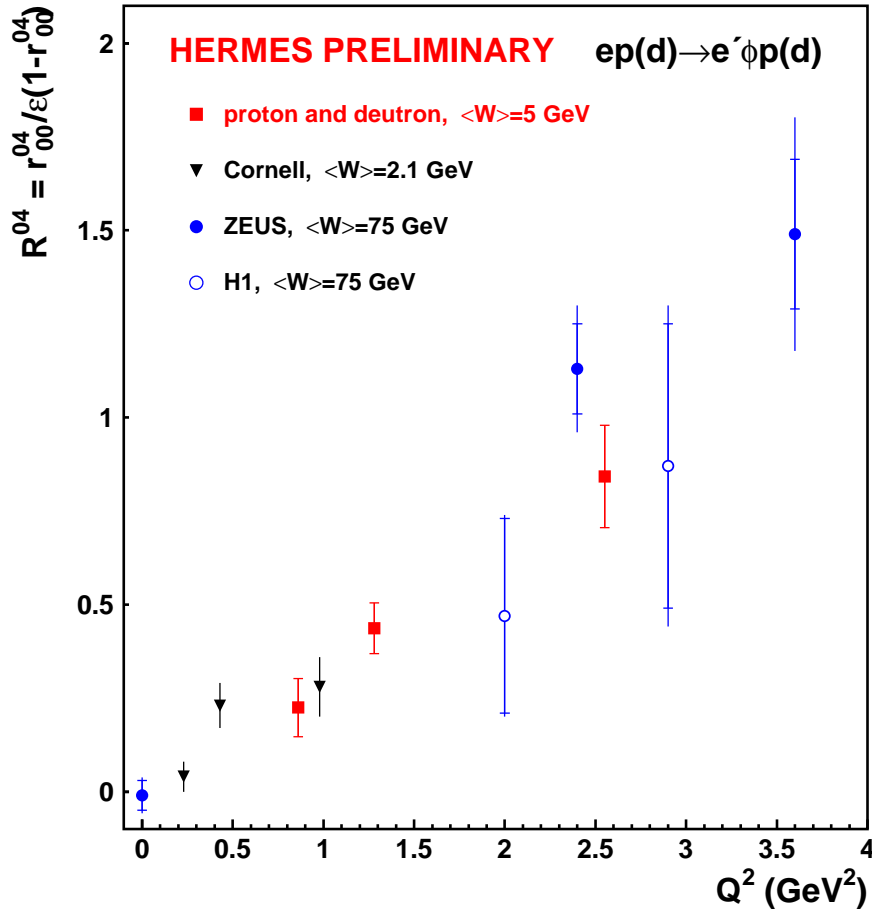


The SDME's dependencies on Q^2 for proton and deuteron data. The outer bars represent the total, the inner ones the statistical errors.

Dependencies of ϕ meson SDME's on t'



The SDME's dependencies on t' for proton and deuteron data. The outer bars represent the total, the inner ones the statistical errors.



Comparison of commonly measured:

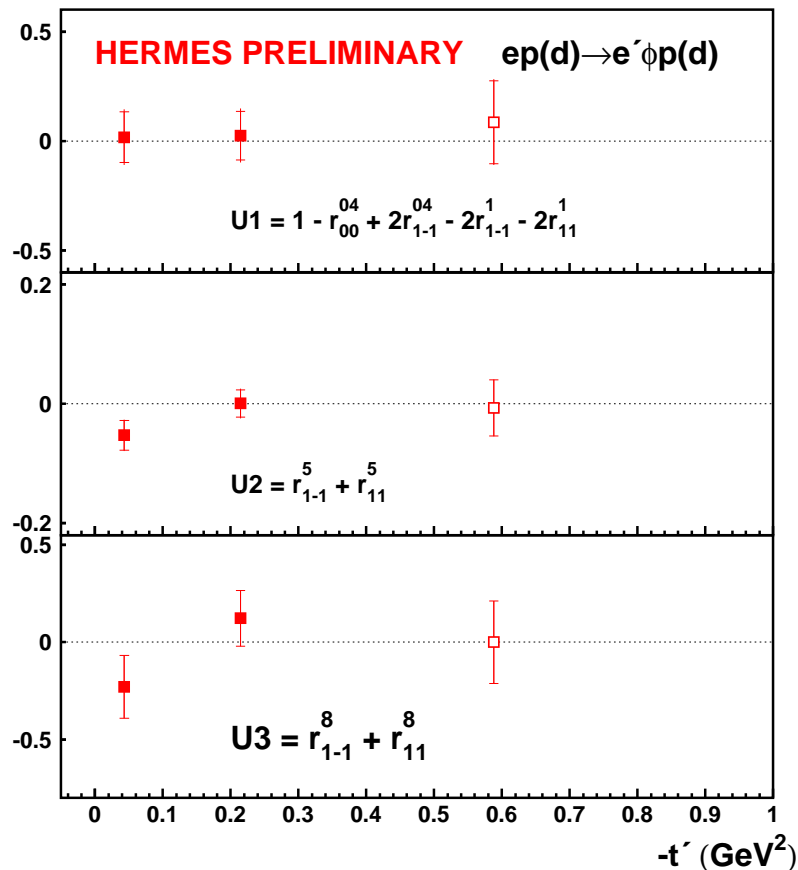
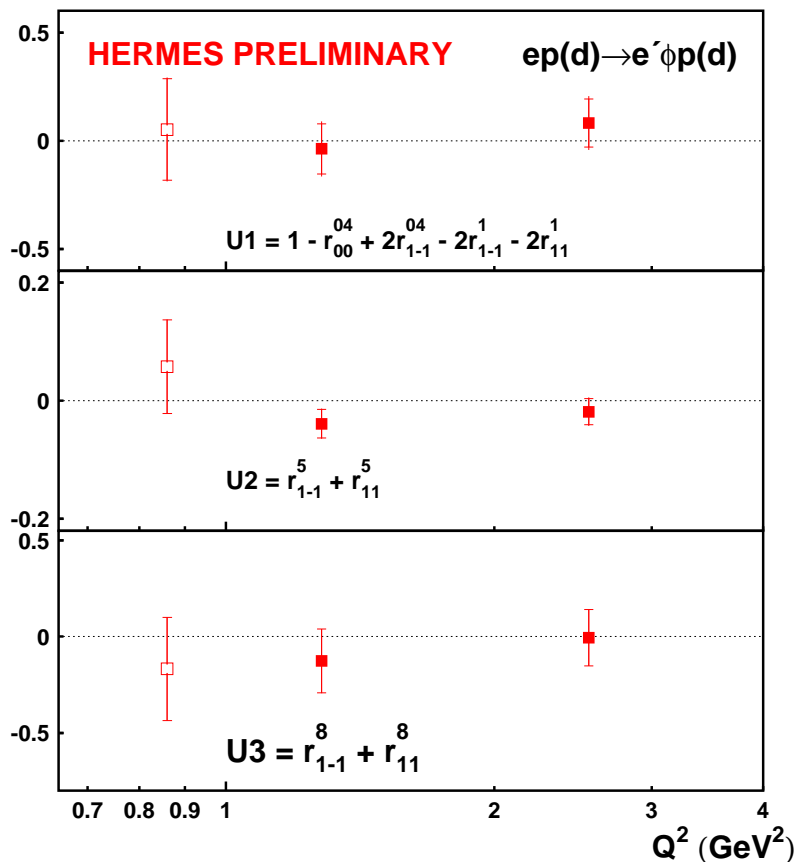
$$R^{04} = \frac{\sigma_L}{\sigma_T} = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}},$$

where:

$$r_{00}^{04} = \sum \{ \epsilon |T_{00}|^2 + \} / \sigma_{tot}$$

$$\sigma_{tot} = \epsilon \sigma_L + \sigma_T$$

$\Rightarrow R^{04}$ for ϕ meson at HERMES is in good agreement with world data.



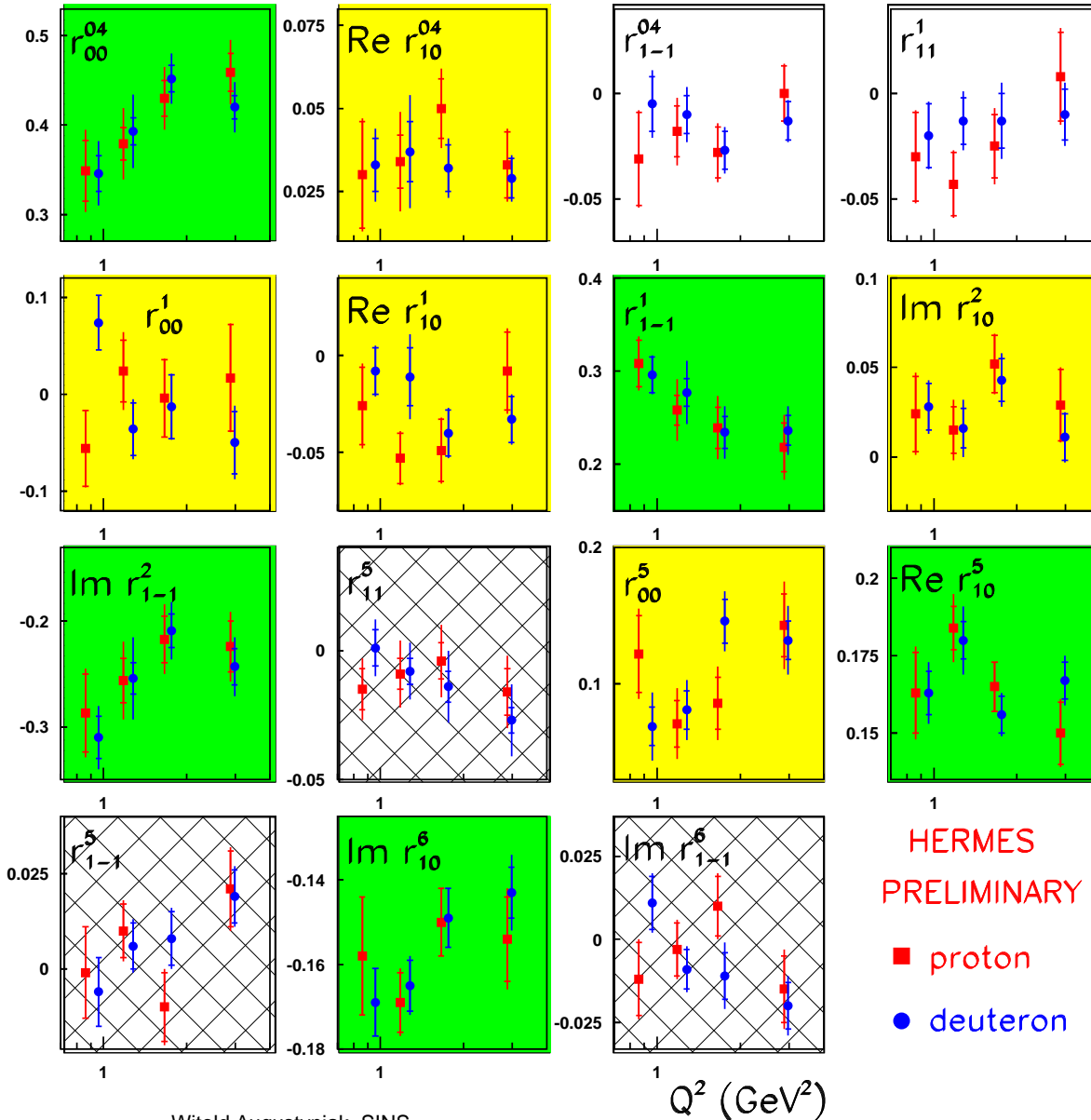
$$U1 = 0.02 \pm 0.07_{stat} \pm 0.16_{syst}$$

$$U2 = -0.03 \pm 0.01_{stat} \pm 0.03_{syst}$$

$$U3 = -0.05 \pm 0.11_{stat} \pm 0.07_{syst}$$

\Rightarrow no UPE for ϕ meson production, as expected

Dependencies of ρ^0 meson SDME's on Q^2



INDICATIONS:

green: SCHHC - $\gamma_L^* \rightarrow V_L, \gamma_T^* \rightarrow V_T$

yellow: Single Flip - $\gamma_T^* \rightarrow V_L$

grid: Single Flip - $\gamma_L^* \rightarrow V_T$

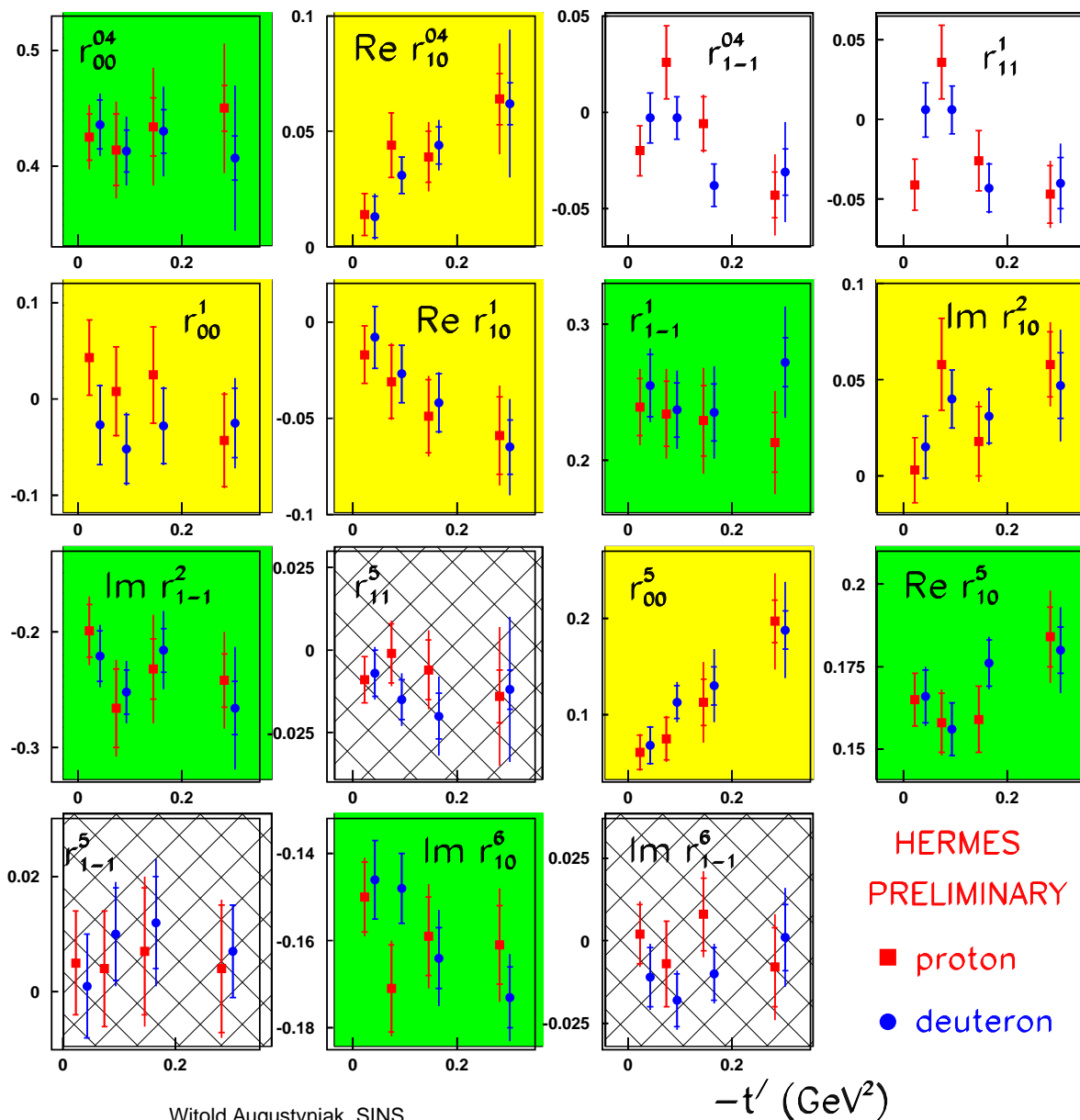
blank: Double Flip - $\gamma_T^* \rightarrow V_{-T}$

HERMES
PRELIMINARY

■ proton

● deuteron

Dependencies of ρ^0 meson SDME's on t'



INDICATIONS:

green: SCHHC - $\gamma_L^* \rightarrow V_L, \gamma_T^* \rightarrow V_T$

yellow: Single Flip - $\gamma_T^* \rightarrow V_L$

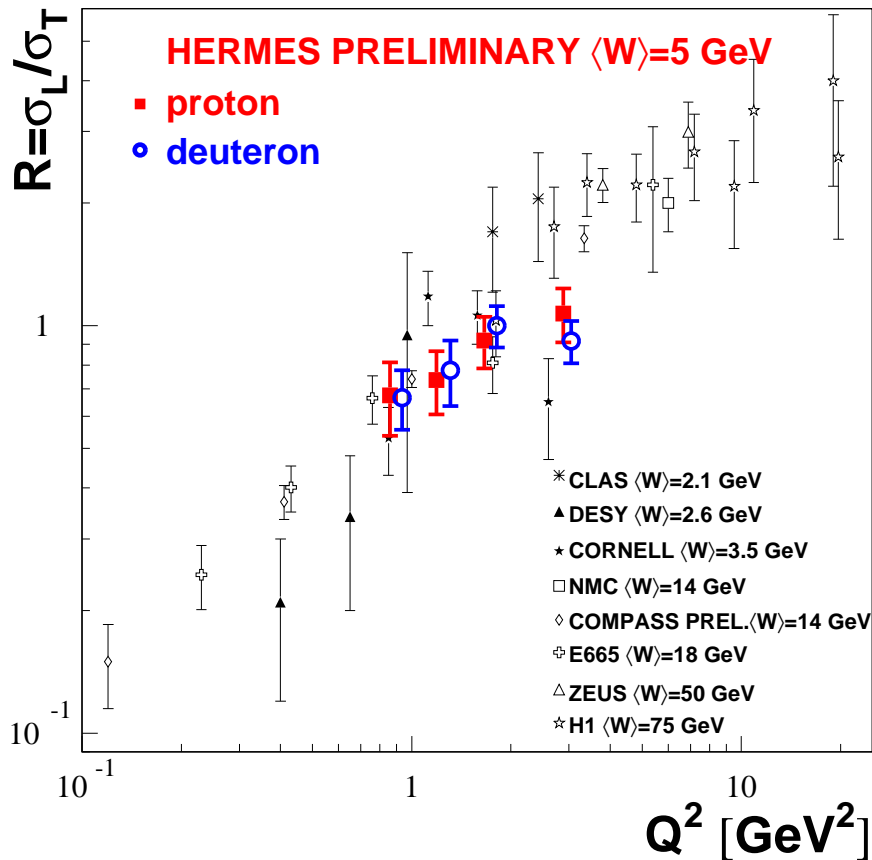
grid: Single Flip - $\gamma_L^* \rightarrow V_T$

blank: Double Flip - $\gamma_T^* \rightarrow V_{-T}$

HERMES
PRELIMINARY

■ proton

● deuteron



Comparison of commonly measured:

$$R^{04} = \frac{1}{\epsilon} \frac{r_{00}^{04}}{1 - r_{00}^{04}},$$

$$r_{00}^{04} = \sum \{ \epsilon |T_{00}|^2 + |T_{01}|^2 + |U_{01}|^2 \} / \sigma_{tot},$$

$$\sigma_{tot} = \epsilon \sigma_L + \sigma_T,$$

$$\sigma_T = \sum \{ |T_{11}|^2 + |T_{01}|^2 + |T_{1-1}|^2 + |U_{11}|^2 \},$$

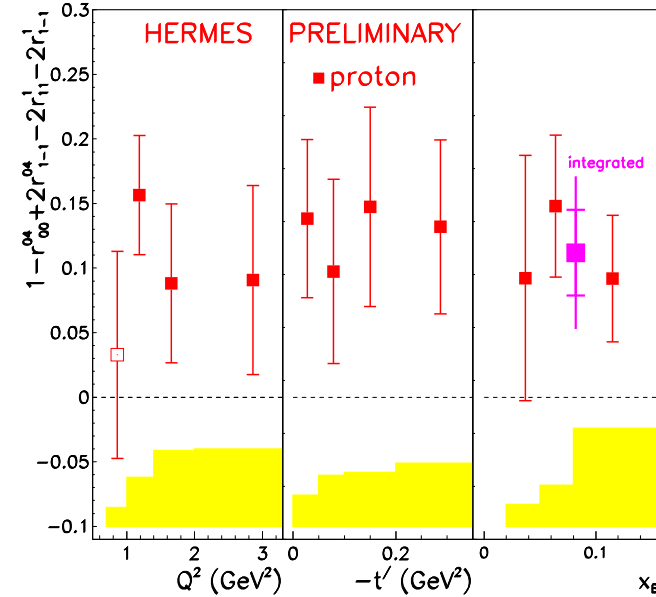
$$\sigma_L = \sum \{ |T_{00}|^2 + 2|T_{10}|^2 \}.$$

Due to the helicity-flip and unnatural parity amplitudes R^{04} depends on kinematic conditions, and is not identical to $R \equiv |T_{00}|^2 / |T_{11}|^2$ at SCHC and NPE dominance.

⇒ Second order contribution of spin-flip amplitudes to R^{04}

⇒ HERMES ρ^0 data on R^{04} are suggestive to R(W)-dependence

- Natural-parity exchange: interaction is mediated by a particle of 'natural' parity: vector or scalar meson:
 $J^P = 0^+, 1^-$ e.g. ρ^0, ω, a_2
- Unnatural parity exchange is mediated by pseudoscalar or axial meson:
 $J^P = 0^-, 1^+$, e.g. $\pi, a_1, b_1 \rightarrow$ only quark-exchange contribution
- No interference between NPE and UPE contributions on unpolarized target
- Extracted from SDMEs:
 $U2 + iU3 \propto (U_{11} + U_{1-1}) * U_{10}$
 $U2 = r_{11}^5 + r_{1-1}^5$
p: $U2 = -0.012 \pm 0.006_{stat} \pm 0.012_{syst}$
d: $U2 = -0.008 \pm 0.0046_{stat} \pm 0.010_{syst}$
 $U3 = r_{11}^5 + r_{1-1}^5$
p: $U3 = -0.020 \pm 0.050_{stat} \pm 0.007_{syst}$
d: $U3 = -0.021 \pm 0.038_{stat} \pm 0.011_{syst}$



- $U1 \propto \epsilon |U_{10}|^2 + 2|U_{11} + U_{1-1}|^2$
 $U1 = 1 - r_{00}^{04} + 2r_{1-1}^{04} - 2r_{11}^1 - 2r_{1-1}^1$
p: $U1 = 2|U_{11}|^2 = 0.132 \pm 0.026_{st} \pm 0.053_{syst}$
d: $U1 = 0.094 \pm 0.020_{st} \pm 0.044_{syst}$
p+d: $U1 = 0.109 \pm 0.037_{tot}$

\Rightarrow Indication on hierarchy of ρ^0 UPE amplitudes:
 $|U_{11}| \gg |U_{10}| \sim |U_{01}|$

-
- The transitions $\gamma_L^* \rightarrow \phi_L^0(\rho_L^0)$ and $\gamma_T^* \rightarrow \phi_T^0(\rho_T^0)$, are dominant for both ρ and ϕ .
SDME's: $(1 - r_{00}^{04})$, r_{1-1}^1 , $\text{Im } r_{1-1}^2$, depend on Q^2 i.e. $\sim 1/(Q^2 + m_v^2)$.
- Helicity-Flip transition $\gamma_T^* \rightarrow \rho_L^0$:
Observed only for ρ . Regular dependence of three elements belonging to this set: $\text{Re } \{r_{10}^{04}\}$, $\text{Re}\{r_{10}^1\}$ and r_{00}^5 are nicely seen. In the case of r_{00}^1 and $\text{Im } \{r_{10}^2\}$ dependencies are less pronounced.
- Helicity-Flip transition $\gamma_L^* \rightarrow \rho_T^0$:
Very weak oscillating near zero.
- Double Helicity-Flip transition $\gamma_T^* \rightarrow \rho_{-T}^0$: Very weak indication of the dependence on Q^2 .
- Dependence on target (H, D): not observed.
- Only natural-Parity Exchange: only in the case of ϕ .
- Unnatural-Parity Exchange: only in the case of ρ .
- Comparisons with other measurements: Good agreement. For measurements with higher W small differences are seen due .

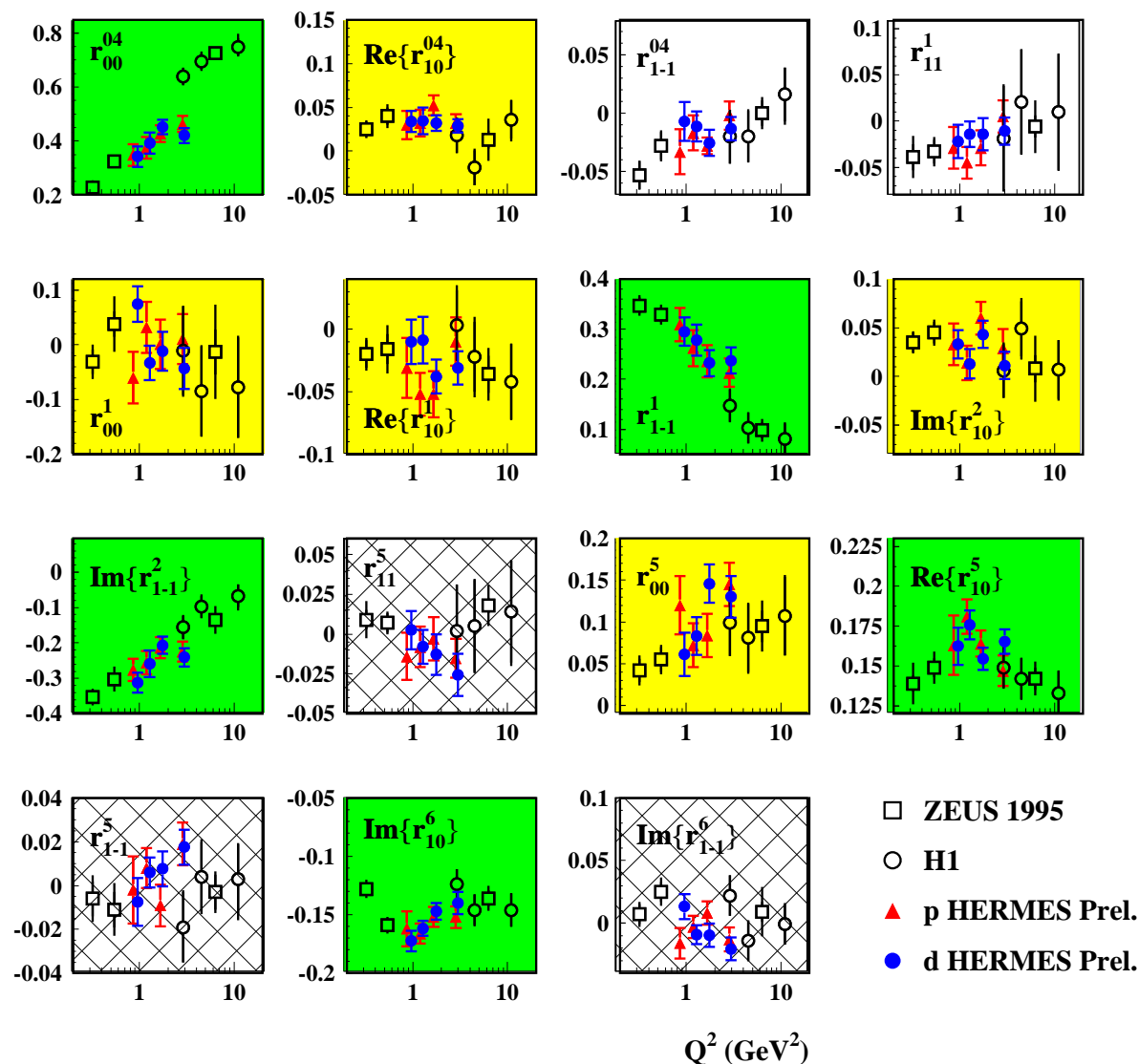
-
- Determine SDME's for samples with polarized target.



END ADDITIONAL SLIDES



Dependencies of ρ^0 meson SDME's on Q^2



The SDME's as function of Q^2 . For HERMES data received at $W=5$.GeV for **proton** and **deuteron** targets as well as H1 and ZEUS data at $W=75$ GeV.

⇒ Several SDMEs ($r_{00}^{04}, r_{1-1}^1, Im(r_{1-1}^2 \dots)$) indicate possible W -dependence, in addition to Q^2 -dependence